

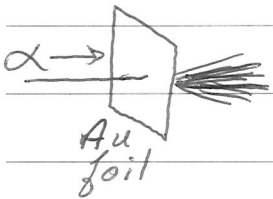
# Particle Physics

Atoms have substructure

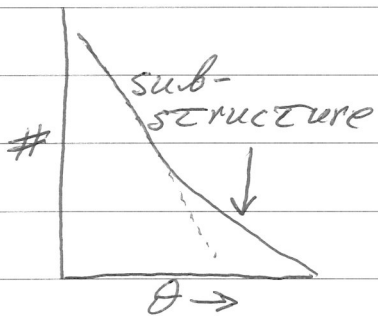
- electrons in shells
  - positively charged, massive nucleus
- } Rutherford Experiment

Rutherford gold foil experiment \*

- if charge evenly spread thru atom
- no strong concentration of charge
- slight deflections of  $\alpha$  particles



- but saw some huge deflections
- <sup>there</sup> must exist a strong repulsive force



- cannot be due to  $e^-$ 's: mass much too small
- implies a nucleus
- small, massive, positive charge
- surrounded by electrons in "orbit"

\* Phil. Mag. xxi, 669 (1911).

# The Nucleus:

②

Total charge

- integral # of charges with same magnitude as electron

→ opposite sign ⇒ protons

But additional mass ⇒ neutrons observed } the Nucleus ("nucleons")

How are like charges kept together in nucleus?

Mass?

Gravity  $10^{-37}$  too weak

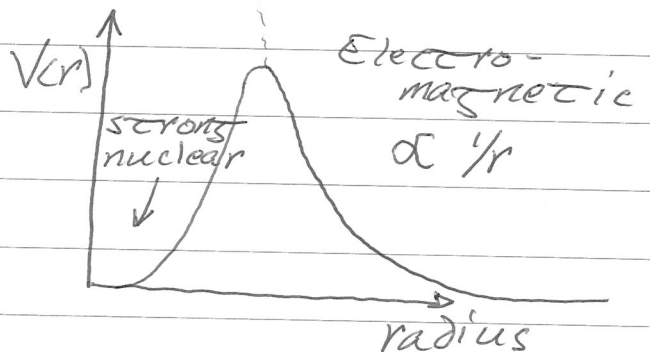


$\alpha$  particle (He nucleus)

New force: "strong nuclear force"

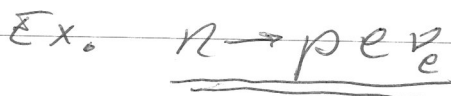
→ affects protons + neutrons the same

→ limited range



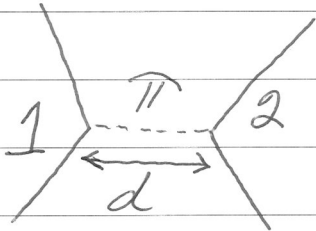
Weak nuclear force

- radioactive  $\beta$  decay



## Yukawa Model

Particles interact via exchange of bosons  
 → momentum transfer involved



Nuclear force has a range  
 of  $\sim 2 \text{ fm}$

Suggests exchange of a  
 massive "meson" ( $\pi$ )

Consider case where 1 is at rest

→ emission of a massive meson  
 requires violation of energy  
 conservation  $\Rightarrow \Delta E \Delta t \approx \hbar/2$

→ assuming velocity of  $\pi \sim c$   
 $\Delta t \sim d/c$

$$\therefore \boxed{m_{\pi} = 130 \text{ MeV}} \quad (1935)$$

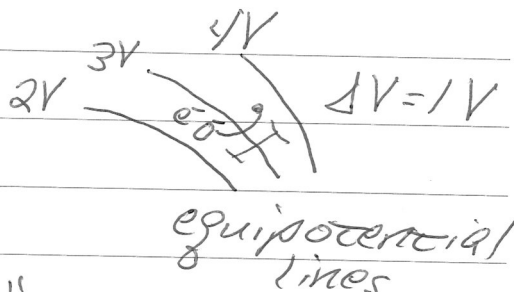
→ first evidence of  $\pi$  in 1947!

