

Q1 (9)

FIND  $U_E = 8 U_B$

$$U_E = \frac{1}{2} \epsilon_0 \int E^2 dV'$$

$$= \frac{1}{2} \epsilon_0 \left( \frac{e}{4\pi\epsilon_0} \right)^2 \int_R^\infty \frac{4\pi r^2 dr}{r^4}$$

$$= \frac{1}{2} \cdot \frac{1}{4\pi} \frac{e^2}{\epsilon_0 R} \left( -\frac{1}{r} \right) \Big|_R^\infty$$

$$\boxed{U_E = \frac{1}{8\pi} \frac{e^2}{\epsilon_0 R}}$$

$$U_B = \frac{1}{2\mu_0} \int_{\text{ALL SPACE}} B^2 dV$$

$$\underline{B}(r < R) = \frac{2}{3} \mu_0 \pi \omega R^2 \hat{z}$$

EX 5.11

$$\frac{1}{r} = \frac{e}{4\pi R^2}$$

$$\underline{B}(r < R) = \frac{2}{3} \mu_0 \frac{e}{4\pi R^2} \omega R^2 \hat{z}$$

$$B^2(r < R) = \frac{\mu_0^2 e^2 \omega^2}{36\pi^2 R^2}$$

$$U_B^{IN} = \frac{1}{2\mu_0} \frac{\mu_0^2 e^2 \omega^2}{36\pi^2 R^2} \frac{4}{3} \pi R^3$$

$$\boxed{U_B^{IN} = \frac{1}{54\pi} \mu_0 e^2 \omega^2 R}$$

Q1 (e) FOR  $r \geq R$ ,  $\vec{B}$  IS LIKE DIPOLE

$$\vec{B} = \frac{\mu_0}{4\pi r^3} \left[ 3 (\vec{m} \cdot \hat{r}) \hat{r} - \vec{m} \right]$$

NEED TO FIND MAGNETIC MOMENT  $\vec{m}$ .

FROM EXAMPLE 6.1

$$\vec{m} = \vec{M} \cdot \frac{4}{3} \pi R^3$$

$$\vec{K} = \sigma \omega R \sin \theta \hat{\varphi}$$

$$\vec{K} = \vec{M} \times \hat{z} = M \sin \theta \hat{\varphi}$$

$$\Rightarrow M = \sigma \omega R$$

$$= \frac{e}{4\pi R^2} \omega R$$

$$\vec{M} = \frac{e \omega}{4\pi R} \hat{z}$$

$$\vec{m} = \frac{e \omega R^2}{3} \hat{z}$$

Q1(b) cont.

$$\begin{aligned}
 B^2 &= \left( \frac{\mu_0}{4\pi r^3} \right)^2 \left[ 9(m_2 \cdot \hat{r})^2 + m^2 - 2 \cdot 3(m_2 \cdot \hat{r})r^{-1}m_2 \right] \\
 &= \left( \frac{\mu_0}{4\pi r^3} \right)^2 \left[ 9m^2 \cos^2 \theta + m^2 - 6m^2 \cos^2 \theta \right] \\
 &= \left( \frac{\mu_0}{4\pi r^3} \right)^2 \left[ 3m^2 \cos^2 \theta + m^2 \right] \\
 &= \left( \frac{\mu_0}{4\pi r^3} \right)^2 m^2 \left[ 3 \cos^2 \theta + 1 \right] \\
 B^2 &= \left( \frac{\mu_0}{4\pi r^3} \right)^2 \cdot \frac{e^2 \omega^2 R^4}{9} \left[ 3 \cos^2 \theta + 1 \right]
 \end{aligned}$$

$$= C_0 \left[ 3 \cos^2 \theta + 1 \right] / r^6 \quad \omega / C_0 = \left( \frac{\mu_0}{4\pi} \right)^2 \frac{e^2 \omega^2 R^4}{9}$$

$$U_B^{\text{OUT}} = \frac{1}{2\mu_0} C_0 \int_R^{\infty} \frac{1}{r^6} \left[ 3 \cos^2 \theta + 1 \right] r^2 dr \sin \theta d\theta d\phi$$

$$= \frac{2\pi}{2\mu_0} C_0 \int_{\theta=0}^{\pi} \int_{r=R}^{\infty} \frac{3 \cos^2 \theta \sin \theta d\theta}{r^6} r^2 dr$$

