THE HIGGS BOSON & BEYOND*

Tao Han Univ. of Pittsburgh June 27, 2018 CTEQ Summer School

Univ. of Puerto Rico - Mayaguez





* Beyond the SM & Beyond the LHC



WHAT WE KNOW NOW?

WHAT IT TELLS US?

WHAT ELSE WE WOULD LIKE TO KNOW?

HOW TO PROCEED FROM HERE?



2013 Nobel Laureate

© The Nobel Foundation. Photo: Lovisa Engblom.

François Englert and Peter W. Higgs "for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"



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The Higgs mechanism (1964)







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The Standard Model (1960-1967, 1972, 1973)



John ELLIS, Mary K. GAILLARD * and D.V. NANOPOULOS ** CERN, Geneva

Received 7 November 1975

A discussion is given of the production, decay and observability of the scalar Higgs boson H expected in gauge theories of the weak and electromagnetic interactions such as the Weinberg-Salam model. After reviewing previous experimental limits on the mass of the Higgs boson, we give a speculative cosmological argument for a small mass. If its mass is similar to that of the pion, the Higgs boson may be visible in the reactions $\pi^-p \rightarrow Hn$ or $\gamma p \rightarrow Hp$ near threshold. If its mass is $\leq 300 \text{ MeV}$, the Higgs boson may be present in the decays of kaons with a branching ratio $O(10^{-7})$, or in the decays of one of the new particles: $3.7 \rightarrow 3.1 + H$ with a branching ratio $O(10^{-4})$. If its mass is $\leq 4 \text{ GeV}$, the Higgs boson may be visible in the reaction $pp \rightarrow H + X$, $H \rightarrow \mu^+\mu^-$. If the Higgs boson has a mass $\leq 2m_{\mu}$, the decays $H \rightarrow e^+e^-$ and $H \rightarrow \gamma\gamma$ dominate, and the lifetime is $O(6 \times 10^{-4} \text{ to} 2 \times 10^{-12})$ seconds. As thresholds for heavier particles (pions, strange particles, new particles) are crossed, decays into them become dominant, and the lifetime decreases rapidly to $O(10^{-20})$ sec for a Higgs boson of mass 10 GeV. Decay branching ratios in principle enable the quark masses to be determined.

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The "EHLQ" (80's)

Supercollider Physics

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The "EHLQ" (80's)





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The Higgs Hunter's Guide

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Sakurai Prize 2017



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Collider

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Moriond 2018 on Higgs: (D. Sperka, ATLAS & CMS) Four production channels with sensitivities; Five decay channels observed; Fermionic & bosonic couplings verified:



Measured mass accuracy < 0.2% : $m_H = 125.26 \pm 0.20(\text{stat}) \pm 0.08(\text{syst}) \text{ GeV}$

June 4, 2018: CERN new release, ATLAS & CMS





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All indications: SM-like Higgs boson, "elementary" at a scale $\Lambda < O(1 \text{ TeV})$ WHAT (ELSE) WE KNOW $m_{\rm H} = 125 \text{ GeV}$ In the SM, the EWSB is parameterized as $V(|\Phi|) = -\mu^2 \Phi^{\dagger} \Phi + \lambda (\Phi^{\dagger} \Phi)^2$ $\Rightarrow \mu^2 H^2 + \lambda v H^3 + \frac{\lambda}{4} H^4$ WHAT (ELSE) WE KNOW $m_{\rm H} = 125 \text{ GeV}$ In the SM, the EWSB is parameterized as $V(|\Phi|) = -\mu^2 \Phi^{\dagger} \Phi + \lambda (\Phi^{\dagger} \Phi)^2$ $\Rightarrow \mu^2 H^2 + \lambda v H^3 + \frac{\lambda}{4} H^4$

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Completion of the SM: The 1st time! A perturbative, renormalizable theory, valid up to a high scale of TeV? ..., M_{Pl} ?

--- Albert Michelson (1894)

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Will History repeat itself (soon)?

