

"Conventional" Neutrino Physics in the Near and Not-so-Near Future

Is there anything left to learn from ν Deeply Inelastic Scattering?

Jorge G. Morfin
Fermilab

1



HERA is not the only place to study Deeply Inelastic Scattering

CTEQ QCD
Summer School
J. G. Morfin
June 2000

- $\sigma_{\nu N}$ is many orders of magnitude smaller than σ_{eN} .
Events are hard to accumulate.
- However, $\nu/\bar{\nu}$ have ability to resolve flavor of nucleon's constituents: ν interacts with $d, s, \bar{u},$ and \bar{c} while $\bar{\nu}$ interacts with u, c, \bar{d} and \bar{s} .
- $\nu/\bar{\nu}$ have definite spin/helicity states, can they be used to determine partonic spin contribution to the nucleon?
- The ν is a very effective partner in determining the partonic and spin structure of the nucleon via measurement of the structure functions.

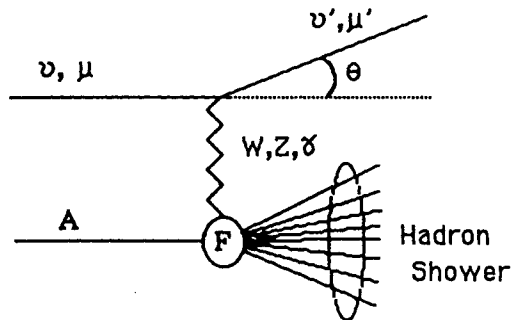
■ Standard Kinematic Variables:

$$Q^2 = 4EE' \sin^2 \theta / 2 = -q^2$$

$$\nu = E - E'$$

$$x_{Bj} = Q^2 / 2M\nu$$

$$y = \nu / E$$



The Vocabulary of Deep Inelastic Lepton-Hadron Scattering

CTEQ QCD
Summer School
J. G. Morfin
June 2000

$$\frac{d^2\sigma^{VN}}{dx dy} = \frac{G_F^2 M E}{\pi(1 + Q^2/M_W^2)} *$$

$$\left[F_2^{VN}(x, Q^2) \left(\frac{y^2 + (2Mxy/Q)^2}{2 + 2R_L^{VN}(x, Q^2)} + 1 - y - \frac{Mxy}{2E_\nu} \right) + xF_3^{VN}(x, Q^2)y \left(1 - \frac{y}{2} \right) \right]$$

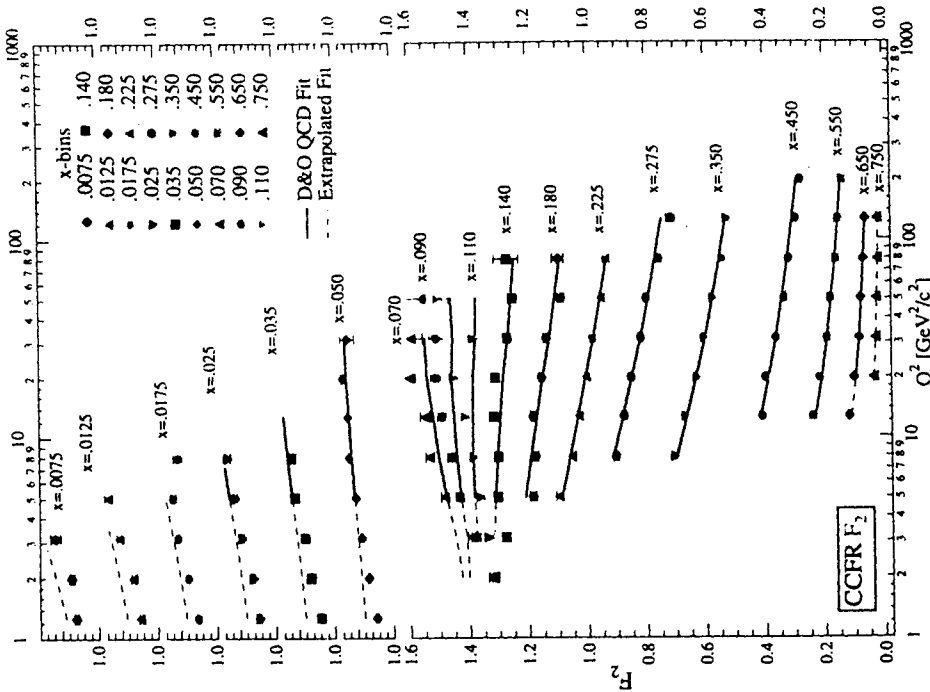
$$R_L(x, Q^2) = \frac{F_2(x, Q^2)}{2xF_1(x, Q^2)} (1 - (2Mx/Q)^2) - 1$$

