

$$F(\omega) = \int_{-\infty}^{+\infty} \frac{1}{2\pi} \left[\int_{-\frac{T}{2}}^{\frac{T}{2}} e^{-i\omega t'} F(t') dt' \right] e^{i\omega t} dt$$

forward

$$\tilde{F}(\omega) \equiv \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-i\omega t'} F(t') dt'$$

notice signs are diff.

$$F(t) = \int_{-\infty}^{+\infty} \frac{1}{\sqrt{2\pi}} e^{i\omega t} \tilde{F}(\omega) d\omega = F(t)$$

reverse

Fourier & Reverse Fourier Transformations

Important bc you can start w/ a difficult time-domain prob. that becomes easy in freq. domain!

Time-Domain

F.T. Frequency-Domain

- F(t) that solves some time-dependent diff. eq.

- now it's just an algebraic eqn (F(ω))

get your answer back in the time-domain

convert back inverse F.T.

$$\tilde{F}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-i\omega t} F(t) dt$$

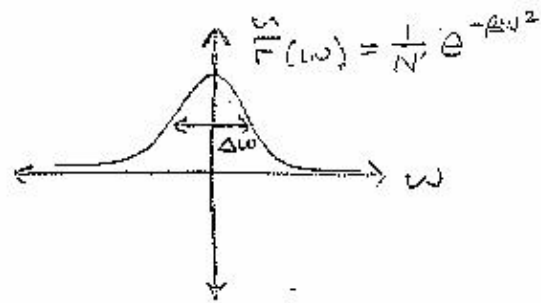
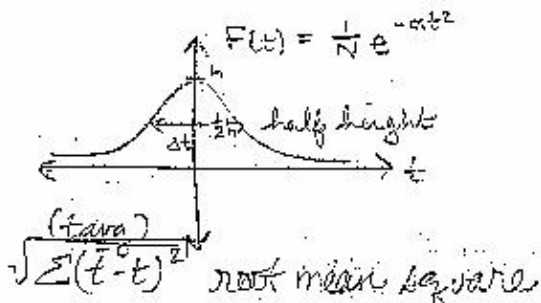
Fourier Transform

$$f(s) = N \int e^{-st} F(t) dt \quad \omega \equiv s$$

Laplace Transform

normalization → no hw or test problems

* dirac delta function introduced next time



as $\Delta t \downarrow, \Delta w \uparrow$
 & as $\Delta t \uparrow, \Delta w \downarrow$ } where $\Delta t \cdot \Delta w = C$
 Some constant

for Gaussians: $\Delta t \cdot \Delta w \geq \frac{1}{2}$

when $t = \text{time}$ $w = \text{angular frequency}$

$hw = \text{quantum mechanical energy}$
 $\therefore \Delta t (\frac{1}{h} \Delta w) \geq \frac{1}{2}$ $\Delta t \cdot \Delta E \geq \frac{h}{2}$

