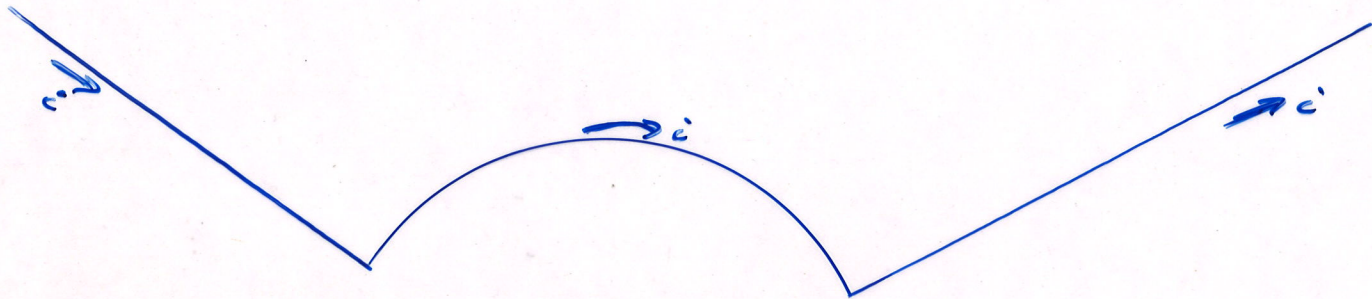


# Biot - Savart Examples

Ex Circular Arc



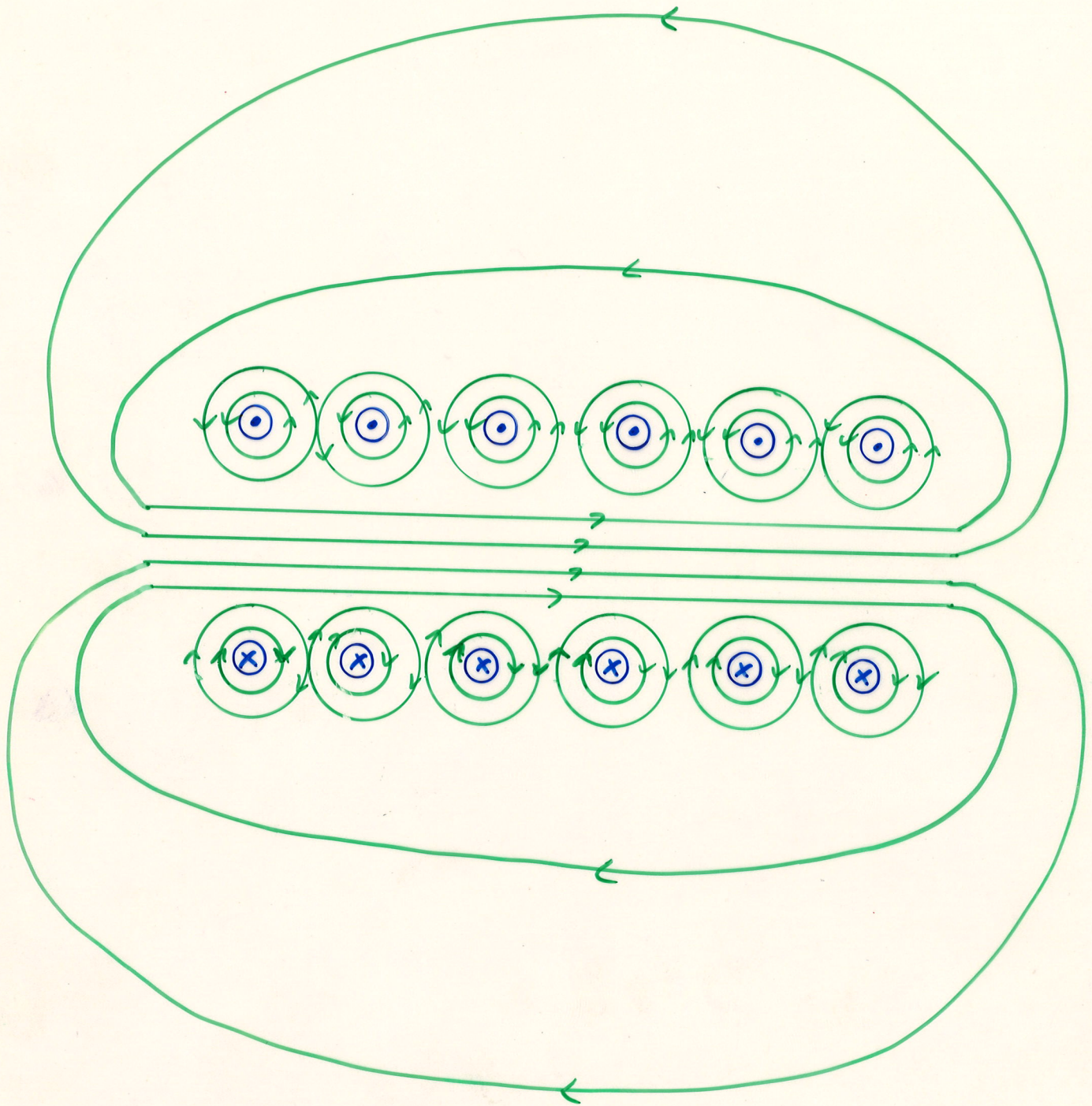
•  $\vec{B}$  "field point"