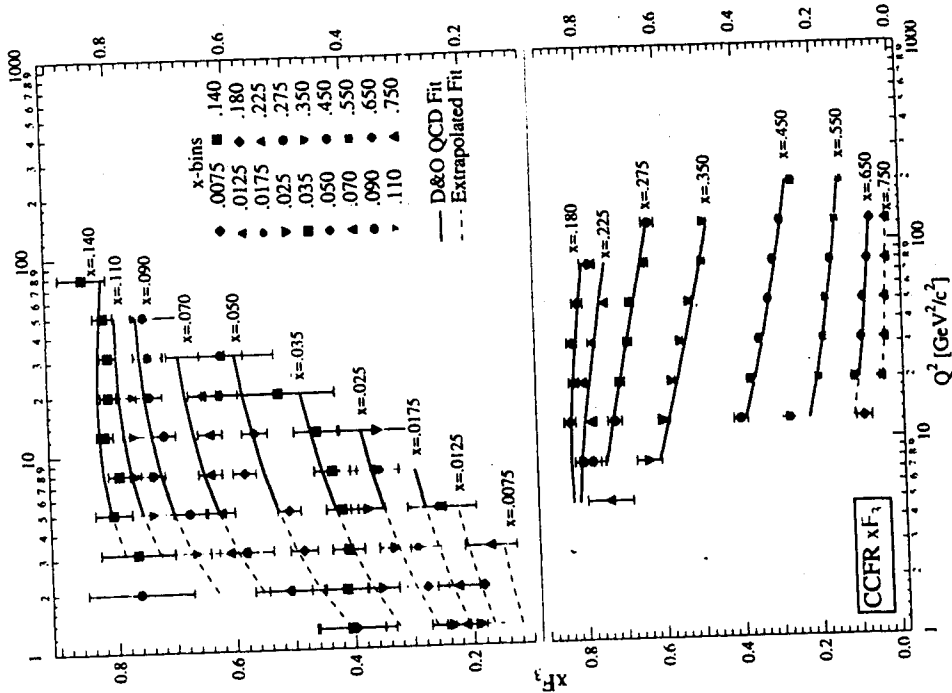
- Start with unpolarized functions.
- Latest measurements of F_2 and xF_3 from CCFR.
Results from E-815/NuTeV "soon".
- Based on 950,000 ν events and 170,000 $\bar{\nu}$ events.

■ Errors:

