

Lecture 4

What do we know so far:

atoms consist of nuclei immersed in a cloud of electrons

particles are small, light, energetic: p, n, e, neutrino

*processes are described by **quantum mechanics** and **relativity***

***electromagnetic interactions** keep atoms together*

***strong interactions** keep nuclei together*

Additional comment:

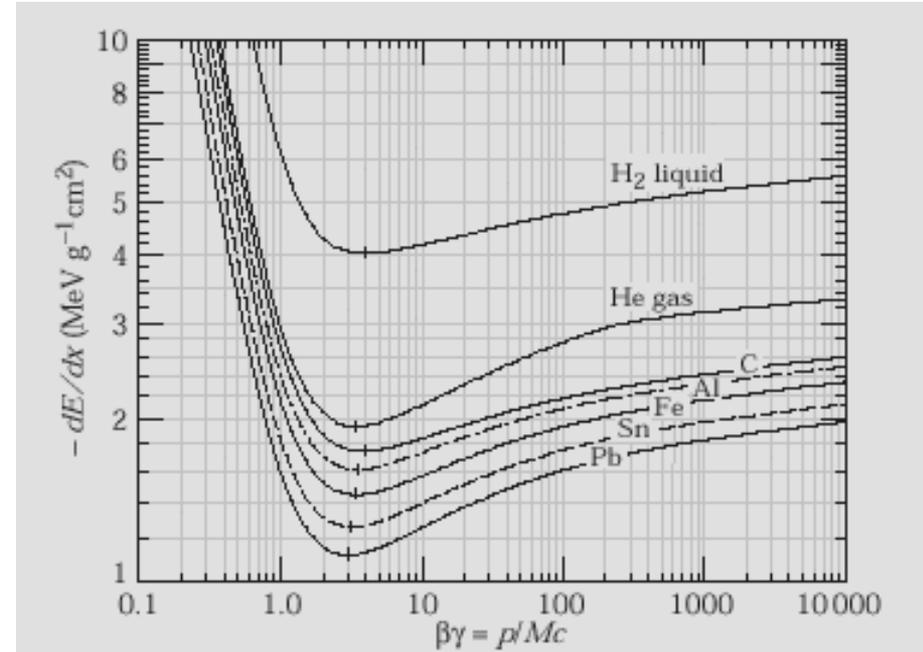
- Strong force has a short range and must act fast*
- Free neutron lives for a while before decaying*
 - > the force responsible for the decay must be weak*
- At quantum level all processes are probabilistic, i.e., we cannot know the outcome of an individual process but only the **probability** of each specific outcome*

Reminder - Energy loss

Ionization + excitation of atomic energy levels \longrightarrow **energy loss**

Mean energy loss rate – dE/dx

- proportional to (electric charge)² of incident particle
- for a given material, function only of incident particle velocity
- typical value at minimum:
 $-dE/dx = 1 - 2 \text{ MeV}/(\text{g cm}^{-2})$



Bremsstrahlung – energy loss due to emission of photons

NOTE: traversed thickness (dx) is given in g/cm^2 to be independent of material density (for variable density materials, such as gases)

\rightarrow multiply dE/dx by density (g/cm^3) to obtain dE/dx in MeV/cm

COSMIC RAYS

- Discovered by V.F. Hess in the 1910' s by the observation of an increase of radioactivity with altitude during a balloon flight
- Until the late 1940' s, the only existing source of high-energy particles

