

MEAS. T :

THERMOMETERS :

1st question - 3.1
t.1

(1) and temp around
(2) what are they

ship { - WANT TO KNOW HOW Q FLOWS FROM SYS A TO B
WITHOUT PUTTING IN CONTACT,

no { - NOT DIRECT → T DEF'D IN TERMS OF MICRO SYS. ; HOW TO GET AT IT VIA MEAS?

THERMOMETER: SYS M WHICH

(1) HAS SOME MICRO "THERMOMETRIC PARAM" $\Theta(\bar{E})$
WHICH CHANGES AS M ABSORBS Q. } will detect diff in T by Q

ex

(a) MERCURY THERM:

no { ABS Q, AVE E OF MOLS INCR → EXPANDS

→ $\underline{\Theta = V}$ (meas. temp in units of vol.)

(b) RESISTANCE THERM:

RES. INCR. AS ABSORBS E ⇒ $\underline{\Theta}$ IN OHMS

(c) CONST VOL GAS THERM ⇒ $\underline{\Theta = \bar{p}}$

(2) SMALL → DOESN'T AFFECT MEAS'D SYS

TO USE :

(a) PUT IN THERM. CONTACT "w/ SYS A

ABSORBS Q UNTIL REQUIL ; THEN $\beta_M = \beta_A$

→ RECORD $\Theta(A)$

(b) REPEAT FOR SYS B

IF $\Theta(A) = \Theta(B)$, THEN $\beta_M = \beta_A, \beta_M = \beta_B$ } Therm. equil depends on single param.

∴ $\beta_A = \beta_B$ → EQUIL IF PUT IN CONTACT



(that doesn't tell me what β is)

DIFF THERMS USEFUL IN DIFF CASES:

EX LOWER T: CONST V GAS BETTER THAN Hg (FREEZES)

- HARD TO COMPARE!

DIFF UNITS

NOT PROPORTIONAL { EX Hg MIGHT DOUBLE WHILE
RESIST. INCR. A LITTLE

NOTE:

IF PARAM Θ WORKS, SO DOES ANY $f(\Theta)$ (IF $\frac{df}{d\Theta} \neq 0$)

VERY USEFUL $f(\Theta)$:

$$\beta_M(\Theta) \equiv \frac{1}{kT}$$

\Rightarrow IF HAVE $\beta_M(\Theta)$ & $\beta_{M'}(\Theta')$, CAN CP. M, M' DIRECTLY.
NOW MEAS. SAME FUND. QTY

TO DO THIS:

NEED 1 SYS. WHERE KNOW $\beta_M(\Theta)$ (OR $T_M(\Theta)$)

THEN:

USE TO CALIBRATE EVERY OTHER $M' \Rightarrow \beta_{M'}(\Theta')$

EX

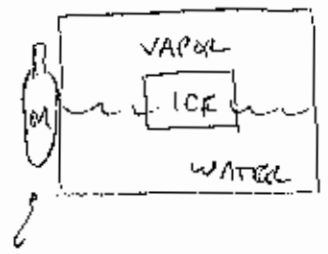
IDEAL GAS

$$\begin{array}{l} \bar{p} V = NkT = (\nu N_A) kT \\ \parallel \quad \uparrow \\ \Theta \quad \text{HELD CONST} \end{array} \quad \left. \begin{array}{l} \text{GIVES } T(\Theta) \\ = T(\bar{p}) \end{array} \right\}$$

USEFUL: ALL GASES OBEY FOR LOW DENSITY
SOME TO VERY LOW T

DEFN. FOR °K (KELVIN): (INTL CONVENTION)

- (1) MAKE IG. FIXED-V THERM
- (2) CHOOSE REF. TEMP $T_E \equiv$ TRIPLE PT. OF H_2O



TEMP. AT WHICH ICE, WATER, VAPOR
IN EQUIL \Rightarrow UNIQUE, AVAIL., EASY TO
PREP. (JUST PUT TOGETHER AND WAIT
FOR EQUIL)

MEAS. \bar{P}_E AT T_E FOR M

- (3) TO MEAS. T FOR NEW SYS:

$$\frac{\bar{P} V}{\bar{P}_E V} = \frac{\nu N_A k T}{\nu N_A k T_E} \Rightarrow \frac{T}{T_E} = \frac{\bar{P}}{\bar{P}_E} \quad \left. \vphantom{\frac{T}{T_E} = \frac{\bar{P}}{\bar{P}_E}} \right\} \text{GIVES RATIO OF } T \text{ TO } T_E$$