Ex.: Magnetic field due to a finite length of current-carrying wire.



•  
 $\vec{B}$  "field point"

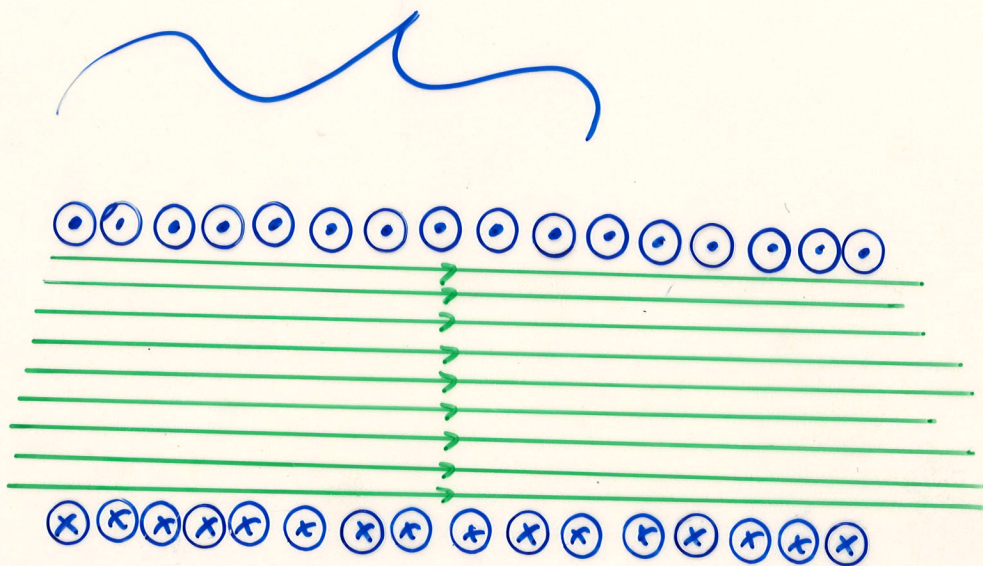
# The Solenoid



$\vec{B}$  field lines are closed loops.