# Units

## Energy

- amount of work to move  $e^-$  thru  
1 Volt of potential

change in Kinetic Energy  
by electron



"1 electron volt"

$$1 \text{ eV} = 1.6 \times 10^{-19} \text{ J}$$

Momentum:  $\text{eV}/c$

Mass:  $\text{eV}/c^2$

In particle physics, often take  $c = \hbar = 1$   
 $\therefore$  mass, energy + momentum:  $\text{eV}$   
length + time:  $\text{eV}^{-1}$

It's also useful to remember

$$\hbar c = 0.2 \text{ GeV fm} \quad (\text{fm} = 10^{-15} \text{ m})$$

"fermi"

Cross section:

$$\text{"1 barn"} = 10^{-28} \text{ m}^2$$

# Energy Scales + Sizes!

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## atomic

- set by electromagnetic interaction

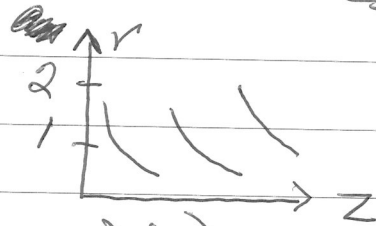
- energy scales:

$$E_0 = -m_e c^2 \alpha^2 / 2 = -13.6 \text{ eV}$$

$\therefore$  in eV range

- size:

$\rightarrow$  decreases as  
each shell fills ( $Z \uparrow$ )



$0.5 - 2 a_0$  (Bohr radius)

$$\therefore \mathcal{O}(10^{-10} \text{ m}) \sim a_0$$

- classical cross section

$$\sigma_{\text{atom}} \sim \pi a_0^2 \sim \underline{3 \times 10^8 \text{ barns}}$$

## nuclear

- residual strong interaction

- energy scale:

$$\sim m_N c^2 \frac{\alpha_s^2}{2} \sim 100 \text{ MeV (many MeV)}$$

- size:

Compton wavelength of proton

$$\lambda_p = \hbar / m_p c \sim 0.2 \text{ fm}$$

Radius of close-packed array of nucleons:  $a_N \sim \lambda_p A^{1/3}$  (fermi)

- classical cross section:  $\sigma_N \sim \pi a_N^2 \sim \underline{31 \text{ mb}}$

# Particle Milestones

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antiparticles:

$e^+$ ,  $\bar{p}$ ... every particle has one

First indication

Dirac Eq.

neutrino: ( $\nu$ )

before

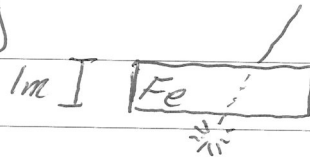
$$\vec{p}_\pi = 0$$

after

$$\vec{p}_e \quad \Sigma \vec{p}_i = ?$$

\* momentum conservation \*

muon: ( $\mu^-$ )



\* cosmic rays

pion (pi meson): ( $\pi^+$ ,  $\pi^0$ )

strong nuclear mediator

higher mass hadrons

- mesons: ( $K, \rho, \eta, \dots$ )

- baryons: ( $\Lambda, \Sigma, \Xi, \dots$ )

\* cosmic rays & accelerators

quarks: (up, down, strange, ~~bottom~~)

\* patterns in observed hadrons

gluons/jets:

QCD model

W & Z bosons:

electroweak model

charm, top quarks,  $\nu_\tau$

" "

# Hadrons, QCD and Jets

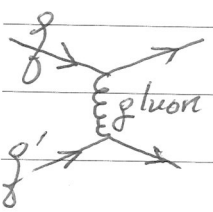
→ quarks <sup>are</sup> fundamental particles bound inside hadrons by strong interaction

mesons:  $q\bar{q}$  pairs (bosons / integer spin)

baryons:  $qqq$  "triplets" (fermions / half-integer spin)

## Strong Interactions

- increasing strength with distance (compare to  $E + m \ll V(r) \propto 1/r$ )



"asymptotic freedom":

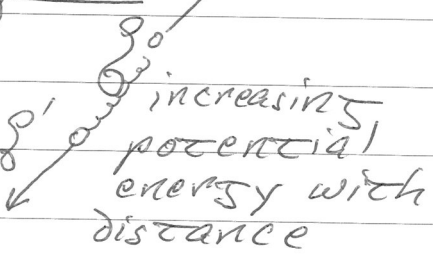
@ small distance (large energy), strength is small:

"confinement":

coupling strength large ( $\alpha_s \sim 1$ )

∴ quarks remain in hadrons

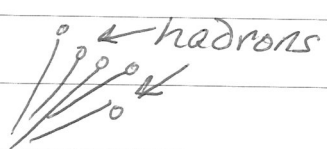
Colored Quarks



Difficulty in calculating "soft" (low energy) strong processes:

- nuclear reactions
- inelastic hadron processes
- jet production

Colorless Jets



Energy goes to particle production

Each will come up in portions of this course.