NEW ERA: UNDER THE HIGGS LAMP POST



NEW ERA: Under the Higgs lamp post

The "Observation" papers: Now 7,250 cites each! 8,350



Vast scope of topics, from interpretations, explorations in & beyond the SM; applications in astronomy, cosmology, CC; strings/branes, to "Philosophical Perspectives"

A REMINDER: A. The Higgs mechanism ≠ a Higgs boson !

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A REMINDER: A. The Higgs mechanism \neq a Higgs boson ! From theoretical point of view, **3 Nambu-Goldstone bosons were all we need!** A non-linear realization of the gauge symmetry: $U = \exp\{i\omega^{i}\}^{i}/v\}, \quad D_{\mu}U = \partial_{\mu}U + igW_{\mu}^{i}\frac{\tau^{i}}{2}U - ig'UB_{\mu}\frac{\tau^{3}}{2}$ $\mathcal{L} = \frac{v^{2}}{2}[D^{\mu}U^{\dagger}D_{\mu}U] \to \frac{v^{2}}{4}(\sum_{i}g^{2}W_{i}^{2} + g'^{2}B^{2})$
A REMINDER: A. The Higgs mechanism ≠ a Higgs boson ! From theoretical point of view, **3 Nambu-Goldstone bosons were all we need!** A non-linear realization of the gauge symmetry: $U = \exp\{i\omega^{i}\}^{i}/v\}, \quad D_{\mu}U = \partial_{\mu}U + igW_{\mu}^{i}\frac{\tau^{i}}{2}U - ig'UB_{\mu}\frac{\tau^{3}}{2}$ $\mathcal{L} = \frac{v^{2}}{2}[D^{\mu}U^{\dagger}D_{\mu}U] \to \frac{v^{2}}{4}(\sum g^{2}W_{i}^{2} + g'^{2}B^{2})$ The theory is valid to a unitarity bound ~ 2 TeV The existence of a light, weakly coupled Higgs boson carries important message for our understanding & theoretical formulation in & beyond the SM – **Ultra-Violet completion / renormalizibility**

WHAT IT TELLS US 1. $V = -\mu^2 |\phi|^2 + \sqrt[3]{\phi} |^4$

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 (due to the loop suppression by design) Already possess challenge to BSM theories.

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The SM can be a consistent perturbative theory up to M_{pl} ! allowing M_N , M_{GUT} , ... Bezrukov et al., arXiv:1205.2893.



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600

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If $\Lambda^2 \gg m_H^2$, then unnaturally large cancellations must occur. Cancelation in perspective:

$$\begin{split} m_{\rm H}^{\ \ 2} &= 36,127,890,984,789,307,394,520,932,878,928,933,023 \\ &\quad -36,127,890,984,789,307,394,520,932,878,928,917,398 \\ &= (125 \ {\rm GeV})^2 \ ! \ ? \end{split}$$

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-36,127,890,984,789,307,394,520,932,878,928,917,398 = (125 GeV)²! ? "Naturalness" -> TeV scale new physics.





Unbelievable! 4 mm² / 20 cm² ~ 10⁻³ fine-tune.



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Weak-scale Supersymmetry:

Extended symmetry between opposite spin & statistics

"s"-particles	symbol	spin	SUSY breaking mass param.
gluino	$ ilde{g}$	1/2	M_3
charginos	$ ilde{\chi}_1^\pm$, $ ilde{\chi}_2^\pm$	1/2	M_2
neutralinos	$ ilde{\chi}_{1}^{0}, \ ilde{\chi}_{2}^{0}, \ ilde{\chi}_{3}^{0}, \ ilde{\chi}_{4}^{0}$	1/2	$M_1, \ \mu, \ B$
			$m_{H_{\pi}}^2, \; m_{H_{d}}^2$
sleptons	$ ilde{e}_L$, $ ilde{ u}_{e_L}$, $ ilde{e}_R$	0	$m_{\ell L}^2$
	$ ilde{\mu}_L$, $ ilde{ u}_{\mu_L}$, $ ilde{\mu}_R$	0	
	$ ilde{ au}_1, ilde{ au}_2,\; ilde{ u}_{ au_L}$	0	$m_{\ell R}^2$
squarks	$ ilde{u}_L$, $ ilde{d}_L$, $ ilde{u}_R$, $ ilde{d}_R$	0	$m_{q_L}^2$
	$\tilde{c}_L, \tilde{s}_L, \ \tilde{c}_R, \tilde{s}_R$	0	
	\tilde{t}_1 , \tilde{t}_2 , \tilde{b}_1 , \tilde{b}_2	0	$m_{q_R}^2$
Higgs	h^0, H^0, A^0, H^{\pm}	0	m_A^2 , tan eta