Composition of cosmic rays at sea level – two main components

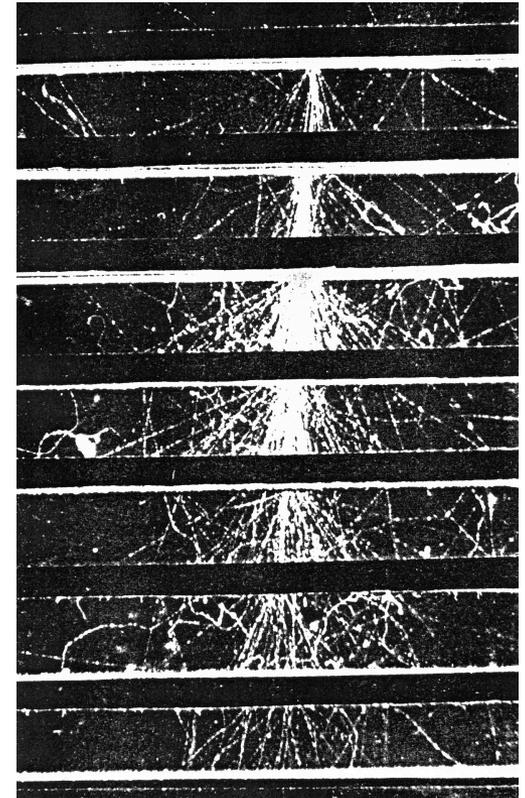
- Electromagnetic “showers”, consisting of many e^- and γ -rays, mainly originating from:
 $\gamma + \text{nucleus} \rightarrow e^+e^- + \text{nucleus}$ (pair production);
 $e^- + \text{nucleus} \rightarrow e^- + \gamma + \text{nucleus}$ (“bremsstrahlung”)

The typical mean free path for these processes (“radiation length”, x_0) depends on Z .

- For Pb ($Z = 82$) $x_0 = 0.56$ cm
- Thickness of the atmosphere $\approx 27 x_0$
- Muons (μ^-) capable of traversing as much as 1 m of Pb without interacting; tracks observed in cloud chambers in the 1930' s.

Determination of the mass by simultaneous measurement of momentum $p = mv(1 - v^2/c^2)^{-1/2}$ (track curvature in magnetic field) and velocity v (ionization):

$$m_\mu = 105.66 \text{ MeV}/c^2 \approx 207 m_e$$



Cloud chamber image of an electromagnetic shower.
Pb plates, each 1.27 cm thick

1937: Theory of nuclear forces (H. Yukawa)

Existence of a new light particle (“meson”) as the carrier of nuclear forces

Relation between interaction radius and meson mass m :

$$R_{\text{int}} = \frac{\hbar}{mc} \quad \longrightarrow \quad mc^2 \approx 200 \text{ MeV} \\ \text{for } R_{\text{int}} \approx 10^{-13} \text{ cm}$$



Hideki Yukawa

Yukawa's meson initially identified with the muon – in this case μ^- stopping in matter should be immediately absorbed by nuclei \rightarrow nuclear breakup (not true for stopping μ^+ because of Coulomb repulsion - μ^+ never come close enough to nuclei, while μ^- form “muonic” atoms)

Experiment of Conversi, Pancini, Piccioni (Rome, 1945):
study of μ^- stopping in matter using μ^- magnetic selection in the cosmic rays

In light material ($Z \sim 10$) the μ^- decays mainly to electron (just as μ^+ to positron)
In heavier material, the μ^- disappears partly by decaying to electron, and partly by nuclear capture (process later understood as $\mu^- + p \rightarrow n + \nu$).
However, the rate of nuclear captures is consistent with the weak interaction.

 **the muon is not Yukawa's meson**

1947: Discovery of the π^- meson (the “real” Yukawa particle)

Observation of the $\pi^+ \rightarrow \mu^+ \rightarrow e^+$ decay chain in nuclear emulsion exposed to **cosmic rays at high altitudes**

Nuclear emulsion: a detector sensitive to ionization with ~ 1 mm space resolution (size of AgBr microcrystals suspended in gelatin)

In all events the muon has a fixed kinetic energy (range of ~ 600 mm in nuclear emulsion \sim i.e. 4.1 MeV) \rightarrow two-body decay

Today's knowledge:

$m_\pi = 139.57 \text{ MeV}/c^2$; spin = 0

Dominant decay mode: $\pi^+ \rightarrow \mu^+ + \nu$

(and $\pi^- \rightarrow \mu^- + \nu$)

