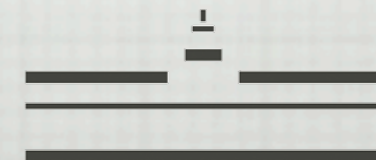


Parton Distribution Functions

for beginners

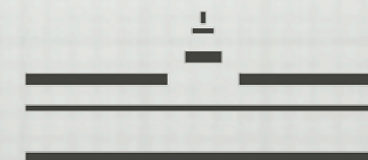
KAROL KOVAŘÍK

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Lecture 2



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Recap of Lecture 1

- Parton Distribution Functions - parameterisation @ initial scale (DGLAP initial condition)

CTEQ 6

hep-ph/0201195

$$x f_k(x, Q_0) = c_0 x^{c_1} (1-x)^{c_2} e^{c_3 x} (1 + e^{c_4 x})^{c_5} \quad k = u_v, d_v, g, \bar{u} + \bar{d}$$

$$\bar{d}(x, Q_0)/\bar{u}(x, Q_0) = c_0 x^{c_1} (1-x)^{c_2} + (1 + c_3 x)(1-x)^{c_4}$$

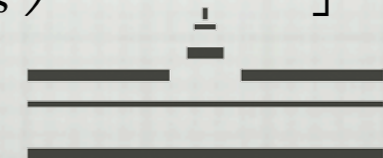
$$s = \bar{s} = 0.2 (\bar{u} + \bar{d}) \quad c(x, Q_0^2) = b(x, Q_0^2) = 0$$

- Parton Distribution Functions - Q dependance & evolution (DGLAP)

- all contributions together mix different PDF via DGLAP equations

$$\frac{d f_q(x, Q^2)}{d \log Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{d\xi}{\xi} \left[P_{qq} \left(\frac{x}{\xi} \right) f_q(\xi, Q^2) + P_{qg} \left(\frac{x}{\xi} \right) f_g(\xi, Q^2) \right]$$

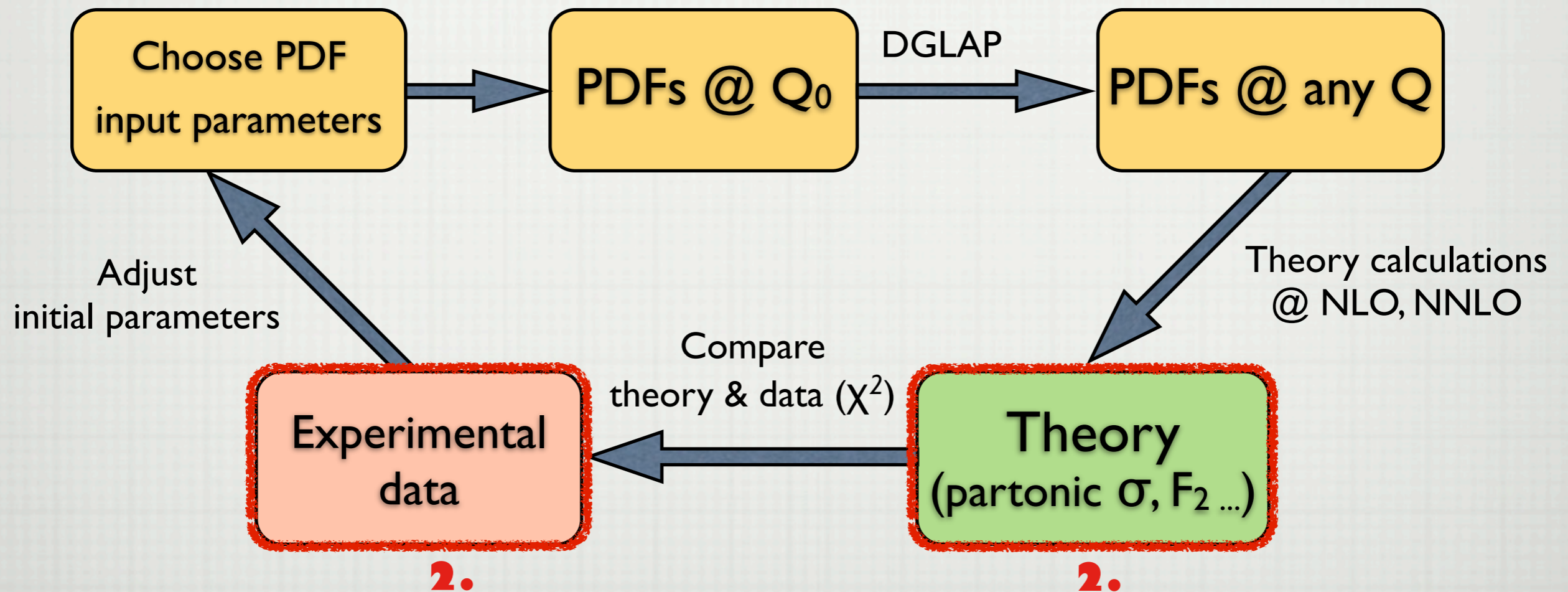
$$\frac{d f_g(x, Q^2)}{d \log Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{d\xi}{\xi} \left[P_{gg} \left(\frac{x}{\xi} \right) f_g(\xi, Q^2) + P_{gq} \left(\frac{x}{\xi} \right) f_q(\xi, Q^2) \right]$$



Fitting PDF

- *Parton Distribution Functions*

- how do we determine them? What are the moving parts in a typical PDF fitting-machine?



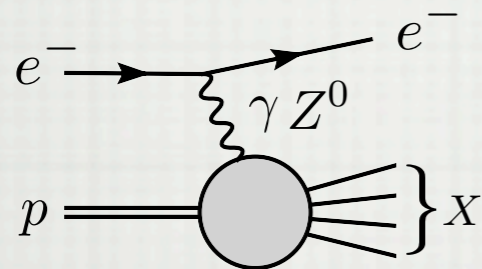
Fitting PDF

- *Parton Distribution Functions - experimental data & theory*

- which data are included in a fit ? Which PDFs do they constrain ?
- kinematic cuts on data
- difference between global fits like CTEQ, MSTW, NNPDF and not so global HERApdf, ABM...

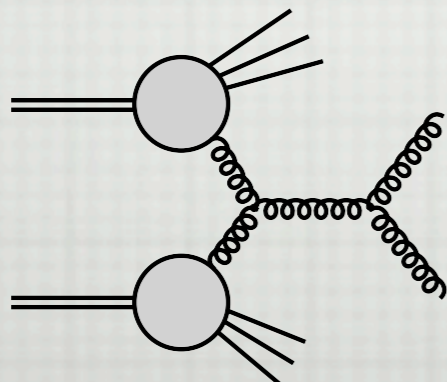
- ☑ *Neutral current DIS*

(HERA, SLAC, NMC, BCDMS)



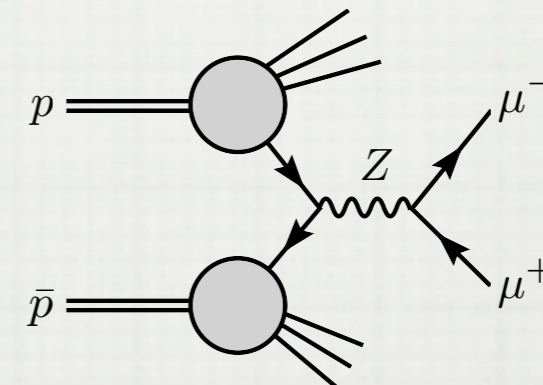
- ☑ *Jet data from Tevatron & LHC*

(D0, CDF, ATLAS, CMS)



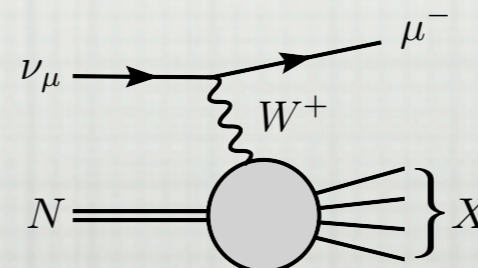
- ☑ *DY data (W,Z production) from Tevatron & LHC*

(E602, E866, D0, CDF, ATLAS, CMS)



- ☑ *Neutrino DIS & di-muon*

(CDHSW, CHORUS, NuTeV)



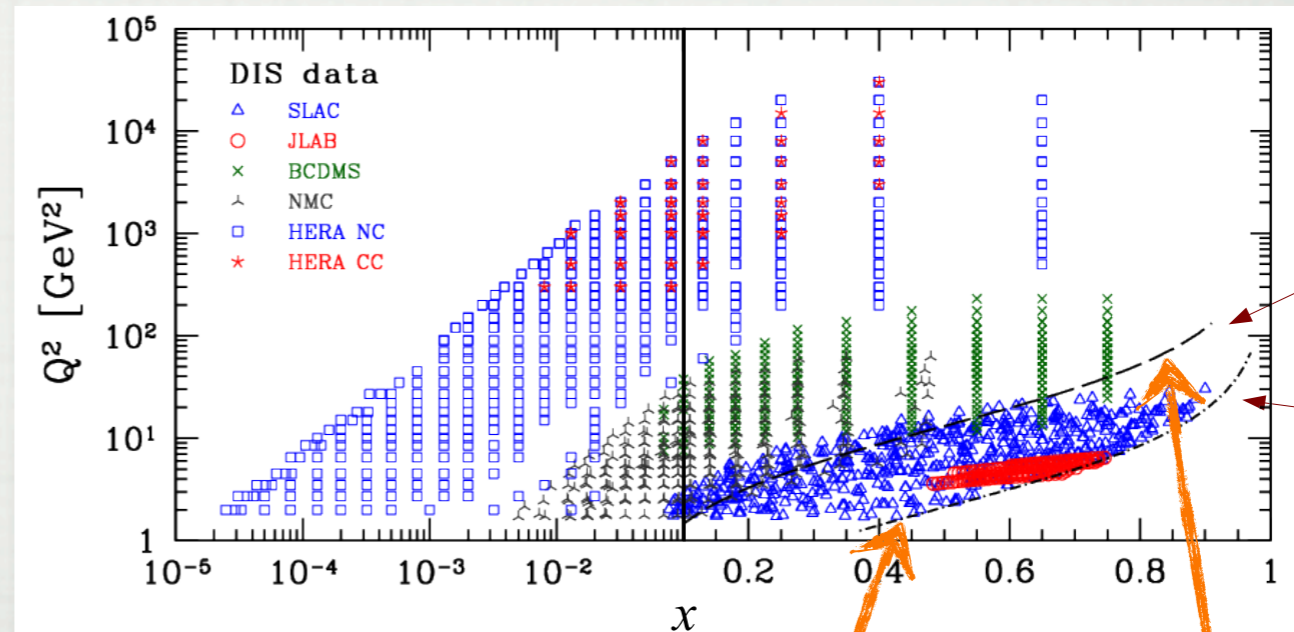
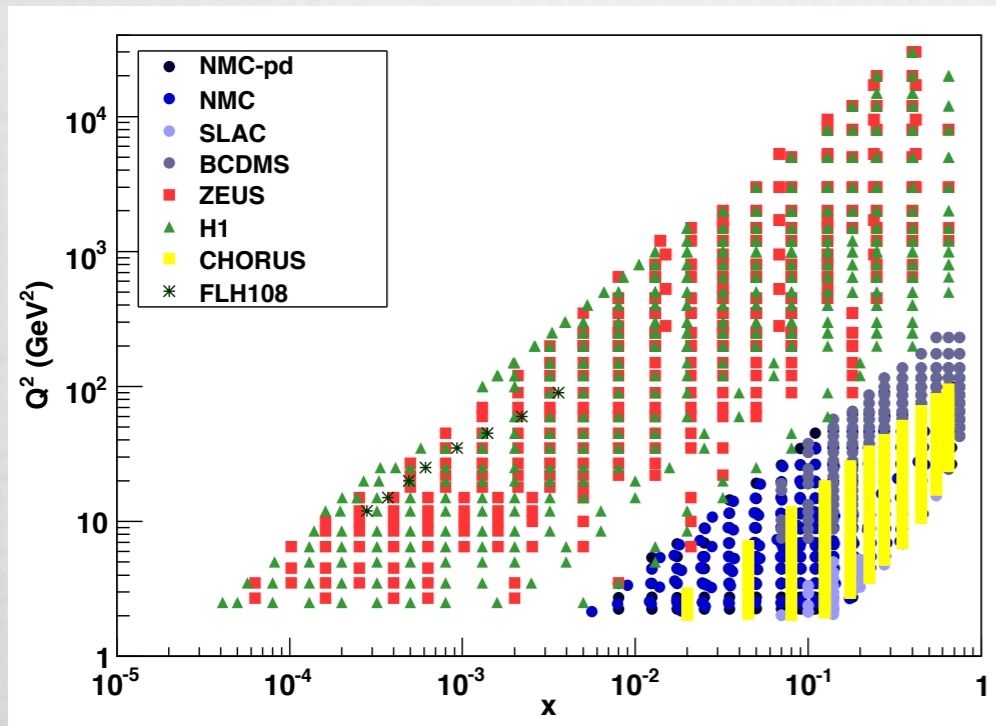
Fitting PDF