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	$\tilde{c}_L, \tilde{s}_L, \ \tilde{c}_R, \tilde{s}_R$	0	
	$ ilde{t}_1$, $ ilde{t}_2$, $ ilde{b}_1$, $ ilde{b}_2$	0	$m_{q_R}^2$
Higgs	h^0, H^0, A^0, H^{\pm}	0	$m_A^2, aneta$

Natural cancellations: \tilde{t} versus t \tilde{W} versus W \tilde{H} versus H H_d versus H_u , $\Delta m_u^2 \sim (M_{excurr}^2 - M_{excurr}^2) - \frac{\lambda_f^2}{f} \ln \left(-\Lambda \right)$

Weak scale SUSY is natural if
$$M_{SUSY} \sim \mathcal{O}(1 \text{ TeV})$$
.

Additional Feature I :

• Unification of forces?

Do the forces E & M/Weak/Strong all unify into a single force ?



Possible if there is a TeV scale new physics threshold !

Additional Feature II :

Natural existence of a lightest, neutral fermion: The "neutralino" as the WIMP (weakly interacting massive particle) cold dark matter candidate!

Additional Feature III :

Natural extension / maximal symmetry in quantum field theory; Elegant mathematical structure; Connection with quantum gravity/string theory.

Dynamical approach for mass generation:

• Technicolor: A lesson from QCD $SU(N_{TC})$ gauge theory, TC fermions Q = U, D, ...EWSB by TC-fermion condendation at Λ_{TC} : $v \sim \langle \overline{Q_L} Q_R \rangle^{1/3} \sim 246$ GeV.

 $\begin{array}{l} \checkmark \text{ no elementary scalar, like Higgs.} \\ \checkmark \text{ theory natural: } \Lambda_{TC} \text{ dynamical.} \\ \checkmark \text{ predicts new strong dynamics at the TeV scale: } \pi_T, \ \eta_T, \ \rho_T, \ \omega_T... \end{array}$

Dynamical approach for mass generation:

• Technicolor: A lesson from QCD $SU(N_{TC})$ gauge theory, TC fermions Q = U, D, ...EWSB by TC-fermion condendation at Λ_{TC} : $v \sim \langle \overline{Q_L} Q_R \rangle^{1/3} \sim 246$ GeV.

√ no elementary scalar, like Higgs. √ theory natural: Λ_{TC} dynamical. √ predicts new strong dynamics at the TeV scale: π_T , η_T , ρ_T , ω_T ...

• Extended Technicolor:* Fermion mass generation G_{ETC} gauge theory, ETC fermions: U, D, ..., u, d...After intrgrating out ETC gauge bosons at the scale Λ_{ETC} , with TC-fermion condensate,SM fermion mass generated:

 $m_f \sim \langle \overline{Q_L} Q_R \rangle / \Lambda_{ETC}^2 \sim \Lambda_{TC}^3 / \Lambda_{ETC}^2.$

Not supported by observation

• "Warped" Extra-dimension Scenario: The Randall-Sundrum model In a 5-dim space, Randall and Sundrum found a static solution of the form:*

 $ds^2 \sim e^{-2ky} \eta_{\mu\nu} dx^{\mu} dx^{\nu} - dy^2,$

where k is the curvature scale in the 5th-dim.

The extra dimension y is "warped".



*L. Randall, R. Sundrum.

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Dual description to strong dynamics/composite theory!

Little Higgs Models: A less ambitious approach

keep the Higgs boson "naturally" light (at 1-loop level).