# Fundamental Interactions + Particles

Four "interactions" (forces renamed)

|                  | Distance             | Strength   | Mediators                 |
|------------------|----------------------|------------|---------------------------|
| strong           | $10^{-14} \text{ m}$ | $\sim 1$   | 8 gluons                  |
| electro-magnetic | $\infty$             | $10^{-2}$  | 1 photon                  |
| weak             | $10^{-18} \text{ m}$ | $10^{-5}$  | 3 bosons ( $W^\pm, Z^0$ ) |
| gravity          | $\infty$             | $10^{-38}$ | (graviton?)               |

## Three Generations of Particles (Fermions)

(mass in MeV)

|                  | 1st            | 2nd           | 3rd             | Q              | L  | B              |
|------------------|----------------|---------------|-----------------|----------------|----|----------------|
| charged leptons  | $e^- (0.5)$    | $\mu^- (105)$ | $\tau^- (1700)$ | -1             | +1 | 0              |
| neutrinos        | $\nu_e (0)$    | $\nu_\mu (0)$ | $\nu_\tau (0)$  | 0              | +1 | 0              |
| up-type quarks   | $u (\sim 100)$ | $c (1400)$    | $t (173k)$      | $+\frac{2}{3}$ | 0  | $+\frac{1}{3}$ |
| down-type quarks | $d (\sim 100)$ | $s (350)$     | $b (4500)$      | $-\frac{1}{3}$ | 0  | $+\frac{1}{3}$ |



# Relativistic Kinematics

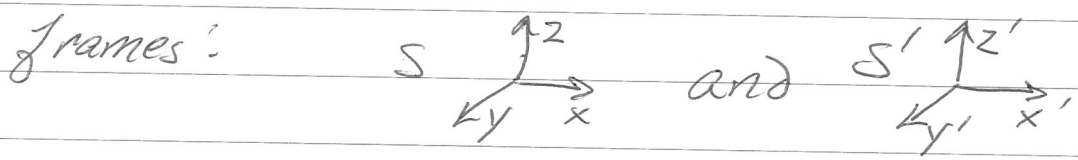
Consider a relativistic particle

$$\vec{v}, \vec{\beta} = \vec{v}/c$$

momentum:  $\vec{p} = mc\gamma\vec{\beta}$

energy:  $E = mc^2\gamma$  ( $\gamma = 1/\sqrt{1-\beta^2}$ )

## Lorentz Transformation



$S$  has velocity  $\vec{\beta}_0$  compared to  $S'$  frame ( $-\vec{\beta}_0$  in  $-z$  direction)

$$p_x = p'_x; \quad p_y = p'_y$$

$$p_z = \gamma (p'_z + \beta_0 \frac{E'}{c})$$

$$\frac{E}{c} = \gamma (\frac{E'}{c} + \beta_0 p'_z)$$

→ interpret as "four-momentum"  
 $\tilde{p} = (E/c, \vec{p})$

## Time Dilation

→ for particle with finite lifetime  $\tau$  in rest frame  $\tau_{LAB} = \gamma \tau$

## Particle Decay:

(10)

Only a few particles appear to have infinite, or nearly infinite, lifetimes:

$$e^{\pm}, \nu_e, p, \gamma$$

For particles with finite lifetime

Distance traveled in Lab

$$\lambda_D = \cancel{\beta} c \tau = \gamma c \tau = \left(\frac{p}{mc}\right) c \tau$$

$\propto p$

If  $N_0$  unstable particles at  $x = \tau = 0$ ,

$N(x)$  = number particles at position  $x$

$dN(x)$  = number in interval  $dx$  around  $x$

$$dN(x) = -N(x) \frac{dx}{\lambda_D}$$

Integrating:  $N(x) = N_0 e^{-x/\lambda_D}$

## Some Important Decaying Particles

(11)

Long-lived particles: travel meters or more in typical lab frame

$$\mu, \pi^\pm, \eta, K_L^0, K^\pm$$

These particles may strike a detector near where they're produced

Short-lived particles:

