

Dispersion Relation $\omega(k)$

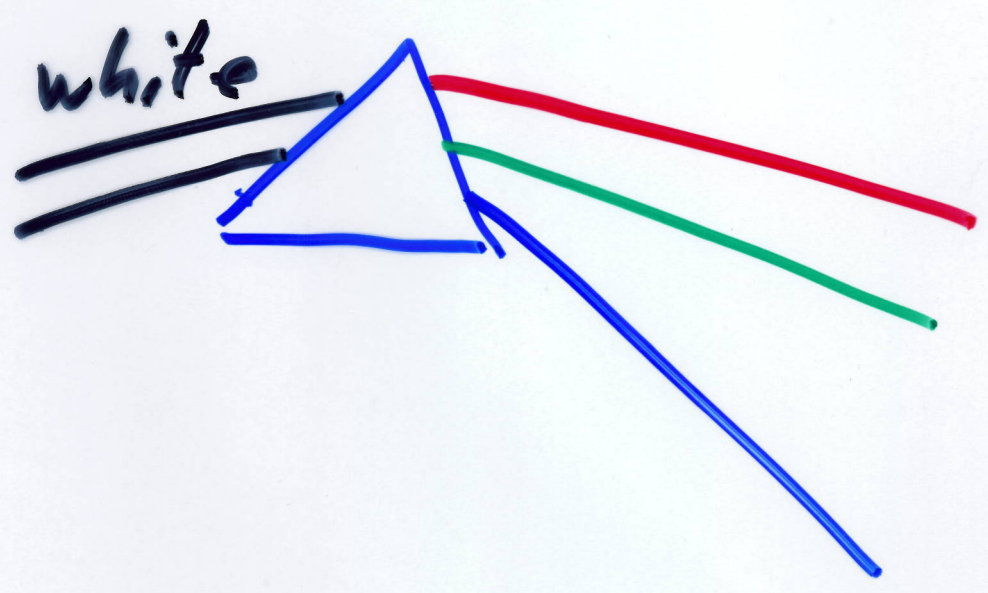
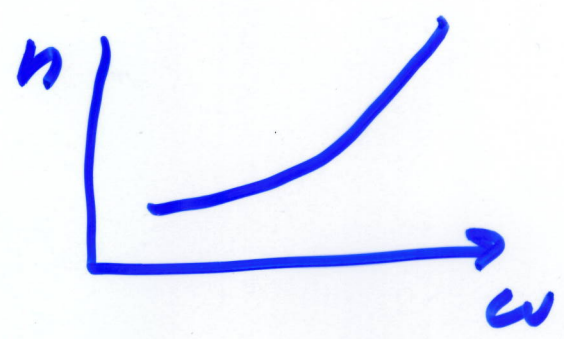
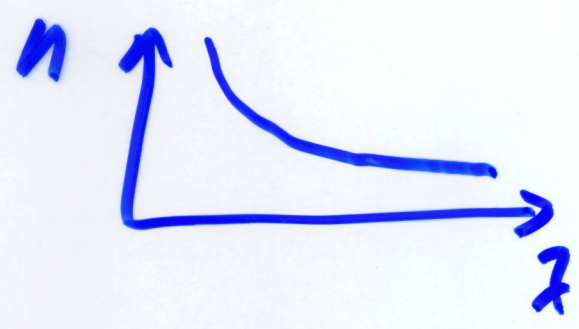
e.g. Light in vacuum