Higgs is a pseudo-Goldstone boson from global symmetry breaking (at scale 4πf)[‡]
Higgs acquires a mass radiatively at the EW scale v, by collective explicit breaking
Consequently, quadratic divergences absent at one-loop level[§]

 $W, Z, B \leftrightarrow W_H, Z_H, B_H; \quad t \leftrightarrow T; \quad H \leftrightarrow \Phi.$ (cancellation among same spin states!)









The "Little hierarchy": Tension with the observation – • In SUSY, $m_{H}^{2} \approx M_{Z}^{2} \cos^{2}2\beta + \Delta m_{SUSY}^{2}$ Tree-level <(80 GeV)² + loop-level: >(45 GeV)²



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Barbieri, Giudice, 1988 Kitano et al, 2005 Giudice, 2007 Feng, 2013

Draper, Shih, Meade, Reece, 2011 Hall, Pinner, Ruderman, 2012 Carena et al., 2012, 2013 S. Heinemeyer et al., 2012-2014 The "Little hierarchy": Tension with the observation –

• In composite/strong dynamics: (dual of extra dimension theory)

The "Little hierarchy": Tension with the observation – In composite/strong dynamics: (dual of extra dimension theory) The Higgs boson as a pseudo-Goldstone boson:

$$m_H^2 \sim \frac{f^2}{(4\pi)^2} \sim \frac{m_t^2 M_T^2}{f^2}.$$

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Pomarol, ICHEP'12

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Both SUSY/Compositeness suffer from some degree of "fine-tune": < 1%.

Too heavy to be light; too light to be heavy!

WHAT WE WISH TO KNOW IN THE LHC ERA 1. THE NATURE OF EWSB

 $V(|\Phi|) = -\mu^2 \Phi^{\dagger} \Phi + \lambda (\Phi^{\dagger} \Phi)^2$ $\Rightarrow \quad \mu^2 H^2 + \lambda v H^3 + \frac{\lambda}{4} H^4$

WHAT WE WISH TO KNOW IN THE LHC ERA **1. THE NATURE OF EWSB** $V(|\Phi|) = -\mu^2 \Phi^{\dagger} \Phi + \lambda (\Phi^{\dagger} \Phi)^2$ $\Rightarrow \mu^2 H^2 + \lambda v H^3 + \frac{\lambda}{4} H^4$ Fully determined at the weak scale: $v = (\sqrt{2}G_F)^{-1/2} \approx 246 \text{ GeV} \quad m_H \approx 126 \text{ GeV}$ $m_H^2 = 2\mu^2 = 2\lambda v^2 \quad \Rightarrow \quad \mu \approx 89 \text{ GeV}, \quad \lambda \approx \frac{1}{8}.$

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WHAT WE WISH TO KNOW IN THE LHC ERA **1. THE NATURE OF EWSB** $V(|\Phi|) = -\mu^2 \Phi^{\dagger} \Phi + \lambda (\Phi^{\dagger} \Phi)^2$ $\Rightarrow \mu^2 H^2 + (\lambda v H^3) + \frac{\lambda}{4} H^4$ Fully determined at the weak scale: $v = (\sqrt{2}G_F)^{-1/2} \approx 246 \text{ GeV} \quad m_H \approx 126 \text{ GeV}$ $m_H^2 = 2\mu^2 = 2\lambda v^2 \quad \Rightarrow \quad \mu \approx 89 \text{ GeV}, \quad \lambda \approx \frac{1}{8}.$ O(1) deviation on λ_{hhh} could make EW All we know: phase transition strong 1st order!

X.M.Zhang (1993); C. Grojean et al. (2005)

"Natural SUSY":

Cohen, Kaplan, Nelson, 1996 Hall, Pinner, Ruderman, 2012 Baer, Barger, Huang, Tata, 2012



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Relevant to the Higgs and the "Most Wanted": $\tilde{H}^{0,\pm}, \tilde{t}, \tilde{b}, (\tilde{g}); S, \tilde{S}...$



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• "Compositeness": the T', current ATLAS limit: $M_{\rm T}$ > 480 GeV, for $M_{\rm A}$ < 100 GeV.