$$+ \frac{4\pi}{2\mu_0} C_0 \int_R^{\infty} \frac{1}{r^4} dr$$

$$= \frac{2\pi}{2\mu_0} C_0 \left( -\frac{r^{-3}}{3} \right) \Big|_R^{\infty} \int_{\theta=0}^{\pi} 3 \cos^2 \theta \sin \theta d\theta$$

$$+ \frac{4\pi}{2\mu_0} C_0 \left( \frac{-r^{-3}}{3} \right) \Big|_R^{\infty}$$

$$U_B^{\text{out}} = \frac{\pi}{\mu_0} C_0 \left( \frac{1}{3R^3} \right) \left( -\cos^3 \theta \Big|_0^{\pi} \right)$$

$$+ \frac{2\pi}{\mu_0} C_0 \left( \frac{1}{3R^3} \right)$$

$$U_B^{\text{out}} = \frac{4\pi}{3\mu_0} \frac{C_0}{R^3} = \frac{4\pi}{3\mu_0 R^3} \left( \frac{\mu_0}{4\pi} \right)^2 \frac{e^2 \omega^2 R^4}{9}$$

$$U_B^{\text{out}} = \frac{\mu_0 e^2 \omega^2 R}{27 \cdot 4\pi}$$

$$U_B^{\text{ttl}} = U_B^{\text{in}} + U_B^{\text{out}}$$

$$U_B^{\text{ttl}} = \frac{1}{36\pi} \mu_0 e^2 \omega^2 R$$

$$U_{\text{ttl}} = U_E + U_B$$

$$U_{\text{ttl}} = \frac{1}{8\pi} \frac{e^2}{\epsilon_0 R} + \frac{1}{36\pi} \mu_0 e^2 \omega^2 R$$

Q1 (b)  $\vec{L} = \epsilon_0 \int \vec{L} \times (\vec{E} \times \vec{B})$

$$\vec{L} = \int_{\text{ALL SPACE}} \vec{L} dV$$

$$\vec{E} (r < R) = 0$$

$$\vec{E} (r \geq R) = \frac{e \hat{r}}{4\pi\epsilon_0 r^2}$$

$$\vec{B} (r \geq R) = \frac{\mu_0 m}{4\pi r^3} [2 \cos\theta \hat{r} + \sin\theta \hat{\theta}]$$

$\vec{B} (r < R)$  IRRELEVANT SINCE  $\vec{E} = 0$  THERE

ALSO,  $m = \frac{e\omega R^2}{3}$  FROM COMPARING  
EXAMPLE 5.11 EXPRESSION FOR  $\vec{B}$   
AND SECTION 5.4.3 FOR  $\vec{B}$  DIPOLE.

Now,

$$\vec{E} \times \vec{B} = \frac{e}{4\pi\epsilon_0 r^2} \frac{\mu_0 m}{4\pi r^3} \hat{r} \times [2 \cos\theta \hat{r} + \sin\theta \hat{\theta}]$$

$$\omega/m = \frac{e\omega R^2}{3}$$

$$= \frac{\mu_0 e^2 \omega R^2 \sin\theta}{3 \epsilon_0 (4\pi)^2 r^5} \hat{r} \times \hat{\theta}$$

Now,

$$\hat{r} \times \hat{\theta} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \sin\theta \cos\varphi & \sin\theta \sin\varphi & \cos\theta \\ \cos\theta \cos\varphi & \cos\theta \sin\varphi & -\sin\theta \end{vmatrix}$$

$$= \hat{x} \left[ -\sin^2\theta \sin\varphi - \cos^2\theta \sin\varphi \right]$$

$$- \hat{y} \left[ -\sin^2\theta \cos\varphi - \cos^2\theta \cos\varphi \right]$$

$$+ \hat{z} \left[ \sin\theta \cos\theta \sin\varphi \cos\varphi - \sin\theta \cos\theta \sin\varphi \cos\varphi \right]$$

$$= -\hat{x} \sin\varphi + \hat{y} \cos\varphi$$

$$\boxed{\hat{r} \times \hat{\theta} = \hat{\varphi}}$$

$$\text{So, } \int_{\nu} = \epsilon_0 \int_{\nu} \hat{r} \times (\hat{E} \times \hat{B}) = \frac{\mu_0 e^2 \omega^2 R^2 \sin\theta}{3 (4\pi)^2 r^4} \hat{r} \times \hat{\varphi}$$

