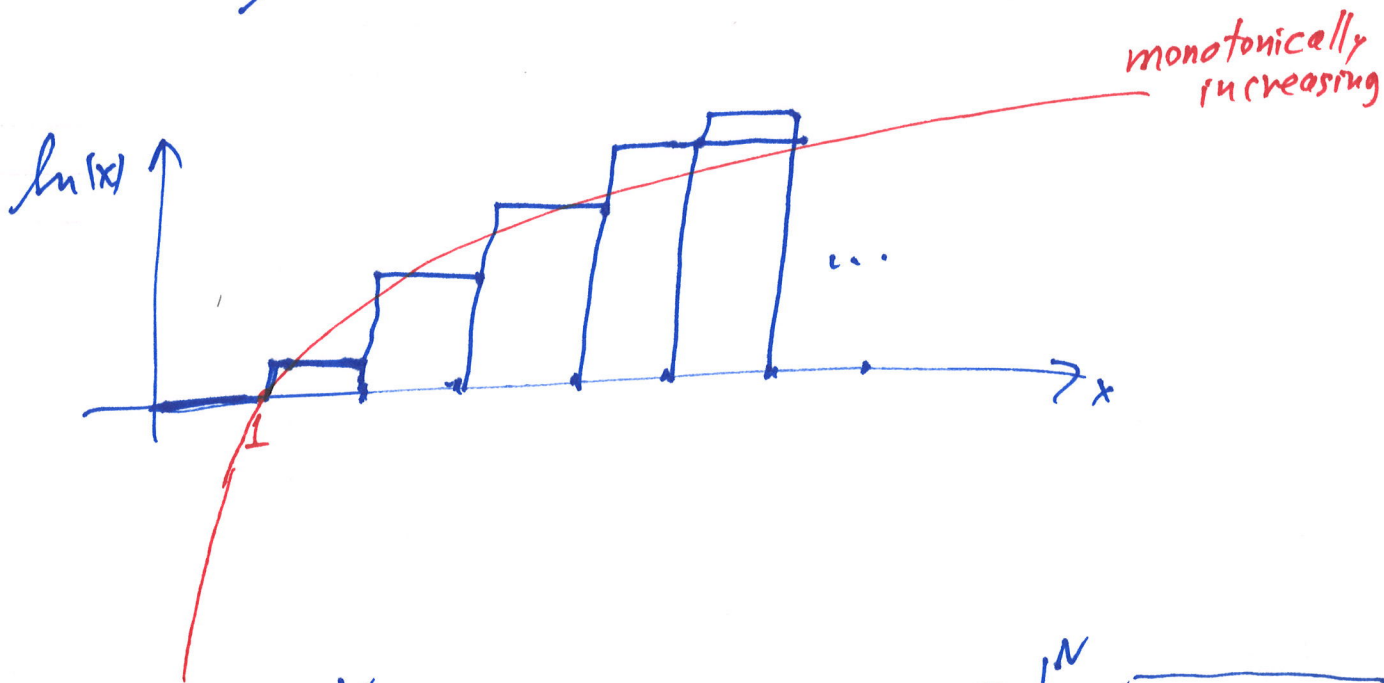


Stirling's Approximation

Need a formula to evaluate $N!$ for large N
eg. $N = N_A = 10^{23}$, $N_A!$ is very large

$$\begin{aligned}\ln(N!) &= \ln[N \cdot (N-1) \cdot (N-2) \cdots 2 \cdot 1] \\ &= \ln(N) + \ln(N-1) + \cdots + \ln(2) + \ln(1)\end{aligned}$$



$$\textcircled{1} \ln(N!) \approx \int_{x=0}^N \ln(x) dx = \left[x \ln(x) - x \right]_0^N = \boxed{N \ln(N) - N}$$