D.R. $\omega = c k$
 c speed of light

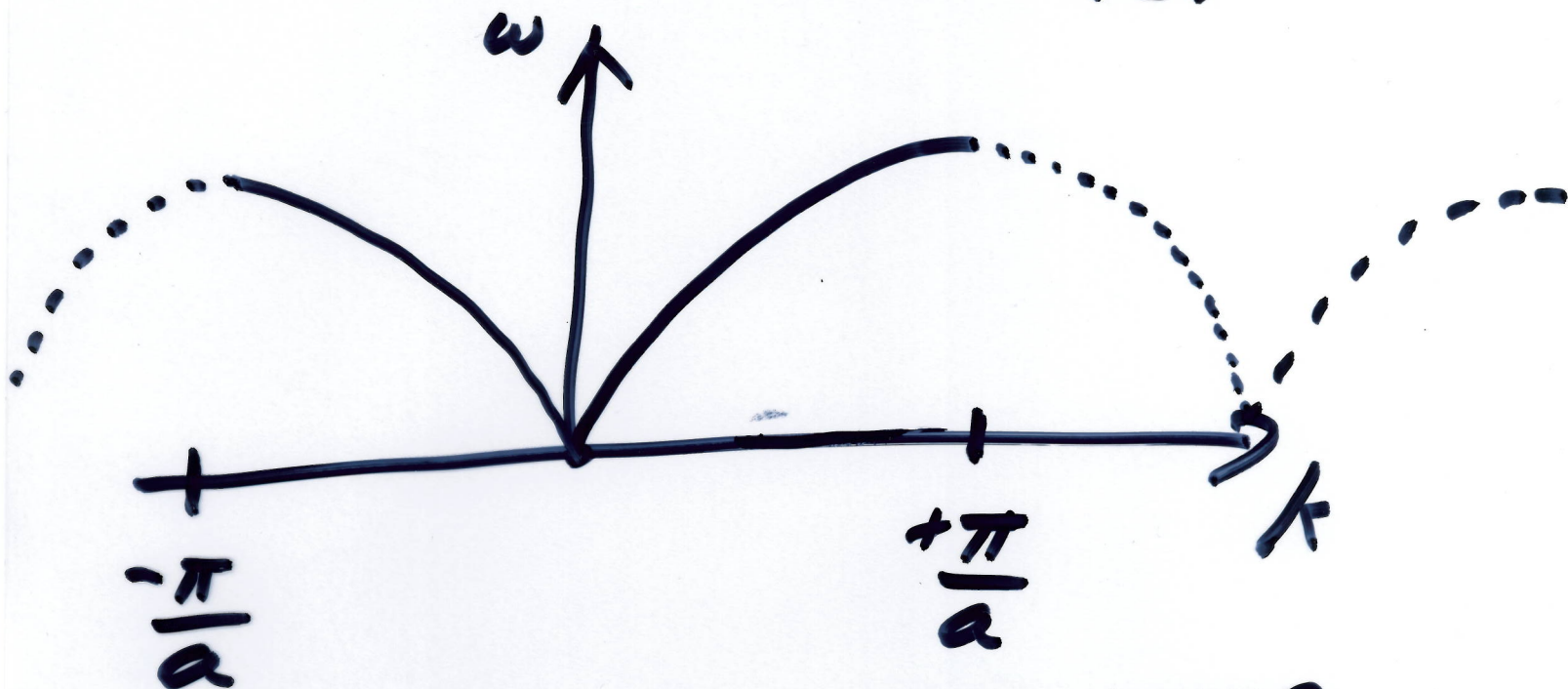
e.g. Light in a material with index of refraction $n(\omega)$

D.R. $\omega = \frac{c}{n(\omega)} k$

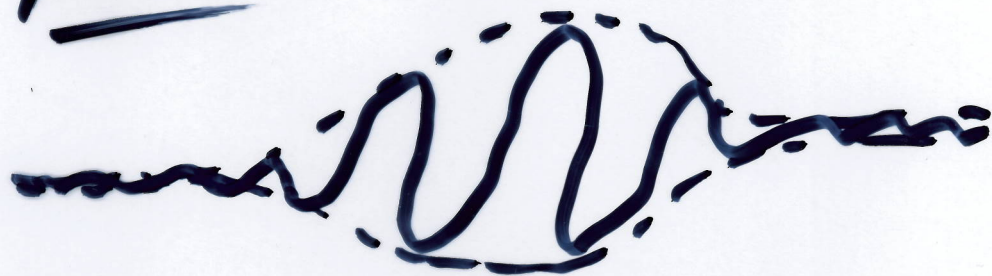
$n=1$ vacuum
 $n \geq 1$



$$\text{D.R. } \omega = 2\sqrt{\frac{\epsilon'}{n}} \left| \sin\left(\frac{ka}{2}\right) \right|$$