Heisenberg Uncertainty Relationship

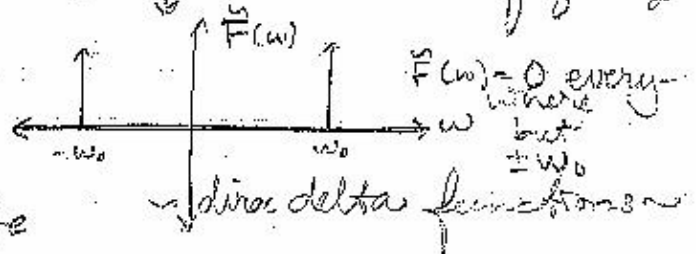
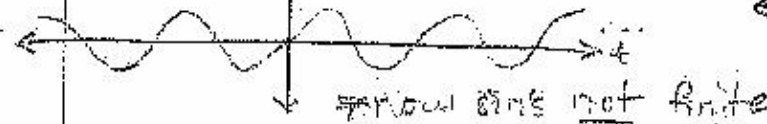
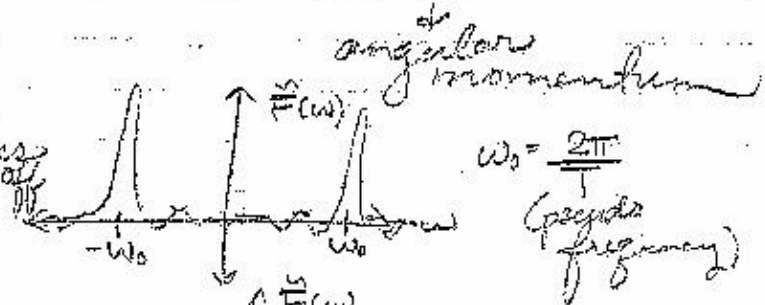
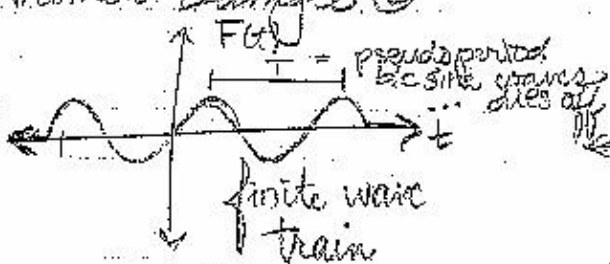
$\{t, E\}$ complimentary observables

* you can only measure one or the other, but not both at the same time

Some other complimentary observables:

$\{x, p_x\}, \{y, p_y\}, \{z, p_z\} \{L, \theta\}$

Another Example (C)



$$F(t) = \frac{1}{2\pi} \int_{\omega=-\infty}^{+\infty} \int_{t'=-\infty}^{+\infty} e^{-i\omega t'} F(t') dt' e^{i\omega t} d\omega$$

$\omega \rightarrow -z$
 $d\omega \rightarrow -dz$

-changing the minus sign \rightarrow taking complex conjugates

$$F(t) = \frac{1}{2\pi} \int_{z=-\infty}^{+\infty} \int_{t'=-\infty}^{+\infty} e^{izt'} F(t') dt' e^{-izt} (-dz)$$

use this negative to flip integration limits \rightarrow then let $z = \omega$ & $dz = d\omega$

Gives You:

$$F(t) = \frac{1}{2\pi} \int_{\omega=-\infty}^{+\infty} \int_{t'=-\infty}^{+\infty} e^{+i\omega t'} F(t') dt' e^{-i\omega t} d\omega$$

rewritten as:

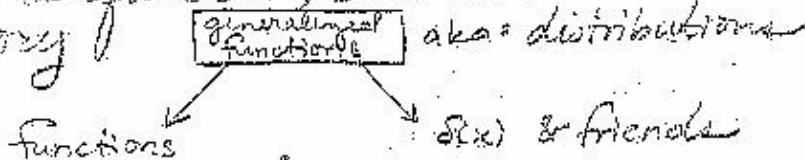
$$F(t) = \int_{t'=-\infty}^{+\infty} \textcircled{?} F(t') dt'$$

acts as delta function for stuff on left

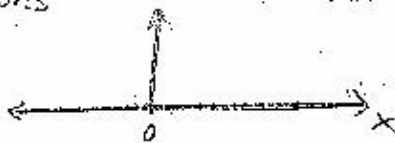
$$\textcircled{?} = \frac{1}{2\pi} \int_{\omega=-\infty}^{+\infty} e^{i\omega(t'-t)} d\omega = \delta(t'-t) = \delta(x)$$

Dirac Delta "Function"

