

Plane Wave solutions to Maxwell's Equations:

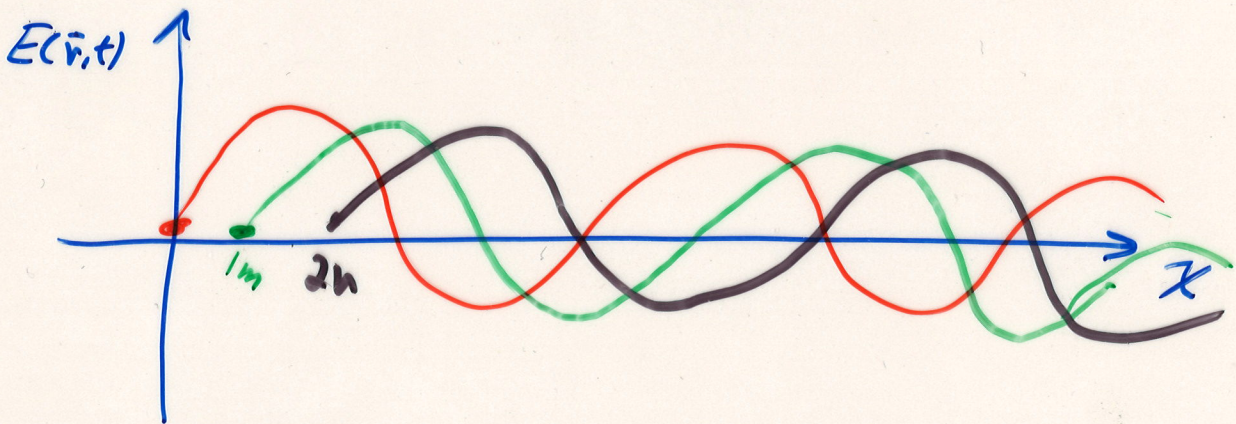
Suppose the wave is travelling along the x -axis from $-\infty$ to $+\infty$.

The function

$$E(\vec{r}, t) = E_{\max} \sin\left[\omega\left(\frac{x}{c} - t\right)\right]$$

satisfies Maxwell's Eq's in free space and describes a wave moving along the x -axis with speed c .

$$t=0 \quad \vec{E} = E_{\max} \sin\left(\frac{\omega x}{c}\right)$$



$$t = \frac{1\text{m}}{c} = \frac{1\text{m}}{3 \times 10^8 \text{ m/s}} = \frac{1}{3} \times 10^{-8} \text{ s}$$

$$\left(\frac{x}{c} - t\right) = 0 \quad \text{when } x = 1\text{m}$$

Notice that there is no y or z dependence in $E(\vec{r}, t)$.

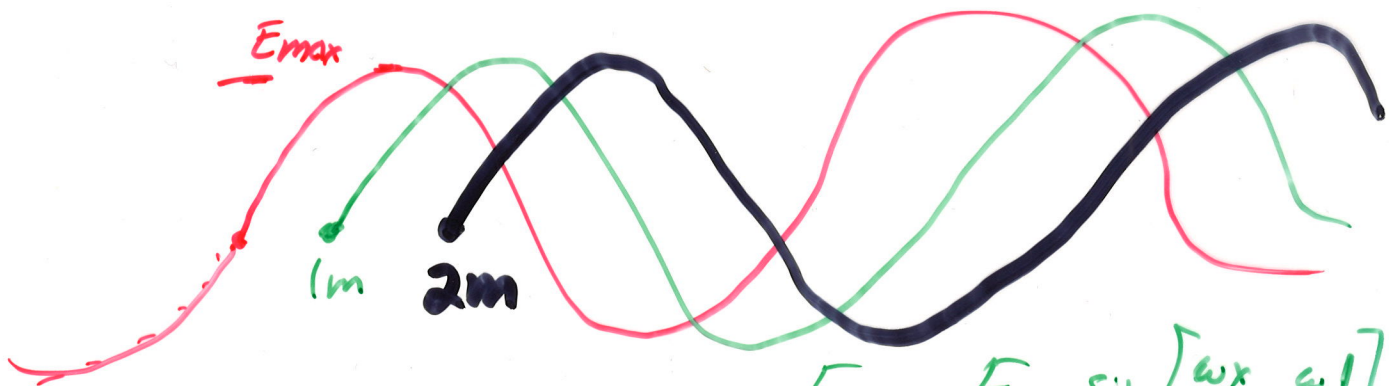
- What does this look like?
- What about the magnetic field?
- $\vec{E}(\vec{r}, t)$ and $\vec{B}(\vec{r}, t)$ are vector fields.

They are perpendicular to each other and to the direction of propagation.

For example: \vec{E} along y
 \vec{B} along z
motion along x

$$t = 0$$

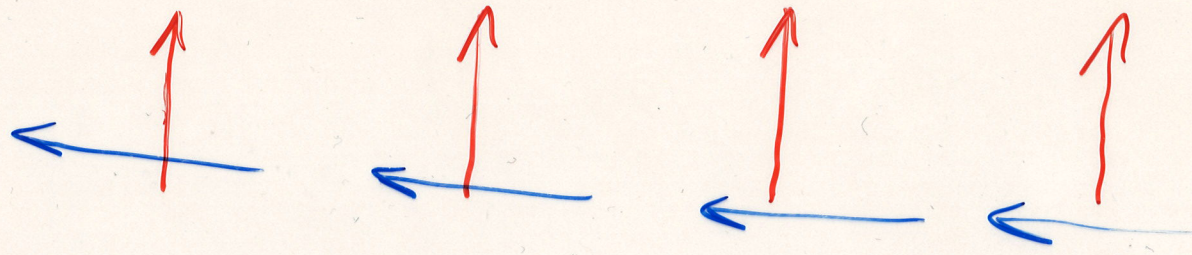
$$E(x) = E_{\max} \sin\left(\frac{\omega x}{c}\right)$$