F_2 : stat. (1-3)%
 xF_3 : stat (2-10)%

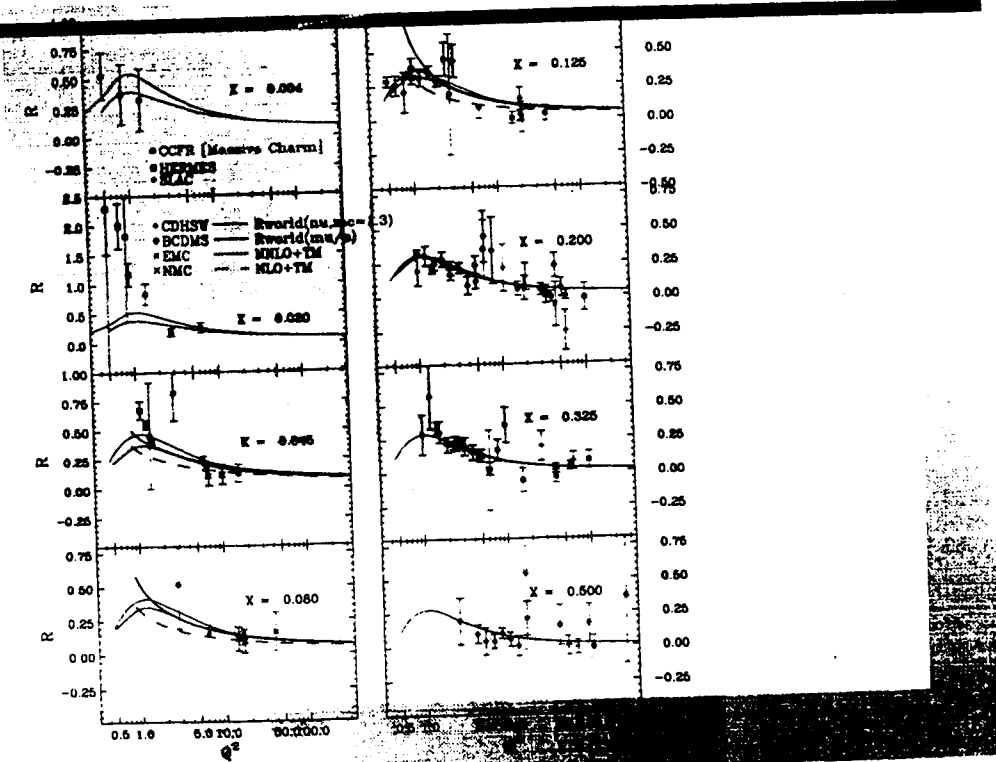
Systematic errors about the same size





Current Status of Neutrino Measurements of $R_L(x, Q^2)$ - on Fe

CTEQ QCD
 Summer School
 J. G. Morfin
 June 2000



- Heavy nuclear targets needed to accumulate sufficient statistics. Measuring ν -Fe not ν -N

- Statistics and systematics not yet at level to allow extraction of all possible information.

Independent ν and $\bar{\nu}$ structure functions
not yet well measured!



Effect of These Compromises: Measured Structure Functions

CTEQ QCD
Summer School
J. G. Morfin
June 2000

$$\frac{d^2\sigma^{\nu N}}{dx dy} = \frac{G_F^2 ME}{\pi(1+Q^2/M_W^2)} *$$

$$\left[F_2^{\nu N}(x, Q^2) \left(\frac{y^2 + (2Mxy/Q)^2}{2 + 2R_L^{\nu N}(x, Q^2)} + 1 - y - \frac{Mxy}{2E_\nu} \right) + xF_3^{\nu N}(x, Q^2)y \left(1 - \frac{y}{2} \right) \right]$$

$$\frac{d^2\sigma^{\bar{\nu} N}}{dx dy} = \frac{G_F^2 ME}{\pi(1+Q^2/M_W^2)} *$$

$$\left[F_2^{\bar{\nu} N}(x, Q^2) \left(\frac{y^2 + (2Mxy/Q)^2}{2 + 2R_L^{\bar{\nu} N}(x, Q^2)} + 1 - y - \frac{Mxy}{2E_\nu} \right) - xF_3^{\bar{\nu} N}(x, Q^2)y \left(1 - \frac{y}{2} \right) \right]$$

$$R_L(x, Q^2) = \frac{F_2(x, Q^2)}{2xF_1(x, Q^2)} (1 - (2Mx/Q)^2) - 1$$

$$F_2^{VN}(\mathbf{x}, Q^2) = \mathbf{x}[u + \bar{u} + d + \bar{d} + 2\bar{s} + 2c]$$

$$F_2^{\bar{V}N}(\mathbf{x}, Q^2) = \mathbf{x}[u + \bar{u} + d + \bar{d} + 2s + 2\bar{c}]$$

$$xF_3^{VN}(\mathbf{x}, Q^2) = \mathbf{x}[u + d - \bar{u} - \bar{d} - 2\bar{s} + 2c]$$

$$xF_3^{\bar{V}N}(\mathbf{x}, Q^2) = \mathbf{x}[u + d - \bar{u} - \bar{d} + 2s - 2\bar{c}]$$

■ What do we gain by measuring separate v and \bar{v} structure functions?