$\delta(x)$ is not quite a function but almost \rightarrow new category

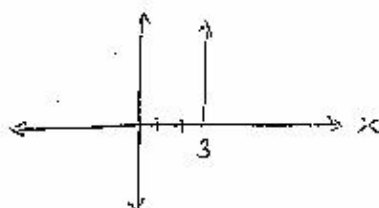


$$\delta(x) = \begin{cases} 0 & x \neq 0 \\ \infty & x = 0 \end{cases}$$

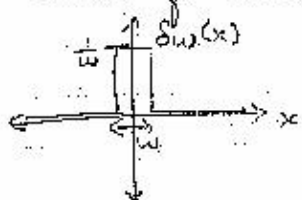


$$\int_{x=-\infty}^{+\infty} \delta(x) dx = 1$$

$$\delta(x-3) =$$



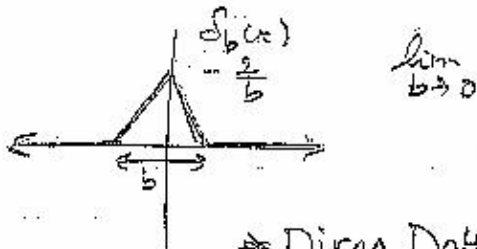
Think of $\delta(x)$ as the limit of:



$$\int_{-w}^w dx = \text{Area} = w(w) = 1$$

then take limit as $w \rightarrow 0$

* you can use any shape that peaks around zero



Dirac Delta function is the Continuum Analog of Kronecker Delta

$$f_i = \sum_{n=1}^{\infty} \delta_{in} f_n = 0f_1 + 0f_2 + \dots + 1f_i + 0f_{i+1} + \dots = \underline{f_i}$$

Kronecker delta

$$f(x) = \int_{-\infty}^{+\infty} \delta(x-y) f(y) dy$$

or just something large.

scales

$$\int_{-\infty}^{+\infty} \delta(ax) f(x) dx \quad \text{let } ax = y \quad a dx = dy$$

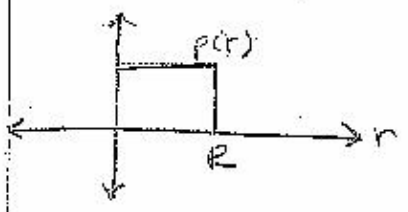
$$\frac{1}{|a|} \int_{-\infty}^{+\infty} \delta(y) f\left(\frac{y}{a}\right) dy = \frac{1}{|a|} f\left(\frac{0}{a}\right)$$

as peak is @ zero now

Delta "Function" was invented by Dirac to describe the charge density of a pt. particle at the origin

(a) (where you don't need the delta function)
 THE PROTON: $R \approx 10^{-15} \text{ m} = 1 \text{ fermi}$
 model of solid sphere - uniform charge
 (only depends on distance from charge)

charge density $\rho(r) = \begin{cases} \frac{e}{\frac{4}{3}\pi R^3} & r < R \\ 0 & r > R \end{cases}$ (outside the proton)



(b) electron (point-like) ~ no volume (up to 10^{-18} m)
 experimental error

$\rho(r) = \begin{cases} 0 & r > 0 \\ \infty & r = 0 \end{cases}$
 b/c it's like $\frac{-e}{0(\text{volume})} = \infty$

but, $\iiint_V \rho(r) dV = \text{total charge} = -e$
 works b/c they're variables you can't name $\delta = \delta(x)$

$\rho(r) = -e \delta^{(3)}(\vec{r}) = -e \delta(x) \delta(y) \delta(z)$
 product of 3 $\delta(r)$'s... 1 for each coordinate

$\iiint_V \rho(r) dV = \int_x \int_y \int_z [-e \delta(x) \delta(y) \delta(z)] dx dy dz$
 $= -e \int_x \delta(x) dx \int_y \delta(y) dy \int_z \delta(z) dz = -e$