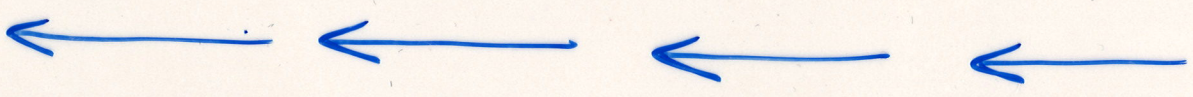
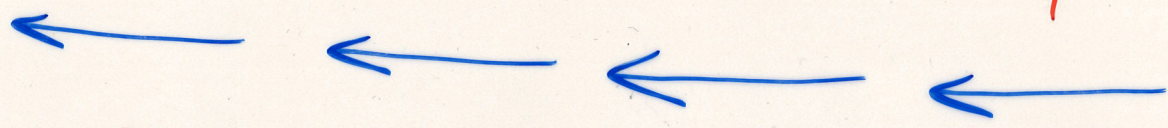
$$E(x) = E_{\max} \sin\left[\frac{\omega x}{c} - \frac{\omega t}{c}\right]$$

$$t = \frac{1m}{c} \sim 3ns$$

$$t = \frac{2m}{c} \sim 6ns$$

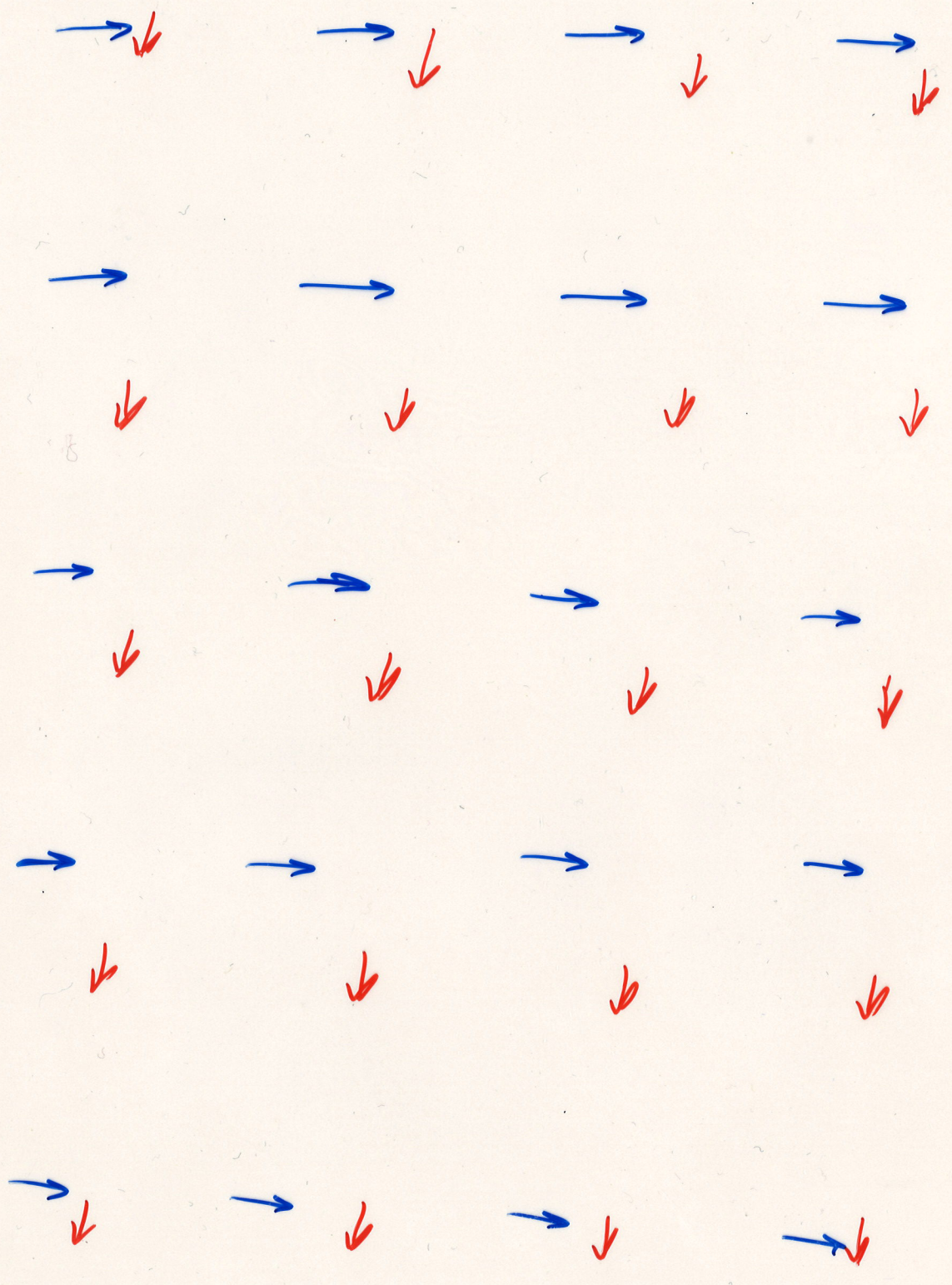


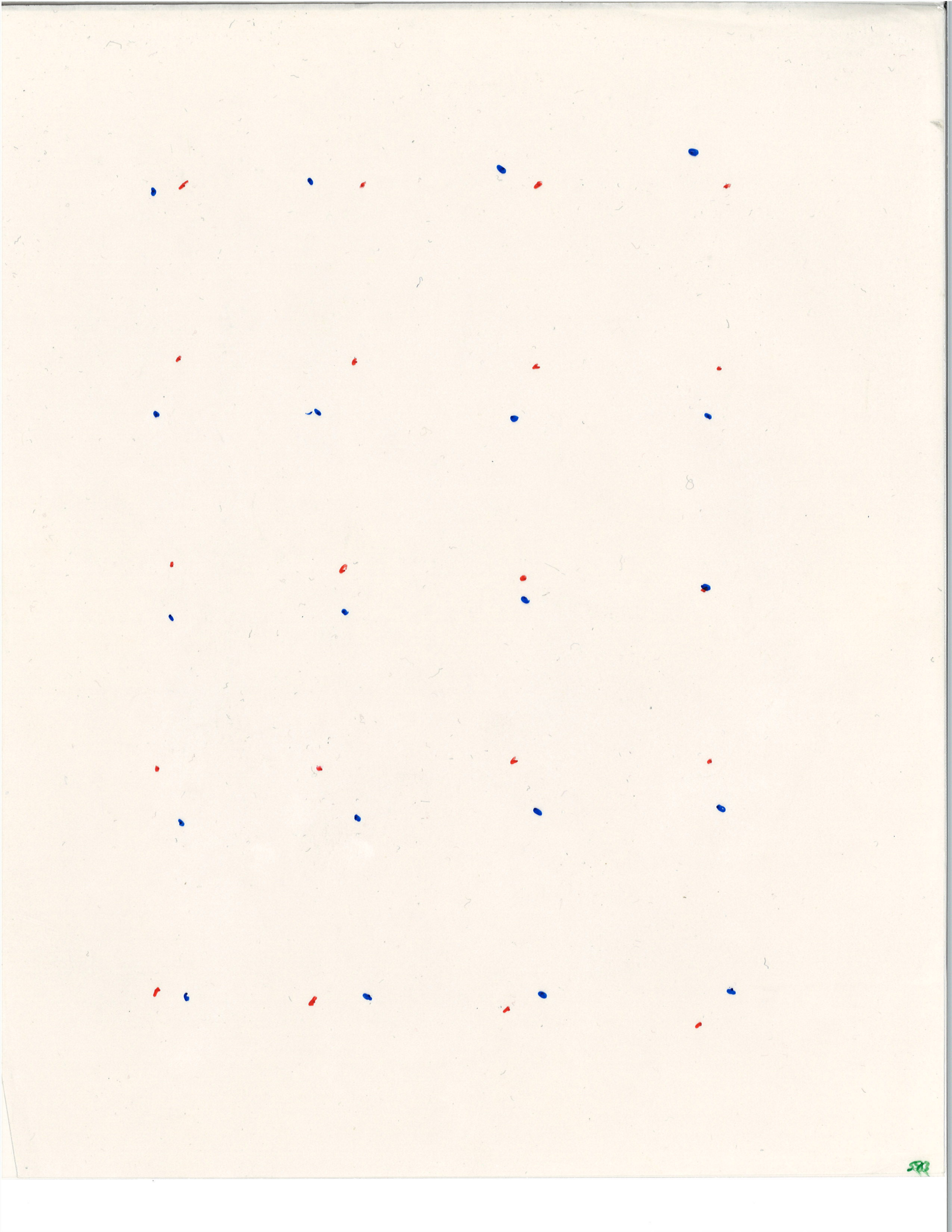
E₁

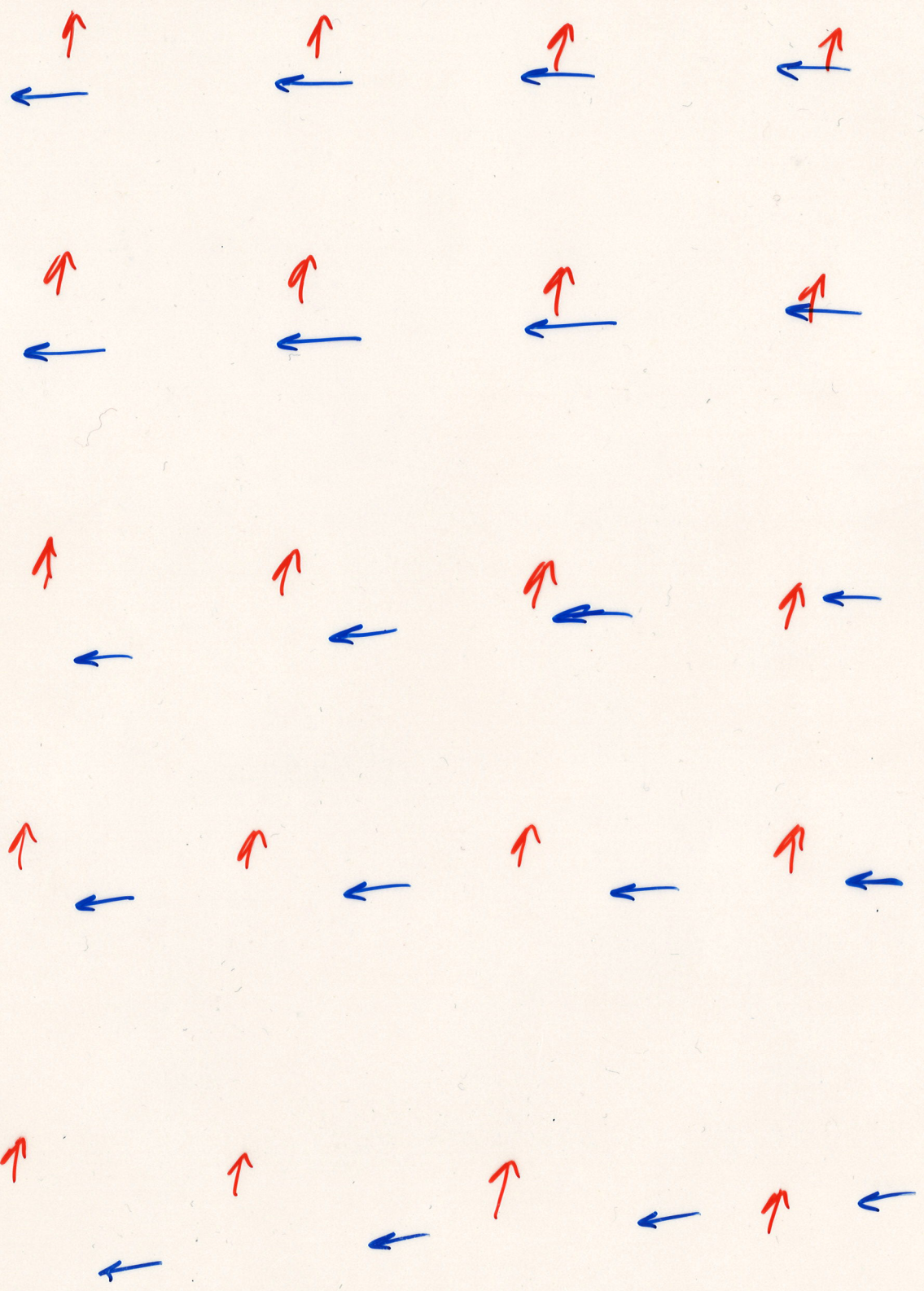


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Maxwell's Equations in vacuum

put a constraint on E_{\max} and B_{\max} .

They are not both arbitrary. You can select one, then the other must satisfy

$$\frac{E_{\max}}{B_{\max}} = c \quad (\text{speed of light})$$

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

In fact, the electromagnetic wave moves along the vector

$$\vec{E} \times \vec{B}$$

The direction of \vec{E} is called the polarization of the wave.

$$\vec{E}(\vec{r}, t) = E_{\max} \sin\left[\omega\left(\frac{x}{c} - t\right)\right] \hat{j}$$

$$\vec{B}(\vec{r}, t) = B_{\max} \sin\left[\omega\left(\frac{x}{c} - t\right)\right] \hat{k}$$

$\frac{\omega}{c}$ is often denoted as k , the wave number.