→ some live long enough that their distance traveled in the lab frame is detectable

$K_S^0, \Lambda$ : several cm in rest frame

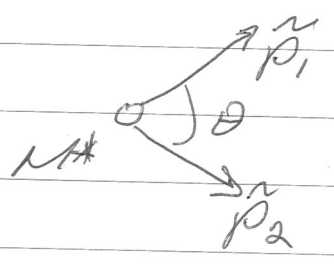
$\tau$  lepton,  $c, b$  quarks: up to several 100  $\mu\text{m}$  in LAB frame

Some important decays:

|  | <u>Branching Fraction</u> |
|--|---------------------------|
| $\pi^0 \rightarrow \gamma\gamma$                     | 98.8%                     |
| $\mu \rightarrow e \bar{\nu}_e \nu_\mu$              | 100%                      |
| $\pi^+ \rightarrow \mu^+ \nu_\mu$                    | 100%                      |
| $\tau \rightarrow \pi^+ \pi^- \pi^+ \nu_\tau$        | 9%                        |
| $b \rightarrow c \mu \bar{\nu}_\mu, c e \bar{\nu}_e$ | 17% each                  |

# Lorentz Invariant Quantities

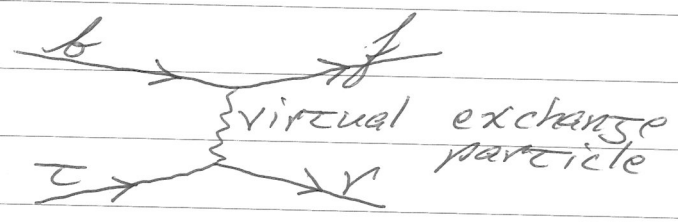
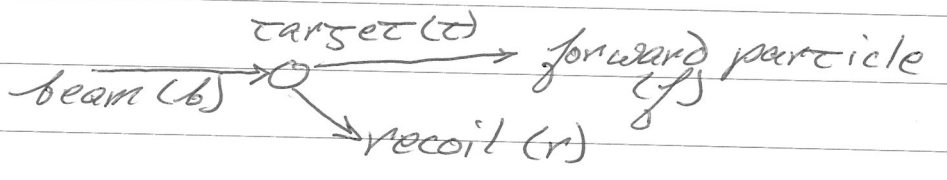
"Invariant mass" of two decay particles



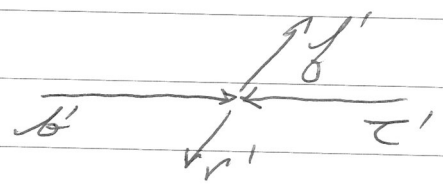
$$\begin{aligned}
 M &= (\vec{p}_1 + \vec{p}_2)^2 \\
 &= (\vec{p}_1^2 + \vec{p}_2^2 + 2\vec{p}_1 \cdot \vec{p}_2) \\
 &= m_1^2 + m_2^2 + 2(E_1 E_2 - p_1 p_2 \cos \theta)
 \end{aligned}$$

Total center-of-momentum (CM) frame energy:

In LAB:



In CM frame:



In either case (primed or unprimed), total energy involved in interaction:

$$\begin{aligned}
 s &= (\vec{p}_b + \vec{p}_t)^2 = m_b^2 + m_t^2 + 2(E_b E_t - \vec{p}_b \cdot \vec{p}_t) \\
 &= W^2 \text{ (total energy)}
 \end{aligned}$$

In symmetric collider

$$\begin{aligned}
S &= (\vec{p}_b + \vec{p}_c)^2 = (\vec{E}_b + \vec{p}_b + \vec{E}_c + \vec{p}_c) \\
&= (\vec{E}'_b + \vec{E}'_c)^2 = W^2
\end{aligned}$$

$$W \propto \sqrt{2} E_b \propto \underline{\underline{E_b}}$$

In fixed target ( $\vec{p}_c = 0$ )

$$S = m_b^2 + m_c^2 + 2m_c E_b$$

At high energy ( $E_b \gg m_b$  or  $m_c$ )

$$\begin{aligned}
W^2 &\propto E_b \\
W &\propto \sqrt{E_b}
\end{aligned}$$

∴ colliders more efficient to turn beam energy into collision energy  
 → and possibly mass for heavy particles