CONVENTION:

$$T_E \equiv 273.16 \text{ } ^\circ\text{K} \text{ (EXACT)} \quad \left. \vphantom{T_E \equiv 273.16 \text{ } ^\circ\text{K}} \right\} \text{DEFINES } ^\circ\text{KELVIN} \dots$$

$$\Rightarrow T = (273.16 \text{ } ^\circ\text{K}) \frac{\bar{P}}{\bar{P}_E}$$

$$(4) \bar{P}_E V = \nu N_A k T_E$$

FIXES $k = 1.38 \times 10^{-16} \text{ ergs/} ^\circ\text{K}$

"BOLTZMANN CONST"

(COULD HAVE ALSO FIXED k AS EXACT, MEAS T_E , TO DEF UNIT)

TORTURED DEFN?

DESIGNED TO MATCH CELSIUS (w/ OFFSET) (OLDRA) (IN WIDE USE)
 i.e. $1 \text{ } ^\circ\text{K} = 1 \text{ } ^\circ\text{C}$

$$\theta \text{ } ^\circ\text{C} = (T - 273.15) \text{ } ^\circ\text{K}$$

DEFN. BASED ON H_2O :

FREEZES AT $0 \text{ } ^\circ\text{C}$ ($p = 1 \text{ ATM}$)

BOILS AT $100 \text{ } ^\circ\text{C}$

\Rightarrow TRIPLE PT: $0.01 \text{ } ^\circ\text{C}$ (NOTE: p DIFF)

LOWEST TEMP. (ABS ZERO): $-273.15 \text{ } ^\circ\text{C}$

MORE RATIONAL:

$$T \equiv \frac{1}{\beta} = \left(\frac{\partial \ln \Omega}{\partial E} \right)^{-1} \quad (\text{w/ } k=1)$$

\Rightarrow UNITS OF E (OK UNLESS)

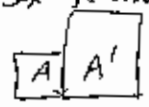
THEN

$$pV = \nu N_a T \quad \Rightarrow \text{DON'T NEED REF.}$$

just say
 WILL SEE: PROVIDES CONVENIENT, MEAS. BLE
 MACRO VAR. WHICH GIVES INFO ON.
 S, WHICH COUNTS MICRO INFO
 ⇒ BEHAVIOR AT LOW T GAVE FIRST
 CLUE THAT CLASS. PHYS WRONG → NEED QM

HEAT CAPACITY:

ANOTHER USEFUL MACRO PARAM RELATING Q, T:
 ex problem



IF $T_A < T_{A'}$

$dQ_A > 0$

T_A INCR $T_{A'}$ DECR UNTIL $T_A = T_{A'}$

(smaller system will usually have
 T chg more)

FIND
 FINAL T? (from properties of materials)

step

HEAT CAPACITY:

not a deriv, just small amt of
 dQ dev. by small chg in T

$C_y \equiv \left(\frac{dQ}{dT} \right)_y$

{ how much heat abs.
 for given chg in T

use small amts because it's not const

PARAMETER HELD FIXED (ex p or V)

⇒ IN GENERAL $C_y = C_y(T, y)$ {OR $C_y(\bar{E}, y)$ }

(mention once give \bar{E}, y for sys, completely
 specifies; can compute T. Could chg var then $\bar{E} \rightarrow T$; easier to measure)



SPECIFIC HEAT: DIV. OUT AMT OF MATERIAL
 SO GIVES PROPERTY OF " , NOT AMT

PER MOLE:

$C_y \equiv \frac{1}{\nu} C_y$
 ↑
 # MOLES

PER GRAM:

$C'_y \equiv \frac{1}{m} C_y$
 ↑
 MASS

(usually obvious
 which one by units)

ASIDE:

EXTENSIVE VS INTENSIVE PARAMS.

USEFUL TO CATEGORIZE MACRO PARAMS BY
DEPEND. ON SYSTEM SIZE:

EXTENSIVE:- INCR \propto SYS. SIZE

ex $V, \bar{E}, S, C_y, M(\text{MASS}), \nu$ (# MOLES)

~ IF DOUBLE SIZE OF SYS, THESE DOUBLE

{ how to know S is extensive?

$$[I] \Omega_{\text{tot}} = \Omega_A + \Omega_{A'} \quad S_{\text{tot}} = S_A + S_{A'}$$

INTENSIVE:- INTRINSIC PROPERTIES OF SYS \rightarrow IND. OF SIZE

ex \bar{p}, T, ρ, c_y

NOTE:

$$\frac{\text{EXT.}}{\text{EXT.}} \rightarrow \text{INT.}$$

$$(\text{ex } \rho = \frac{M}{V})$$

$$c_y = \frac{G_A}{\nu}$$

etc

(like units, eqns must balance)