Suppose we send a wave packet into the crystal



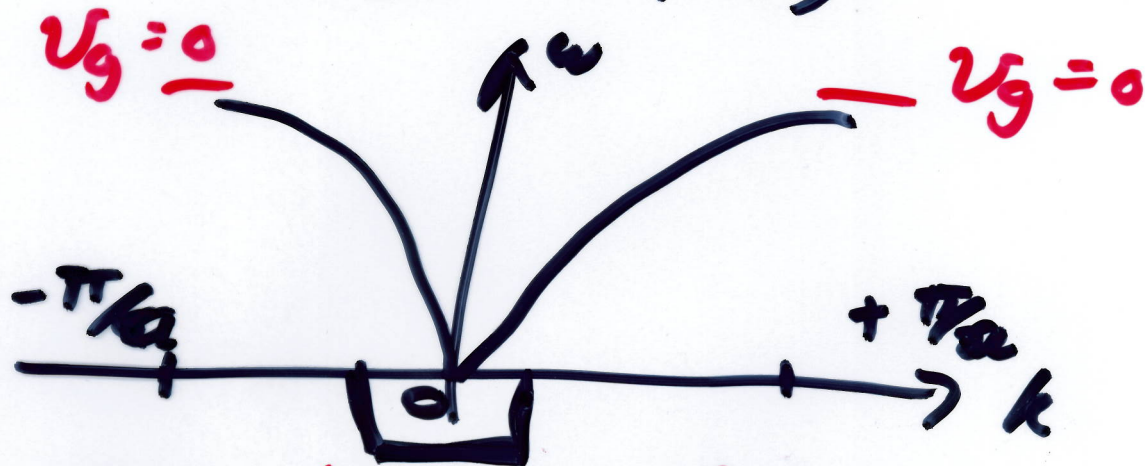
Two velocities

phase velocity $v_p = \frac{\omega}{k}$

group velocity $v_g = \frac{d\omega}{dk}$

$$v_g = \frac{d\omega}{dk} = \frac{d}{dk} \left[2\sqrt{\frac{E}{m}} \sin\left(\frac{ka}{2}\right) \right]$$

$$= a\sqrt{\frac{E}{m}} \cos\left(\frac{ka}{2}\right)$$



v_g is tangent slope



first Brillouin zone

At the zone boundaries $k = \pm \frac{\pi}{a}$



standing wave
 $v_g = 0$

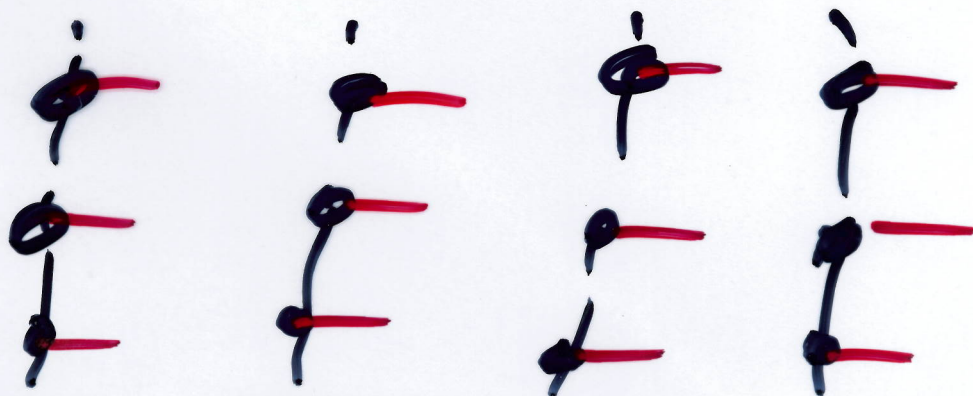
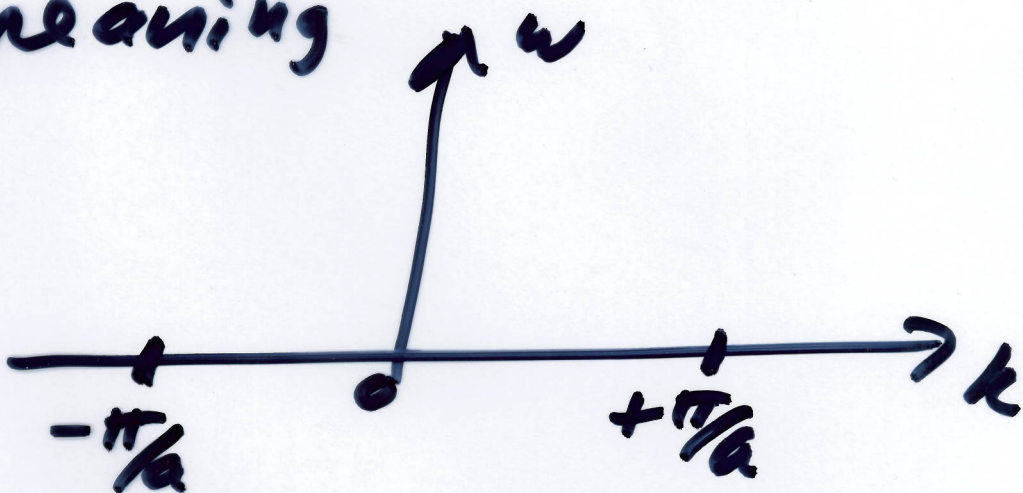
$$\frac{u_{s+1}}{u_s} = \frac{A e^{ika(s+1)}}{A e^{ikas}} = e^{ika}$$