3. EXTENDED HIGGS SECTOR? The Higgs boson should have not only relatives: $\tilde{t}, \ \tilde{b}, \ \tilde{H}^{\pm,0}; \ T',$ But also siblings: $H_i^0, A_i^0, H^{\pm}, H^{\pm\pm}, \dots$ Haber, 2012 Branco, Ferreira, Rebelo, Sher, Silva, arXiv:1106.0034; • Two Higgs Doublet Model (2HDM): Coleppa, Kling, Su, arXiv:1305.0002. rich phenomenology, Type II SUSY option ... Ellwanger, Gunion et al., 2012 S. King et al., 2012 • Plus a singlet: R. Barbieri et al., 2013, NMSSM, solve the µ-problem, relax fine-tune, light DM... • Triplet Model: m, L-R symmetric theories, Little Higgs ...

neutrino mass connection via Type II seesaw.

• Type II Two Higgs Doublet Model

$$H_u = \begin{pmatrix} H_u^+ \\ H_u^0 \end{pmatrix} \xrightarrow{} v_u / \sqrt{2} \qquad H_d = \begin{pmatrix} H_d^0 \\ H_d^- \end{pmatrix} \xrightarrow{} v_d / \sqrt{2}$$
$$v_u^2 + v_d^2 = v^2 = (246 \text{GeV})^2 \qquad \tan \beta = v_u / v_d$$

after EWSB 5 physical Higgses CP-even Higgses: h⁰, H⁰ CP-odd Higgs: A⁰ Charged Higgses: H[±]

\odot tree level masses determined by mA, tan β

$$m_{h^0,H^0}^2 = \frac{1}{2} \left((m_A^2 + m_Z^2) \mp \sqrt{(m_A^2 - m_Z^2)^2 + 4m_A^2 m_Z^2 \sin^2 2\beta} \right)$$
$$m_{H^{\pm}}^2 = m_A^2 + m_W^2, \quad \cos^2(\beta - \alpha) = \frac{m_{h^0}^2 (m_Z^2 - m_{h^0}^2)}{m_A^2 (m_{H^0}^2 - m_{h^0}^2)}.$$

 $m_{h^0} \approx \min\{m_A, m_Z\} |\cos 2\beta|, \quad m_{H^0} \approx \max\{m_A, m_Z\}, \quad m_{H^{\pm}} \approx \max\{m_A, m_W\}.$

Tao Han

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Example: MSSM Two Higgs-Doublet Model after the discovery: Arbey et al., 2011,

5 Higgs bosons: h^0 , H^0 , A^0 , H^{\pm}

Arbey et al., 2011, 2012 Baer et al., 2012 Heinemeyer et al., 2012 Carena 2012, 2013,

Tree-level masses given by M_A , $\tan \beta$ Collider bounds:



4. THE HIGGS PORTALS TO COSMOS? (a). Dark Matter $H^{\dagger}H$ is the only bi-linear SM gauge singlet.

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 $k_s H^{\dagger} H S^* S, \quad \frac{k_{\chi}}{\Lambda} H^{\dagger} H \bar{\chi} \chi.$

4. THE HIGGS PORTALS TO COSMOS?

(a). Dark Matter

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 $k_s H^{\dagger} H S^* S, \quad \frac{k_{\chi}}{\Lambda} H^{\dagger} H \bar{\chi} \chi.$ Missing energy at LHC













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TH, Z.Liu, A.Natarajan, arXiv:1303.3040





100 m_{χ1} (GeV)

10-5

TH, Z.Liu, A.Natarajan, arXiv:1303.3040



1000



Fowlie, Roszkowsky et al., arXiv:1306.1567

(b). Baryon – anti-baryon Asymmetry For $M_H = 125 \text{ GeV}$, EW baryogenesis needs light sparticles: $m_{stop} \approx 150 \text{ GeV}$ plus a light neutralino, singlets ... Chung et al., 2011; Chung et al., 2011;

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The existence of a fundamental scalar encourages the consideration of scalar fields in cosmological applications.

5. FLAVOR & YUKAWA COUPLINGS

- Particle mass hierarchy
- Patterns of quark, neutrino mixings
- New CP-violation sources?

Higgs Yukawa couplings as the pivot!