# The Solenoid

$n$  turns per unit length  
(100 wires per inch)



If the solenoid is very long compared to its radius and if the coils are closely spaced then:

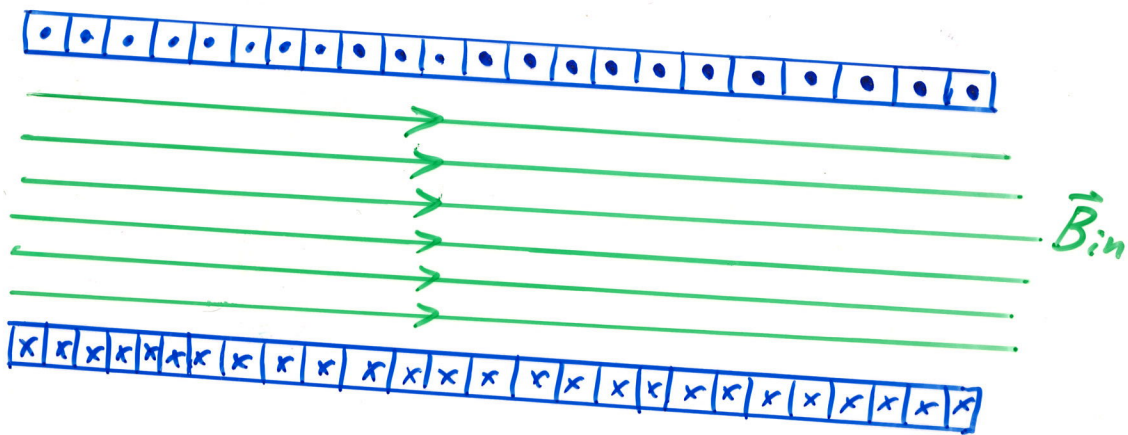
$$\vec{B}_{\text{inside}} \approx \text{constant}$$

$$\vec{B}_{\text{outside}} \approx 0$$

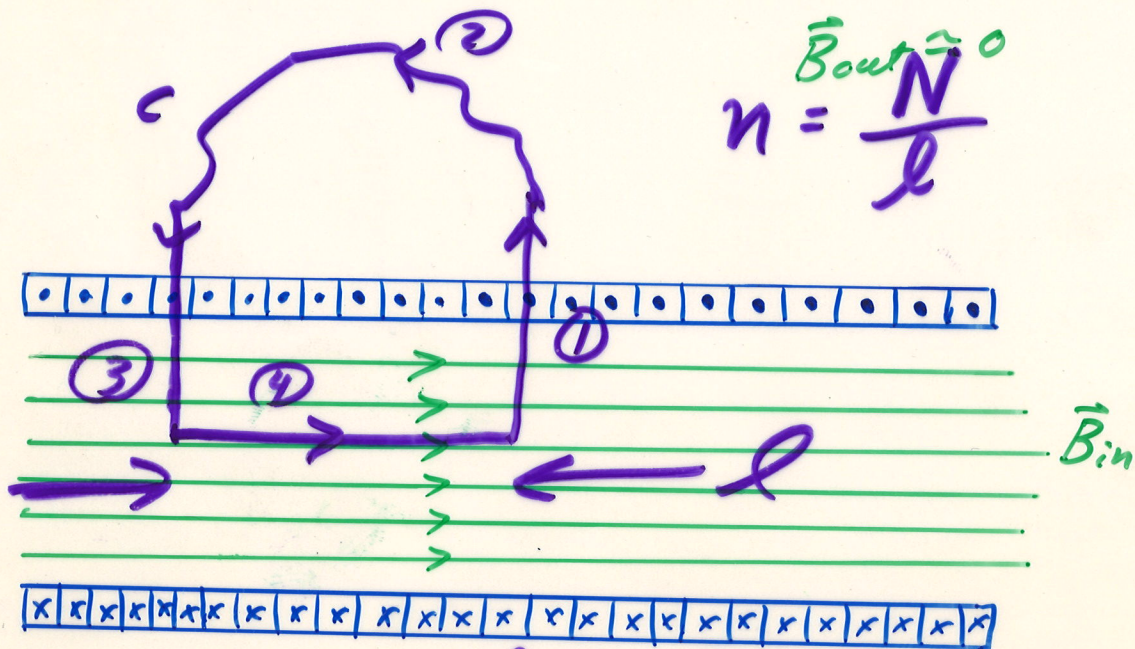
Well, not really, but the  $\vec{B}$  field is much less dense outside.