At zone boundary $ka = \pm \pi$

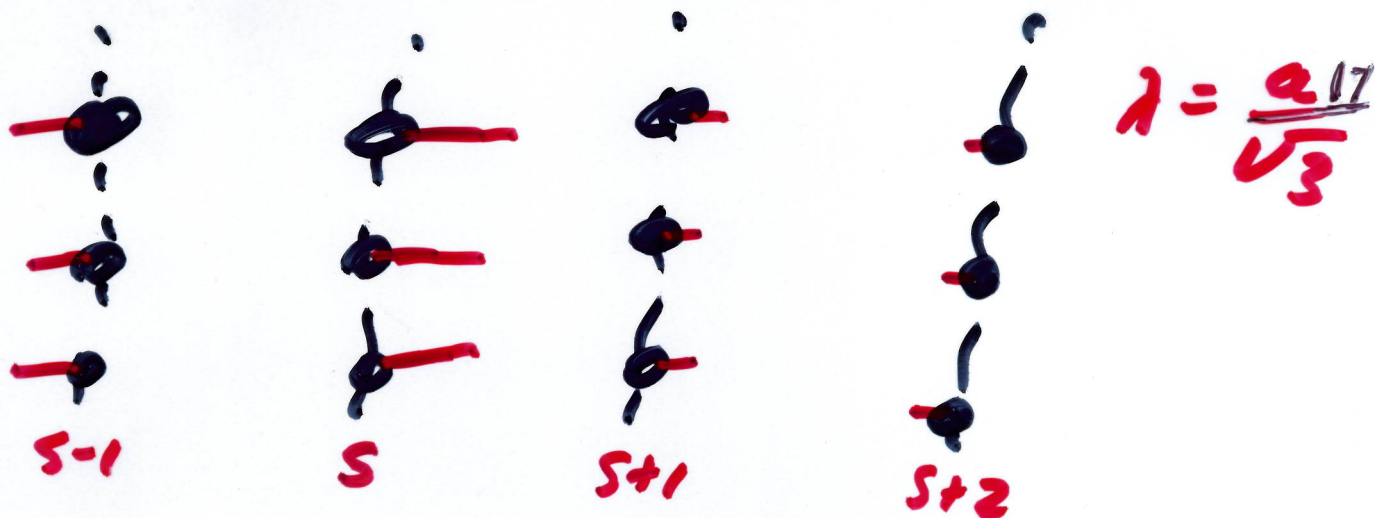
$$\frac{u_{s+1}}{u_s} = e^{\pm i\pi} = -1$$