Mean life at rest: $t_\pi = 2.6 \times 10^{-8} \text{ s} = 26 \text{ ns}$

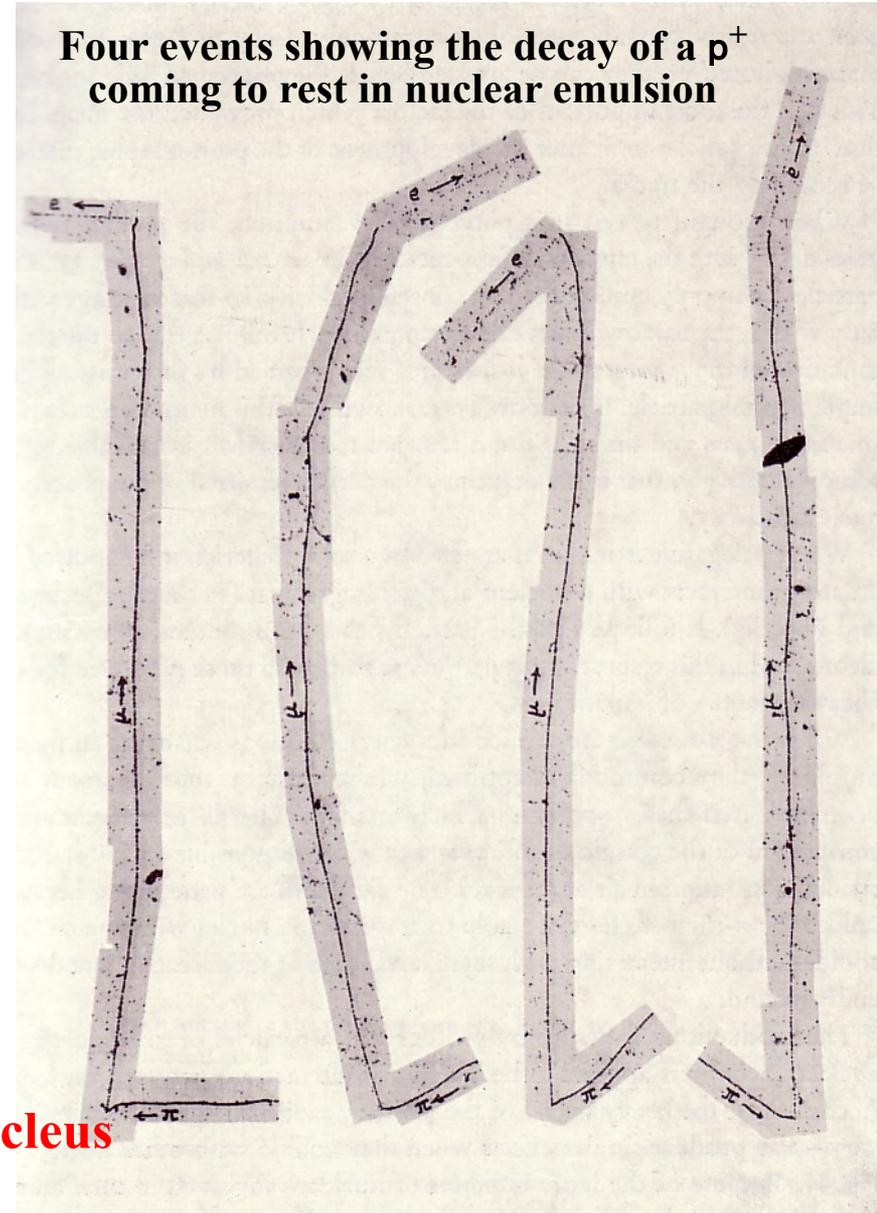
π^- at rest undergoes nuclear capture, as expected for the Yukawa particle

A neutral π – meson (π^0) also exists:

$m(\pi^0) = 134.98 \text{ MeV}/c^2$

Decay: $\pi^0 \rightarrow \gamma + \gamma$, mean life = $8.4 \times 10^{-17} \text{ s}$

π mesons are the most copiously produced particles in proton – proton and proton – nucleus collisions at high energies



The pion is the most common particle produced in nuclear collisions. Most of the tracks you have seen in the photographs of cosmic rays interactions with nuclear emulsion were pions.