☑ DIS data from HERA & fixed target experiments (SLAC, NMC, BCDMS)

- different observables used in the fit (HERA - $\frac{d\sigma}{dx dy}$, fixed target - $F_2(x, Q^2)$)

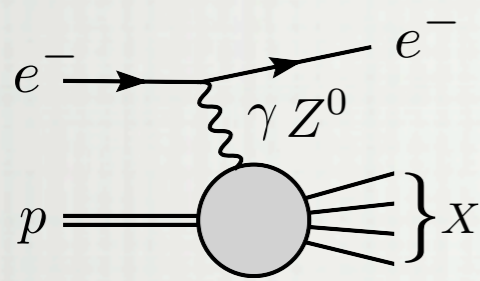
$$\frac{d\sigma}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} [Y_+ F_2 - y^2 F_L \pm Y_- x F_3],$$

- motivation to still use fixed target is different coverage in x - Q^2 plane & sensitivity to large- x



Fitting PDF

☑ DIS data from HERA & fixed target experiments (SLAC, NMC, BCDMS)



- DIS @ low Q - dominated by photon exchange
- DIS @ high Q - dominated by photon-Z interference

$$F_{1,2} = F_{1,2}^{\gamma\gamma} + \frac{2g_V}{4s_W^2 c_W^2} \frac{Q^2}{Q^2 + M_Z^2} F_{1,2}^{\gamma Z} + \frac{g_V^2 + a_V^2}{16s_W^4 c_W^4} \frac{Q^4}{(Q^2 + M_Z^2)^2} F_{1,2}^{ZZ}$$

$$F_3 = \frac{2a_V}{4s_W^2 c_W^2} \frac{Q^2}{Q^2 + M_Z^2} F_3^{\gamma Z} + \frac{2g_V a_V}{16s_W^4 c_W^4} \frac{Q^4}{(Q^2 + M_Z^2)^2} F_3^{ZZ},$$

sensitive to quark & anti-quark PDF @ LO

- quarks & anti-quarks enter together with different weights depending on the exchange vector boson
- access interference comparing different helicity leptons & electrons/positrons

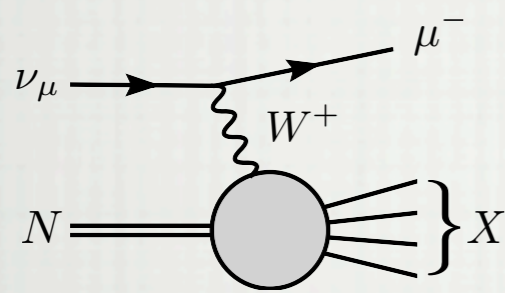
$$F_2^{\gamma\gamma}(x, Q^2) = x \sum_q e_q^2 [q(x, Q^2) + \bar{q}(x, Q^2)] \quad \text{photon}$$

$$\text{NC DIS} \quad F_2^{\gamma Z}(x, Q^2) = x \sum_i B_i [q_i(x, Q^2) + \bar{q}_i(x, Q^2)] \quad \text{photon-Z interference}$$

$$xF_3^{\gamma Z}(x, Q^2) = x \sum_i D_i [q_i(x, Q^2) - \bar{q}_i(x, Q^2)] \quad \text{photon-Z interference}$$

Fitting PDF

☑ Neutrino DIS & di-muon (CDHSW, CHORUS, NuTeV)



- neutrino DIS contributes to $F_2(x, Q^2)$ and $F_3(x, Q^2)$

- different PDF combinations contribute to flavor separation together with NC DIS

$$F_2(x, Q^2) = x \sum_q [q(x, Q^2) + \bar{q}(x, Q^2)]$$

$$xF_3(x, Q^2) = x \sum_q [q(x, Q^2) - \bar{q}(x, Q^2)]$$

- neutrino DIS data on protons are scarce and hard to come by (WA21/22)

- neutrino DIS typically taken on nuclei - need for nuclear corrections

☑ Charge current DIS on proton (HERA)

- neutrino DIS can be replaced by CC DIS on protons (still experimentally challenging)

$$F_2^{W^\pm}(x, Q^2) = x(\bar{u} \pm d \pm s + \bar{c})$$

Fitting PDF

✓ DIS @ NLO

sensitive to gluon PDF @ NLO

- breaking of Bjorken scaling at small x driven by gluon PDF

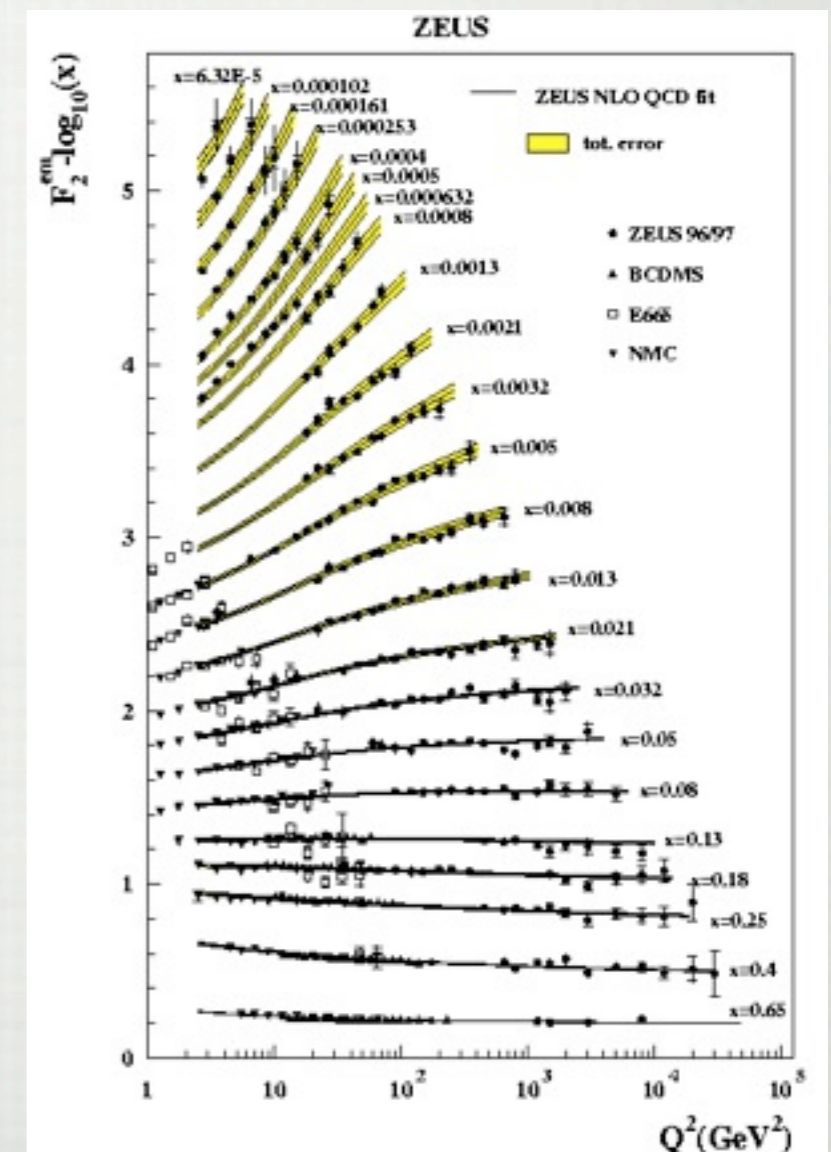
$$Q^2 \frac{dF_2}{dQ^2} = \frac{\alpha_S}{2\pi} \sum e_i^2 \int_x^1 \frac{dy}{y} P_{qg}(y) f_g(x/y, Q^2)$$

- longitudinal structure function $F_L(x, Q^2)$

$$F_L(x, Q^2) = F_2(x, Q^2) - 2xF_1(x, Q^2) = 0$$

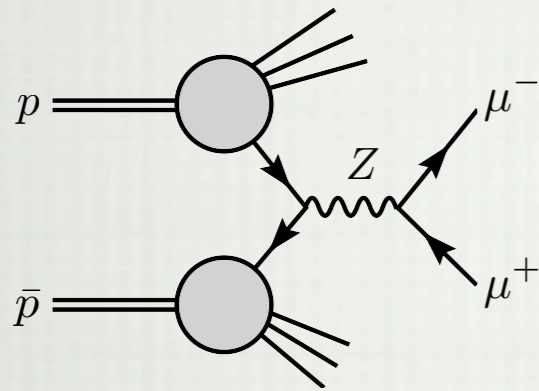
Callan-Gross relation @ LO

- gluon not subdominant in $F_L(x, Q^2)$ as in $F_2(x, Q^2)$
- experimental separation of $F_2(x, Q^2)$ and $F_L(x, Q^2)$ requires measurements at different \sqrt{s} - lower statistics
- using differential cross-section in the fit - effective separation of structure functions



Fitting PDF

☑ Drell-Yan lepton pair production (W,Z production)



- DY dominated by photon exchange away from W & Z resonances

$$\frac{d\sigma}{dQ^2 dy} = \frac{4\pi\alpha^2}{9Q^2 s} \sum_i e_i^2 [q_i(x_a, Q^2)\bar{q}_i(x_b, Q^2) + a \leftrightarrow b]$$

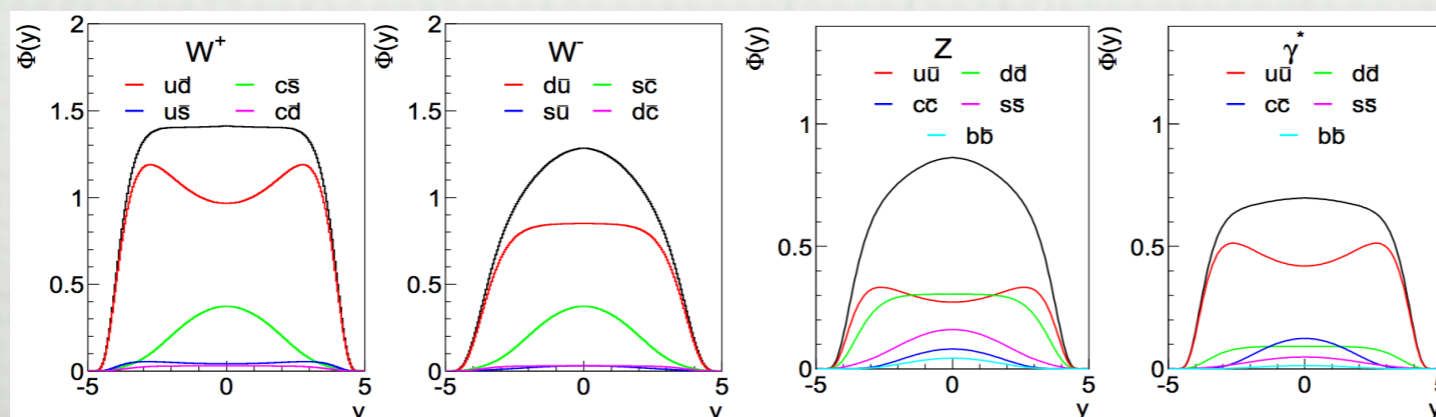
$$Q^2 = (p_l + p_{\bar{l}})^2 \quad x_{a,b} = \frac{Q}{\sqrt{s}} \exp(\pm y)$$

- DY at the W & Z resonances - different PDF combinations

$$\frac{d\sigma^W}{dy} = \frac{\sqrt{2}\pi G_F m_W^2}{3s} \sum_{i,j} |V_{ij}^{\text{CKM}}| [q_i(x_a, Q^2)\bar{q}_j(x_b, Q^2) + a \leftrightarrow b]$$

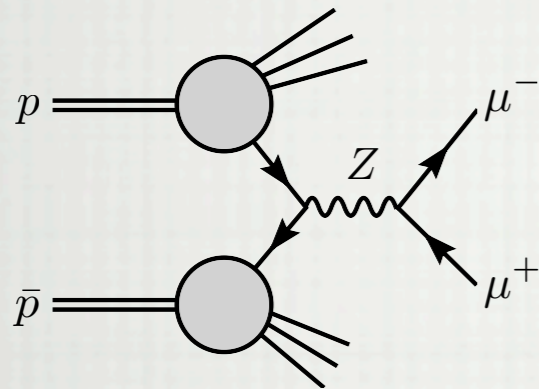
$$\frac{d\sigma^Z}{dy} = \frac{\sqrt{2}\pi G_F m_Z^2}{3s} \sum_i (V_i^2 + A_i^2) [q_i(x_a, Q^2)\bar{q}_i(x_b, Q^2) + a \leftrightarrow b]$$