Only wave vectors k in the first Brillouin have physical meaning

$$k = \frac{2\pi}{\lambda}$$



$\lambda = a$
not allowed



For k 's outside the Brillouin zone
 $|k| > \frac{\pi}{a}$

$$\begin{aligned}
 u_s &= A e^{i k a s} = A e^{i a s (k_n + 2\pi n)} \\
 &= A e^{i a s k_n}
 \end{aligned}$$

Long wavelength limit
 (small k limit)

$$\text{D.R. } \omega^2 = \frac{2C}{m} [1 - \cos(ka)]$$

$$\lambda \gg a, \quad ka \ll 1$$

$$\cos(x) = 1 - \frac{x^2}{2!} + \frac{x^4}{4!} - + \dots$$

$$\omega^2 = \frac{2C}{m} \left[1 - \left(1 - \frac{k^2 a^2}{2} \right) \right] + \dots$$

$$\omega^2 = \frac{2C}{m} \frac{k^2 a^2}{2}$$

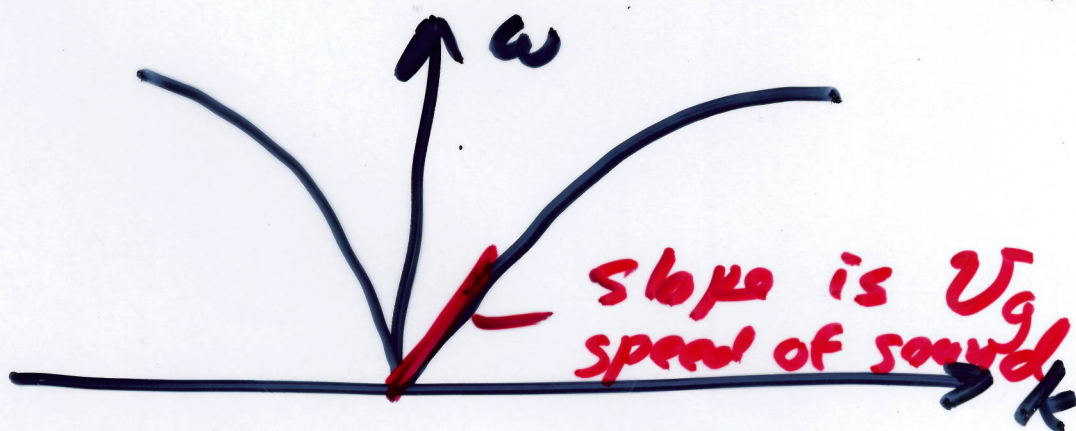
$$\omega = a \sqrt{\frac{C}{m}} k$$

$\omega \propto k$
no dispersion

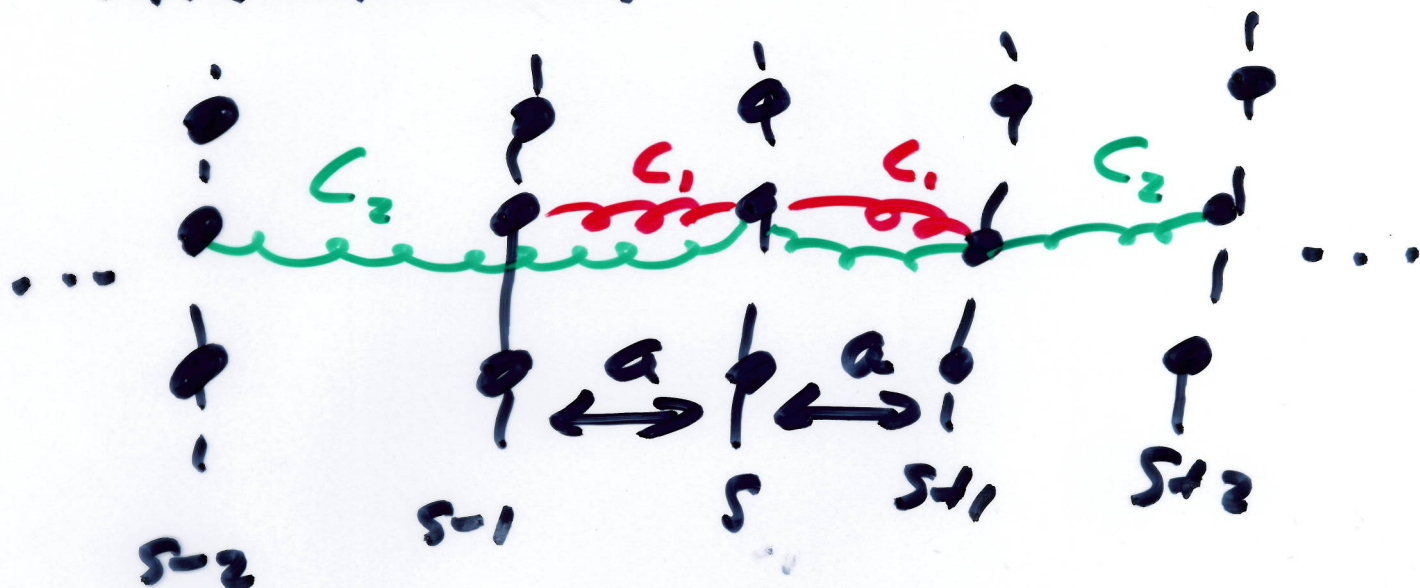
Speed of sound (long λ 's)

