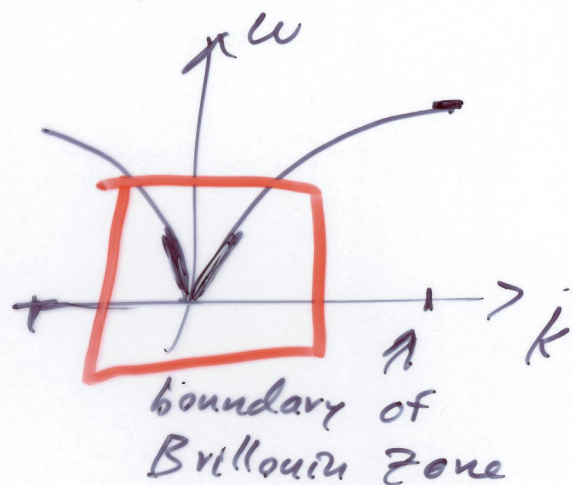


Debye Model

Try to get the T^3 behavior of C_V valid at low temperature

Dispersion Relation (monatomic basis)



Speed of phonons is constant = v

Debye dispersion relation

$$\omega = vk$$

$$\frac{d\omega}{dk} = v$$

Density of states

$$D(\omega) = \frac{V k^2}{2\pi^2} \frac{1}{v_{group}} = \frac{V \omega^2}{2\pi^2 v^3}$$

Constraint

$$\int_0^{\omega_D} D(\omega) d\omega = \int \frac{dN}{d\omega} d\omega = \int dN = N$$

ω_D - Debye frequency - cutoff frequency

$$\int_0^{\omega_D} \frac{V \omega^2}{2\pi^2 v^3} d\omega = \frac{V}{2\pi^2 v^3} \cdot \frac{1}{3} \omega^3 \Big|_0^{\omega_D} = \frac{V \omega_D^3}{6\pi^2 v^3} = N$$

$$\omega_D = \left[\frac{6\pi^2 N}{V} \right]^{1/3} v$$

$$\text{Debye wave number } k_D = \frac{\omega_D}{v} = \left[\frac{6\pi^2 N}{V} \right]^{1/3}$$

Debye Temperature: $\Theta_D \equiv T_D$

$$k_B T_D = \text{energy} = \hbar \omega_D \Rightarrow T_D = \frac{\hbar \omega_D}{k_B}$$

$$T_D = \frac{\hbar v}{k_B} \left[\frac{6\pi^2 N}{V} \right]^{1/3}$$

$$U = \int \underline{D(\omega)} \underline{\langle n \rangle} \underline{\hbar \omega} d\omega$$

$$U = \int_0^{\omega_D} \left(\frac{V \omega^2}{2\pi^2 v^3} \right) \frac{1}{e^{\frac{\hbar \omega}{k_B T}} - 1} \hbar \omega d\omega$$

only temperature dependence

Assume v is the same for all 3 polarizations
 1 longitudinal 2 transverse

$$U = \frac{3Vh}{2\pi^2v^3} \int_0^{\omega_D} \frac{\omega^3 d\omega}{e^{\frac{h\omega}{k_B T}} - 1}$$

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V = \frac{3Vh}{2\pi^2v^3} \int_0^{\omega_D} \frac{\omega^3 d\omega (-1) e^{\frac{h\omega}{k_B T}} \cdot \frac{h\omega}{k_B T^2} (-1)}{\left[e^{\frac{h\omega}{k_B T}} - 1 \right]^2}$$

$$= \frac{3Vh^2}{2\pi^2v^3 k_B T^2} \int_0^{\omega_D} \frac{\omega^4 e^{\frac{h\omega}{k_B T}} d\omega}{\left[e^{\frac{h\omega}{k_B T}} - 1 \right]^2}$$

Low T behavior. should be $\propto T^3$ - check

$$U = ? \quad \text{Define } x = \frac{h\omega}{k_B T}, \quad x_D = \frac{h\omega_D}{k_B T} = \frac{T_D}{T} = \frac{\Theta_D}{T}$$

$$U = 9Nk_B T \left(\frac{T}{T_D} \right)^3 \int_0^{x_D} \frac{x^3}{e^x - 1} dx$$

for low $T \ll T_D$, replace x_D by ∞

$$U = 9Nk_B T \left(\frac{T}{T_D} \right)^3 \int_0^{\infty} \frac{x^3}{e^x - 1} dx = 6 \zeta(4) = \frac{\pi^4}{15}$$

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V = \frac{\partial}{\partial T} \left[\frac{3}{5} \frac{N k_B T^4 \pi^4}{T_D^3} \right]$$

$$= \frac{12 \pi^4}{5} N k_B \left(\frac{T}{T_D} \right)^3 \leftarrow \text{Debye } T^3 \text{ Law}$$