Magnetic field inside a solenoid by  
Ampere's Law:

$$\vec{B}_{out} = 0$$



Magnetic field inside a solenoid by  
 Ampere's Law:  $n$  turns of wire per  
 unit length



Ampere's Law  $\oint_C \vec{B} \cdot d\vec{l} = I_{enc} \mu_0 =$

$$= \int_1 \vec{B} \cdot d\vec{l} + \int_2 \vec{B} \cdot d\vec{l} + \int_3 \vec{B} \cdot d\vec{l} + \int_4 \vec{B} \cdot d\vec{l}$$

$\begin{matrix} 1 & \vec{B} \perp d\vec{l} \\ 2 & B=0 \\ 3 & \vec{B} \perp d\vec{l} \end{matrix}$

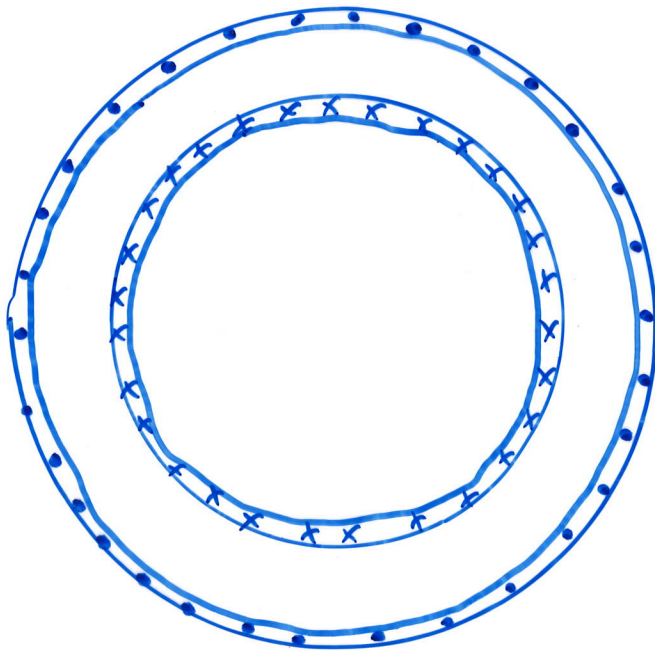
$$= \int_4 B dl = B \int_4 dl = B l = \mu_0 I (n l)$$

$$B = \begin{cases} \mu_0 n I, & \text{inside} \\ 0, & \text{outside} \end{cases}$$

# The Toroid

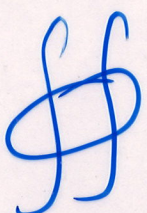
Instead of making the solenoid infinitely long to get a very small  $\vec{B}$  field outside, one can attach the open ends to each other to make a doughnut shape.

Total of  
 $N$  turns



# Gauss' Law

infinitesimal area vector that points out


$$\oint_S \vec{E} \cdot d\vec{A} = \frac{q_{enc}}{\epsilon_0}$$

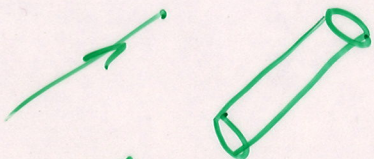


closed  
Surface

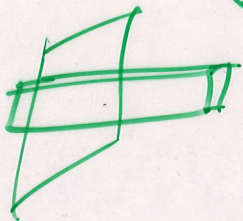
evaluated  
on the closed  
surface S



spherical Gaussian  
surface



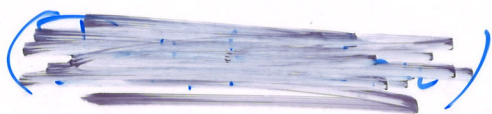
cylinder



pillbox



## Gauss' Law for Magnetism



$$\oint_S \vec{B} \cdot d\vec{A} = 0$$

# Ampere's Law

infinitesimal line element points  
along the curve C



$$\oint_C \vec{B} \cdot d\vec{l} = \mu_0 i_{enc}$$

current  
within the  
closed curve C

evaluated  
on closed  
curve C

closed  
~~curve~~  
curve

"Amperean  
Loops"