- DY helps better determine (anti-)quark PDFs



Fitting PDF

☑ Drell-Yan lepton pair production (W,Z production)



- different asymmetries & ratios sensitive to different combinations of PDFs

$$\frac{\sigma_{W^+} - \sigma_{W^-}}{\sigma_{W^+} + \sigma_{W^-}} \sim \frac{u_v(x_1) - d_v(x_1)}{u(x_1) + d(x_1)} \quad \frac{\sigma_{W^+}}{\sigma_{W^-}} \sim \frac{u(x_1)\bar{d}(x_2)}{d(x_1)\bar{u}(x_2)} \sim \frac{u(x_1)}{d(x_1)}$$

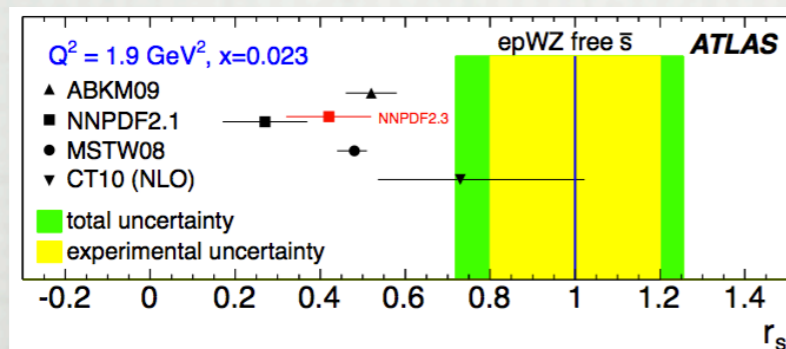
$$\frac{\sigma_{W^+} + \sigma_{W^-}}{\sigma_Z} \sim \frac{u(x_1) + d(x_1)}{0.29u(x_1) + 0.37d(x_1)}$$

- Z cross-section + rapidity distribution used to extract strange quark PDF

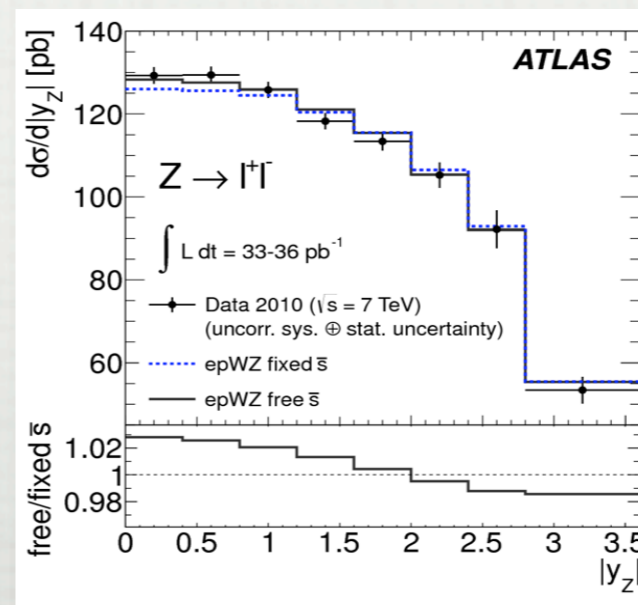
typical assumptions related to strange PDF (motivated by di-muon DIS data)

$$s(x, Q) = \bar{s}(x, Q) \quad r_s \sim 0.75$$

$$s(x, Q) = r_s d(x, Q)$$

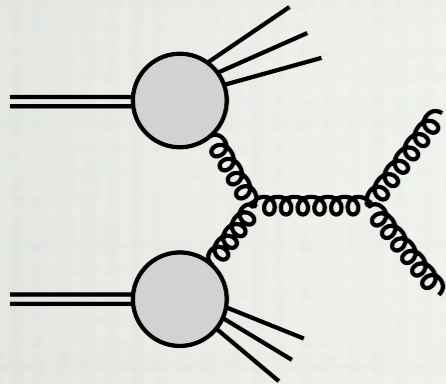


arXiv:1203.4051 [hep-ex]



Fitting PDF

✓ Jet data from Tevatron & LHC (D0, CDF, ATLAS, CMS)

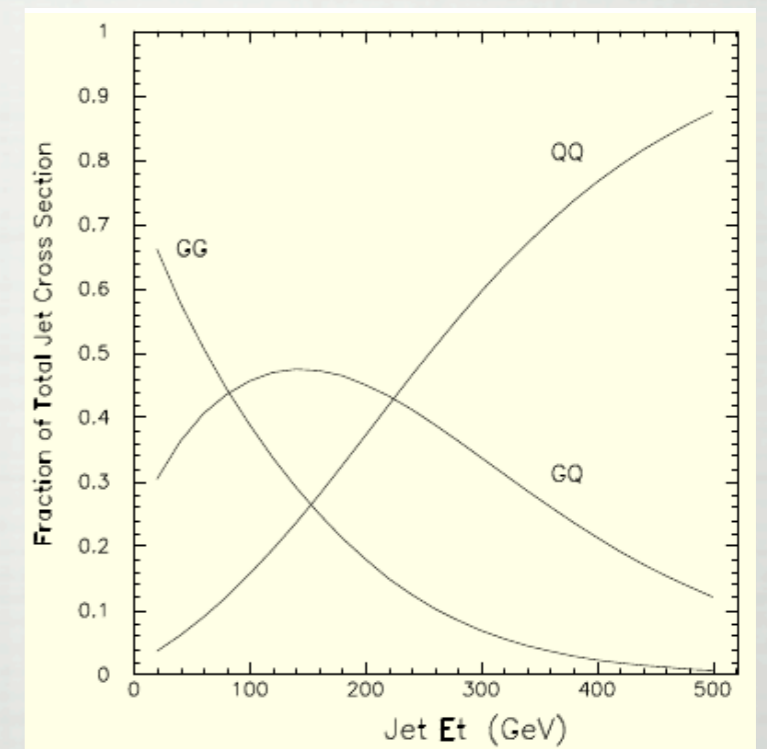


- hadronic jet production at leading order proceeds through

$$qq \rightarrow qq \quad qg \rightarrow qg \quad gg \rightarrow gg$$

- qq subprocess dominates high-Et jets but gluon important enough to allow jet data to put constraints on large-x gluon PDF

- combined with low-x constraints on gluon PDF from DIS and with sum rules one has strong constraints on the gluon PDF
- additional direct probes of gluon PDF needed to constrain the gluon PDF at mid-x and large-x for future searches e.g. SUSY @ LHC



Fitting PDF

Using LHC data in Parton Distribution Functions

Motto: "Yesterday's signal is today's background"

Top pair production

sensitive to gluon PDF at high-x

- very precise top pair production expected from LHC top-factory
- ratios of top/anti-top cross-sections sensitive also to u/d

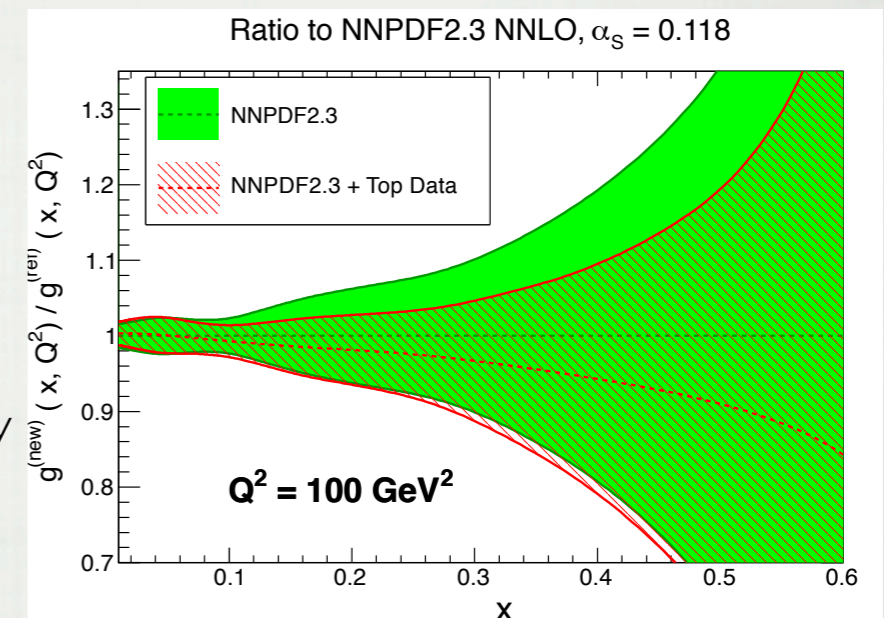
Direct photons

additional, complementary probe of gluon PDF (same x as gg Higgs production)

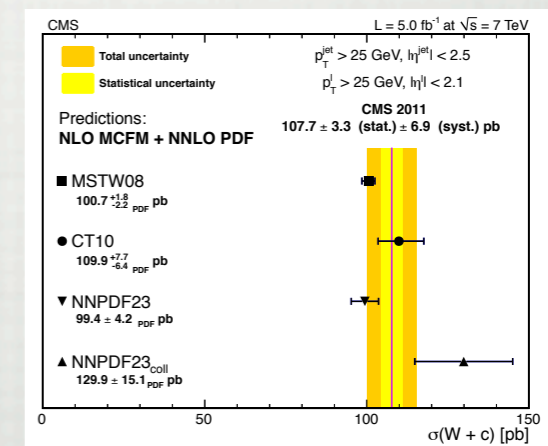
W+charm production

sensitivity to strange quark PDF (difficult to extract elsewhere)

arXiv:1301.7215



CMS-SMP-12-002



Fitting PDF

- *Theory predictions for relevant observables*

- Theory predictions can be done at LO, NLO or NNLO (at the moment)

- at each order PDF has modified meaning - need to use LO PDF with LO predictions etc.

- *leading-order*

- hard scattering results simple but with no scale dependence (no cancellation - large scale dependence)
- LO PDF useful for some MC applications where only LO exists
- data descriptions unsatisfactory

- *next-to-leading-order*

- scale cancellation between hard scattering and PDF - lesser scale dependence
- hard scattering matrix elements complicated & need to be evaluated many times in a fit
- NLO hard matrix elements together with NLO DGLAP - current state-of-the-art

- *next-to-next-to-leading-order*

- NNLO splitting functions known but not all relevant hard scattering matrix element known
e.g. jets or other additional processes

Fitting PDF

- *Theory predictions for relevant observables*

- NLO & NNLO theory predictions complicated functions - need a way how to evaluate them quickly & efficiently
- Full cross-section is a convolution of hard matrix elements with PDFs (or two convolutions)

$$\sigma(\mu_r, \mu_f) = \sum_{n,i} c_{n,i}(x_a, x_b, \mu_r, \mu_f) \otimes [\alpha_s^n(\mu_r) F_i(x_a, x_b, \mu_f)]$$

- Old (but effective) method - use K-factors - lose some (N)NLO information about shape
- Alternative - decouple PDFs and strong coupling dependence by putting everything on a (x,Q) grid & pre-compute complicated matrix element once and for all

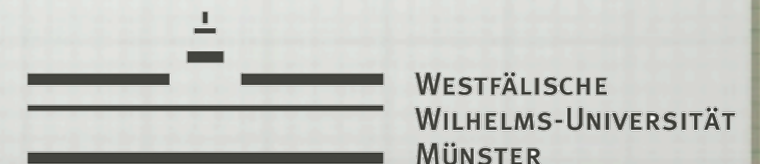
$$\sigma(\mu) \simeq \sum_{n,i,k,l,m} \tilde{\sigma}_{n,i,k,l,m}(\mu) \alpha_s^n(\mu^{(m)}) F_i(x_a^{(k)}, x_b^{(l)}, \mu^{(m)})$$

