

# Why is $\hat{H}$ the Hamiltonian?

①  $\left[ \begin{array}{l} \frac{\partial}{\partial x} \rightarrow \text{momentum} \\ \frac{\partial}{\partial t} \rightarrow \text{energy} \end{array} \right. \quad \hat{p} = \frac{\hbar}{i} \frac{\partial}{\partial x} \quad \hat{H} = i\hbar \frac{\partial}{\partial t}$

$g = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & +1 & 0 \\ 0 & 0 & 0 & +1 \end{pmatrix}$

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②  $\hat{H}$  plays the same role as classical Hamiltonian from mechanics.

$$\frac{d}{dt} \langle B \rangle \sim [H, B] \xrightarrow{\text{classical}} \{H, B\}$$

Poisson brackets

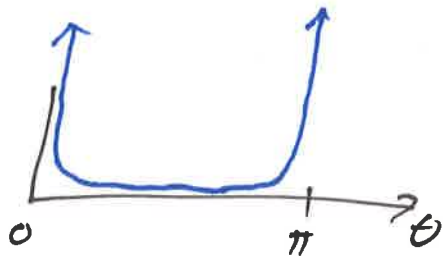
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$\psi(x) = \langle x | \psi \rangle$  projection of ket  $|\psi\rangle$  on  $|x\rangle$  basis

must be single-valued (function)

$\psi(x)$  need not be finite

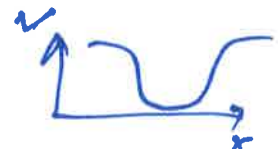
$$dP = |\psi|^2 dV \quad \uparrow \text{goes to zero}$$





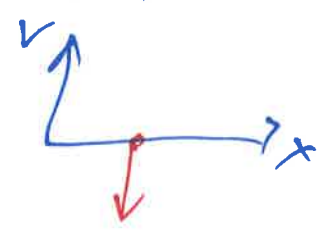
# Continuity of $\psi(x)$

$$-\frac{\hbar^2}{2m} \nabla^2 \psi(x) + V(x) \psi(x) = E \psi(x)$$

$\psi(x)$  is continuous for most common potentials

-  $V(x)$  is continuous   $\psi', \psi$  continuous

-  $V(x)$  has discontinuities   
but finite (e.g. finite square well)  
 $\psi''$  has discont., [ $\psi', \psi$  continuous]

-  $V(x)$  has discont. and is infinite   
• (e.g. infinite square well)  
 $\psi'$  is discontinuous,  $\psi$  is continuous  
• (e.g. delta function well) 

- If  $V(x)$  contains derivatives of delta functions  
then  $\psi(x)$  will not be continuous.

# Georg Cantor's Diagonalization Proof

There are more real numbers than integers.

Try to list every real number between 0 and 1

• 0.7154925...

• 1.000000...

• 0.555555...

• 0.2147999279...

⋮

• 0.3142

aleph-null

Cardinality of the set of integers  $\equiv \aleph_0$

“ “ real  $\equiv \mathbb{C}$

$\mathbb{C} = \aleph_1$  continuum hypothesis

Zermelo-Frankel set theory

+ Axiom of choice.  $\boxed{\aleph_1 = \mathbb{C}}$

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Fourier Series — denumerably infinite number of  $\omega$ 's

“ Transform — non-denumerably “ “  $a$ 's

$|\psi\rangle$  ket contains everything that can be known about the quantum system

$\langle x|\psi\rangle \equiv \psi(x)$  old friend the wavefunction  
The ket  $|\psi\rangle$  expressed in the  $|x\rangle$  basis

$|x\rangle$  is an eigenstate of the  $\hat{X}$  operator

$$\hat{X}|\lambda_0\rangle = \lambda_0|\lambda_0\rangle, \quad \langle x|\lambda_0\rangle = \delta(x-x_0)$$

↑ op    ↑ ket                      ← eigenvalue, real number                      ↑ ket

$|p\rangle$  is an eigenstate of the  $\hat{P}$  momentum op.

$$\hat{P}|p_0\rangle = p_0|p_0\rangle, \quad \langle p|q\rangle = \delta(p-q)$$

↑ op    ↑ ket                      ↑ ket                      ← eigenvalue, real number

$\langle p|\psi\rangle \equiv \tilde{\psi}(p)$  Fourier transform of  $\psi(x)$

the ket  $|\psi\rangle$  expressed in the  $|p\rangle$  basis.

$$\langle x|p\rangle = ?$$

$|p\rangle$  momentum eigenstate  
 $\Rightarrow$  plane wave

$$= \frac{1}{\sqrt{2\pi\hbar}} e^{\frac{ipx}{\hbar}}$$

$$\langle p|x\rangle = \langle x|p\rangle^* = \frac{1}{\sqrt{2\pi\hbar}} e^{-\frac{ipx}{\hbar}}$$

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$$\tilde{\psi}(p) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} dx \psi(x) e^{-\frac{ipx}{\hbar}}$$

$$\psi(x) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{+\infty} dp \tilde{\psi}(p) e^{+\frac{ipx}{\hbar}}$$