$$\lim_{x \rightarrow 0} x \ln(x) = 0$$

check with L'Hôpital

$$a \ln(b) \neq \ln(b^a)$$

$$\ln(N!) \approx N \ln(N) - N$$

exponentiate both sides

$$e^{-\ln(N!)} = N! = e^{N \ln(N) - N} = e^{N \ln N - N} = e \cdot e^{-N}$$

$$= e^{\ln(N^N)} \cdot e^{-N} = N^N \cdot e^{-N} = \left(\frac{N}{e}\right)^N$$

$$N! \approx \left(\frac{N}{e}\right)^N$$

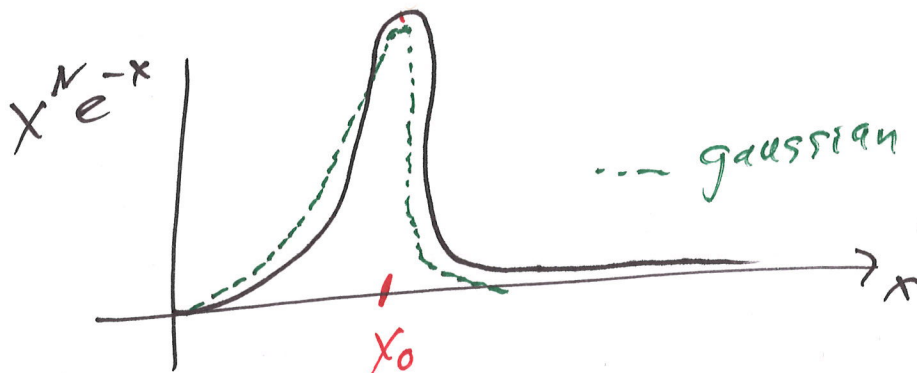
0th approximation
usually good enough

First correction - use gaussian approximation to the integral

$$N! = \int_{x=0}^{\infty} \underline{x^N \cdot e^{-x}} dx \equiv \Gamma(N+1) \quad \leftarrow \begin{array}{l} \text{Gamma} \\ \text{function} \end{array}$$

$$0! = 1, \quad \left(-\frac{1}{2}\right)! = \sqrt{\pi} = \Gamma\left(\frac{1}{2}\right)$$

Integrand $x^N e^{-x}$ $\left\{ \begin{array}{l} x^N \text{ grows rapidly with } x \\ e^{-x} \text{ shrinks } \textit{more} \text{ rapidly with } x \end{array} \right.$



$$\frac{d}{dx} [x^N e^{-x}] \Big|_{x=x_0} \stackrel{!}{=} 0 = N x^{N-1} e^{-x} + x^N e^{-x} \Big|_{x=x_0}$$

$$= e^{-x_0} x_0^{N-1} (N - x_0) = 0 \Rightarrow \boxed{x_0 = N}$$

$$\ln(x^N e^{-x}) = \ln(x^N) + \ln(e^{-x}) = N \ln(x) - x$$

change variables $y = x - N$ distance from peak
 $dy = dx$

$$N \ln(x) - x = N \ln(y + N) - y - N$$

$$= N \ln \left[N \left(1 + \frac{y}{N} \right) \right] - y - N$$

$$= N \ln(N) - N + N \ln \left(1 + \frac{y}{N} \right) - y$$

Near the peak, $x \approx N$, $\frac{y}{N} \ll 1$

Taylor Expansion: $\ln \left(1 + \frac{y}{N} \right) = \frac{y}{N} - \frac{1}{2} \left(\frac{y}{N} \right)^2 + \frac{1}{3} \left(\frac{y}{N} \right)^3 - \dots$

$$N \ln(x) - x = N \ln(N) - N - \frac{1}{2} \frac{y^2}{N} + \mathcal{O} \left(\frac{y^3}{N} \right)$$

$$x^N e^{-x} = \exp(N \ln(x) - x) \approx \exp\left[N \ln(N) - N - \frac{y^2}{2N} + \dots\right]$$

$$N! = \int_{x=0}^{\infty} x^N e^{-x} dx \approx \int_{y=-N}^{\infty} N^N e^{-N} e^{-\frac{y^2}{2N}} dy + \dots$$

\nearrow Replace $-N$ by $-\infty$

$$N! \approx \int_{-\infty}^{+\infty} N^N e^{-N} e^{-\frac{y^2}{2N}} dy = N^N e^{-N} \int_{-\infty}^{+\infty} e^{-\frac{y^2}{2N}} dy$$

$\underbrace{\hspace{10em}}_{\text{Gaussian integral}}$
 $= \sqrt{\pi 2N}$

$$N! \approx N^N e^{-N} \sqrt{2\pi N} \left(1 + \frac{1}{12N}\right)$$

Gaussian Integrals

$$I = \int_{x=-\infty}^{+\infty} e^{-x^2} dx = \text{number (not a function of } x)$$

$$I^2 = \left(\int_{x=-\infty}^{+\infty} e^{-x^2} dx \right) \left(\int_{y=-\infty}^{+\infty} e^{-y^2} dy \right)$$

Now consider $x+y$ as Cartesian coordinates.