C_y DEP. ON WHICH PARAM FIXED

ex SPS MACRO SYSTEM SPECIFIED BY \bar{E}, V (or T, V):

- (a) V FIXED: $dQ = d\bar{E}$
(CLAMPED PISTON) (no work)



- (b) \bar{p} FIXED: PISTON (w/ FIXED WT) FREE

$$dQ = d\bar{E} + \bar{p} dV$$

$\Rightarrow T$ dep. ^{strongly} in gen. on \bar{E} (ideal gas \Rightarrow only on T)

- (a) ~~GIVEN~~ dQ ALL GOES TO $d\bar{E} \Rightarrow$ RAISES T

- (b) SOME dQ TO $d\bar{E}$
SOME TO WORK

$d\bar{E}$ LESS \Rightarrow RAISES T LESS

EXPECT

\Rightarrow

$$C_p > C_v$$

CAN RELATE TO ENTROPY:

$$dS = \frac{dQ}{T} \quad (9-5)$$

$$\Rightarrow C_y = \left(\frac{dQ}{dT} \right)_y = T \left(\frac{\partial S}{\partial T} \right)_y$$

is a deriv;
partial since y
held fixed

THINKING OF $S(T, y)$
RATHER THAN $S(\bar{E}, y)$

ALSO,
FOR

$$C_v, \quad dQ = d\bar{E}$$

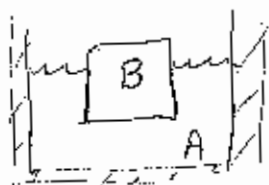
$$\Rightarrow C_v = T \left(\frac{\partial S}{\partial T} \right)_v = \left(\frac{\partial \bar{E}}{\partial T} \right)_v$$

USING C:

(1) FIND T_f

~~JK~~ PUT BLOCK OF CU IN H_2O

\Rightarrow LET V CHG, BUT DO IN ROOM $\Rightarrow \bar{p} = 1 \text{ ATM}$



INIT TEMPS: T_A, T_B

AMT: m_A, m_B

(GRAM) SP. HT: $C_{p,A}', C_{p,B}'$

(properties of H_2O, Cu)

(easier for Eq, solids to keep \bar{p} fixed, not V.)

$$Q_A + Q_B = 0$$

$$Q = \int \delta Q = \int C(T) dt = \int m C'(T) dt$$

$$0 = \int_{T_A}^{T_f} m_A C_{p,A}'(T) dT + \int_{T_B}^{T_f} m_B C_{p,B}'(T) dT$$

\Rightarrow CAN SOLVE FOR T_f

(stop writing !!)

SPECIAL CASE: $C_{p,A}' \sim \text{CONST}$ (FOR BOTH) (over this range)

$$m_A C_{p,A}' (T_f - T_A) + m_B C_{p,B}' (T_f - T_B) = 0$$

$$\Rightarrow T_f = \frac{m_A C_{p,A}' T_A + m_B C_{p,B}' T_B}{m_A C_{p,A}' + m_B C_{p,B}'}$$

(note if m_B or $C_{p,B}'$ large
 $T_f \sim T_B$)

(2) COMPUTE ΔS :

not a well-defined integral; more a meas. prescription
 ↙ ↘
 now it's well-def'd

Q-S PROCESS

$$S(T_b) - S(T_a) = \int_a^b \frac{dQ}{T} = \int_{T_a}^{T_b} \frac{C_y(T) dT}{T}$$

SPECIAL CASE:

IF $C_y(T) \sim \text{CONST}$ OVER THIS RANGE (NOT GENERALLY TRUE)
 $\sim C_y \ln(T_b/T_a)$ (why not? consider $T \rightarrow 0$)

something fishy here:

THIS EX: NOT Q-S

- ONLY IN EQUIL AT BEG & END
- FLOW OF E DEPENDS ON ^{MICRO} DETAILS
- NO WELL-DEFINED T, C(T) DURING HEAT TRANS.