$$\vec{E}(\vec{r}, t) = E_{\max} \sin(kx - \omega t) \hat{j}$$

Energy in Electromagnetic Radiation

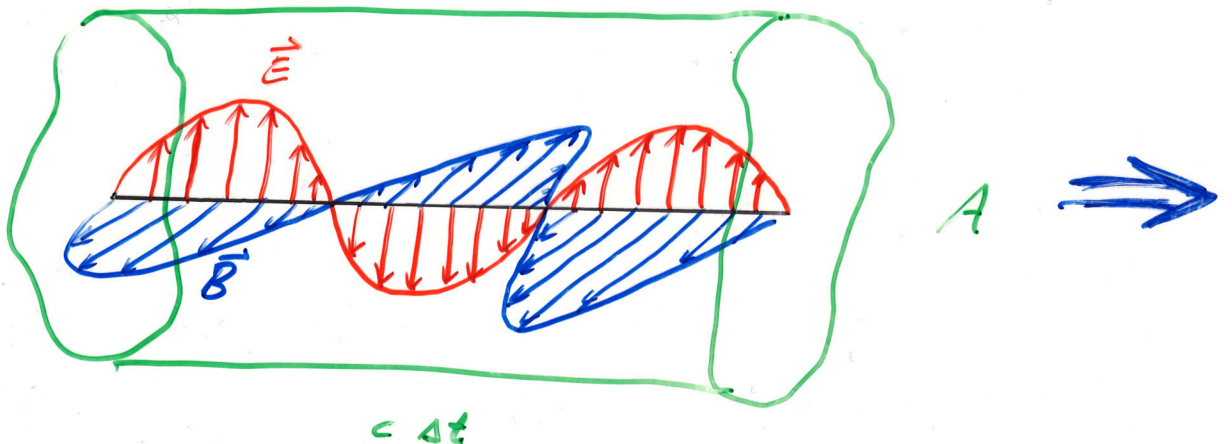
$$E(\vec{r}, t) = E_{\max} \sin(kx - \omega t) = E_{\max} \sin\left[\omega\left(\frac{x}{c} - t\right)\right]$$

$$B(\vec{r}, t) = B_{\max} \sin(kx - \omega t) = \frac{E_{\max}}{c} \sin\left[\omega\left(\frac{x}{c} - t\right)\right]$$

$$B_{\max} = \frac{E_{\max}}{c}$$

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

Consider a box of cross-sectional Area A and length $c\Delta t$. In a time Δt , all the energy in this box will fall on a screen (detector).



electric energy
density

$$u_E(\vec{r}, t) = \frac{1}{2} \epsilon_0 E^2(\vec{r}, t)$$

magnetic energy
density

$$\begin{aligned} u_B(\vec{r}, t) &= \frac{1}{2\mu_0} B^2(\vec{r}, t) = \frac{1}{2\mu_0} \frac{E^2(\vec{r}, t)}{c^2} \\ &= \frac{1}{2} \epsilon_0 E^2(\vec{r}, t) = u_E(\vec{r}, t) \end{aligned}$$

energy in box:

$$\begin{aligned} \Delta U &= (u_E + u_B) A c \Delta t \\ &= \epsilon_0 E^2(\vec{r}, t) A c \Delta t \end{aligned}$$

rate of energy hitting the detector per unit area:

$$\begin{aligned} \frac{\Delta U}{\Delta t} \frac{1}{A} &= \epsilon_0 c E^2(\vec{r}, t) = \frac{1}{c\mu_0} E^2(\vec{r}, t) \equiv S(\vec{r}, t) \\ &= \frac{1}{\mu_0} \vec{E}(\vec{r}, t) \cdot \vec{B}(\vec{r}, t) \end{aligned}$$

$S(\vec{r}, t)$ is the magnitude of the
Poynting vector.

$$\vec{S}(\vec{r}, t) = \frac{1}{\mu_0} \vec{E}(\vec{r}, t) \times \vec{B}(\vec{r}, t)$$

direction of travel, direction of energy flow

Intensity

Intensity I is the time-averaged Poynting vector.

$$I = \frac{1}{T} \int_0^T S(\vec{r}, t) dt = \frac{1}{c\mu_0} \frac{1}{T} \int_0^T E^2(\vec{r}, t) dt$$

$$= \frac{1}{c\mu_0} E_{\max}^2 \left(\frac{1}{T} \int_0^T \sin^2 \left[\omega \left(\frac{x}{c} - t \right) \right] dt \right) = \frac{1}{2}$$

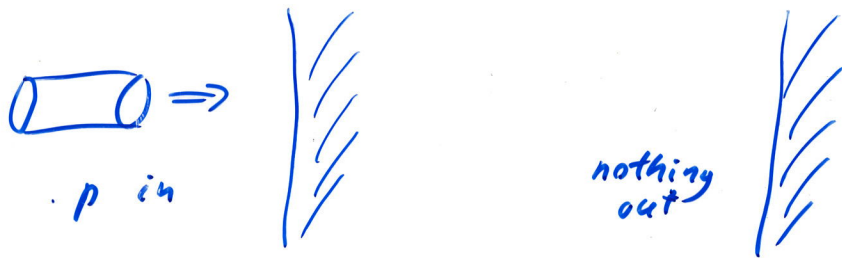
$$= \frac{1}{c\mu_0} \frac{E_{\max}^2}{2} = \frac{1}{c\mu_0} E_{\text{rms}}^2$$

$$E_{\max} = \sqrt{2} E_{\text{rms}}$$

Light carries momentum

If the energy in the box is ΔU ,
then the momentum in the box is $p = \frac{\Delta U}{c}$.

Case 1: Total absorption



momentum transferred to screen : $\frac{\Delta U}{c} = \Delta p$
 $= p_f - p_i$

Case 2: total reflection



momentum transferred to screen : $\frac{2\Delta U}{c} = \Delta p$

Radiation Pressure

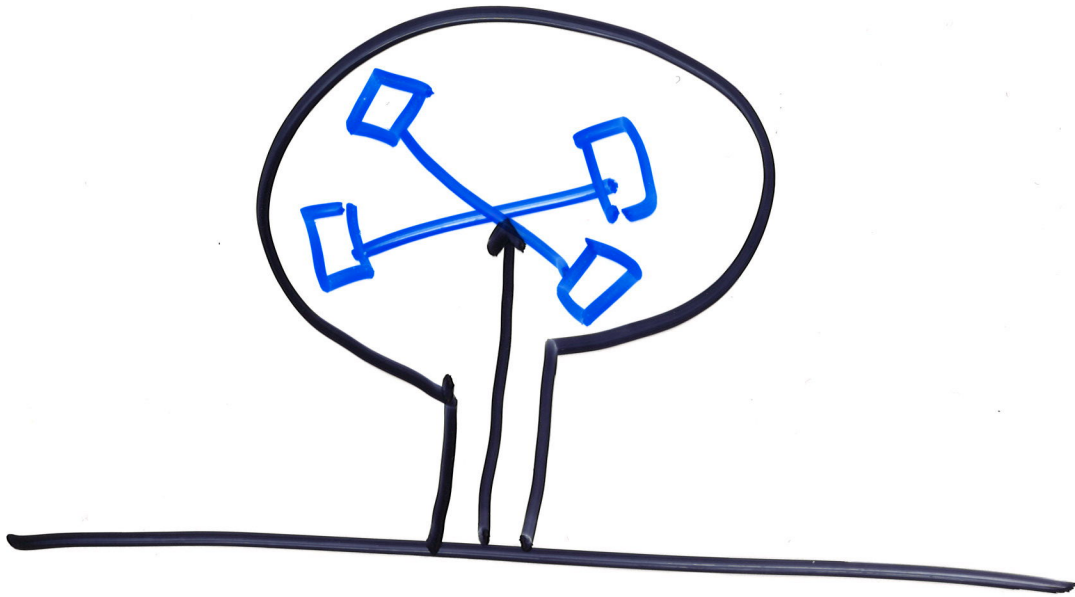
Average force : $F = \frac{\Delta p}{\Delta t}$