FastNLO

hep-ph/0609285

APPLGRID

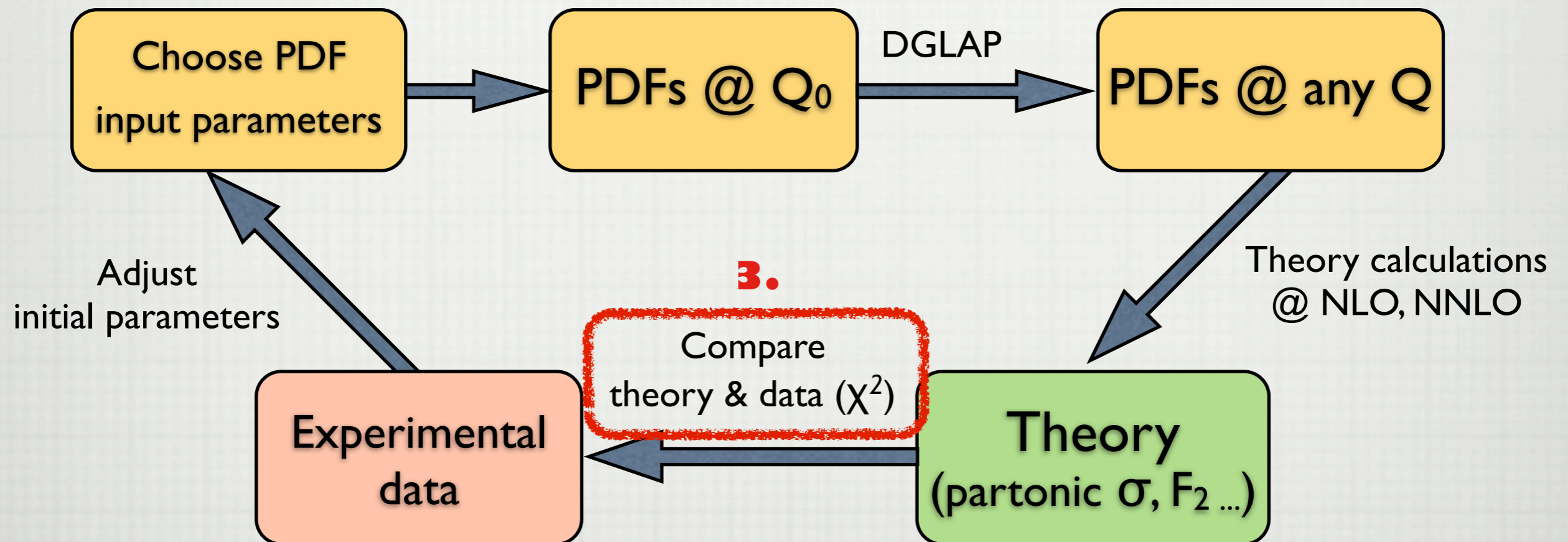
arXiv:0911.2985



Fitting PDF

- *Parton Distribution Functions*

- how do we determine them? What are the moving parts in a typical PDF fitting-machine?



Fitting PDF

- Parton Distribution Functions - χ^2 -fit & errors

- Most PDF fitters use χ^2 -function to measure the goodness of the fit

standard definition

$$\chi^2 = \sum_{i=1}^{N_{\text{dat}}} \left(\frac{D_i - T_i}{\sigma_i} \right)^2$$

definition with correlated errors

$$\chi^2 = \sum_{i=1}^{N_{\text{dat}}} \sum_{j=1}^{N_{\text{dat}}} (D_i - T_i) (V^{-1})_{ij} (D_j - T_j)$$

$$V_{ij} = \delta_{ij} (\sigma_i^{\text{uncorr}})^2 + \sum_{k=1}^{N_{\text{corr}}} \sigma_{k,i}^{\text{corr}} \sigma_{k,j}^{\text{corr}}$$

covariance matrix

- Try to use all possible experimental information available
 - statistical errors
 - systematic errors - (un)correlated
 - normalisation uncertainty (might be multiplicative)

Fitting PDF

● Parton Distribution Functions - what do we get ?

- What are the features we see in a typical PDF result ?

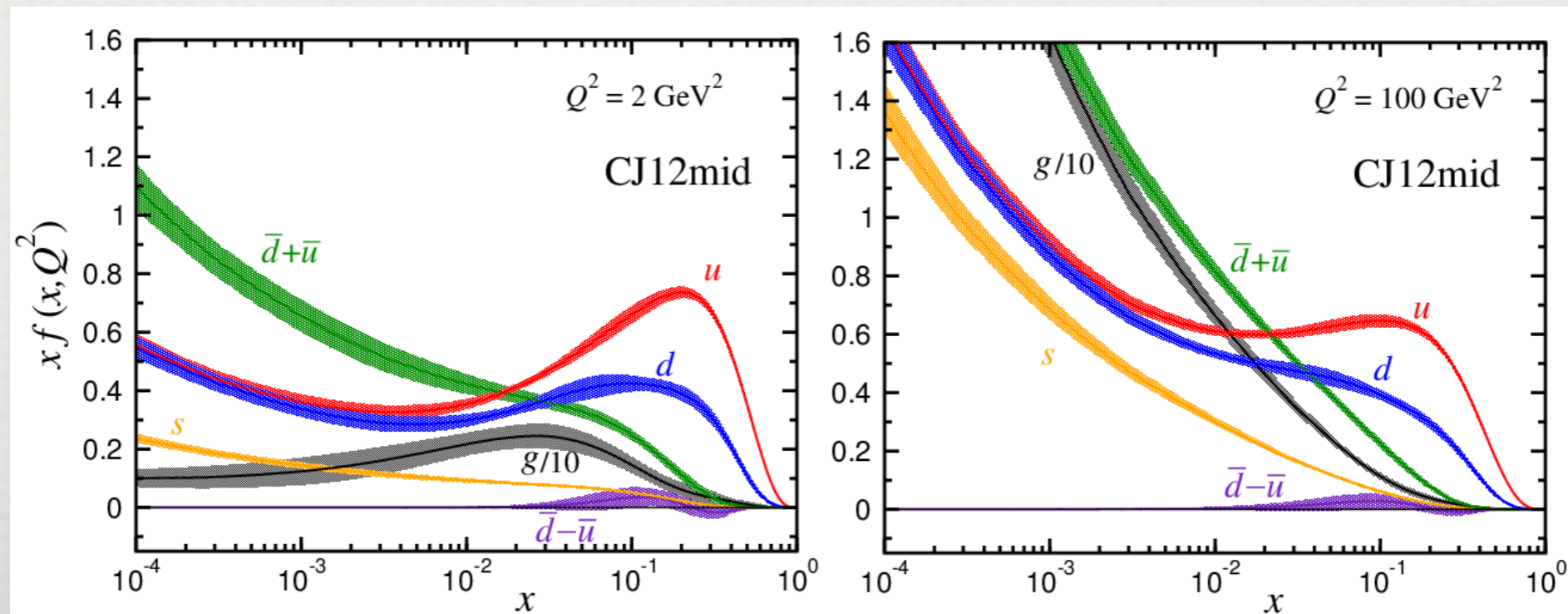
☑ *u&d quarks - valence & sea*

- valence part causes u&d dominate all other PDF at large-x where $u > d$
- symmetric sea-quark: q & anti- q comparable at low-x
- at high Q - contribution of the sea component increases through gluon radiation (DGLAP)

☑ *strange quarks*

- strange quark PDF suppressed at initial scale but enhanced at high- Q

arXiv:1212.1702



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Fitting PDF

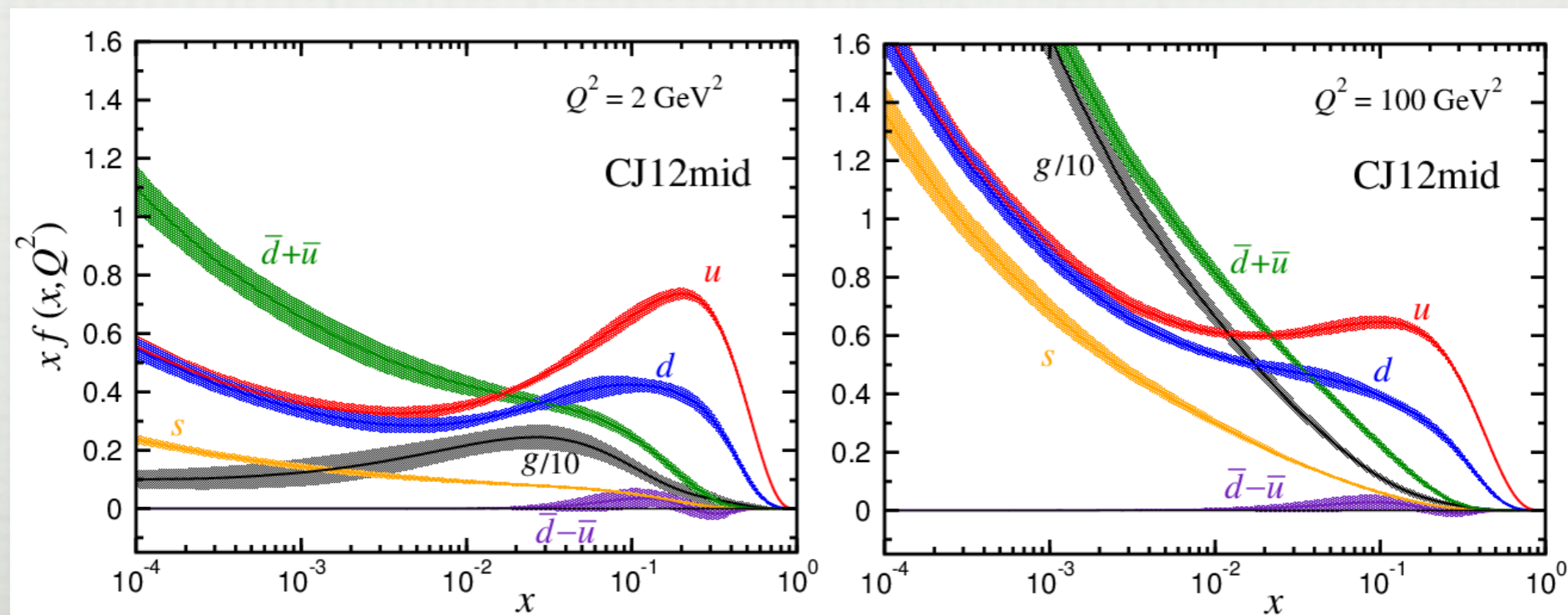
● Parton Distribution Functions - what do we get ?

- What are the features we see in a typical PDF result ?

☑ gluon

- dominate at small- x but fall off steeply as x increases
- going to high- Q - gluon radiation reduces momenta of partons - everything shifts to smaller x
- gluon radiates q - q bar pairs or additional gluons - at small- x gluon PDF and sea quark PDF get steeper
- gluon can radiate even heavy quarks at high- Q so charm and bottom PDF are non-zero

arXiv:1212.1702



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PDF uncertainties

● Parton Distribution Functions - what are PDF uncertainties

☑ error PDF

- uncertainty of experimental data can be interpreted as uncertainty of the underlying PDF parameters
- different approaches how to translate experimental uncertainties to PDFs

☑ other uncertainties (not in error PDFs)

- choice of data sets or observables (include neutrino DIS or not, LHC or not ...)
- choice of kinematic cuts (looser cuts might constrain PDF better but ...)
- parameterisation bias
- pQCD choices (NLO vs NNLO, strong coupling)
- heavy-quark schemes (FFS, ZM-VFNS, VF-VFNS)
- higher-twist terms, nuclear corrections etc...