Since we need to integrate ~~this~~  $\int_{\nu}$ ,  
we need to get its DIRECTION.  
We need  $\hat{r} \times \hat{\varphi}$

AWAY WE GO:

$$\vec{r} \times \vec{\varphi} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ r \sin\theta \cos\varphi & r \sin\theta \sin\varphi & r \cos\theta \\ -r \sin\varphi & r \cos\varphi & 0 \end{vmatrix}$$

$$= -\hat{x} r \cos\theta \cos\varphi - \hat{y} r \sin\theta \cos\theta + \hat{z} r \sin\theta$$

$$\vec{r} \times \vec{\varphi} = -\hat{\theta}$$

BECAUSE OF  $\sin\theta$  TERM IN MAGNITUDE OF  $\hat{\theta}$ , INTEGRATING OVER  $\theta$  YIELDS  $\int_0^{\pi} \cos\theta \sin\theta d\theta = 0$  FOR  $\hat{x}$  INTEGRATION  $\theta=0$

LIKE WISE, FOR  $\hat{y}$  INTEGRATION OVER  $\theta$

$$\int \sin^2\theta \cos\theta d\theta = 0$$

ONLY  $\hat{z}$  TERM WILL SURVIVE.

FOR  $\hat{z}$  INTEGRATION

$$L = \hat{z}$$

$$= - \frac{\mu_0 e^2 \omega R^2}{3 (4\pi)^2} \int_R^R \frac{r^2 dr}{r^4} \int_0^{2\pi} d\phi \int_0^\pi \sin^2 \theta d\theta \hat{z}$$

$$= - \frac{\mu_0 e^2 \omega R^2}{3 (4\pi)^2} \left( -\frac{1}{R} \Big|_R^R \right) \cdot 2\pi \cdot \frac{4}{3} \hat{z}$$

$$= - \frac{\mu_0 e^2 \omega R^2}{3 (4\pi)^2} \frac{2\pi}{R} \cdot \frac{4}{3}$$

$$L = - \frac{\mu_0 e^2 \omega R}{10\pi} \hat{z}$$



Q2 (a)  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

$$\vec{E} = 0 \Rightarrow \frac{\partial \vec{B}}{\partial t} = 0$$

(b)  $\vec{E} = -\frac{d\Phi}{dt}$

$$\oint \vec{E} \cdot d\vec{l} = -\frac{d\Phi}{dt}$$

SINCE  $\vec{E} = 0$ ,  $\frac{d\Phi}{dt} = 0 \Rightarrow \Phi = \text{const.}$

(c)  $\nabla \times \vec{B} = \mu_0 \vec{J} + \epsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t}$

$$\vec{B} = 0 \Rightarrow \nabla \times \vec{B} = 0$$

$$\vec{E} = 0 \Rightarrow \frac{\partial \vec{E}}{\partial t} = 0$$

SUBSTITUTING INTO EQN  $\Rightarrow \vec{J} = 0$

(d) WE KNOW  $\vec{B} = 0$  &  ~~$\vec{B} = 0$~~   $\vec{J} = 0$ , CURRENT FLOWS ONLY AT SURFACE, SURFACE CURRENT MUST FLOW IN SUCH A WAY THAT THE  $B$ -FIELD  $\vec{B}_{\text{ind}}$  IT PRODUCES EXACTLY CANCELS THE STATIC FIELD  $\vec{B}_0$

$$\vec{B}_{\text{ind}} + \vec{B}_0 = 0$$

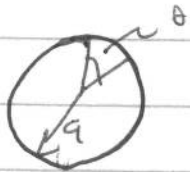
Q2

From EXAMPLE 5.11 we know that a spherical shell with constant charge density  $\sigma$  spinning with angular speed  $\omega$  produces a steady B-field inside shell.

$$\text{From EX 5.11} \quad \vec{B} = \frac{\mu_0}{3} \sigma \omega a \hat{z}$$

$a = \text{SHELL RADIUS}$

RECAST IN TERMS OF  $\vec{K}$



FIRST NOTE THAT WE NEED

$$\sigma \omega a = -\frac{3}{2} \frac{B_0}{\mu_0}$$

$$\text{NOW, } \vec{K} = \sigma \vec{v} = \sigma \omega a \sin \theta \hat{\phi}$$

$$\Rightarrow \vec{K} = -\frac{3}{2} \frac{B_0}{\mu_0} \sin \theta \hat{\phi}$$