Average Pressure : $P = \frac{F}{A}$ force per unit area

pressure \rightarrow $P = \frac{\Delta p}{A \Delta t} = \begin{cases} \frac{\Delta U}{c A \Delta t} & \text{total absorption} \\ \frac{2 \Delta U}{c A \Delta t} & \text{total reflection} \end{cases}$

momentum \swarrow

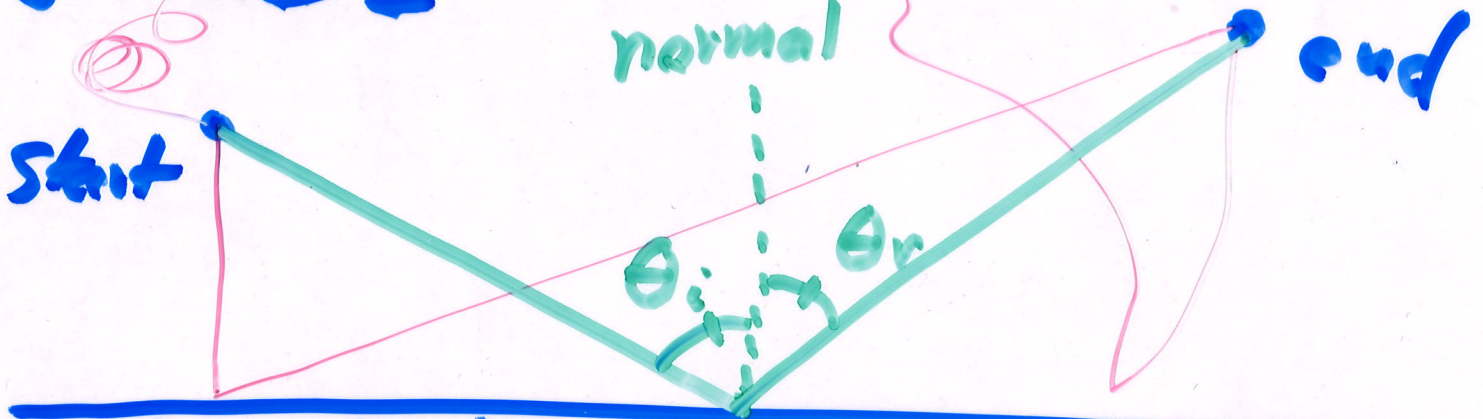
$$P = \begin{cases} \frac{I}{c} & \text{total absorption} \\ \frac{2I}{c} & \text{total reflection} \end{cases}$$



white / black
→

Reflection + Refraction ↑ (Bending)

Game #1



specular reflection : angle of incidence
= angle of reflection



Fermat law of least time

Game # 2

start

normal

θ_i

sand $v_s = \frac{c}{n_s}$

water $v_w = \frac{c}{n_w}$

fastest running speed: c

$$1 \leq n_i \leq \infty$$

Snell's Law

Snell ~~Law~~

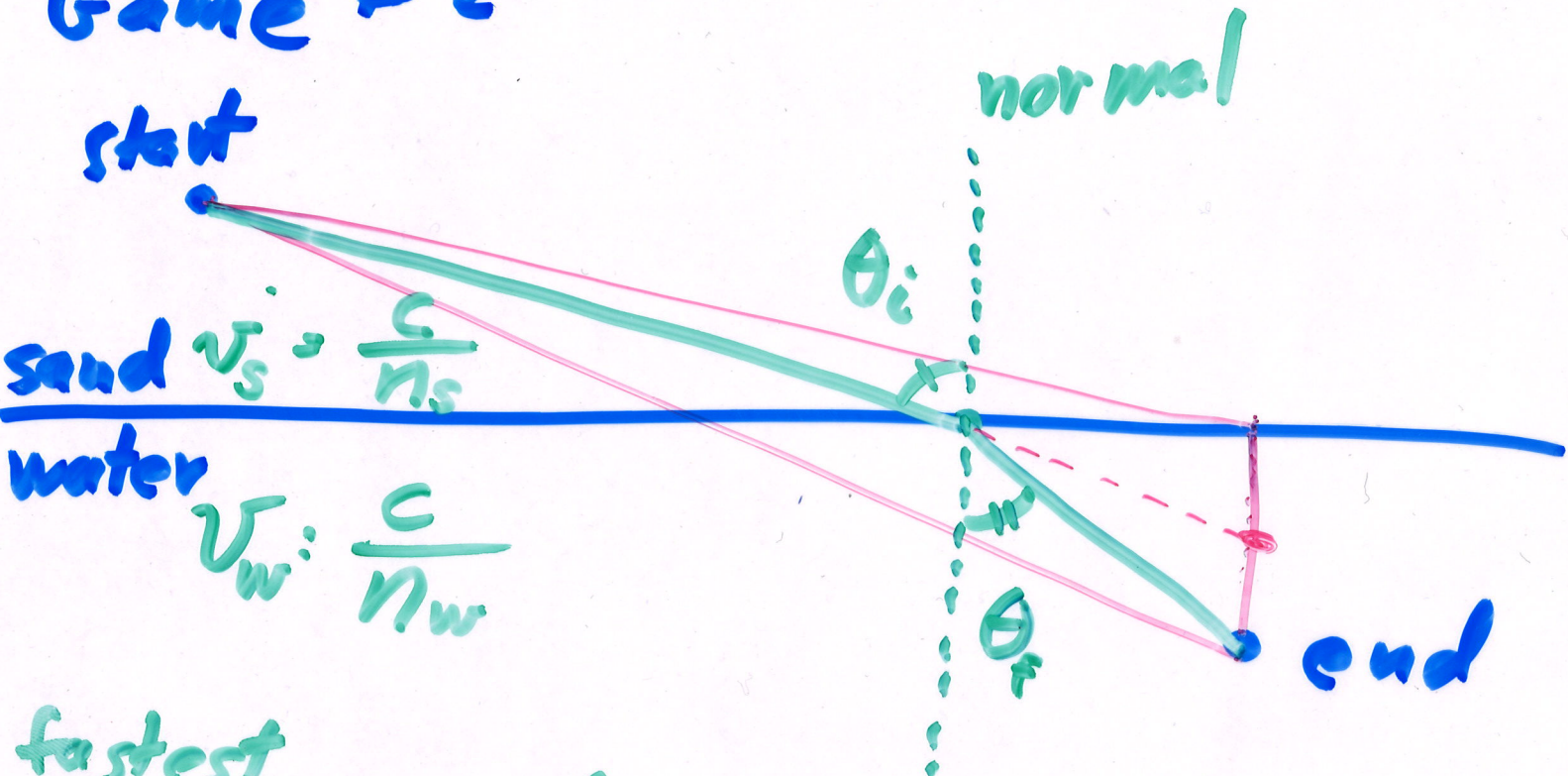
$$n_i \sin \theta_i = n_f \sin \theta_f$$