Pions come with three different charges: π^+ , π^- , π^0

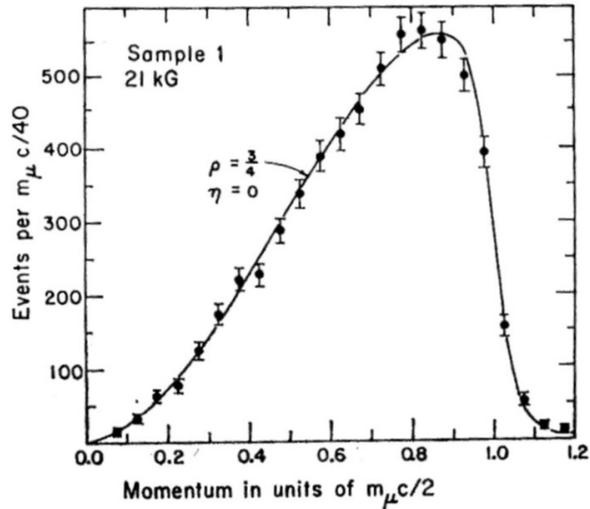
They are unstable and decay after $\sim 10^{-8}s$

As the detection techniques improved, many other short-lived particles were discovered. These we will discuss later

Muon decay

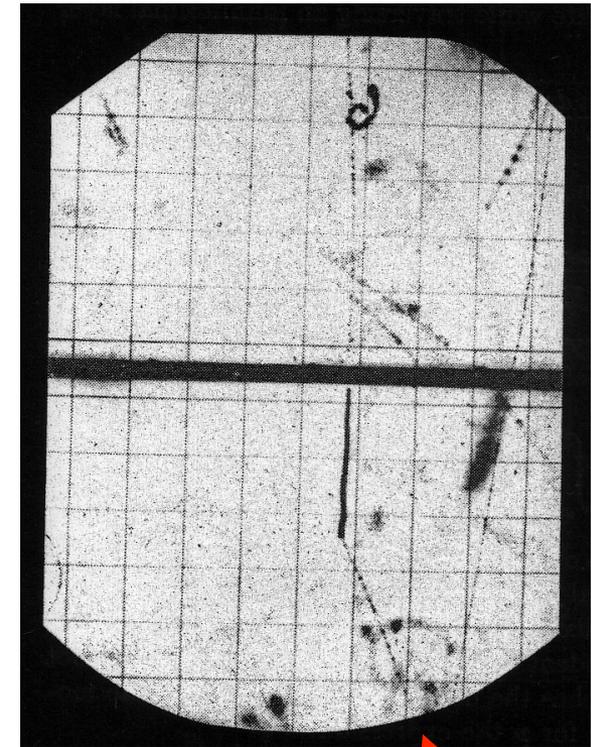


Decay electron momentum distribution



Cosmic ray muon stopping in a cloud chamber and decaying to an electron

Muon spin = $\frac{1}{2}$



decay electron track

Muon lifetime at rest: $t_\mu = 2.197 \times 10^{-6} \text{ s} \rightarrow 2.197 \text{ ms}$

Muon decay mean free path in flight:

$$\lambda_{decay} = \frac{v \tau_\mu}{\sqrt{1-(v/c)^2}} = \frac{p \tau_\mu}{m_\mu} = \frac{p}{m_\mu c} \tau_\mu c$$

$\left[\begin{array}{l} p : \text{muon momentum} \\ t_m c \sim 0.66 \text{ km} \end{array} \right]$

→ muons can reach the Earth surface after a path $\approx 10 \text{ km}$ because the decay mean free path is stretched by the relativistic time expansion