$\delta(x)$ has dimension $\frac{1}{\text{length}}$

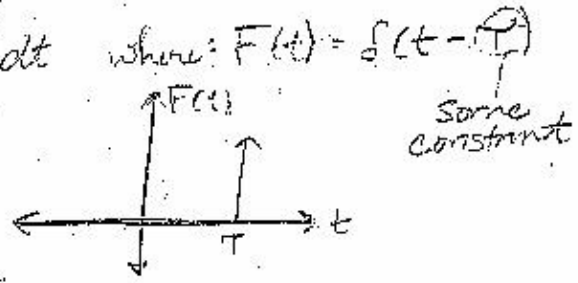
$\int \delta(x) dx = 1$ (as long as integration limits include 0)
 length \times dimensionless = 1
 $\therefore \delta(x)$ must have units of $\frac{1}{\text{length}}$
 ~ don't make 0 a limit of integration though

$\lim_{\epsilon \rightarrow 0} \int_{-\epsilon}^{\epsilon} f(x) dx = 0$ for a real function

Fourier Transform of a Delta Function:

$\hat{F}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{i\omega t} F(t) dt$ where: $F(t) = \delta(t - \tau)$

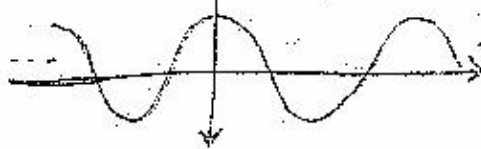
$\hat{F}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{i\omega t} \delta(t - \tau) dt$



Sifting Property:

$\hat{F}(\omega) = \frac{1}{\sqrt{2\pi}} e^{i\omega \tau}$
 or also:

$\hat{F}(\omega) = \frac{1}{\sqrt{2\pi}} [\cos(\omega \tau) + i \sin(\omega \tau)]$
 $\uparrow \text{Re}[\hat{F}(\omega)]$

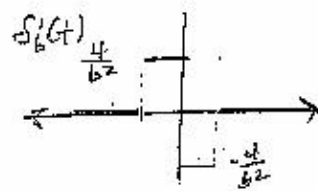
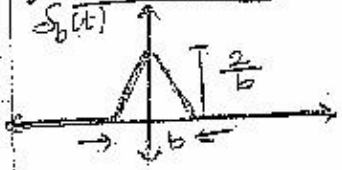


As $\tau \rightarrow 0$, the cos flattens out and eventually forms a line @ 1

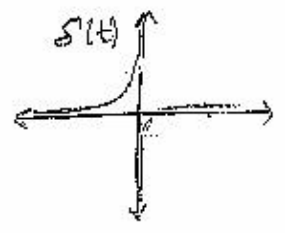
Derivative of Dirac Delta Function:

$S'(t) = \frac{d}{dt} S(t)$

remember:



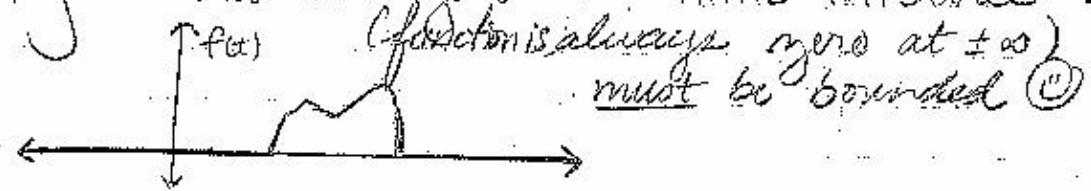
then:



$$\int_{-\infty}^{+\infty} \delta'(t) f(t) dt = - \int_{-\infty}^{+\infty} \delta(t) f'(t) dt = -f'(0)$$

Special Property:

$f(t)$ must have "bounded support", meaning: $f(t)$ is nonzero on a finite interval $[a, b]$



$$d(uv) = v(du) + u(dv)$$

$$\int_a^b d(uv) = \int_a^b v du + \int_a^b u dv$$

& integrate both sides

→ now, let $u = \delta(t)$ & $dv = f(t) dt$ integration by parts

$$\int_{-\infty}^{+\infty} \delta'(t) f(t) dt = \underbrace{\delta'(t) f(t)}_{\text{surface term} = 0} \Big|_a^b - \int_a^b \delta(t) f'(t) dt$$

divergence theorem = Gauss's Law

$$= -f'(0) \equiv - \frac{df}{dt} \Big|_{t=0}$$

EXAMPLE → another distribution

$\sin(\alpha t) \delta'(t)$ is a generalized distribution

SIMPLIFY THIS

→ distributions only make physical sense inside an integral & multiplied by a "test" function, $f(t)$ w/ bounded support

$$\int_{-\infty}^{+\infty} \sin(\alpha t) \delta'(t) f(t) dt$$

↓
test function

$$= \sin(\alpha t) \delta(t) f(t) \Big|_{-\infty}^{+\infty} - \int_{-\infty}^{+\infty} \frac{d}{dt} [\sin(\alpha t) f(t)] \delta(t) dt$$

$$= - \int_{-\infty}^{+\infty} [\alpha \cos(\alpha t) f(t) + \sin(\alpha t) f'(t)] \delta(t) dt \Rightarrow$$

plug in 0 for all t's...

$$= \int_{-\infty}^{\infty} [\alpha \cos(\omega t) f(t) + \sin(\omega t) f'(t)] \delta(t) dt$$

$$= \int_{-\infty}^{\infty} \alpha f(t) \delta(t) dt = \alpha f(0)$$

undergraduate

original line: $\int_{-\infty}^{\infty} \sin(\omega t) \delta'(t) f(t) dt$

can't see the distribution if you go that far... b/c it didn't start just a function

$$\sin(\omega t) \delta'(t) = -\alpha \delta(t)$$

graduate

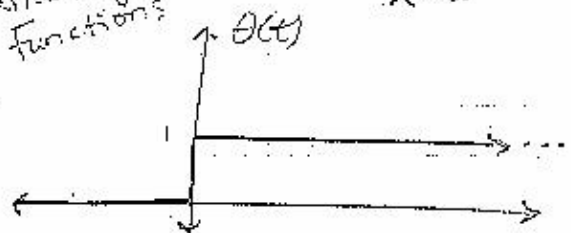
don't report ans. w/ f'(t) b/c you didn't start w/ f'(t)...

→ SIMILAR TO BONUS PROBLEMS ←

Integration of Dirac Delta Function

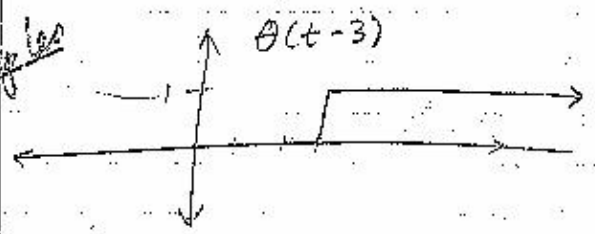
Heavy-Side Functions
Unit Step Functions

$$\theta(t) \equiv \int_{-\infty}^t \delta(x) dx = \begin{cases} 0, & t < 0 \\ 1, & t > 0 \end{cases}$$



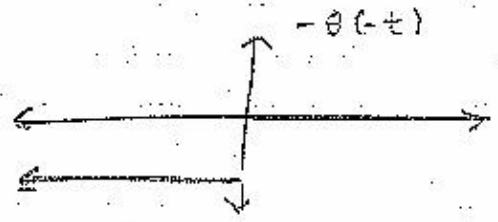
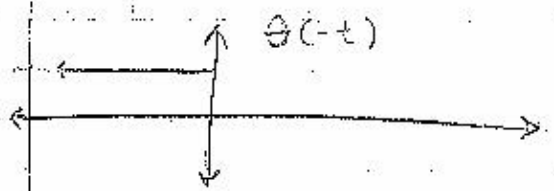
at t=0, you pick what theta(t) equals → some popular choices are ±, 0 & 1/2