↑
incident

↑
refracted

$n_{air} \approx 1.00$?

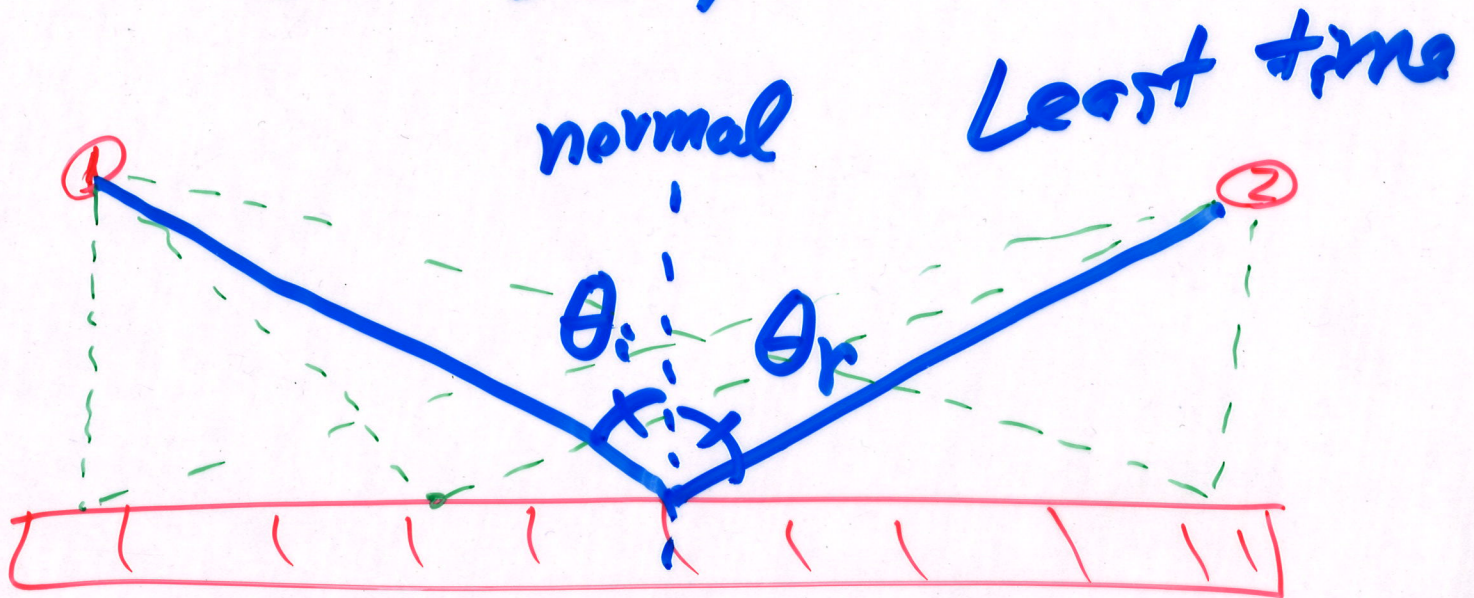
$n_{water} \approx \frac{4}{3} \approx 1.33$



Fermat's Principle of Least time

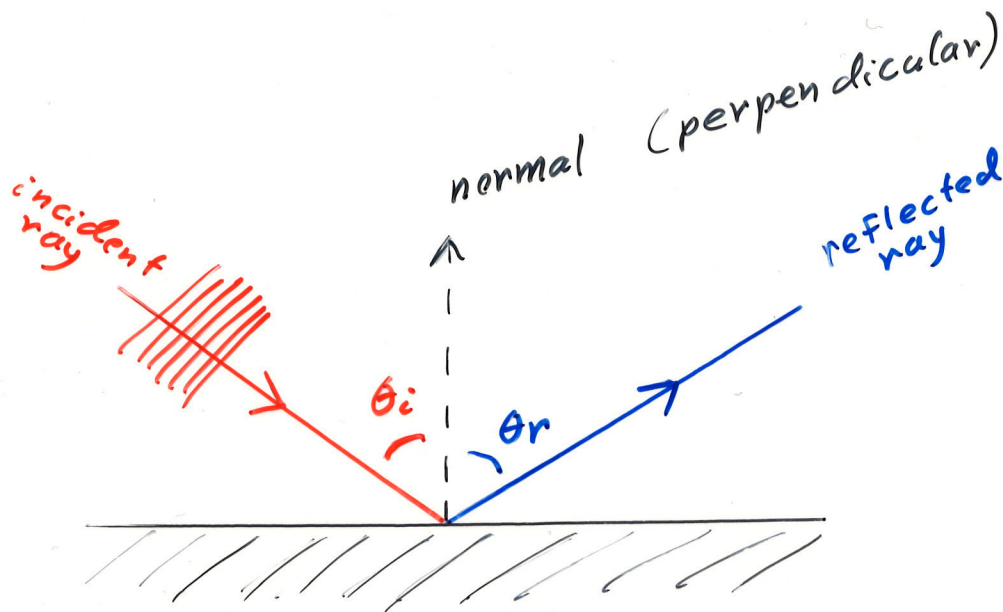
$$c^2 = a^2 + b^2$$

$$s^2 = z^2 + y^2$$



$$\theta_i = \theta_r \text{ for reflection}$$

Specular Reflection



$$\theta_i = \theta_r$$

angle of incidence = angle of reflection

The incident and reflected rays are in the same plane as the normal vector.

Refraction (Bending)

- The speed of light in a medium is less than the speed of light in vacuum.

$$v \leq c$$

- In fact, $v = \frac{c}{n}$ where $n \geq 1$. n is called the index of refraction.