change variables to polar coordinates

$$I^2 = \int_{x=-\infty}^{+\infty} \int_{y=-\infty}^{+\infty} e^{-(x^2+y^2)} dx dy$$

$$x = r \cos \varphi$$

$$y = r \sin \varphi$$

$$r = \sqrt{x^2 + y^2}$$

$$\varphi = \arctan\left(\frac{y}{x}\right) (+180^\circ)$$

↑
maybe

$$I^2 = \int_{r=0}^{\infty} \int_{\varphi=0}^{2\pi} e^{-r^2} r dr d\varphi$$

↑
Jacobian

$$= \left(\int_{r=0}^{\infty} e^{-r^2} r dr \right) \left(\int_{\varphi=0}^{2\pi} d\varphi \right)$$

2π

$$I^2 = 2\pi \left[-\frac{e^{-r^2}}{2} \right]_0^\infty = 2\pi \left[0 - \left(-\frac{1}{2}\right) \right]$$

$$\frac{d}{dr} \left(-\frac{e^{-r^2}}{2} \right) = -\frac{1}{2}(-2r)e^{-r^2} = r e^{-r^2}$$

$$I^2 = \pi \Rightarrow I = \sqrt{\pi}$$

$$I = \int_{-\infty}^{+\infty} e^{-x^2} dx = \sqrt{\pi}$$

change variable

$$x = \sqrt{a} z$$

$$x^2 = a z^2$$

$$dx = \sqrt{a} dz$$

$$= \sqrt{a} \int_{-\infty}^{\infty} e^{-a z^2} dz = \sqrt{\pi}$$

Rename $z \rightarrow x$

$$\int_{-\infty}^{+\infty} e^{-ax^2} dx = \sqrt{\frac{\pi}{a}}$$

$\leftarrow \frac{d}{da}$ both sides

$$\int_{-\infty}^{+\infty} (-x^2) e^{-ax^2} dx = \sqrt{\pi} \frac{d}{da} (a^{-1/2}) = \sqrt{\pi} \left(-\frac{1}{2}\right) a^{-3/2}$$

$$\int_{-\infty}^{+\infty} x^2 e^{-ax^2} dx = \frac{1}{2} \left(\frac{\pi}{a^3} \right)^{1/2}$$

later, reuse, repeat

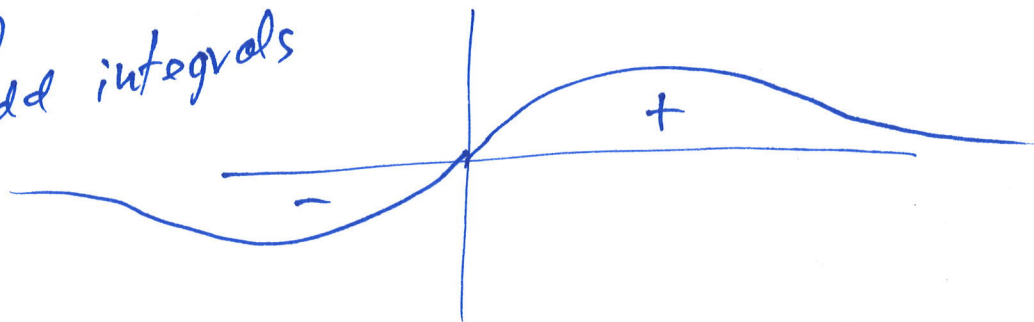
$$\frac{d^n}{da^n}$$

⋮

$$\int_{-\infty}^{+\infty} x^{2n} e^{-ax^2} dx = ?$$

$$\int_{-\infty}^{+\infty} x^{2n+1} e^{-ax^2} dx = \emptyset$$

↑
odd integrals



$$\int_0^{\infty} x e^{-x^2} dx = \frac{1}{2}$$

$$x \rightarrow \sqrt{a} z$$

$$\frac{d^n}{da^n} \text{ under } \int \text{ sign.}$$

Hyper area + hypervolumes

$$\int_{x_1=-\infty}^{+\infty} e^{-x_1^2} dx_1 \cdot \int_{x_2=-\infty}^{+\infty} e^{-x_2^2} dx_2 \cdots \int_{x_n=-\infty}^{+\infty} e^{-x_n^2} dx_n = \prod_{k=1}^n \int_{x_k=-\infty}^{+\infty} e^{-x_k^2} dx_k$$

$\underbrace{\hspace{10em}}_{\sqrt{\pi}} \quad \underbrace{\hspace{10em}}_{\sqrt{\pi}} \quad \dots \quad \underbrace{\hspace{10em}}_{\sqrt{\pi}}$

$$= (\sqrt{\pi})^n = \pi^{n/2} = \int d\Omega \int_{r=0}^{\infty} e^{-r^2} r^{n-1} dr$$