BUT USE YOUR IMAGINATION: EQUIL

- $S(T_a, b)$ PROPERTY OF MACRO SYS. AT ENDS — $T = \frac{\partial S}{\partial E}$ (same $\frac{\partial S}{\partial E}$ every instance, so all have same T)

- ΔS SAME IF FOLLOW Q-S ROUTE ^{FOR EACH SEPARATELY} (VIA SEQUENCE OF HT. RESERVOIRS) \Rightarrow pretend each A, B got there via Q-S process

PREVIOUS EXAMPLE:

$$\Delta S_A = S_A(T_f) - S_A(T_A) = \int_{T_A}^{T_f} \frac{m_A C_A' dT}{T} = m_A C_A' \ln \frac{T_f}{T_A}$$

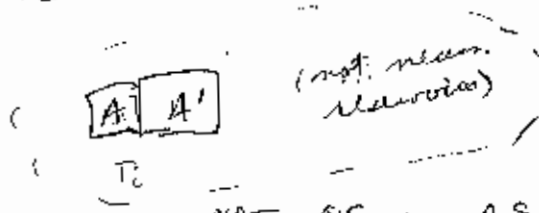
$$\Delta S_A + \Delta S_B = m_A C_A' \ln \frac{T_f}{T_A} + m_B C_B' \ln \frac{T_f}{T_B}$$

= REIF SHOWS $\Delta S_A + \Delta S_B > 0$ FOR ANY T_A, T_B (AS IT MUST BE FOR ANY Q-S PROCESS ON ISOLATED SYS)
 (have T_f above)



only if necessary

NOTE



CLOSED SYS

compute ΔS separately for each as if Q-S

NOT QS ; $\Delta S \geq 0$

BUT FOR A, IF FINAL STATE GIVEN BY SAME E, V (OR T, V) DOESN'T MATTER HOW GOT THERE \rightarrow SAME S

FOCUS ON A alone:

CLOSED SYS



$\Delta S = 0$

Then do same for A'



(3) FIND HEAT OF TRANSFORMATION:

OR WHITE / GRAY TIN (Sn)

$T > T_0 = 292^{\circ}K$ (~ RM TEMP)

WHITE \Rightarrow METAL

$T < T_0$

GRAY \Rightarrow SEMICONDUCTOR

AT $T \gtrsim T_0$ GRAY \rightarrow WHITE, REQUIRING $Q_D \equiv$ "HEAT OF TRANSFORMATION" (PHASE TRANS) (like melting ice)

BUT: VERY SLOW

- CAN TREAT BOTH AS \approx EQUIL TO LOW T
- CAN MEASURE $C_p^{(w)}(T)$ & $C_p^{(g)}(T)$ BOTH DOWN TO $T \sim 0$

(note: statistical reasoning works either if system reacts very quickly to changes a part of it very slowly; either case, stay close to equil)

2 very diff T scales:
 (1) usual trans of E \rightarrow fast \rightarrow reaches equil
 (2) change of phase structures \rightarrow slow

$T \sim 0$ (BUT NOT LOW ENOUGH TO AFFECT NUCL SPINS) (can use our guide $kT \lesssim E_A$ to know when ok)

$S \rightarrow S_0 = k \ln \Omega_s$

$\Omega_s \neq$ NUCL. SPIN STATES

\Rightarrow SAME NUCLEI $\therefore \Omega_s$ SAME FOR BOTH

} MAIN PT

$\Rightarrow S^{(w)}(T \rightarrow 0) = S^{(g)}(T \rightarrow 0)$

CAN PREDICT Q_D :

1 MOLE; \bar{p} FIXED: $C_p = \nu \bar{C}_p = C_p$

(a) TAKE WHT. FROM 0 TO T_0

$S^{(w)}(T_0) = S^{(w)}(0) + \int_0^{T_0} \frac{C_p^{(w)}(T)}{T} dT$ (HEATS)

(b) TAKE GR. FROM 0 TO T_0 , THEN QS TRANS TO WHT. AT T_0

$S^{(w)}(T_0) = S^{(g)}(0) + \int_0^{T_0} \frac{C_p^{(g)}(T)}{T} dT + \frac{Q_D}{T_0}$

DURING TRANS., T FIXED AT T_0 $\therefore \Delta S = \frac{Q_D}{T_0}$

$$\begin{array}{c} \text{PREDICT} \\ \downarrow \\ \Rightarrow \frac{Q_0}{T_0} = \int_0^{T_0} \frac{[C_p^{(w)}(T) - C_p^{(g)}(T)] dT}{T} \end{array}$$

\Rightarrow WORKS : $Q_0 = 2130 \text{ J}$ (1 MOLE)
(VS 2240 J MEAS DIRECTLY)

\Rightarrow RELATED Q_0 , WHICH DEPENDS ON DIFF IN STRUCTURE OF 2 STATES, TO C_p FOR INDIV. STATES VIA ENTROPY CONSIDERATIONS (ESP: SAME AT $T \sim 0$)