The fermion mass/mixing is a muchⁿ bigger puzzle! What controls the mixing structure: "Minimal Flavor Violation" for BSM? The fermion mass/mixing is a muchⁿ bigger puzzle!

What controls the mixing structure: "Minimal Flavor Violation" for BSM? The b rare decays are pushing the limits: b -> s γ, Bs -> μ⁺ μ⁻ BR(Bs) ~ tan⁶β / M⁴_A

> TH, Liu, arXiv:1303.3040 Carena et al., arXiv:1305.5761.



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Most recent LHCb+CMS: arXiv:1411.4413

B(B ⁰ s	$\mu^+\mu^-)$	=	$2.8^{+0.7}_{-0.6}$	×	10^{-9} and
B(B ⁰	$\mu^+\mu^-)$	=	3.9 ^{+1.6} _{-1.4}	×	10 ^{− 10} ,
$S^{B^0_s}_{SM} =$	0.76^{+0}_{-0}).20).18	and $S^{B^0}_{SM}$	=	$3.7^{+1.6}_{-1.4}$.

⁴⁰ ⁵⁰ With Belle 2 as well, likely a surprise to breakthrough! The Higgs as pivot for "seesaw": $m_{\nu} \sim \kappa \frac{\langle H^0 \rangle^2}{M}$

The Higgs as pivot for "seesaw": $m_{\nu} \sim \kappa \frac{\langle H^0 \rangle^2}{M}$ Type I seesaw: $M = M_{N_i}$ right-handed (sterile) N_R^{i} $H \rightarrow NN, N \rightarrow H\nu, ...$ Yanagida; Ramond et al.; Mohapatra ... The Higgs as pivot for "seesaw": $m_{\nu} \sim \kappa \frac{\langle H^0 \rangle^2}{M}$ Type I seesaw: $M = M_{N,}$ right-handed (sterile) N_R^i $H \rightarrow NN, N \rightarrow H\nu, ...$ Yanagida; Ramond et al.; Mohapatra ... Type II seesaw: $M = M_{H++}$, a Higgs triplet Φ_3 $H^{++} \rightarrow l_i^+ l_j^+$ Mohapatra, Senjanovic, ...




Sype III seesaw: $M = M_T$, a termionic triplet T_3 : $T^+ \rightarrow H l_i$, $T^0 \rightarrow W^{\pm} l$ Senjanovic et al., arXiv:0904.2309.

The Higgs as pivot for "seesaw": $m_{\nu} \sim \kappa \frac{\langle H^0 \rangle^2}{M}$ Type I seesaw: $M = M_N$ right-handed (sterile) N_R^i Yanagida; Ramond et al.; Mohapatra ... $H \rightarrow NN, N \rightarrow Hv, \dots$ Type II seesaw: $M = M_{H++}$, a Higgs triplet $\Phi_3 H^{++} \rightarrow l_i^+ l_i^+$ IH 1.000 E Mohapatra, Senjanovic, ... NH^{1.000} $\mu\mu$ $\tau\tau$ $\mu\mu$ Fileviez-Perez et al., 2008. 0.100 0.100 Chaudhuri, Grimus, $H^{++} = \sum_{0,010}^{10} [l^{+}_{i} l^{+}_{j}]$ Mukhopadyaya, arXiv:1305.5761 BR Chun et al., arXiv:1305.0329 ee 0.010 0.001 m1 (eV) $m_3 (eV)$ Type III seesaw: $M = M_T$, a fermionic triplet T_3 : $T^{+} \rightarrow H l_{+}, T^{0} \rightarrow W^{\pm} l$ Senjanovic et al., arXiv:0904.2309. Watch out: $H^0 \rightarrow \mu \tau (l_i l_j)$ for BSM flavor physics! TH, Marfatia, PRL (2001) 28 Harnik, Kopp, Zupan, 2013

6. COUPLINGS & WIDTH

Higgs boson couplings encode its properties:



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In a pessimistic scenario, the LHC does not see a new particle associated with the Higgs sector, then the effects of a heavy state on Higgs coupling g_i at the scale M:

