

$$\hat{H} \psi_n(x) = E_n \psi_n(x)$$

for any 2 states  $\psi_1(x, 0) = \sum_n c_n \psi_n(x)$   
 $\psi_2(x, 0) = \sum_n d_n \psi_n(x)$  } Energy Eigenbasis

$$\langle \psi_1(x) | \psi_2(x) \rangle = \sum_n c_n^* d_n$$

$$= \int \psi_1^*(x) \psi_2(x) dx$$

$\hat{Q} \phi_n = q_n \phi_n$  : the  $q$ 's will also form a complete basis.

$$\Rightarrow \psi_1(x) = \sum_n f_n^{(1)} \phi_n \quad \psi_2(x) = \sum_n f_n^{(2)} \phi_n$$
 } Eigenbasis of  $\hat{Q}$

$$\langle \psi_1(x) | \psi_2(x) \rangle = \sum_n f_n^{(1)*} f_n^{(2)}$$

$$= \sum_n c_n^* d_n$$

The Inner product is independent of basis!

Furthermore,  $\psi_1(x)$  &  $\psi_2(x)$  are characterized by either  $c_n$  &  $d_n$ 's or  $f_n^{(1)}$ ,  $f_n^{(2)}$ 's expansion coefficients.

Yet another example,

$$\psi_1(x, 0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi_1(k) e^{ikx} dk$$

$$\psi_2(x, 0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi_2(k) e^{ikx} dk$$
 } Eigenbasis of the Momentum operator

$$\langle \psi_1(x) | \psi_2 \rangle = \int_{-\infty}^{\infty} \psi_1^*(x) \psi_2(x) dx = \frac{1}{2\pi} \iiint \phi_1^*(k) \phi_2(k') e^{-ikx} e^{ik'x} dk dk' dx$$

$$= \frac{1}{2\pi} \int dk dk' \phi_1^*(k) \phi_2(k') \int dx e^{i(k'-k)x}$$

(...)

$$= \int \phi_1^*(k) \phi_2(k) dk$$

$$= \langle \phi_1(k) | \phi_2(k) \rangle$$

In Linear Algebra a vector  $\vec{V}$  has components  $(V_x, V_y, V_z)$

$$\vec{V} = \sum_i \hat{e}_i V_i = \hat{i} V_x + \hat{j} V_y + \hat{k} V_z$$

$$= \sum_i \hat{e}_i V_i$$

∴ Dirac found it convenient to invent a formalism where quantum states are defined indep. of basis.

To each quantum state  $\leftrightarrow$   $|\psi\rangle$  "Ket Vector"  
 $\langle\psi|$  "Bra Vector"

$$\hat{H} |E_n\rangle = E_n |E_n\rangle$$

$$\langle E_n | E_m \rangle = \delta_{mn}$$

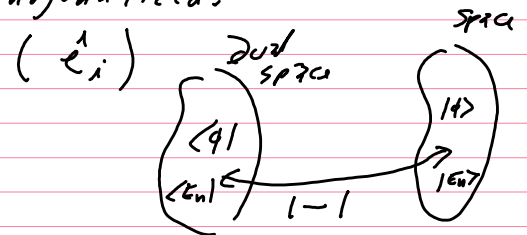
$$|\psi\rangle = \sum c_n |E_n\rangle$$

$$\langle E_m | E_n \rangle = \delta_{mn}$$

$$c_n = \langle E_n | \psi \rangle$$

$|c_n|^2 = \text{prob of obtaining } E_n$

$|E_n\rangle$  "Form a basis"  
 this basis can always be orthogonalized.



$$\langle \hat{H} \rangle = \langle \psi | \hat{H} | \psi \rangle = \sum_{m,n} c_m^* c_n \langle E_m | \hat{H} | E_n \rangle$$

$$= \sum_n |c_n|^2 E_n$$

For any operator  $\hat{Q}$ :

$$\hat{Q} |\psi\rangle = |\psi'\rangle$$

$$\hat{Q} \sum_n c_n |E_n\rangle = \sum_m c_m' |E_m\rangle$$

$$\sum_n c_n \hat{Q} |E_n\rangle = \sum_m c_m' |E_m\rangle$$

$$\sum_n c_n \langle E_1 | \hat{Q} | E_n \rangle = \sum_m c'_m \langle E_1 | E_m \rangle$$

$$\sum_n \underbrace{\langle E_1 | \hat{Q} | E_n \rangle}_{Q_{1n}} c_n = c'_1$$

$$\sum_n Q_{1n} c_n = c'_1$$

$$\begin{pmatrix} Q_{11} & Q_{12} & \dots \\ Q_{21} & Q_{22} & \dots \\ \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \\ \vdots \end{pmatrix} = \begin{pmatrix} c'_1 \\ c'_2 \\ \vdots \end{pmatrix}$$