Theory of β -decay (E. Fermi, 1932-33)

$$\beta^- \text{ decay: } n \rightarrow p + e^- + \bar{\nu}$$

$$\beta^+ \text{ decay: } p \rightarrow n + e^+ + \nu \quad (\text{e.g., } ^{14}\text{O}_8 \rightarrow ^{14}\text{N}_7 + e^+ + \nu)$$

n : the particle proposed by Pauli
(named “neutrino” by Fermi)

$\bar{\nu}$: its antiparticle (antineutrino)



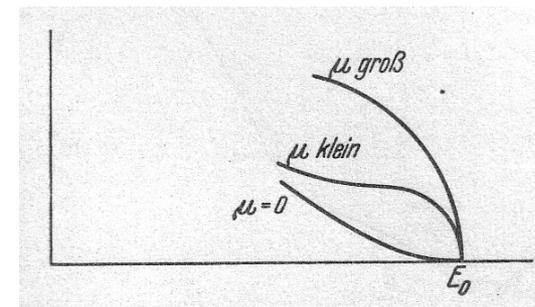
Enrico Fermi

Fermi's theory: a point interaction among four spin $\frac{1}{2}$ particles, using the mathematical formalism of creation and annihilation operators invented by Jordan
 \rightarrow particles emitted in β^- decay need not exist before emission – they are “created” at the instant of decay

Prediction of β^- decay rates and electron energy spectra as a function of only one parameter: Fermi coupling constant G_F (determined from experiments)

Energy spectrum dependence on neutrino mass m
(from Fermi's original article, published in German on *Zeitschrift für Physik*, following rejection of the English version by *Nature*)

Measurable distortions for $m > 0$ near the end-point (E_0 : max. allowed electron energy)

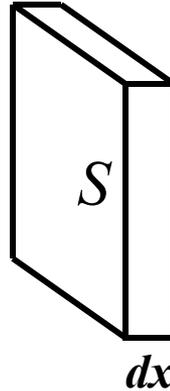


Neutrino detection

Prediction of Fermi's theory: $\bar{\nu} + p \rightarrow e^+ + n$

$\bar{\nu} - p$ interaction probability in thickness dx of hydrogen-rich material (e.g., H_2O)

Incident $\bar{\nu}$:
Flux Φ [$\bar{\nu} \text{ cm}^{-2} \text{ s}^{-1}$]
(uniform over surface S)



Target:
surface S , thickness dx
containing n protons cm^{-3}

$\bar{\nu} p$ interaction rate = $\Phi S n s dx$ interactions per second

σ : $\bar{\nu} - p$ cross-section (effective proton area, as seen by the incident $\bar{\nu}$)

$\bar{\nu} p$ interaction probability = $n \sigma dx = dx/\lambda$

Interaction mean free path: $\lambda = 1/n \sigma$

Interaction probability for finite target thickness $T = 1 - \exp(-T/\lambda)$

$\sigma(\bar{\nu} p) \approx 10^{-43} \text{ cm}^2$ for 3 MeV $\bar{\nu} \rightarrow \lambda \approx 150$ light-years of water !

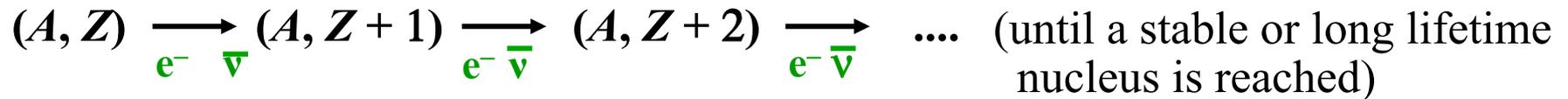
Interaction probability $\sim T/\lambda$ very small ($\sim 10^{-18}$ per meter of H_2O)

\rightarrow need very intense sources for antineutrino detection

Nuclear reactors: very intense antineutrino sources

Average fission: $n + {}^{235}\text{U}_{92} \rightarrow \underbrace{(A_1, Z) + (A_2, 92 - Z)}_{\text{nuclei with large neutron excess}} + 2.5 \text{ free neutrons} + 200 \text{ MeV}$