 $\Delta_i \equiv \frac{g_i}{g_{SM}} - 1 \sim \mathcal{O}(v^2/M^2) \approx \text{a few \% for } \mathbf{M} \approx 1 \text{ TeV}$

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Higgs coupling deviations:Δ:VVHbbH,ττHggH,γγH

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LHC 14 TeV, 3ab⁻¹: 8% 15% few% 50%

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If no deviations, I'd DEFINE it THE SM Higgs! Higgs coupling deviations:

LHC 14 TeV, 3ab⁻¹: 8% 15% few% 50%











Particle physicists around the world are designing colliders that are much larger in size than the Large Hadron Collider at CERN, Europe's particle-physics laboratory. Nature News, July '14 LHC Leads the Way (2015-2030)



Table 1-1. Proposed running periods and integrated luminosities at each of the center-of-mass energies for each facility. Snowmass 1310.8361

Facility	HL-LHC	ILC	ILC(LumiUp)	CLIC	TLEP (4 IPs)	HE-LHC	VLHC
\overline{s} (GeV)	14,000	250/500/1000	250/500/1000	350/1400/3000	240/350	33,000	100,000
$\mathcal{L}dt$ (fb ⁻¹)	3000/expt	250 + 500 + 1000	1150 + 1600 + 2500	500 + 1500 + 2000	10,000+2600	3000	3000
$dt (10^7 s)$	6	3+3+3	(ILC 3+3+3) + 3+3+3	3.1 + 4 + 3.3	5+5	6	6

HL-LHC: The Energy & Precision Frontier

Sasha Valishev



Levelled Luminosity $L = 5 \times 10^{34} cm^{-2} s^{-1}$

HE-LHC: The New Energy Frontier

Vladimir Shiltsev, this workshop

HE-LHC design goals and basic choices

physics goals:

- 2x LHC collision energy with FCC-hh magnet technology
- c.m. energy = 27 TeV ~ 14 TeV x 16 T/8.33T
- target luminosity $\geq 4 \times HL-LHC$ (cross section $\propto 1/E^2$)

key technologies:

- FCC-hh magnets (curved!) & FCC-hh vacuum system
- HL-LHC crab cavities & electron lenses

beam:

• HL-LHC/LIU parameters (25 ns baseline, also 5 ns option)

Fastest Possible Technical Schedules

M. Benedikt, F. Zimmermann '17



HE-LHC design & construction

technical schedule defined by magnets program and by CE \rightarrow earliest possible physics starting dates:

- HE-LHC: 2040 (with HL-LHC stop at LS5 / 2034)
- Options: FCC-ee @ 2039; FCC-hh @ 2043.

THE NEXT ENERGY FRONTIER: 100 TEV HADRON COLLIDER



Higgs Production @ SPPC



Snowmass QCD Working Group: 1310.5189

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Higgs Self-couplings:



Higgs Self-couplings:



Triple coupling sensitivity:

Test the shape of the Higgs potential, and the fate of the EW-phase transition!

$ \begin{array}{c c c c c c c c c c c c c c c c c c c $		HL-LHC	ILC500	ILC500-up	ILC1000	ILC1000-up	CLIC1400	CLIC3000	HE-LHC	VLHC
$ \frac{\int \mathcal{L}dt \ (\text{fb}^{-1})}{\lambda} = \frac{3000/\text{expt}}{50\%} = \frac{500}{83\%} = \frac{1600^{\ddagger}}{1600^{\ddagger}} = \frac{500+1000}{1600+2500^{\ddagger}} = \frac{1500}{1500} + \frac{2000}{3000} = \frac{3000}{3000} = \frac{3000}{3000} = \frac{3000}{3000} = \frac{1000^{\ddagger}}{10\%} = \frac{1000}{10\%} = \frac{1000}$	$\sqrt{s} \; ({\rm GeV})$	14000	500	500	500/1000	500/1000	1400	3000	33,000	100,000
λ 50% 83% 46% 21% 13% 21% 10% 20% 8%	$\int \mathcal{L}dt \; (\mathrm{fb}^{-1})$	3000/expt	500	1600^{\ddagger}	500 + 1000	$1600 + 2500^{\ddagger}$	1500	+2000	3000	3000
	λ	50%	83%	46%	21%	13%	21%	10%	20%	8%

Snowmass 1310.8361

Adequate to test EW phase transition strong 1st order: $\rightarrow O(1)$ deviation on λ_{hhh}

HE-LHC: Higgs self-coupling

D. Goncalves, TH, F. Kling, T. Plehen, M. Takeuchi, arXiv:1802.04319.



Cross section increases $\sim 3x$. hh \rightarrow bb $\gamma\gamma$: $\delta\lambda \sim 30\%$ @ 2σ

Also, S. Homiller, P. Meade, this workshop.