If  $\hat{Q}$  is Hermitian

$$\Rightarrow \langle E_1 | \hat{Q} | E_n \rangle^* = \langle E_n | \hat{Q}^\dagger | E_1 \rangle = \langle E_n | \hat{Q} | E_1 \rangle$$

$$\hat{Q}_{1n}^* = \hat{Q}_{n1} \Rightarrow \hat{Q}^{*T} = \hat{Q}$$

Any operator can be written as the outer product of two ket vectors: of the form  $|\phi\rangle\langle\psi|$

In linear,  
 $\hat{e}_i \otimes \hat{e}_j$

e.g.  $\hat{Q}|\psi\rangle = |\psi'\rangle \Rightarrow \hat{Q} = |\psi'\rangle\langle\psi|$

$$\hat{P} \equiv |\psi\rangle\langle\psi| \quad \text{"Projection operator"}$$

If this operator acts on any ket vector  $|\phi\rangle$ :

$$\hat{P}|\phi\rangle = |\psi\rangle\langle\psi|\phi\rangle = c|\psi\rangle$$

$$\hat{P}^2 = (|\psi\rangle\langle\psi|)(|\psi\rangle\langle\psi|) = \hat{P}$$

Projection operators for basis vectors:

$$\hat{P}_n = |E_n\rangle\langle E_n| \quad \hat{P}_n^2 = \hat{P}_n$$

$$\begin{aligned} \hat{P}_n |\psi\rangle &= |E_n\rangle\langle E_n|\psi\rangle \\ &= c_n |E_n\rangle \end{aligned}$$

$$\sum_n \hat{P}_n |\psi\rangle = \sum_n c_n |E_n\rangle = |\psi\rangle$$

$$\sum_n \mathbb{P}_n = \mathbb{1} \quad \text{"Completeness"}$$

$$\hookrightarrow \sum_n \psi_n(x) \psi_n(x') = \delta(x-x')$$

$$\mathbb{P}_n \mathbb{P}_m = \mathbb{P}_n \delta_{mn}$$

$|\psi\rangle$  &  $|\varphi\rangle$  the inner product  $\langle\varphi|\psi\rangle = \langle\varphi|\mathbb{1}|\psi\rangle$

$$\begin{aligned} \langle\varphi|\psi\rangle &= \sum_n \langle\varphi|\mathbb{P}_n|\psi\rangle = \sum_n \underbrace{\langle\varphi|E_n\rangle}_{d_n^*} \underbrace{\langle E_n|\psi\rangle}_{c_n} \\ &= \sum_n d_n^* c_n \end{aligned}$$

$$\mathbb{P}_n^{(x)} = |q_n\rangle\langle q_n|$$

Basis Labeled by Continuous Index:

$$\hat{X}|x\rangle = x|x\rangle \quad \langle x|x'\rangle = 0 \quad x \neq x'$$

$$\mathbb{P}_x = |x\rangle\langle x| \quad \mathbb{P}_x^2 = \mathbb{P}_x$$

$$\int \mathbb{P}_x dx = \int dx |x\rangle\langle x| = \mathbb{1}$$

$$|\psi\rangle = \int dx c(x) |x\rangle$$

$$\langle x'|\psi\rangle = \int dx c(x) \langle x'|x\rangle \Rightarrow \langle x'|x\rangle = \delta(x-x')$$

$$c(x') = \langle x'|\psi\rangle$$

$$\text{c.f. } c_n = \langle E_n|\psi\rangle$$

$$\int dx \langle y'|x\rangle \langle x|y\rangle = \langle y'|y\rangle$$

$$\Rightarrow \int dx \delta(x-y') \delta(y-x) = \delta(y'-y) \quad \checkmark$$

$$|\varphi\rangle = \int dy f(y) |y\rangle$$

$$\langle \varphi| = \int dy \langle y| f^*(y)$$

$$\langle \varphi|\psi\rangle = \int dx dy f^*(y) c(x) \underbrace{\langle y|x\rangle}_{\delta(x-y)}$$

$$\left. \begin{array}{l} \text{"Coordinate} \\ \text{basis"} \end{array} \right\} \begin{array}{l} = \int dx f^*(x) c(x) \\ = \int dx \varphi^*(x) \psi(x) \end{array} \right\} \Rightarrow \begin{array}{l} f(x) = \varphi(x) \\ c(x) = \psi(x) \end{array}$$

### Momentum Basis

$$\hat{P} |p\rangle = p |p\rangle$$

$$P_p = |p\rangle \langle p|$$

$$\int P_p dp = 1$$

$$\langle x|\psi\rangle = \int dp \varphi(p) \langle x|p\rangle$$

$$\psi(x) = \int dp \varphi(p) \langle x|p\rangle \longleftrightarrow \psi(x) = \frac{1}{\sqrt{2\pi}} \int dk \varphi(k) e^{ikx}$$

$$\langle x|p\rangle \sim e^{ikx} \quad p = \hbar k$$