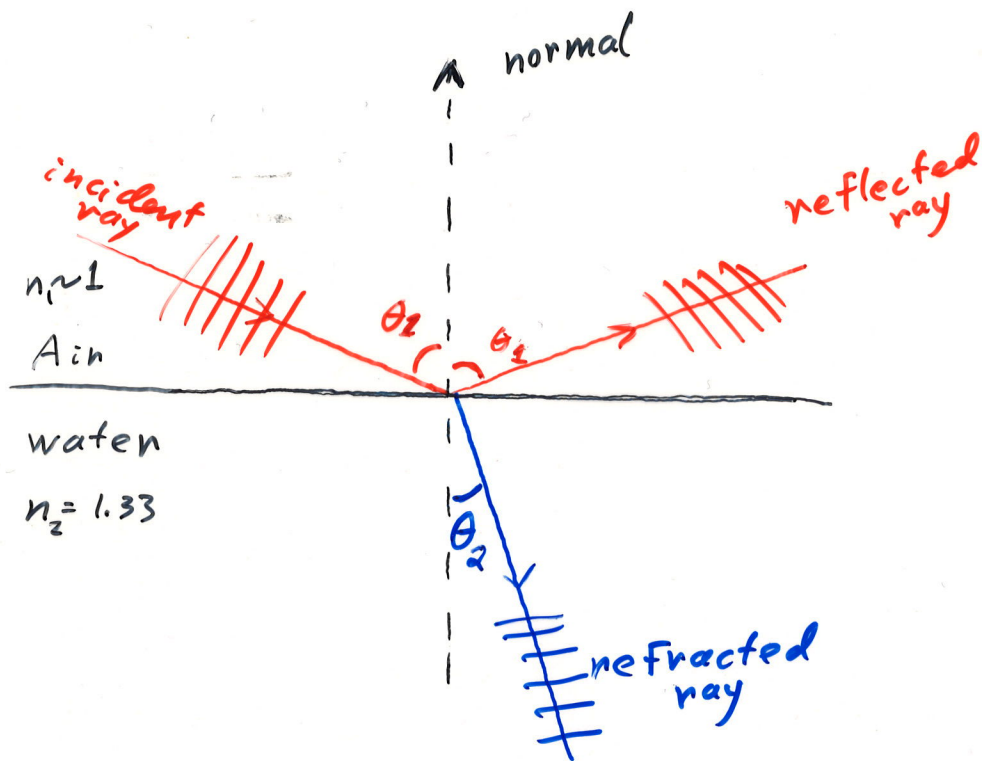
$n_{\text{air}} \approx 1$
 $n_{\text{vac}} = 1$
 $n_{\text{H}_2\text{O}} = 1.33$
 $n_{\text{Glass}} = 1.5$

- The frequency of light does not change in a medium.

- Because $\lambda f = v$, the wavelength changes. λ is shorter in a medium than in vacuum.

big density $\Rightarrow n$ large $\Rightarrow v$ small

All of this has the effect of bending (refracting) a light ray.

