

#1. EXPRESS / EXPAND TRITIUM W. F.
IN TERMS OF EIGENSTATES OF
 He^3 ($Z=2$):

$$\psi(\text{TRITIUM}) = \sum_{n=1}^{\infty} \psi_n(\text{He}^3)$$

PROBABILITY OF REMAINING IN GROUND
STATE AFTER TRANSITION IS

$$P_{\text{GND}} = |\langle \psi(\text{TRITIUM GND STATE}) | \psi(\text{He}^3 \text{ G.S.}) \rangle|^2$$

$$\psi_{100}(\text{TRITIUM}) = \frac{2}{\sqrt{\pi}} \left(\frac{1}{a}\right)^{3/2} e^{-r/a}$$

$$\psi_{100}(\text{TRITIUM}) = \frac{1}{\sqrt{\pi a^3}} e^{-r/a}$$

$$r/a = \frac{4\pi\epsilon_0 \hbar^2}{m e^2}$$

$$\psi_{100}(\text{He}^3) = \sqrt{\frac{8}{\pi a^3}} e^{-2r/a} \quad (2 \text{ protons})$$

NOTE THAT FOR $Z \neq 1$, $a = \frac{4\pi\epsilon_0 \hbar^2}{m Z e^2}$

$$\text{SO, } P = \left[\frac{1}{\sqrt{\pi a^3}} \int_0^{\infty} e^{-r/a} \cdot \sqrt{\frac{8}{\pi a^3}} e^{-2r/a} r^2 dr d\Omega \right]^2$$

#7.

$$P = \left[\frac{\sqrt{8}}{\pi a^3} \cdot 4\pi \int_0^\infty e^{-3r/a} r^2 dr \right]^2$$

$$= \left[\frac{4\sqrt{8}}{a^3} \cdot \int_0^\infty e^{-\eta} \eta^2 d\eta \cdot \frac{a^3}{27} \right]^2$$

$$\text{w/ } \eta \equiv 3r/a$$

$$\text{Now, } \int_0^\infty e^{-\eta} \eta^2 d\eta = 2 \quad \text{via WOLFRAM INTEGRATOR}$$

Hence,

$$P = \left[\frac{4\sqrt{8} \cdot 2}{27} \right]^2$$

$$P = 0.70$$

$$\#2. \int_0^a u^2 dr = 1 \quad \text{w/ } u = A \sin kr + B \cos kr$$

$$= A \sin kr$$

$$\int_0^a A^2 \sin^2 kr dr = 1 \quad \text{since } u(a) = 0$$

$$\text{w/ } k = \sqrt{2mE}/\hbar$$

$$\int_0^a A^2 \frac{(1 - \cos kr)}{2} dr = 1$$

$$A^2 \left[\frac{r}{2} - \frac{\sin kr}{kr} \right] \bigg|_0^a = 1$$

$$A^2 \frac{a}{2} = 1 \quad \text{since } \sin ka = 0$$

$$A = \sqrt{2/a}$$

#3. For H-Atom

$$E_n = - \left[\frac{m_e}{2\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0} \right)^2 \right] \cdot \frac{1}{n^2}$$

$$= -13.6 \text{ eV} / n^2 \quad m/m_e = e^- \text{ mass}$$

MORE GENERALLY

$$E_n = - \frac{\mu}{m_e} \left[\frac{m_e}{2\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0} \right)^2 \right] \cdot \frac{1}{n^2}$$

$$E_n = - \frac{\mu}{m_e} \frac{13.6 \text{ eV}}{n^2} \quad \mu = \text{REDUCED MASS}$$

For ionization, $n=1 \rightarrow \infty$

$$\text{AND } \mu = \frac{m_\mu m_p}{m_\mu + m_p}$$

$$= \frac{206 m_e}{\frac{m_\mu}{m_p} + 1}$$

$$\Rightarrow \mu \approx 182 m_e$$

So,

$$E_{\text{ionization}} \approx 182 \times 13.6 \text{ eV}$$

$$\approx 2.5 \text{ keV}$$

$\mu^- p$ "ATOM" VERY TIGHTLY BOUND?

$$\#4. E_{e^+e^-} = - \left[\frac{\mu}{2\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0} \right)^2 \right] \cdot \frac{1}{n^2}$$

$$= - \frac{m_e/2}{2\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0} \right)^2 \cdot \frac{1}{n^2}$$

$$E_{e^+e^-} = -\frac{1}{2} 13.6 \text{ eV} / n^2$$

BALMER SERIES $n_f = 2$
 LOWEST E TRANSITIONS: $3 \rightarrow 2, 4 \rightarrow 2, 5 \rightarrow 2$.

$$E_f(3 \rightarrow 2) = \frac{13.6 \text{ eV}}{2} \left(\frac{1}{4} - \frac{1}{9} \right)$$

$$E_f(3 \rightarrow 2) = 0.94 \text{ eV}$$

$$\lambda (\text{nm}) = \frac{1240 \text{ eV} \cdot \text{nm}}{E_f(\text{eV})} \quad \text{FROM LECTURE}$$

$$\lambda(3 \rightarrow 2) = 1313 \text{ nm} \quad \underline{\text{NOT VISIBLE!}}$$

NOTE:

$$\sim 390 \text{ nm} \leq \text{VISIBLE LIGHT} \leq \sim 700 \text{ nm}$$

$$E_f(4 \rightarrow 2) = \frac{13.6}{2} \left(\frac{1}{4} - \frac{1}{16} \right) = 1.28 \text{ eV}$$

$$E_f \lambda = \frac{1240 \text{ eV} \cdot \text{nm}}{1.28 \text{ eV}}$$

$$\lambda = 970 \text{ nm} \quad \underline{\text{NOT VISIBLE!}}$$

#4. $E_g(5 \rightarrow 2) = \frac{13.6 \text{ eV}}{Z^2} \left(\frac{1}{4} - \frac{1}{25} \right)$

$$E_g(5 \rightarrow 2) = 1.43 \text{ eV}$$

$$\lambda(5 \rightarrow 2) = 868 \text{ nm} \quad \underline{\text{NOT}} \text{ VISIBLE!}$$