

Any operator \hat{Q} is Hermitian if: $\int \Psi^*(\vec{x}, t) \hat{Q} \Psi(\vec{x}, t) d^3x$

$$= \int (\hat{Q} \Psi(\vec{x}, t))^* \Psi(\vec{x}, t) d^3x$$

$$= \langle \hat{Q} \rangle$$

$$= \langle \hat{Q} \rangle^* \quad \text{"Real"}$$

Properties of Hermitian Operators

In general for any function Φ , $\hat{Q}\Phi = \Phi'$
 e.g. $\hat{p}\Phi = \Phi' = \nabla\Phi$

but if it happens that $\hat{Q}\Phi_n = \lambda_n\Phi_n$ "Eigensvalue Equ"
 $\hookrightarrow \lambda$ "eigenvalue"
 then Φ is said to be an eigenfunction of \hat{Q}
 with eigenvalue λ .

- Claim:
- 1) Eigenvalues of Hermitian Operators are Real.
 - 2) Eigenfunctions of Hermitian Operators corresponding to distinct eigenvalues are orthogonal.

In Linear Algebra:

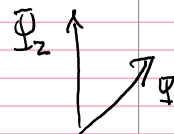
$$M \vec{v} = \lambda \vec{v} = (M - \lambda) \vec{v} = 0 \quad (\text{Aside})$$

for 2x2: $\begin{pmatrix} m_{11} - \lambda & m_{12} \\ m_{21} & m_{22} - \lambda \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = 0$

for non-trivial soln we require that

$$\det \begin{pmatrix} m_{11} - \lambda & m_{12} \\ m_{21} & m_{22} - \lambda \end{pmatrix} = 0$$

$$(m_{11} - \lambda)(m_{22} - \lambda) - m_{12}m_{21} = 0$$



Inner Product of Two wave functions:

$$\langle \Psi_1(\vec{x}, t) | \Psi_2(\vec{x}, t) \rangle = \int \Psi_1^*(\vec{x}, t) \Psi_2(\vec{x}, t) d^3x \neq \langle \Psi_2(\vec{x}, t) | \Psi_1(\vec{x}, t) \rangle$$

"Inner Product" ↖ "Complex"

$$\vec{v} \cdot \vec{w} = v_1 w_1 + v_2 w_2 + v_3 w_3 = \sum_i v_i w_i \rightarrow \int v(x) w(x) dx$$

$$\vec{V} \cdot \vec{W} = V_1 W_1 + V_2 W_2 + V_3 W_3 = \sum_{i=1}^3 V_i W_i \rightarrow \int V(x) W(x) dx$$

Proof of statement (1)

Lets compute the Expectation Value of \hat{Q} using its own eigenfunctions:

$$\langle \hat{Q} \rangle = \int d^3x \left(\Phi_a^*(\vec{x}, t) \hat{Q} \Phi_a(\vec{x}, t) \right)$$

$$= q_a \int d^3x \Phi_a^*(\vec{x}, t) \Phi_a(\vec{x}, t)$$

$$= \langle \hat{Q} \rangle^* \quad \text{"because } \hat{Q} \text{ is Hermitian"}$$

$$\Rightarrow q_a = q_a^* \quad \text{Q.E.D.}$$

Proof of statement (2):

$$(i) \hat{Q} \Phi_a = q_a \Phi_a \quad ; \quad (ii) \hat{Q} \Phi_b = q_b \Phi_b \quad (q_a \neq q_b)$$

$$(i) \int \Phi_b^* \hat{Q} \Phi_a d^3x = q_a \int \Phi_b^* \Phi_a d^3x$$

$$(ii) \int \Phi_a^* \hat{Q} \Phi_b d^3x = q_b \int \Phi_a^* \Phi_b d^3x$$

$$\int (\hat{Q} \Phi_a)^* \Phi_b d^3x = q_b \int \Phi_a^* \Phi_b d^3x \quad \text{by Hermiticity of } \hat{Q}$$

$$(ii) \int \Phi_b^* \hat{Q} \Phi_a d^3x = q_b^* \int \Phi_b^* \Phi_a d^3x \quad \text{"take complex conjugate"}$$

$$0 = (q_a - q_b^*) \int \Phi_b^* \Phi_a d^3x$$

$$\Rightarrow (q_a - q_b) \int \Phi_b^* \Phi_a d^3x = 0$$

$$\int \Phi_b^*(\vec{x}, t) \Phi_a(\vec{x}, t) d^3x = 0 \quad \text{for } q_a \neq q_b$$

Q.E.D.

General Properties of Solutions of the TDSE:

$$\hat{H} \Psi(\vec{x}, t) = i \hbar \frac{\partial}{\partial t} \Psi(\vec{x}, t)$$

$$\hat{H} = -\frac{\hbar^2}{2m} \nabla^2 + V(\vec{x})$$

↳ time independent

$$\hat{H} \neq \hat{H}(t)$$

Formal Particular Solution

$$\Psi(\vec{x}, t) = e^{-\frac{i}{\hbar} \hat{H} t} \psi(\vec{x})$$

Use:

$$e^x = \sum_{n=1}^{\infty} \frac{x^n}{n!}$$

$$\frac{\partial \Psi}{\partial t} = -\frac{i}{\hbar} \hat{H} e^{-\frac{i}{\hbar} \hat{H} t} \psi(\vec{x})$$

$$i \hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi$$

Since \hat{H} is a Hermitian Operator $\hat{H} \psi_E(\vec{x}) = E \psi_E(\vec{x})$

Since the TDSE & TISE are linear in Ψ or ψ respectively it means they obey a superposition principle:

i.e. $\Psi_1(\vec{x}, t)$ & $\Psi_2(\vec{x}, t)$ are solutions

then $\Psi(\vec{x}, t) = (a \Psi_1(\vec{x}, t) + b \Psi_2(\vec{x}, t))$ is also a solution.

We can $\psi(\vec{x}) = \sum_E a_E \psi_E(\vec{x})$

$$\Psi(\vec{x}, t) = e^{-\frac{i}{\hbar} \hat{H} t} \sum_E a_E \psi_E(\vec{x})$$

$$\Psi(\vec{x}, t) = \sum_E a_E e^{-\frac{i}{\hbar} E t} \psi_E(\vec{x}) \quad \hat{H} \psi_E(\vec{x}) = E \psi_E(\vec{x})$$

for a test
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