Some Examples

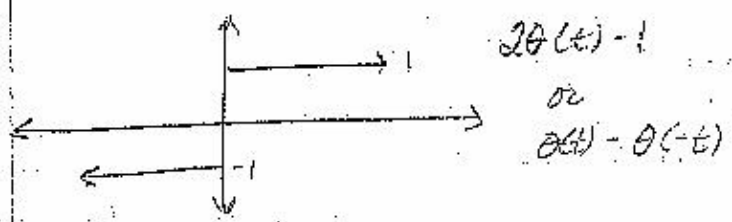


$$t-3 < 0, \theta = 0$$

$$t-3 > 0, \theta = 1$$



homework problem starts at 3 → 6 & at height of 4



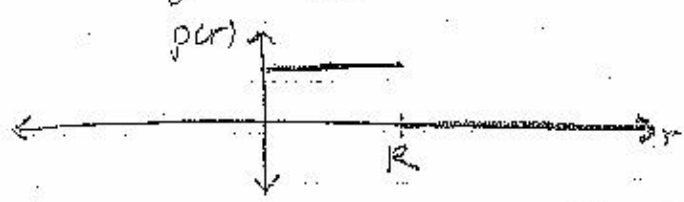
electron: $\rho(\vec{r}) = -e \delta(x) \delta(y) \delta(z)$
 later to be written in spherical coordinates

proton: model (solid ball of charge)
 proton radius / cubical

$$\rho(r) = \begin{cases} \frac{+e}{\frac{4}{3}\pi R^3} & r < R \\ 0 & r > R \end{cases}$$

like $\theta(t)$ function \rightarrow on and then it turns off

$\rho(r) = \frac{+e}{\frac{4}{3}\pi R^3} \theta(R-r)$ rewritten w/ θ , rather than piecewise function



Algebraic Equations: $5-x=2, x^2=4$

Goal: Find the number(s) for x that make the equations true

w/ complex numbers $i^2 = -1$ $E-i, i^3 = \sqrt{-1}$

* good book for history of i

homogeneous ~

Differential Equations: first-order, ordinary, linear

types of diff's ~

$$\frac{df(x)}{dx} = \cos x \quad f'(x) = \cos x \quad f(x) = \sin x + C$$

Goal \rightarrow find function(s), $f(x)$ that make the equation true

* no guaranteed way to do this - bit of guesswork ☺

- Differential Equations -

$$\frac{dy(x)}{dx} - y(x) = 0$$

Goal: find function, $y(x)$ that makes this true $\forall x$ (for all x)

$$y'(x) - y(x) = 0$$

GUESSING... $y(x) = x^2$ $\frac{dy(x)}{dx} = 2x$

$\rightarrow 2x - x^2 \neq 0 \therefore$ it was a bad guess

A better guess... $y(x) = e^x$ $y'(x) = e^x$

$y'(x) - y(x) = e^x - e^x = 0 \forall x$
not most general equation $\rightarrow Ae^x$ is the most general

where A is any constant
Highest derivative tells you # of arbitrary constants
 $\therefore y(t) = Ae^x$ has 1 arbitrary constant
 \therefore it's a first-order differential equation

Order = highest derivative in function, not variable
 $y'(x) - y(x) = 0$ is first-order, linear (in y) & homogeneous

Linear - derivatives of the function $y(x)$, including the 0th derivative occur to the first power.

$\frac{d^2y(x)}{dx^2} + y(x) = 0$ $y'' + y = 0$ 2nd order - linear & homogeneous

$y''(x) + y^2(x) = 0$ non-linear

$y'(x) + \sin(y(x)) = 0$ also non-linear

$[y''(x)]^2 - y'(x) = 0 \rightarrow \left[\frac{dy(x)}{dx}\right]^2 - y(x) = 0$ also non-linear