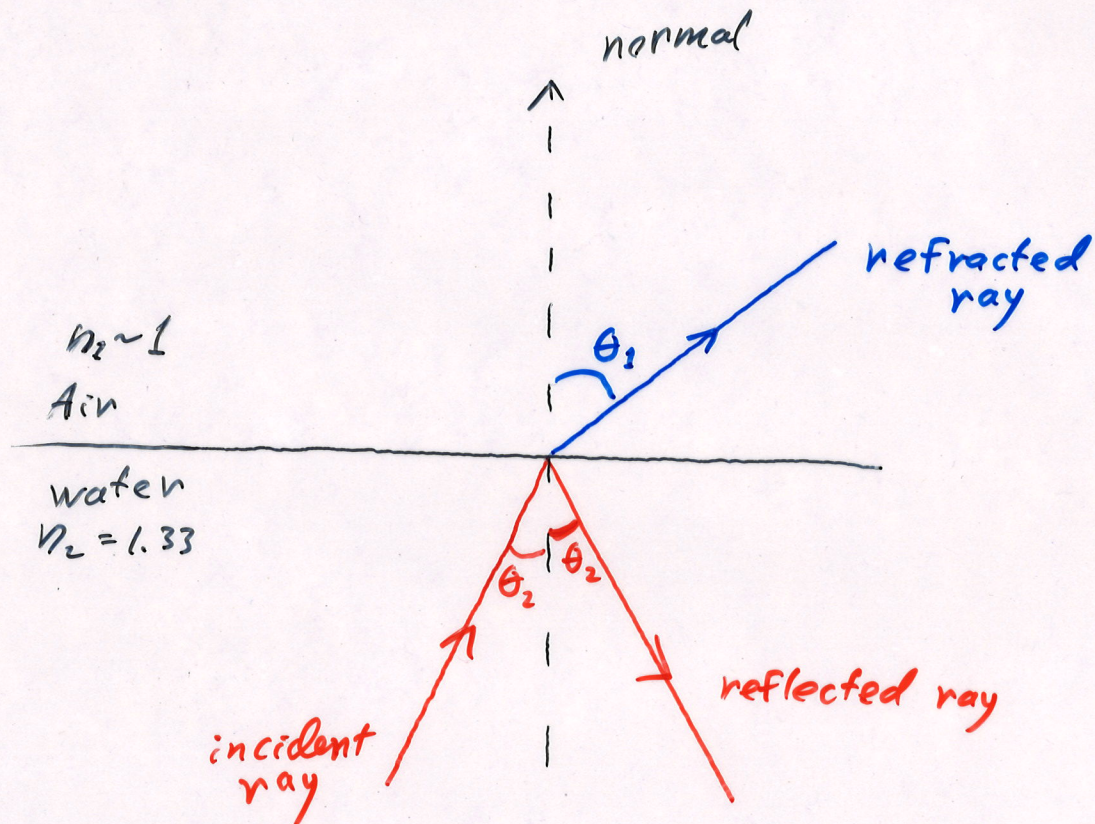


Snell's Law

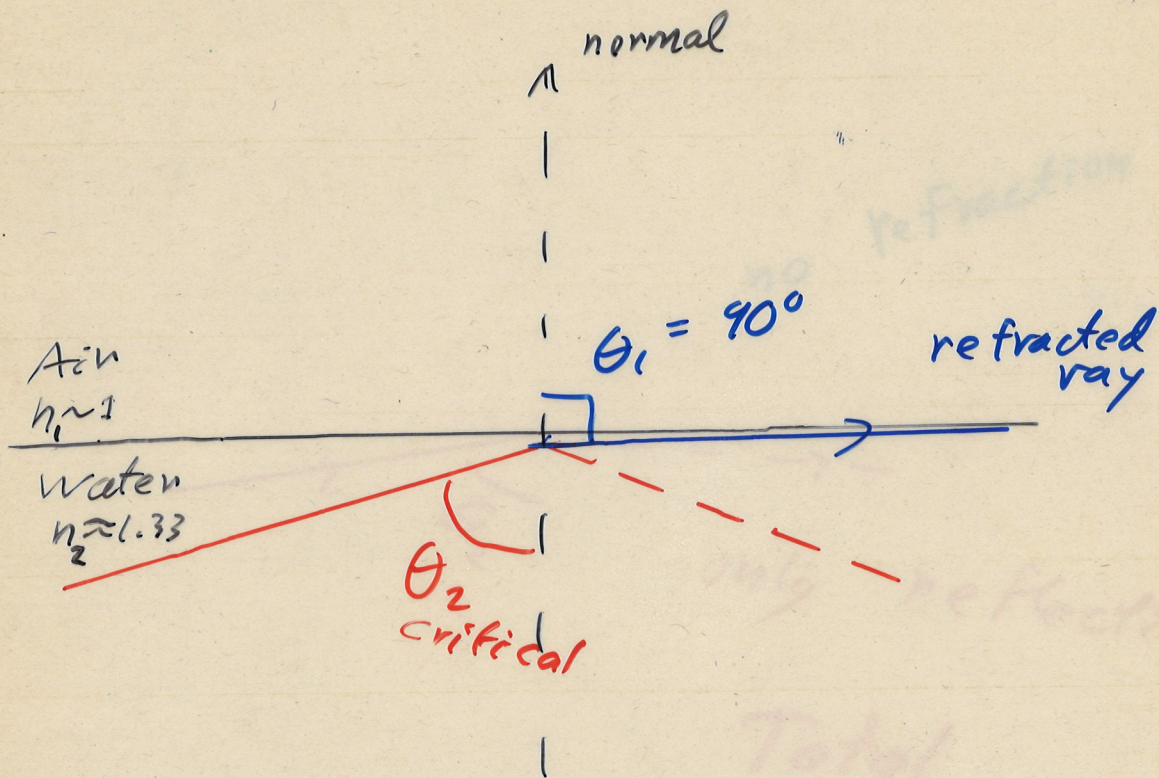
$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

$$\begin{aligned} n_1 &< n_2 \\ \sin \theta_2 &< \sin \theta_1 \\ \theta_2 &< \theta_1 \end{aligned}$$

The incident ray can also originate in the denser medium (the one with larger n).



$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$



Total internal reflection

— no refracted ray
Snell's Law

$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

Ex: What is the critical angle
for total internal reflection
at a water \rightarrow air interface?

$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

$$n_1 \sim 1 \text{ (Air)} \quad n_2 \sim 1.33 \text{ (Water)}$$

when $\theta_2 = \theta_2 \text{ critical}$ then $\theta_1 = 90^\circ$

$$\sin \theta_1 = \sin 90^\circ = 1$$

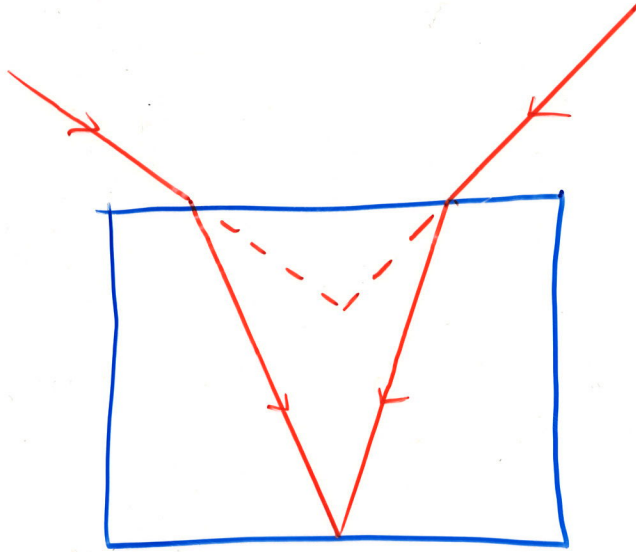
$$n_1 (1) = n_2 \sin (\theta_2 \text{ critical})$$

$$\frac{n_1}{n_2} = \sin (\theta_2 \text{ crit.})$$

$$\sin^{-1} \left(\frac{n_1}{n_2} \right) = \theta_2 \text{ critical}$$

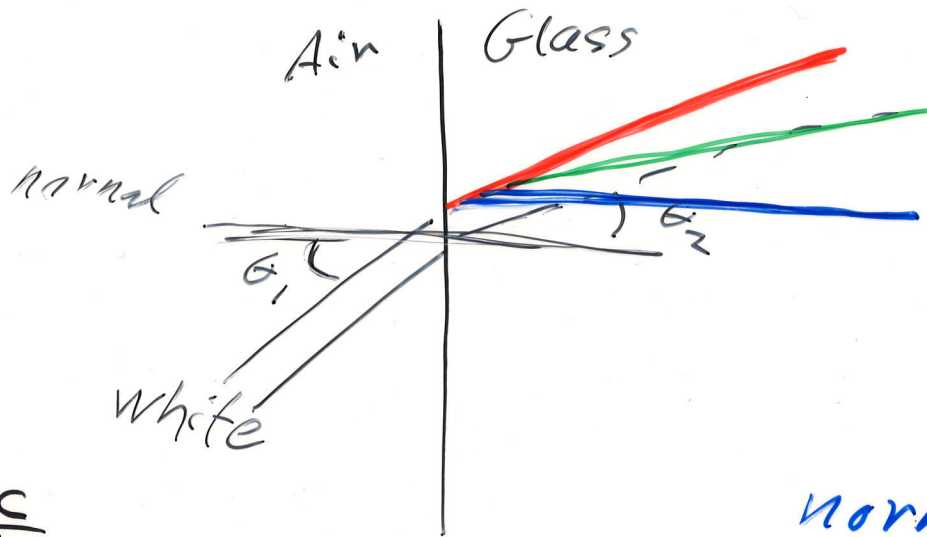
$$\sin^{-1} \left(\frac{1}{1.33} \right) \approx 49^\circ$$

The Glass Block



Objects appear higher than the bottom of the block.

Dispersion (spreading)



$$v = \frac{c}{n}$$

normal dispersion

$$\lambda f = v$$

$$v_{red} > v_{blue}$$

$$n_{red} < n_{blue}$$

Anomalous dispersion

$$v_{red} < v_{blue}$$