DM Searches



DM Searches



40

DM Searches



HE-LHC: WIMP Dark Matter

TH, S. Mukhopadhyay, X. Wang, this workshop; M. Low, D. Egana-Ugrinovic, J. Ruderman, L.-T. Wang, to appear.





 2σ @HL-LHC \rightarrow HE-LHC \rightarrow FCC Wino monojet: 300 GeV 700 GeV 1400 GeV Wino disappearing track: 2000 GeV 800 GeV 5000 GeV Higgsino monojet: 500 GeV 900 GeV 200 GeV Higgsino disappearing track: 250 GeV 500 GeV 1100 GeV

e+e- colliders: Energy/Lumi projection



TLEP Report: 1308.6176

e+e- colliders: Energy/Lumi projection



Ecm	running time	statistics (FCC-ee)
	b,c,т	10 ¹¹ b,с,т
90 GeV	1-2 yrs	10 ¹² Z (Tera Z)
160 GeV	1-2 yrs	10 ⁸ - 10 ⁹ WW(Oku W)
240 GeV	4-5 yrs	2x10 ⁶ ZH (Mega H)
350 GeV	4-5 yrs	10 ⁶ tt (Mega top)

Higgs-Factory: Mega (106) Higgs Physics



- Model-independent measurement: ILC Report: 1308.6176 $\Gamma_{\rm H} \sim 6\%$, $\Delta m_{\rm H} \sim 30 \text{ MeV}$ (HL-LHC: assume SM, $\Gamma_{\rm H} \sim 5-8\%$, $\Delta m_{\rm H} \sim 50 \text{ MeV}$)
- TLEP 10⁶ Higgs: $\Gamma_{\rm H} \sim 1\%$, $\Delta m_{\rm H} \sim 5$ MeV. TLEP Report: 1308.6176

Higgs-Factory: Mega (106) Higgs Physics



- Model-independent measurement: ILC Report: 1308.6176 $\Gamma_{\rm H} \sim 6\%$, $\Delta m_{\rm H} \sim 30 \text{ MeV}$ (HL-LHC: assume SM, $\Gamma_{\rm H} \sim 5-8\%$, $\Delta m_{\rm H} \sim 50 \text{ MeV}$)
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Higgs Total Width & Invisibale BR: TH, Z.Liu, J.Sayre, arXiv:1311.7155



Also see, Peskin, arXiv:1312.4974 including ILC luminosity upgrade.
A GRAND PICTURE:



Summary: The Higgs boson is a new class, at a pivotal point of energy, intensity, cosmic frontiers.



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LHC ROCKS!



HE-LHC: Extended SUSY Reach:



Pushing the "Naturalness" limit



WIMP DM: $M_{\rm DM} < 1.8 \,\,{\rm TeV}\left(\frac{g_{\rm eff}^2}{0.3}\right)$



New Particle Searches

Electroweak Resonances: Z',W'

Colored Resonances:



Z-Factory: Tera (1012) Z Physics

TLEP Report: 1308.6176

- Clean environment, $\Delta E_{cm} < 1 \text{ MeV}$, 10^5 x LEP-I
- possible longitudinal polarization
- Precision measurements (statistical):
 - Z-ploe: ΔM_Z , $\Delta \Gamma_Z < 0.1$ MeV, $\Delta \sin^2 \theta_w < 10^{-6}$;

 $\Delta M_W \sim O(1 \text{ MeV}), \Delta m_t \sim O(10 \text{ MeV}), \Delta m_H \sim O(10 \text{ MeV}).$

