

## The ground state & the Uncertainty Relations for the H.O.

First the Virial Theorem:

$$\frac{d}{dt} \langle A \rangle = \frac{i}{\hbar} \langle [H, A] \rangle$$

Claim:  $2 \langle T \rangle = \langle \vec{r} \cdot \vec{\nabla} V(r) \rangle$

$$= 0 \quad \text{for stationary states}$$

For stationary states of the H.O.

Proof  $\frac{d}{dt} \langle \vec{r} \cdot \vec{p} \rangle = 0$

"For Stationary State"

$$= \frac{i}{\hbar} \langle [H, \vec{r} \cdot \vec{p}] \rangle$$

$$\Rightarrow \langle [H, \vec{r} \cdot \vec{p}] \rangle = 0$$

$$\begin{aligned} [H, \vec{r} \cdot \vec{p}] &= \left[ \frac{p^2}{2m}, \vec{r} \cdot \vec{p} \right] + [V(r), \vec{r} \cdot \vec{p}] \\ &= \frac{1}{2m} [p^2, r_i] p_i + r_i [V(r), p_i] \end{aligned}$$

$$[p^2, r_i] = p_x [p_x, r_i] + [p_x, r_i] p_x = -2i\hbar p_i$$

$$\Rightarrow \langle -i\hbar \frac{p^2}{m} \rangle + \langle i\hbar \vec{r} \cdot \vec{\nabla} V \rangle = 0$$

$$\underline{2 \langle T \rangle = \langle \vec{r} \cdot \vec{\nabla} V \rangle}$$

Q.E.D.  
(For stationary states)

For the 1-d H.O.:  $V(x) = \frac{1}{2} m \omega^2 x^2$

$$\vec{r} \cdot \vec{\nabla} V = x \frac{\partial}{\partial x} V = m \omega^2 x^2 = 2V(x)$$

$$\Rightarrow \langle T \rangle = \langle V \rangle$$

$$\begin{aligned} \hat{H} = T + V &\Rightarrow \langle \hat{H} \rangle = 2 \langle T \rangle \\ &= 2 \langle V \rangle \end{aligned}$$

$$\Rightarrow \langle T \rangle = \frac{1}{2} \hbar \omega (n + 1/2)$$

$$\langle V \rangle = \frac{1}{2} \hbar \omega (n + \frac{1}{2})$$

$$T = \frac{p^2}{2m} \Rightarrow \langle p^2 \rangle = \hbar m \omega (n + \frac{1}{2})$$

$$V(x) = \frac{1}{2} m \omega^2 x^2 \quad \langle x^2 \rangle = \frac{2}{m \omega^2} \langle V \rangle = \frac{2}{m \omega^2} \frac{1}{2} \hbar m \omega (n + \frac{1}{2})$$

$$\langle x^2 \rangle = \frac{\hbar}{m \omega} (n + \frac{1}{2})$$

$$\text{Also, } \left. \begin{array}{l} \langle p \rangle = 0 \\ \langle x \rangle = 0 \end{array} \right\} \text{ since } \psi_n(-x) = (-1)^n \psi_n(x) \\ |\psi_n(-x)|^2 = |\psi_n(x)|^2$$

$$\begin{aligned} (\Delta p \Delta x)^2 &= \langle x^2 \rangle \langle p^2 \rangle \\ &= \frac{\hbar}{m \omega} (n + \frac{1}{2}) \hbar m \omega (n + \frac{1}{2}) \end{aligned}$$

$$\Rightarrow \Delta p \Delta x = \hbar (n + \frac{1}{2})$$

"for stationary states of the H.O."

$$\therefore (\Delta x \Delta p)_{\min} = \frac{\hbar}{2} \quad \text{"Minimum Uncertainty Product"}$$

But  $n=0 \Rightarrow$  The ground state:

$$\psi_0(x) = A_0 e^{-\frac{x^2}{2}} \quad \text{"Gaussian in } x \text{"}$$

$$E_0 = \frac{\hbar \omega}{2} \quad \text{"Zero Point Energy"}$$

Correspondance w/ the Classical H.O.

$$V(x) = \frac{1}{2} K x^2, \quad E = \frac{p^2}{2m} + V(x)$$

$$\Rightarrow x(t) = x(0) \cos(\omega t)$$

$$E = \frac{1}{2} K x^2(0) = \frac{1}{2} m \omega^2 x_0^2$$

= constant in time

$$\begin{array}{c} \text{piecewise} \\ \begin{array}{ccc} \vdots & & \vdots \\ -x(0) & x_0 & x(0) \\ \vdots & & \vdots \end{array} \\ F = -Kx; \quad \omega = \sqrt{K/m} \\ = m \frac{dx}{dt} \end{array}$$

$$x_0 = \sqrt{\frac{2E}{m \omega^2}}$$

Probability distribution in  $x$  : clearly  $p(-x) = p(x)$

Probability distribution in  $x$  : clearly  $\rho(-x) = \rho(x)$

$$\rho(x) dx = \rho(t) dt$$

$$\rho(t) = \text{constant} = \frac{1}{T}$$

$$\rho(x) dx = \frac{1}{T} dt$$

$$\rho(x) = \frac{1}{T} \frac{1}{\frac{dx}{dt}} = \frac{1}{T \dot{x}(t)}$$

$$\dot{x}(t) = -x(t) \omega \cos \omega t$$

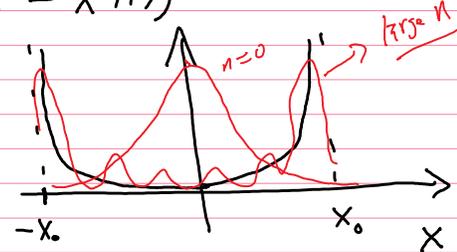
$$\dot{x}^2 = x_0^2 \omega^2 \cos^2 \omega t$$

$$= x_0^2 \omega^2 (1 - \sin^2 \omega t)$$

$$= (x_0^2 - x^2(t)) \omega^2$$

$$\dot{x}(t) = \sqrt{x_0^2 - x^2(t)}$$

$$\rho(x) = \frac{1}{\omega T x_0} \frac{1}{[1 - (x^2/x_0^2)]^{1/2}}$$



$$T = \frac{2\pi}{\omega}$$

$$\rho(x) = \frac{1}{2\pi x_0} \frac{1}{[1 - (x^2/x_0^2)]^{1/2}}$$

$\pm n$  Q.M.  $\rho(x) = |\psi_n(x)|^2 = \frac{A_0^2}{2^n n!} (H_n(\xi))^2 e^{-\xi^2}; \xi = \sqrt{m\omega} x$

$\hookrightarrow$  "n+1 order polynomial"

The classical limit is obtained in the limit  $n \rightarrow \infty$

$$\hbar \omega \ll E_n = \hbar \omega (n + 1/2)$$

Well, not quite. One must also work w/ a superposition of stationary states. One single stationary state will not yield classical limit! The particular superposition which does is called a coherent state.