➔ a chain of β decays with very short lifetimes:



On average, $6 \bar{\nu}$ per fission

$$\bar{\nu} \text{ production rate} = \frac{6P_t}{200 \text{ MeV} \times \underbrace{1.6 \times 10^{-13}}_{\text{conversion factor MeV} \times \text{J}}} = 1.87 \times 10^{11} P_t \bar{\nu}/\text{s}$$

P_t : reactor thermal power [W]

For a typical reactor: $P_t = 3 \times 10^9 \text{ W} \approx 5.6 \times 10^{20} \bar{\nu}/\text{s}$ (isotropic)

Continuous $\bar{\nu}$ energy spectrum – average energy $\sim 3 \text{ MeV}$

*Comanche Peak nuclear Power Plant
Somervell County, 60 miles southwest of Dallas*

2 reactors 1.25 GW each = 2.5 GW

$$\rightarrow 1.87 \times 10^{11} \times 2.5 \times 10^6 = 4.67 \times 10^{20} \text{ v/s}$$

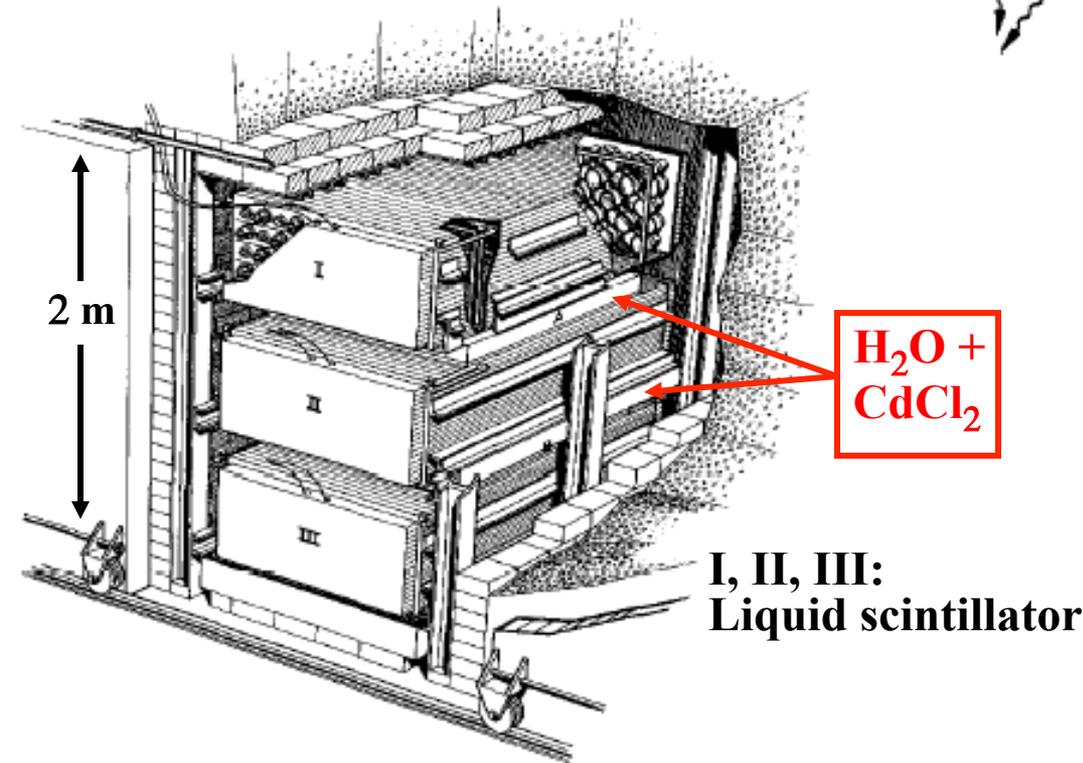
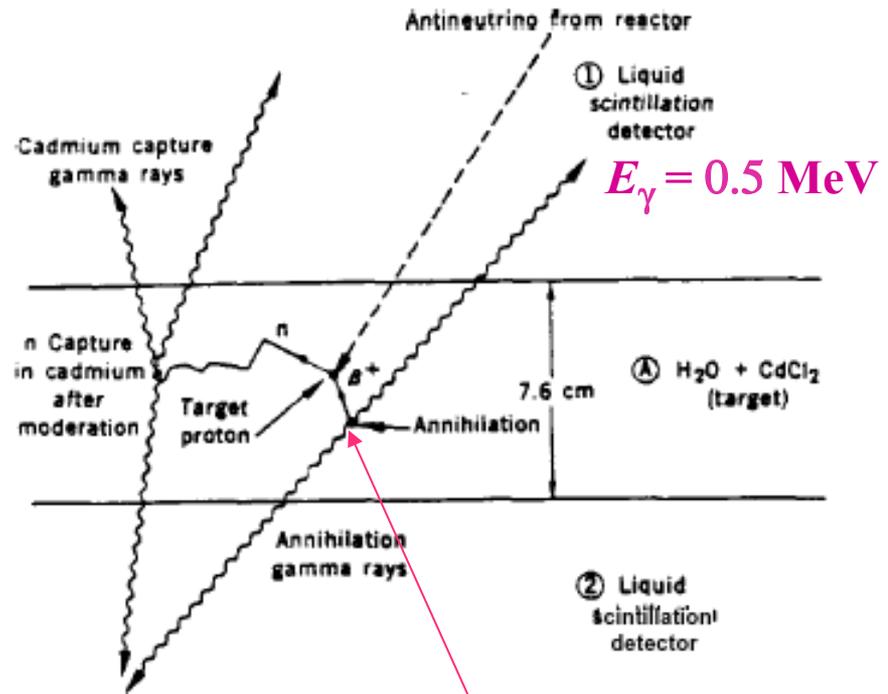
*If you assume that the cross section of your body is 1 m²
and that the flux of reactors neutrinos is distributed uniformly
on a surface of the sphere centered at the reactor then
there are $\sim 3.7 \times 10^9$ v from that reactor crossing your body every
second.*