$\underbrace{\hspace{10em}}_{\text{solid angle}} \quad \underbrace{\hspace{10em}}_{\frac{1}{2} \Gamma(\frac{n}{2})}$

$$r^2 = x_1^2 + x_2^2 + \dots + x_n^2$$

$$\Omega_{n-1} = \text{(n-1)-dimensional solid angle} = \frac{2\pi^{n/2}}{\Gamma(\frac{n}{2})} = \frac{2\pi^{n/2}}{(\frac{n}{2}-1)!}$$

Surface area of an (n-1)-dimensional sphere.




$$S_{n-1} = \Omega_{n-1} R^{n-1} \rightarrow S_m = \Omega_m R^m = \frac{2\pi^{\frac{m+1}{2}} R^m}{(\frac{m-1}{2})!}$$




n -Sphere = n -dimensional manifold
 all points the same distance R from origin

e.g. Soap bubble is 2-sphere.

n -ball = n -dimensional manifold
 all points at or less than distance R from origin.

e.g. Earth is 3-ball

n	S_n	
0	2	 line segment = 1-ball
1	$2\pi R$	 circle = 1-sphere disk = 2-ball
2	$4\pi R^2$	 soap bubble = 2-sphere interior of 2-sphere is 3-ball
3	$2\pi^2 R^3$	

n	V_n	Volume of n -ball
0	1	•
1	$2R$	 line segment <u>length</u>
2	πR^2	 disk <u>area</u>
3	$\frac{4}{3}\pi R^3$	 baseball <u>volume</u>
4	$\frac{1}{2}\pi^2 R^4$	hypersphere hypervolume — glome hypervolume

$$V_n = \int_{r=0}^R S_{n-1}(r) dr$$

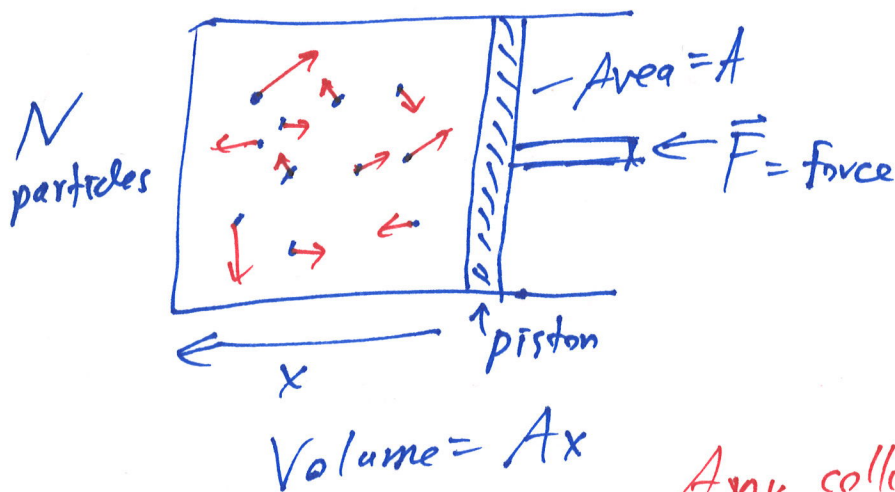


$$= \frac{S_{n-1} R}{n} = \frac{2\pi^{n/2} R^n}{n(\frac{n}{2}-1)!}$$

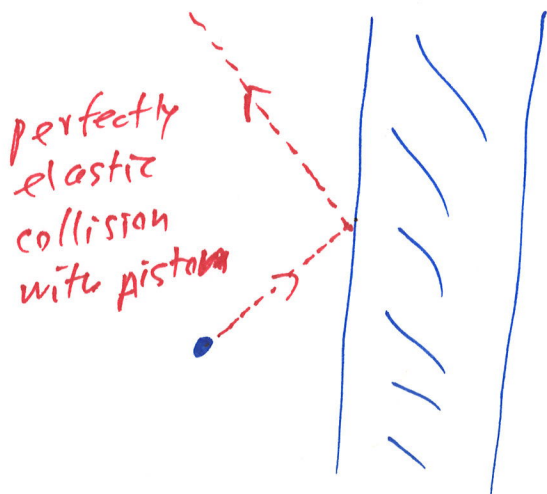
Kinetic Theory of Gases

Atoms + Newton's Laws + Maxwell's Eqs

Not Quantum Mechanics



$$\text{Pressure } P = \frac{F}{A}$$



Any collision conserves linear momentum
Elastic collision conserves kinetic energy

⇒ same speed before and after the collision.

$$p_y \rightarrow p_y, \quad p_z \rightarrow p_z, \quad p_x \rightarrow -p_x$$

$$\begin{aligned} \text{Momentum delivered to piston} &= 2p_x \\ &= 2mv_x \end{aligned}$$