Valid for low T , also correct high T limit - Dulong-Petit Law $C_V = 3Nk_B$

Einstein Model

Independent harmonic Oscillator

one frequency ω_E

$$D(\omega) = A \delta(\omega - \omega_E)$$



$$\text{Total \# of modes} = \int D(\omega) d\omega = N$$

$$D(\omega) = N \delta(\omega - \omega_E)$$

$$U = \int_0^{\infty} D(\omega) \langle n \rangle \hbar \omega d\omega = \int_0^{\infty} N \delta(\omega - \omega_E) \frac{\hbar \omega}{e^{\frac{\hbar \omega}{k_B T}} - 1} d\omega$$

$$= \frac{N \hbar \omega_E}{e^{\frac{\hbar \omega_E}{k_B T}} - 1}$$

$$U = \frac{N \hbar \omega_E}{e^{\frac{\hbar \omega_E}{k_B T}} - 1}$$

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V = \frac{N \hbar \omega_E (-1) e^{\frac{\hbar \omega_E}{k_B T}} \frac{\hbar \omega_E}{k_B T^2} (-1)}{\left[e^{\frac{\hbar \omega_E}{k_B T}} - 1 \right]^2}$$

$$= N k_B \left(\frac{\hbar \omega_E}{k_B T} \right)^2 \frac{e^{\frac{\hbar \omega_E}{k_B T}}}{\left[e^{\frac{\hbar \omega_E}{k_B T}} - 1 \right]^2}$$

multiply by 3 for 3 polarization

high T limit

$$\lim_{x \rightarrow 0} \frac{x^2 e^x}{\left[e^x - 1 \right]^2} = 1$$

$$C_V \rightarrow 3 N k_B \quad (\text{Dulong + Petit})$$

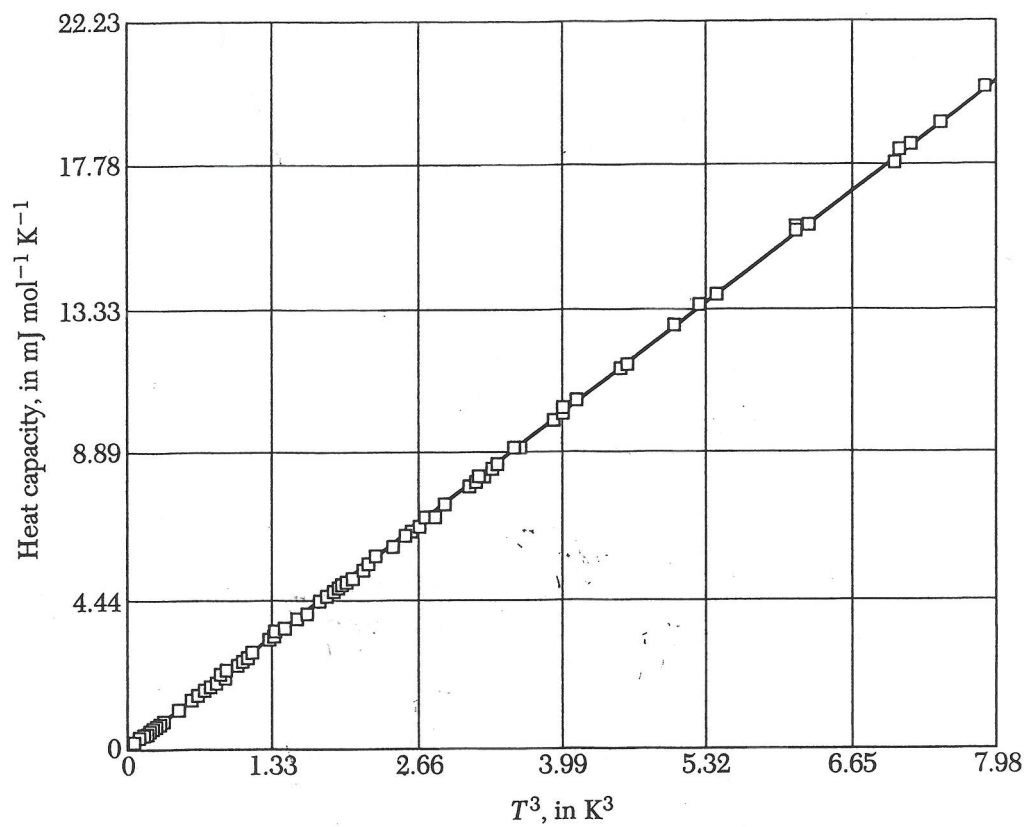


Figure 9 Low temperature heat capacity of solid argon, plotted against T^3 . In this temperature region the experimental results are in excellent agreement with the Debye T^3 law with $\theta = 92.0$ K. (Courtesy of L. Finegold and N. E. Phillips.)