PDF uncertainties

- *Parton Distribution Functions - χ^2 -fit & errors*

- error PDFs are experimental errors translated to errors of free PDF parameters
- all approaches to determine error PDFs give approx. the same results in regions with data

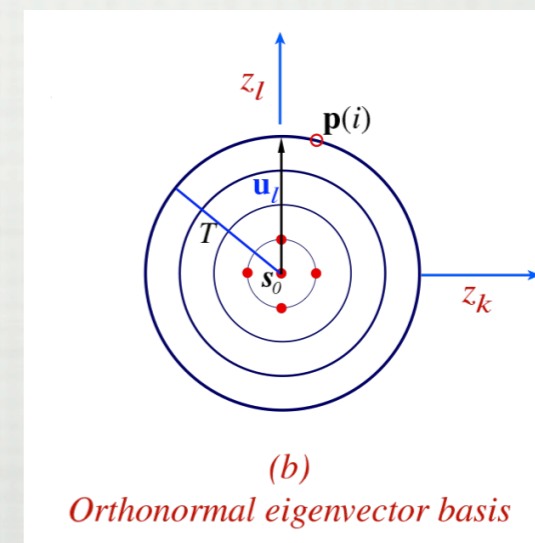
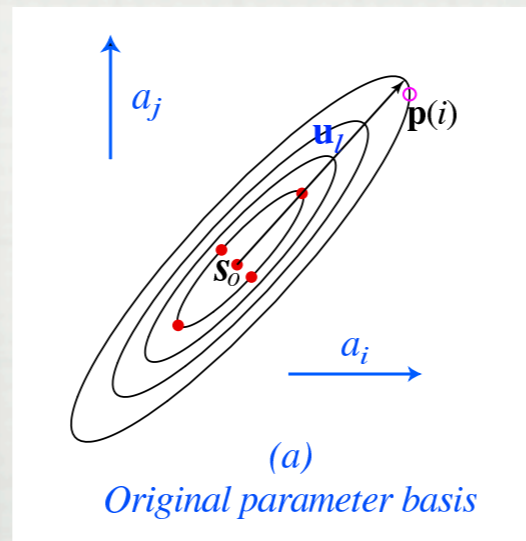
- ☑ **Hessian method**

the most widely used technique to determine error PDFs

Expansion of χ^2 :

$$\chi^2(a) = \chi_0^2 + \frac{1}{2} \frac{\partial^2 \chi^2}{\partial a_i \partial a_j} (a - a_0)_i (a - a_0)_j + \dots \rightarrow \chi_0^2 + \sum_i z_i^2$$

Hessian diagonal Hessian



hep-ph/0201195

PDF uncertainties

- *Parton Distribution Functions - χ^2 -fit & errors*

- error PDFs are experimental errors translated to errors of free PDF parameters
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- ☑ *Hessian method*

the most widely used technique to determine error PDFs

Expansion of χ^2 :

$$\chi^2(a) = \chi_0^2 + \frac{1}{2} \frac{\partial^2 \chi^2}{\partial a_i \partial a_j} (a - a_0)_i (a - a_0)_j + \dots \rightarrow \chi_0^2 + \sum_i z_i^2$$

Hessian diagonal Hessian

Choice of $\Delta\chi^2 = \chi^2 - \chi_0^2$:

ideal choice $\Delta\chi^2 = 1$ pragmatic choice $\Delta\chi^2 \gg 1$ $\Delta\chi^2 \sim 50 - 100$

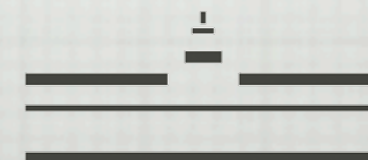
error PDFs

Construct error PDFs for each parameter in 2 directions:

$$z_i = \pm \sqrt{\Delta\chi^2} \qquad X_i^\pm(z) = X_i^\pm(0, 0, \dots, \pm \sqrt{\Delta\chi^2}, \dots, 0, 0)$$

Calculate PDF uncertainty of cross-section

$$(\Delta\sigma)^2 \approx \frac{1}{4} \sum_i^{N_p} \left(\sigma(X_i^+) - \sigma(X_i^-) \right)^2$$



PDF uncertainties

- Parton Distribution Functions - χ^2 -fit & errors

- ☑ Hessian method - dynamical tolerance criterion

- Ideal case would require using $\Delta\chi^2 = 1$ for one sigma (68% CL) or $\Delta\chi^2 = 2.71$ for 90% CL
 - BUT we are fitting data from multiple not necessarily compatible experiments

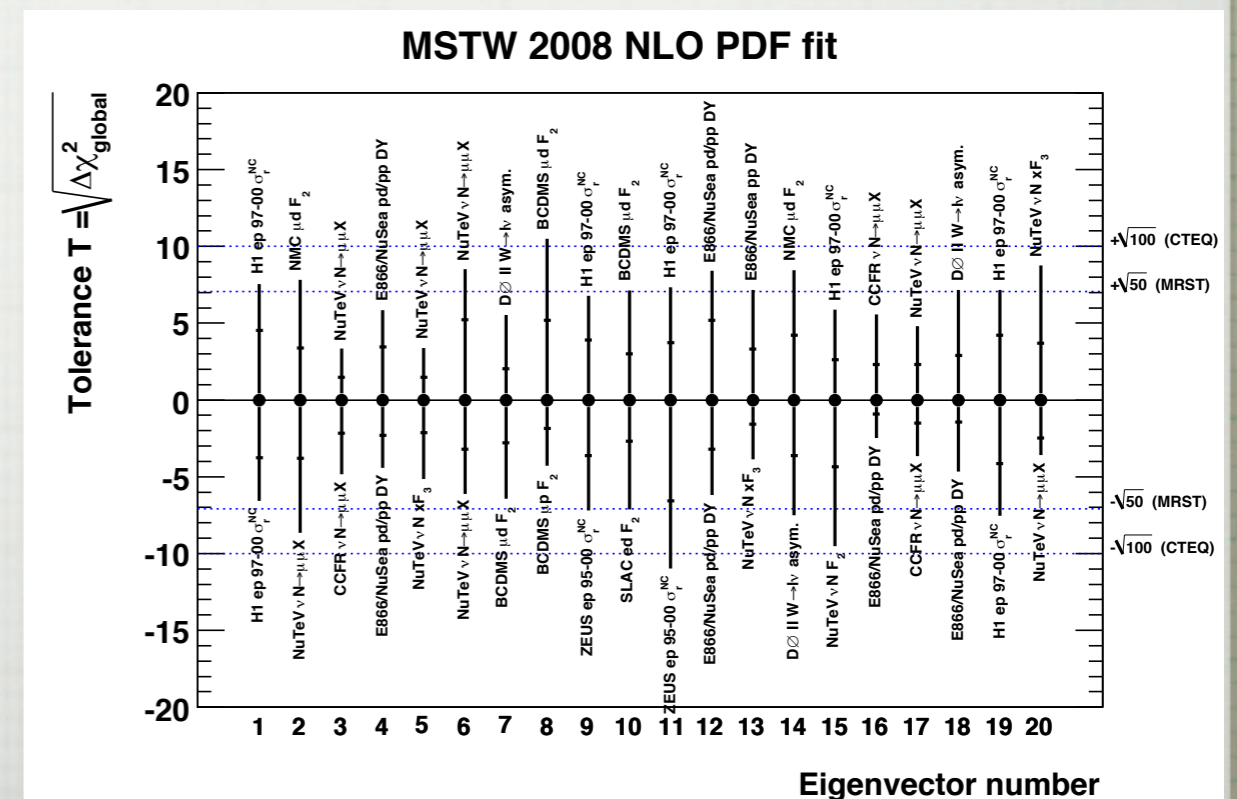
dynamical tolerance by MSTW

- Each experiment (N data points) have to be described up to 90%CL for variations along one eigenvector

$$P_N(\chi^2) = \frac{(\chi^2)^{N/2-1} e^{-\chi^2/2}}{2^{N/2} \Gamma(N/2)}$$

$$\int_0^{\xi_{90}} d\chi^2 P_N(\chi^2) = 0.90$$

- For each eigenvector, take $\Delta\chi_n^2$ where all experiments are described within 90% CL
 - Translate $\Delta\chi_n^2$ to (different) shifts in each parameter



PDF uncertainties

● Parton Distribution Functions - χ^2 -fit & errors

- error PDFs are experimental errors translated to errors of free PDF parameters
- all approaches to determine error PDFs give approx. the same results in regions with data

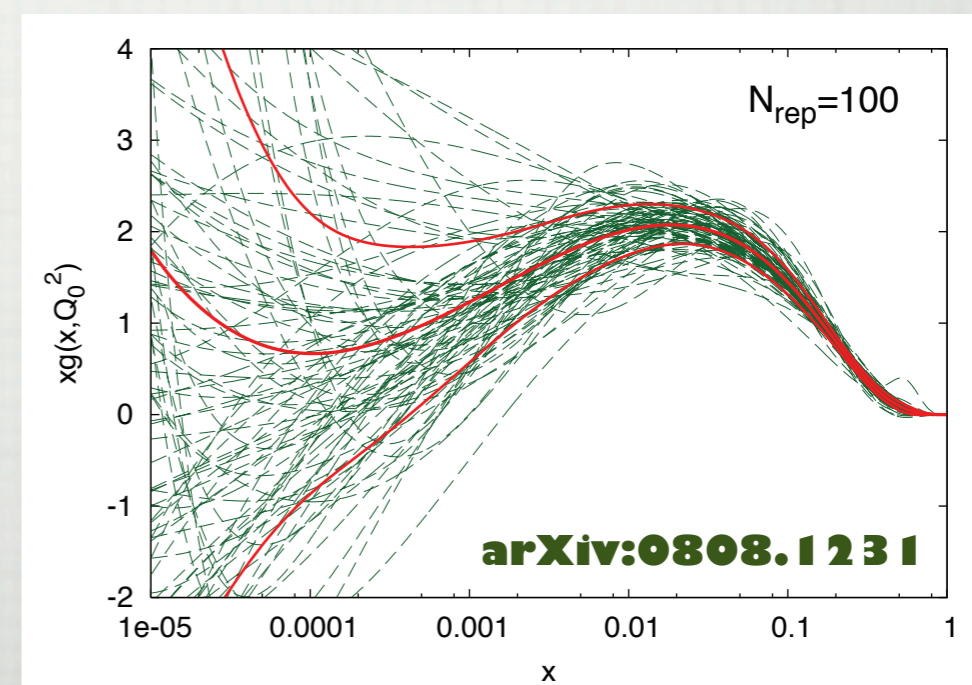
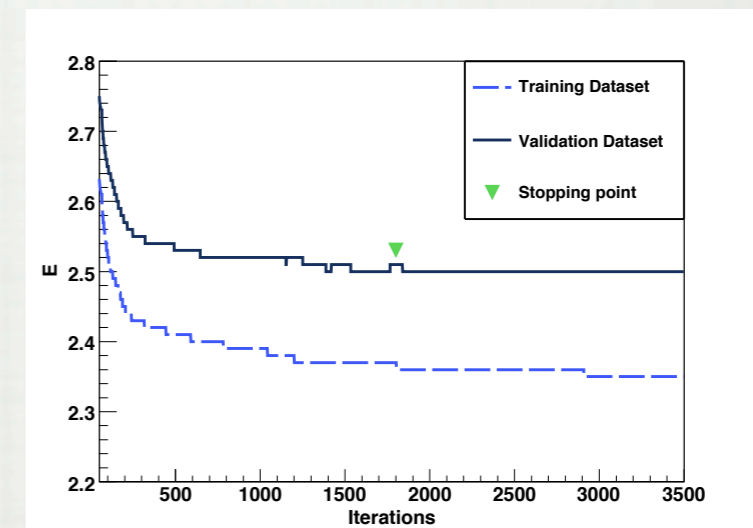
☑ Monte Carlo method (as used by NNPDF)

new technique which allows for more flexible PDF x-shapes

- neural network is used to (over-)parametrize PDFs @ Q_0
- N_{set} artificial replicas of data points generated assuming multi-Gaussian probability distribution
- random separation into training & validation data subsets
- minimize error function (not χ^2) for training set
- stop before overlearning

$$\langle \sigma(X) \rangle = \frac{1}{N_{\text{set}}} \sum_{i=1}^{N_{\text{set}}} \sigma(X^i) \quad \leftarrow \text{error/replica PDFs}$$

$$\Delta \sigma(X) = \left(\sum_{i=1}^{N_{\text{set}}} \left[\sigma(X^i) - \langle \sigma(X) \rangle \right]^2 \right)^{1/2}$$



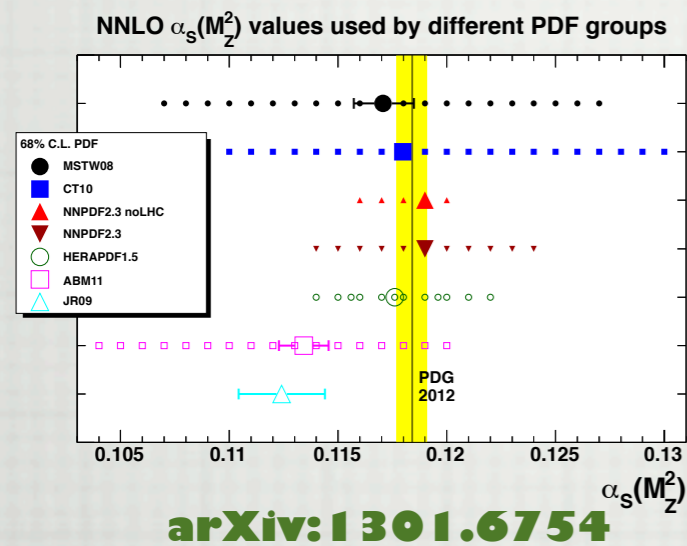
Misc topics in PDF

● Parton Distribution Functions - theory & related issues

- theory calculations @ LO, NLO or NNLO include several important constants which can have large impact on PDFs