no dispersion: $v_g = v_p$

$$v_g = \frac{\omega}{k} = a \sqrt{\frac{C}{m}} = v_p = \frac{d\omega}{dk}$$



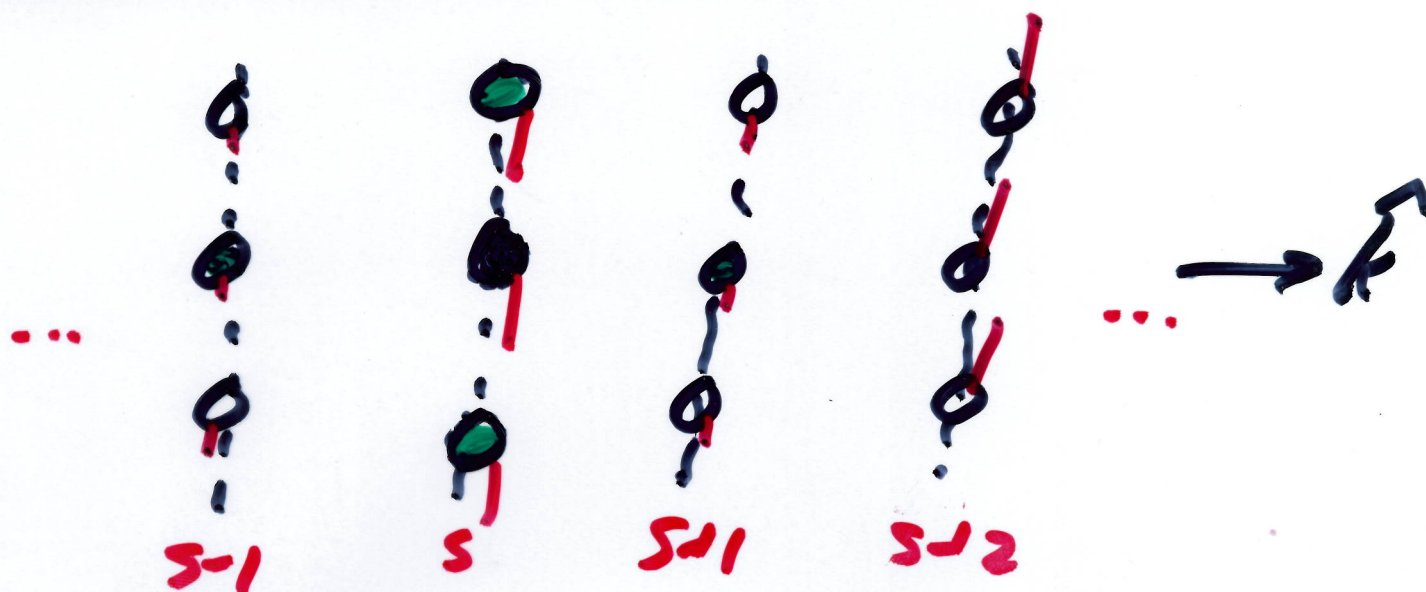
Longitudinal waves
more than nearest neighbor
interactions