Useful in cases of high  
symmetry

# The Laws

of Electricity + Magnetism so far

Gauss' Law:  $\oint_S \vec{E} \cdot d\vec{A} = \frac{q_{enc}}{\epsilon_0}$   $\leftarrow$  by  $S$   
closed surface

Ampere's Law:  $\oint_C \vec{B} \cdot d\vec{l} = \mu_0 i_{enc}$   $\leftarrow$  by  $C$   
closed curve

Faraday's Law:

$$-\frac{d}{dt} \iint_S \vec{B} \cdot d\vec{A} = \oint_C \vec{E} \cdot d\vec{l}$$

↑ open surface      ↑ closed curve

bounded by  $C$

# Faraday's Law

$$\oint_C \vec{E} \cdot d\vec{\ell} = -\frac{d}{dt} \iint_S \vec{B} \cdot d\vec{A}$$



$$\mathcal{E} = -\frac{d}{dt} \Phi_B$$

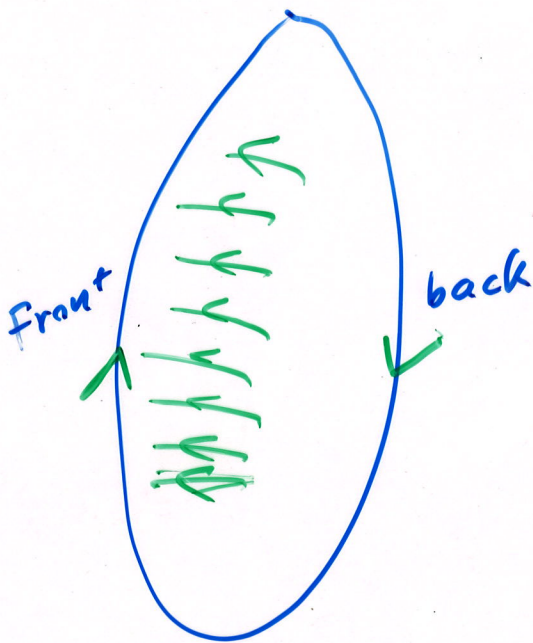


e.m.f. = electro-motive force  
(voltage)

$$\Delta V_{A \rightarrow B} = V_B - V_A = -\int_A^B \vec{E} \cdot d\vec{s}$$

# Lenz's Law

An induced current in a closed conducting loop will create a magnetic field that opposes the change in the external magnetic field.

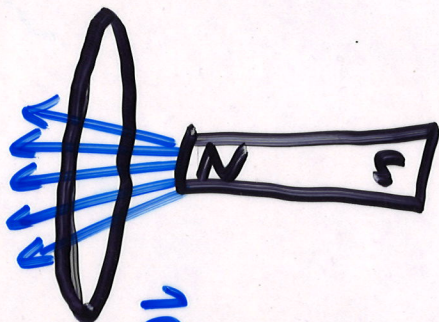


# Lenz's Law

"Lenz's law tries to maintain the status quo."

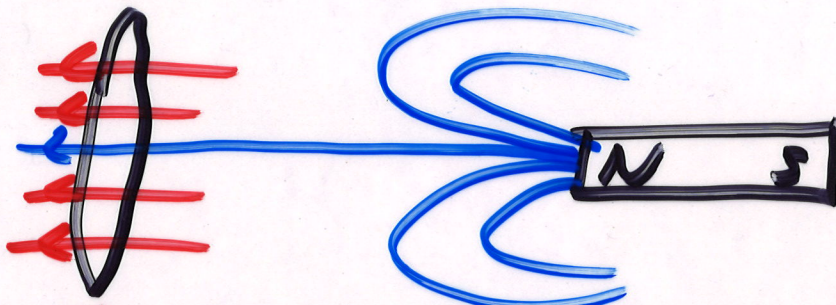
$I_{ind} = \begin{cases} \text{up in front} \\ \text{down in back} \end{cases}$