First neutrino detection

(Reines, Cowan 1953)



- detect 0.5 MeV γ -rays from $e^+e^- \rightarrow \gamma\gamma$ ($t = 0$)
- neutron "thermalization" followed by capture in Cd nuclei \rightarrow emission of delayed γ -rays (average delay ~ 30 ms)



Event rate at the Savannah River nuclear power plant:

3.0 ± 0.2 events/hour

(after subtracting event rate measured with reactor OFF)

in agreement with expectations

Particle interactions (as known until the mid 1960' s)

In order of increasing strength:

■ Gravitational interaction (all particles)

Totally negligible in particle physics

Example: static force between electron and proton at distance D

$$\text{Gravitational: } f_G = G_N \frac{m_e m_p}{D^2} \qquad \text{Electrostatic: } f_E = \frac{1}{4\pi\epsilon_0} \frac{e^2}{D^2}$$

$$\text{Ratio } f_G / f_E \approx 4.4 \times 10^{-40}$$

■ Weak interaction (all particles except photons)

Responsible for β decay and for slow nuclear fusion reactions in the star core

Example: in the core of the Sun ($T = 15.6 \times 10^6$ °K) $4p \rightarrow {}^4\text{He} + 2e^+ + 2n$

Solar neutrino emission rate $\sim 1.84 \times 10^{38}$ neutrinos/s

Flux of solar neutrinos on Earth $\sim 6.4 \times 10^{10}$ neutrinos $\text{cm}^{-2} \text{s}^{-1}$

Very small interaction radius R_{int} (max. distance at which two particles interact)

($R_{\text{int}} = 0$ in the original formulation of Fermi' s theory)

■ Electromagnetic interaction (all charged particles)

Responsible for chemical reactions, light emission from atoms, etc.

Infinite interaction radius

(example: the interaction between electrons in transmitting and receiving antennas)