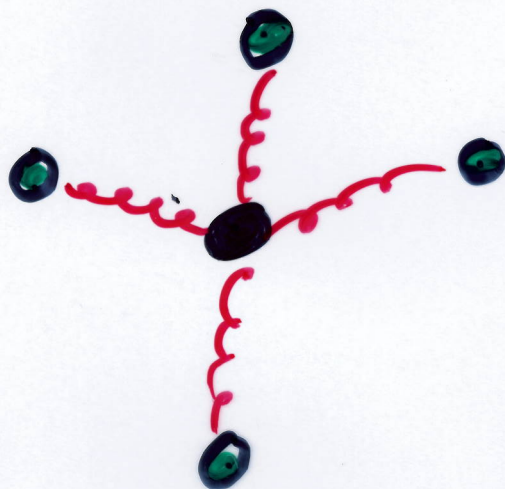
$$\omega^2 = \frac{2C}{M} [1 - \cos(ka)]$$

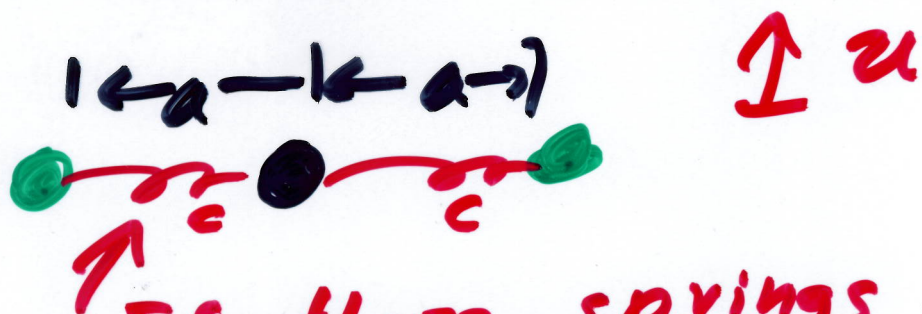
$$\omega^2 = \frac{2}{M} \sum_{p=1}^{\infty} C_p [1 - \cos(kpa)]$$

Transverse waves



No contribution from nearest
neighbor interactions

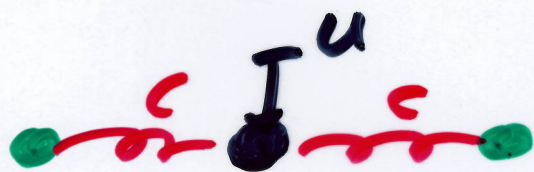




If these springs are under tension (rest length a) there is a restoring force proportional to displacement.

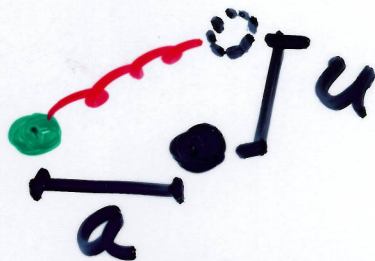
$$F \propto C u_{\text{eff}}$$

Our problem



Rest length of springs = a

$$F \propto u^3$$



new length is $\sqrt{a^2 + u^2}$

amount spring stretched is

$$\sqrt{a^2 + u^2} - a$$

$$u \ll a$$

$$= a \sqrt{1 + \frac{u^2}{a^2}} - a$$

$$= a \left(1 + \frac{u^2}{a^2}\right)^{1/2} - a$$

$$= a \left[1 + \frac{u^2}{2a^2} + \dots\right] - a$$

$$= \frac{u^2}{2}$$



$$\text{Force} = F = C \frac{u^2}{a}$$

$$F_y = F \sin \theta \approx F \tan \theta = F \frac{u}{a}$$

$$2F_y = F_{\text{net}} = C \frac{u^3}{a^2} \propto u^3$$