☑ strong coupling α_s

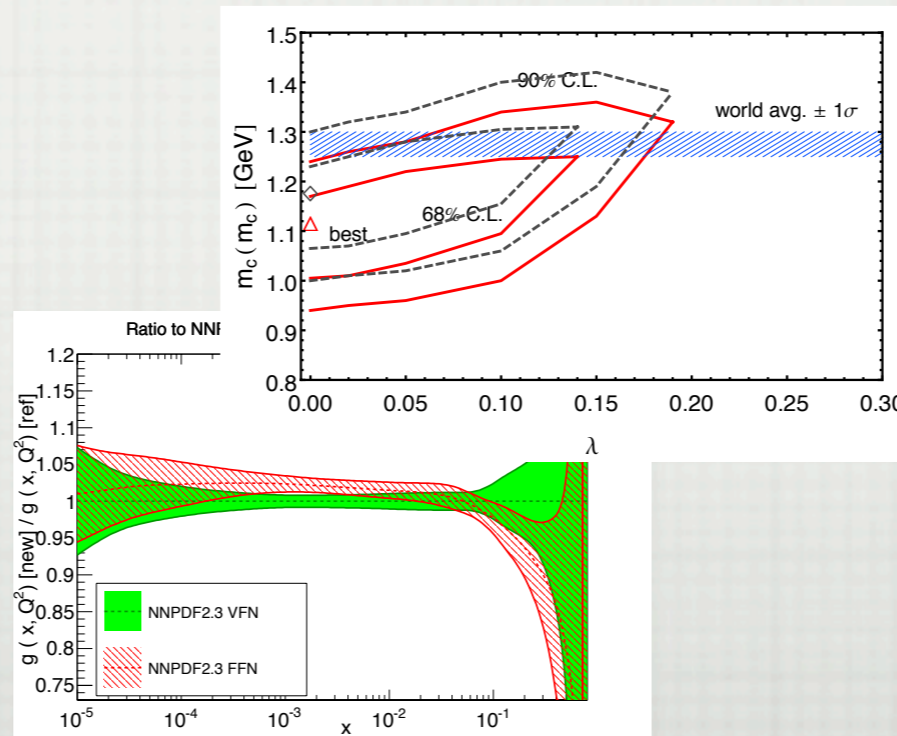
treated as an external parameter
or fitted together with PDFs



☑ quark masses m_c, m_b & heavy quark treatment

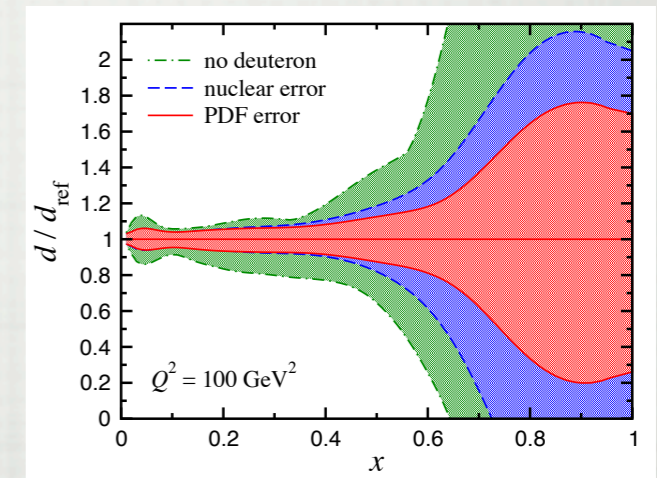
quark masses enter the evolution
& their treatment influences gluon PDFs

arXiv:1304.3494



☑ treatment of deuterium

fixed target DIS experiments provide
important high-x constraints but done
not on proton but deuterium



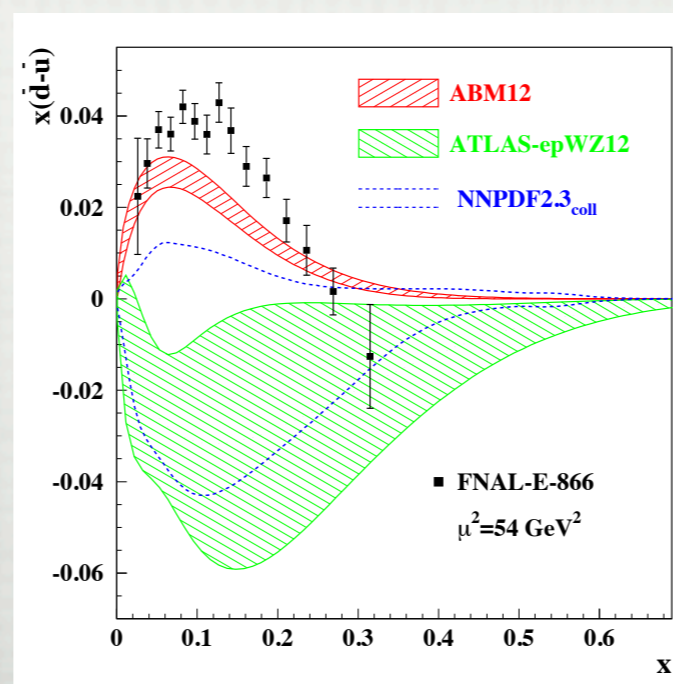
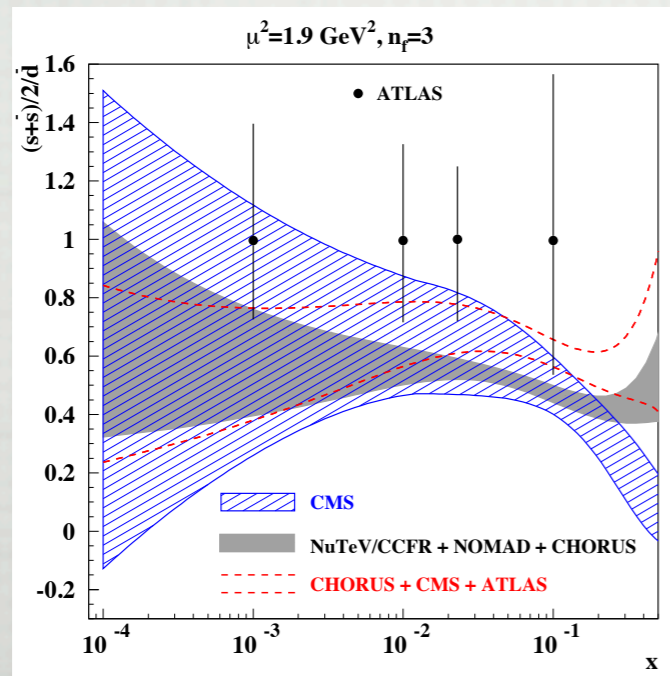
Misc topics in PDF

☑ strange quark PDF

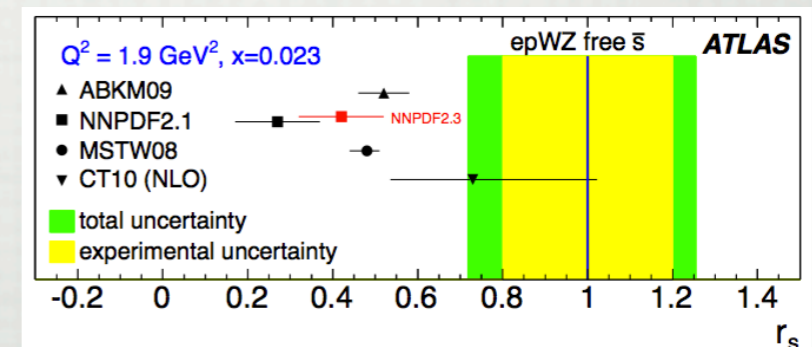
- not many data constrain strange quark PDF - the best pre-LHC constraint comes from neutrino DIS and di-muon subset of DIS (CHORUS, NuTeV, NOMAD)
- new measurements of $W+c$ production from ATLAS & CMS provide new constraints
- with no handle on the strange quark - assumptions were introduced which tied strange PDF to other sea quarks - now these assumptions can be tested

$$s = \bar{s} = \frac{\kappa}{2} (\bar{u} + \bar{d}) \quad r_s = \frac{1}{2} (s + \bar{s}) / \bar{d}$$

arXiv:1404.6469



arXiv:1203.4051 [hep-ex]



Current PDFs

- Parton Distribution Functions - summary of existing PDFs

	DATA	Order	HQ	α_s	Params.	Uncert.
CT14	global	LO,NLO, NNLO	GM-VFNS (s-ACOT)	external	6 indep. PDFs (26 params)	Hessian ($\Delta\chi^2\sim 100$)
MSTW08	global	LO,NLO, NNLO	GM-VFNS (TR)	fit	7 indep. PDFs (20 params)	Hessian ($\Delta\chi^2\sim 25$)
NNPDF	global	LO,NLO, NNLO	GM-VFNS (FONLL)	external	7 indep. PDFs (259 params)	Monte Carlo
CJ12	global	LO,NLO	ZM-VFNS	external	5 indep. PDFs (22 params)	Hessian ($\Delta\chi^2=100$)
HERApdf	DIS (HERA)	NLO NNLO	GM-VFNS (TR)	external	5 indep. PDFs (14 params)	Hessian ($\Delta\chi^2=1$)
ABM11	DIS+DY	NLO NNLO	FFN	fit	6 indep. PDFs (25 params)	Hessian ($\Delta\chi^2=1$)

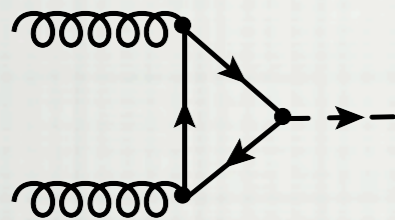
Current PDFs

- *Parton Distribution Functions in Higgs production*

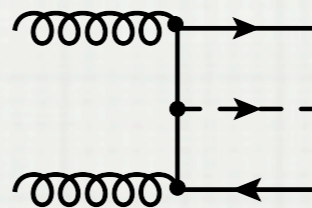
- Higgs is pre-dominantly produced through gluon fusion - gluon PDFs at $x=M_H/\sqrt{s} \sim 0.02$ are crucial

- sub-leading Higgs production via VBF is sensitive to quark & anti-quark PDFs

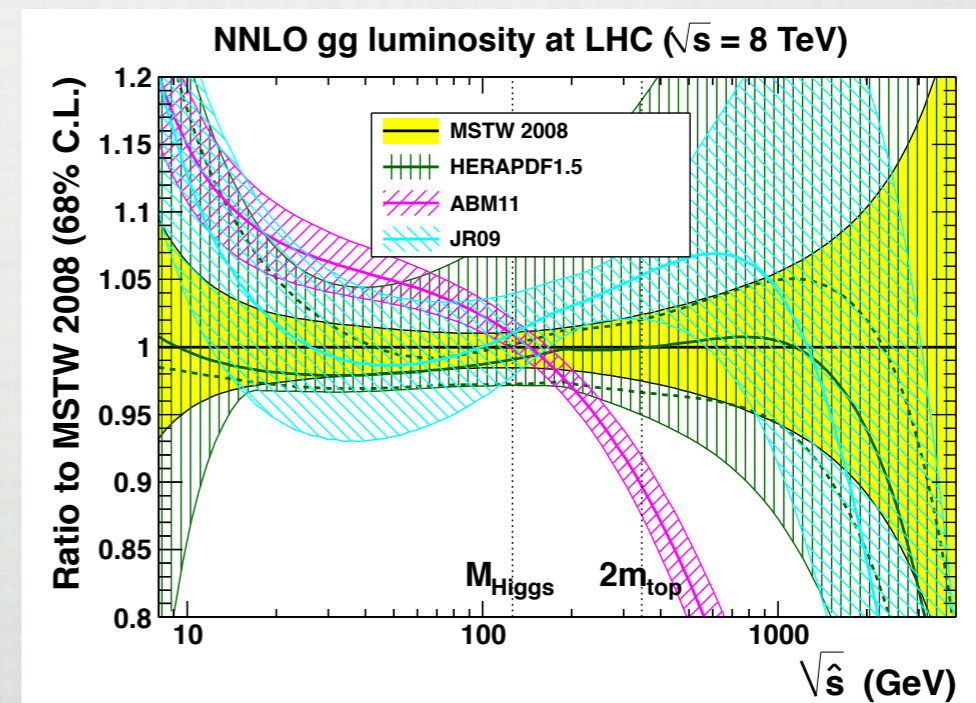
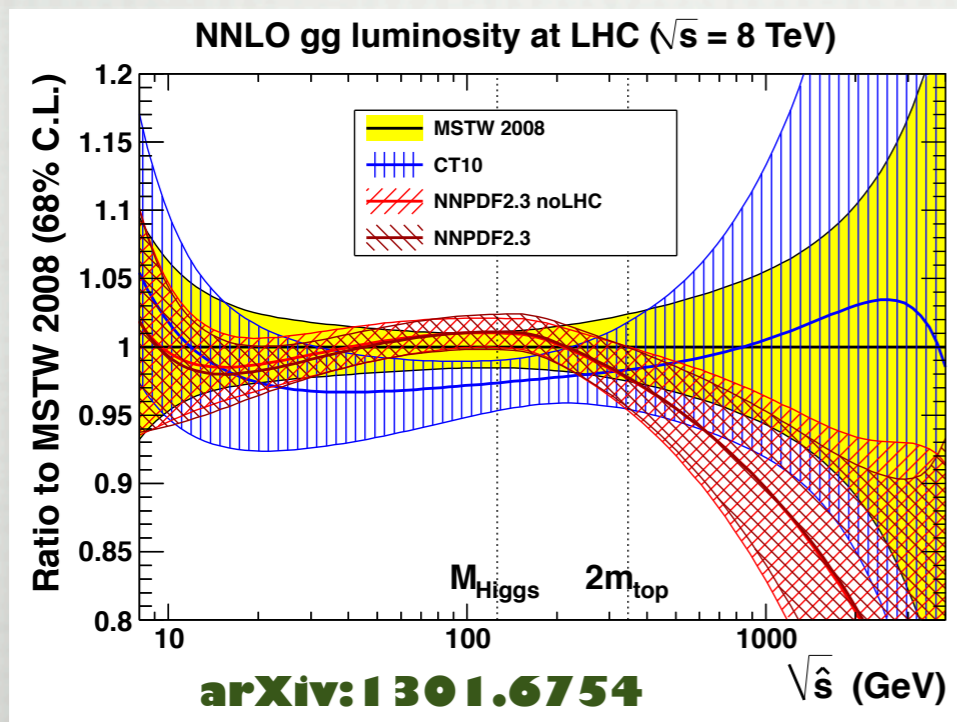
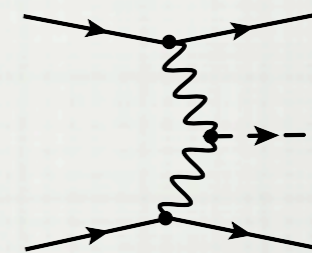
GLUON FUSION