Before



$\vec{B}_{external}$

After



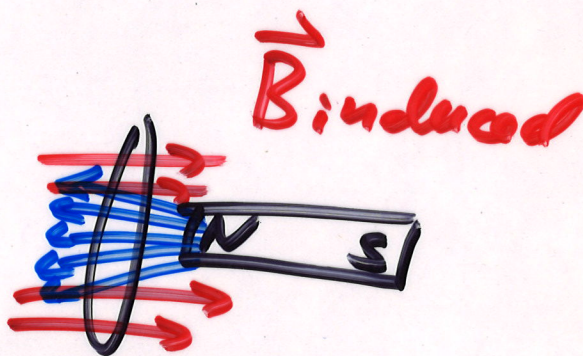
$\vec{B}_{induced}$

Before



$\vec{B}_{external}$

After



$\vec{B}_{induced}$

$\vec{B}_{ext}$

$I_{ind} = \begin{cases} \text{up in back} \\ \text{down in front} \end{cases}$

# Self Inductance (L)



Suppose that you want to start current flowing in a circuit.

When you close the switch, the external voltage  $V$  will begin to push charges around the loop, a current  $i$ , time dependent.

But this current will produce a magnetic field,  $\vec{B}$ . Some of the  $\vec{B}$  lines will penetrate the loop giving a magnetic flux  $\Phi_B$ . Because the current changes in time, so does the flux.

The changing magnetic flux gives rise to a back e.m.f.  $\mathcal{E}'$  by Faraday's Law

$$\mathcal{E}' = - \frac{d\Phi_B}{dt} = - \frac{d}{dt} \iint_{\substack{\text{Area} \\ \text{of loop}}} \vec{B} \cdot d\vec{A}$$

This is a little push against the battery voltage.

Since  $\Phi_B \propto B$  and  $B \propto i$   
we have

$$\mathcal{E} = - \frac{d\Phi_B}{dt} \propto \frac{di}{dt}$$

The constant of proportionality is  
called the (self) inductance

$$\mathcal{E} = -L \frac{di}{dt}$$

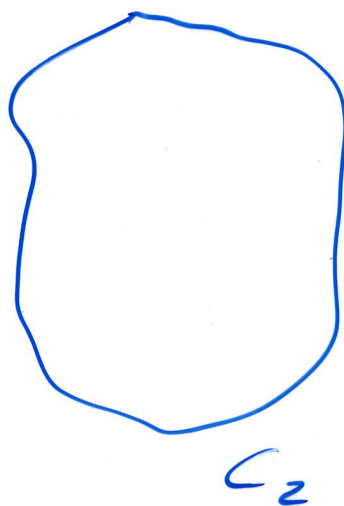
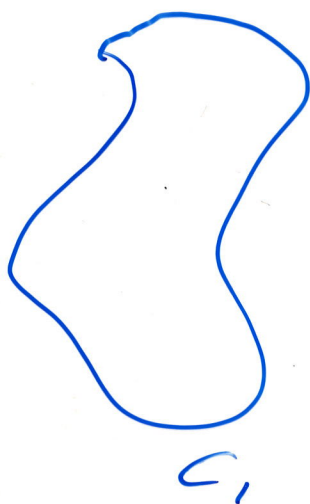
$$Q = CV$$

The MKS unit of inductance is  
the henry (H).

$$H = \frac{V \cdot s}{A}$$



# Mutual Inductance (M)



changing the current in loop 1  
induces an emf. in loop 2

$$\mathcal{E}_2 = -M \frac{di_1}{dt}$$

changing the current in loop 2  
induces an e.m.f. in loop 1

$$\mathcal{E}_1 = -M \frac{di_2}{dt}$$