ASSOCIATED PRODUCTION



VECTOR-BOSON FUSION

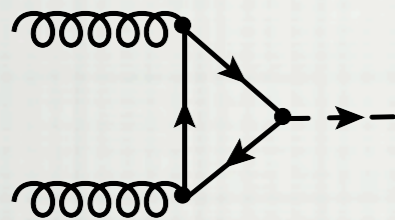


Current PDFs

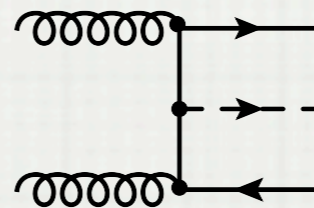
Parton Distribution Functions in Higgs production

- Higgs is pre-dominantly produced through gluon fusion - gluon PDFs at $x=M_H/\sqrt{s} \sim 0.02$ are crucial
- sub-leading Higgs production via VBF is sensitive to quark & anti-quark PDFs

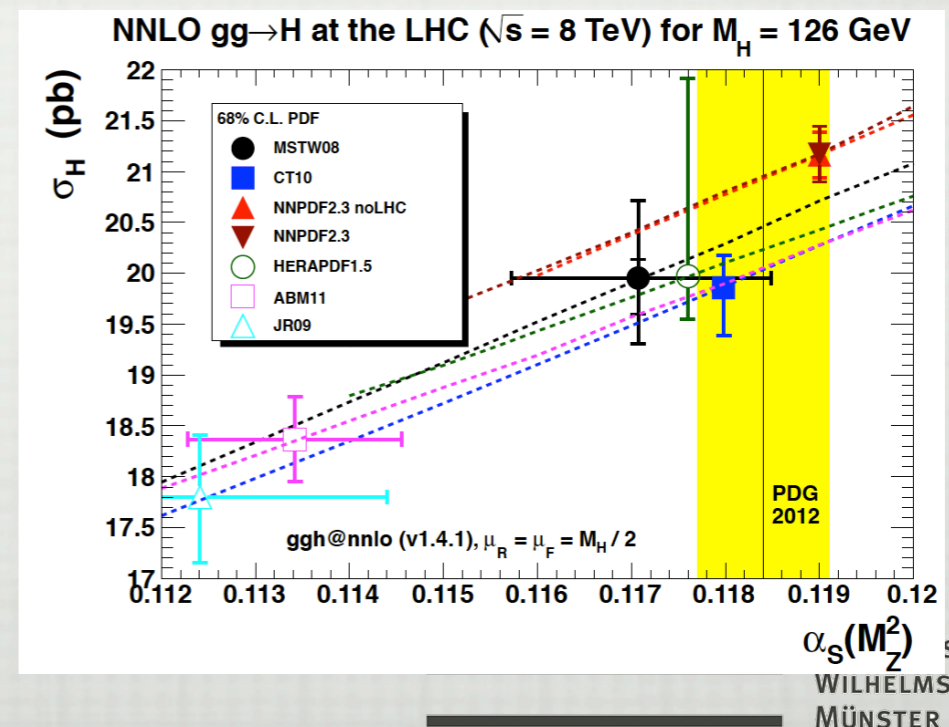
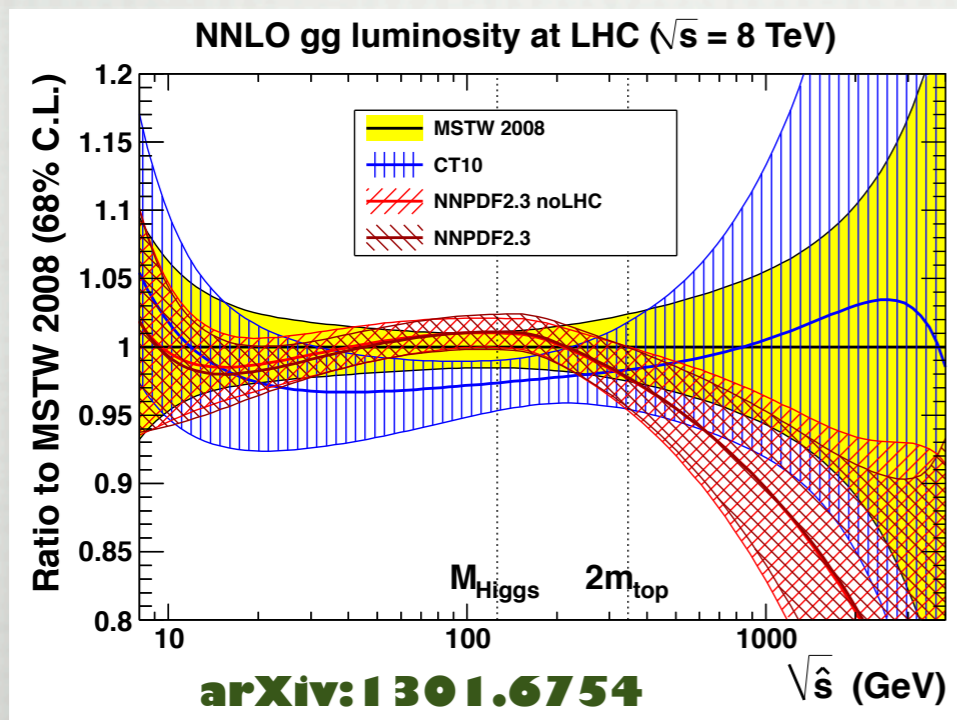
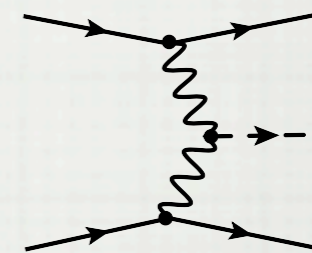
GLUON FUSION



ASSOCIATED PRODUCTION



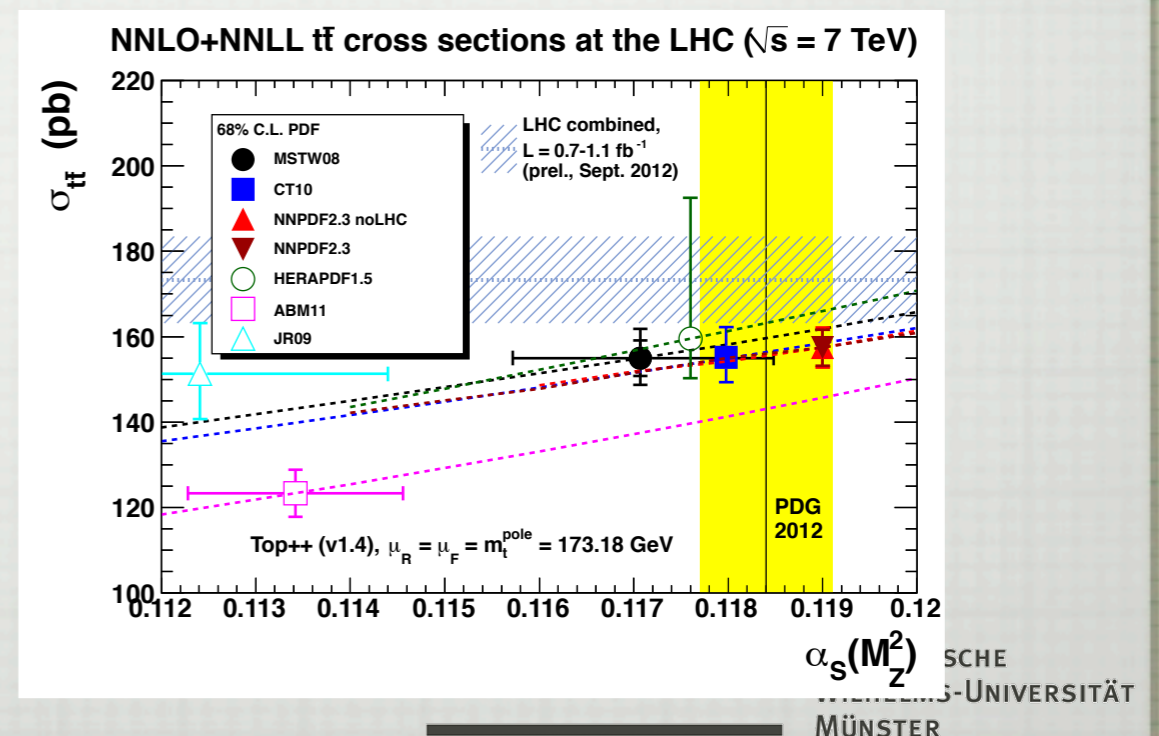
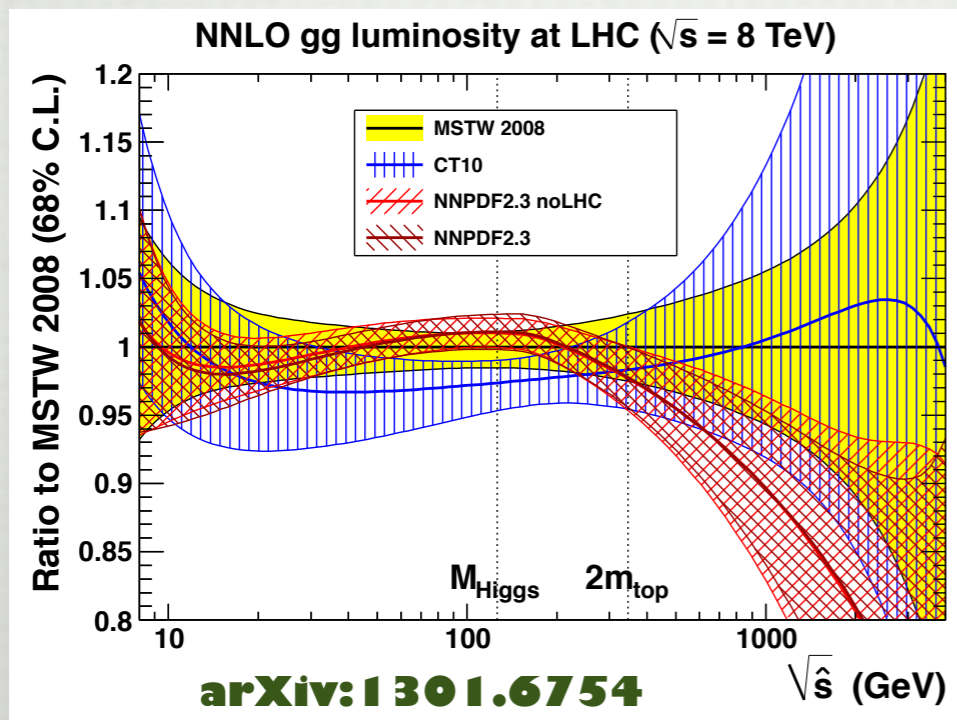
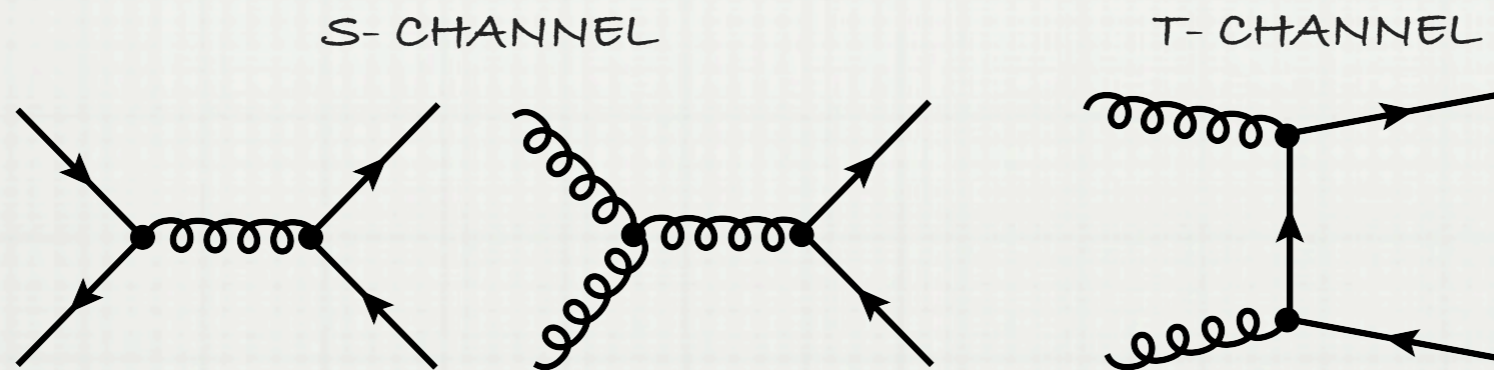
VECTOR-BOSON FUSION



Current PDFs

- Parton Distribution Functions in top quark pair production

- Top quark pair production is dominated by s-channel diagrams where valence quarks & gluons are important at $x=2m_t/\sqrt{s} \sim 0.05$

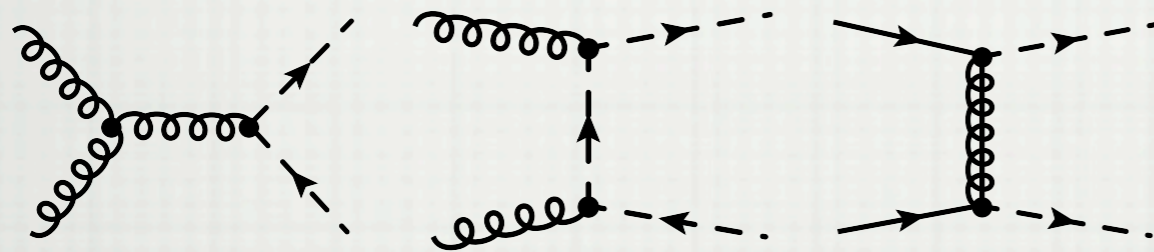


Current PDFs

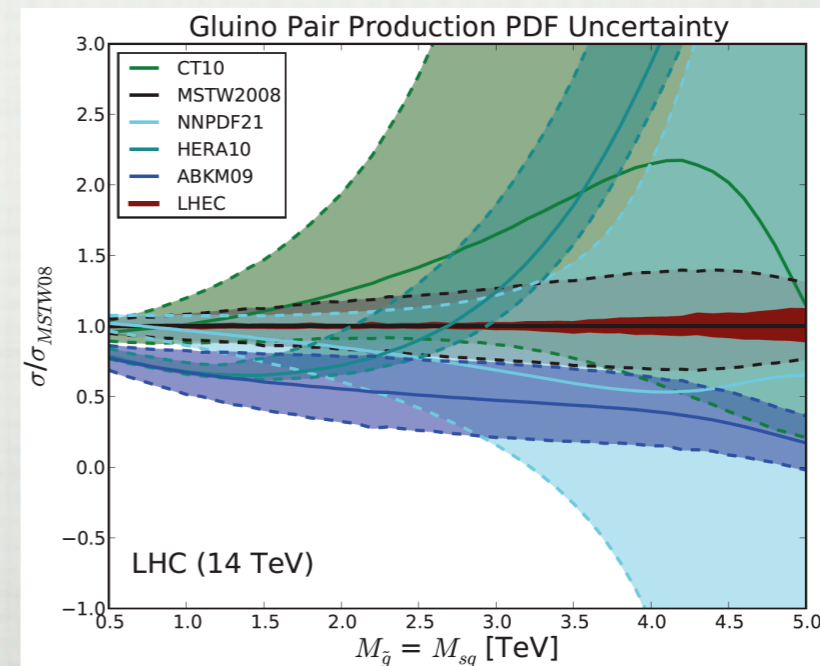
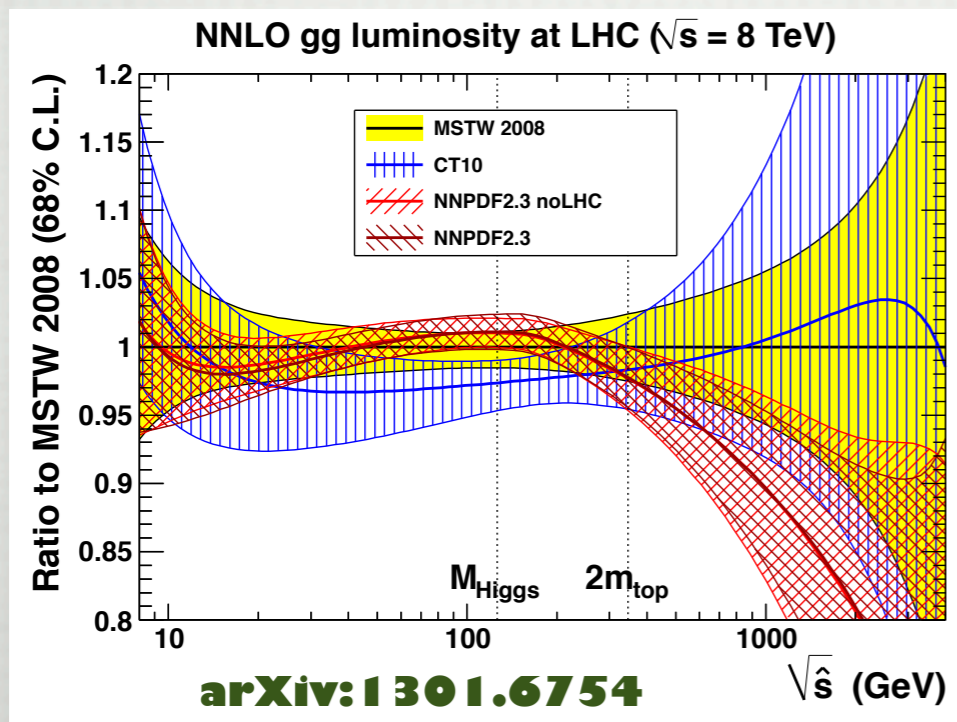
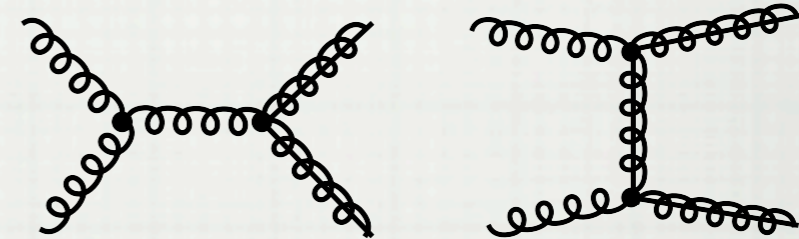
- Parton Distribution Functions in SUSY production

- production of SUSY coloured particles (squarks & gluinos) very sensitive to gluon PDF at very high $x=2m_X/\sqrt{s} \sim 0.2-0.7$ ← very problematic

SQUARK PRODUCTION



GLUINO PRODUCTION



arXiv:1206.2913

PDF after LHC

● Parton Distribution Functions - new dedicated data

- new projects with large possible impact on PDFs

LHeC

colliding electrons / positrons with LHC protons / nuclei

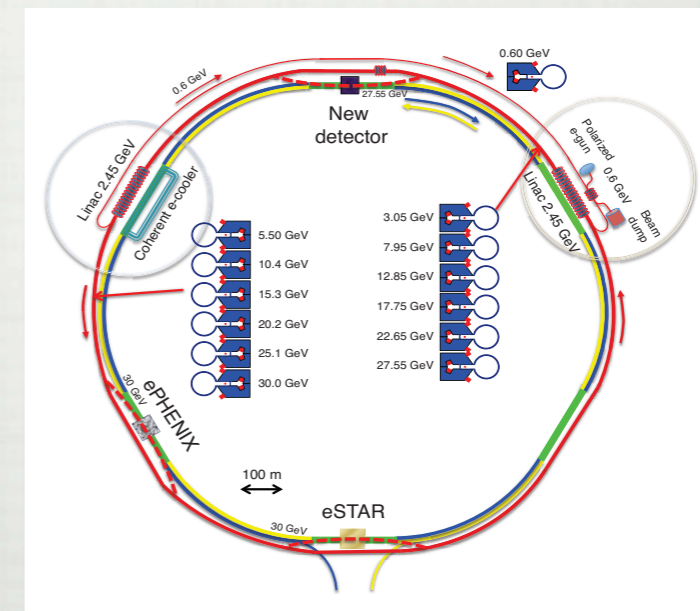
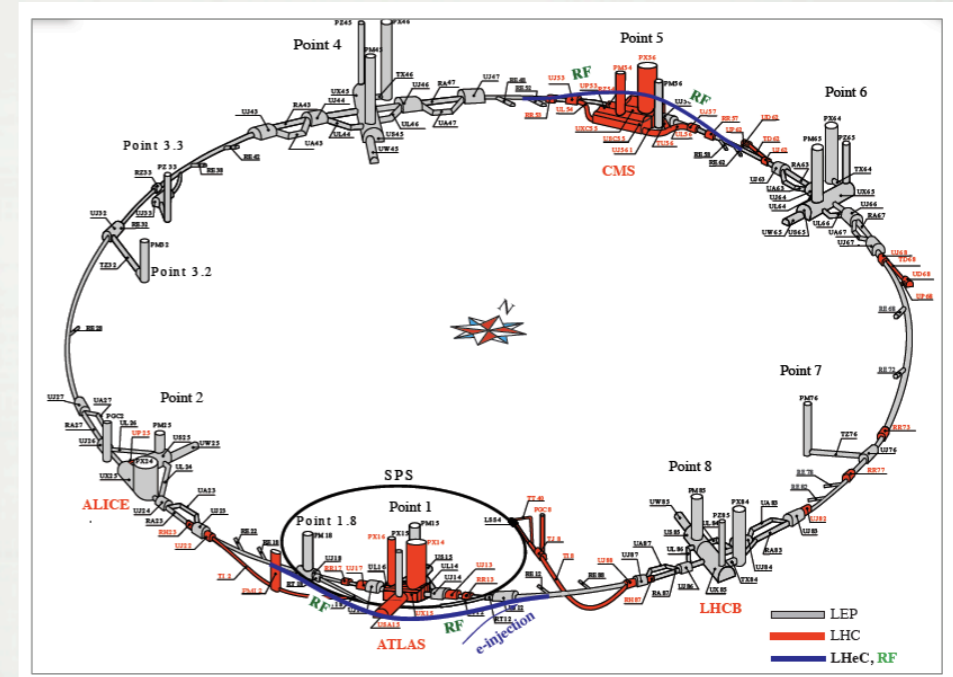
- unprecedented coverage in x - Q^2 plane
- precise determination of the gluon PDF
- interesting also for Higgs & BSM physics programs
- breakthrough machine for nuclear PDFs

EIC

electron ion collider

- high-intensity precision machine with polarized beams
- good coverage in x - Q^2 plane (down to $x \sim 10^{-4}$)
- precise determination of the gluon PDF
- breakthrough machine for nuclear PDFs, saturation, polarized PDFs...

arXiv:1206.2913



arXiv:1212.1701

THANK YOU