■ Does $s = \bar{s}$ and $c = \bar{c}$ over all x ? $F_2^v - F_2^{\bar{v}} = 2[(s - \bar{s}) + (\bar{c} - c)]$

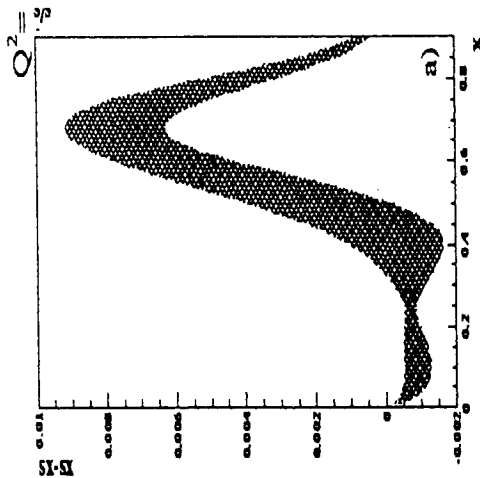
■ If so.....

$$F_2^v - xF_3^v = 2(\bar{u} + \bar{d} + 2\bar{c}) = 2U + 4\bar{c}$$

$$F_2^{\bar{v}} - xF_3^{\bar{v}} = 2(\bar{u} + \bar{d} + 2\bar{s}) = 2U + 4\bar{s}$$

$$xF_3^v - xF_3^{\bar{v}} = 2[(s + \bar{s}) - (\bar{c} + c)] = 4\bar{s} - 4\bar{c}$$

$\overline{s(x)}$ from CDHSW data?



F. Zomer, Eur.Phys.J C12:243,2000

– the CDHSW ν data at low y are larger than both CCFR data and the QCD prediction while their $\bar{\nu}$ data agree with CCFR data

$$\triangleleft \frac{d^2\sigma}{Edxdy}(\nu, y=0) \sim [q + \bar{q}] = (u + \bar{u}) + (d + \bar{d}) + s,$$

$$\frac{d^2\sigma}{Edxdy}(\bar{\nu}, y=0) \sim [q + \bar{q}] = (u + \bar{u}) + (d + \bar{d}) + \bar{s}$$

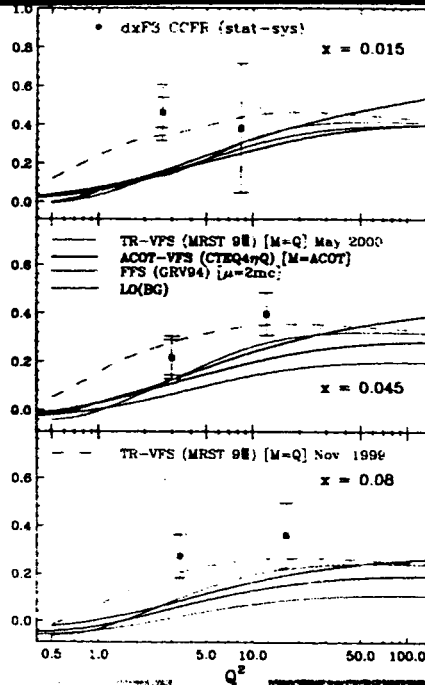
– Thus, the CDHSW data could introduce a large asymmetry between s and \bar{s} ($s > \bar{s}$) because u and d are well constrained

- ΔxF_3 sensitive to s and c (model dependent). Extract directly from data-model independent.

- Extraction of F_2 also dependent on ΔxF_3 .

$$d^2\sigma(v+\bar{v})C \approx (1-y+y^2)F_2 + (y-y^2)\Delta xF_3$$

- Newly extracted F_2^v at low x (< 0.1) now agrees with F_2^μ , a long-standing problem. has been eliminated!
- NuTeV results with increased $\bar{\nu}$ data "soon".



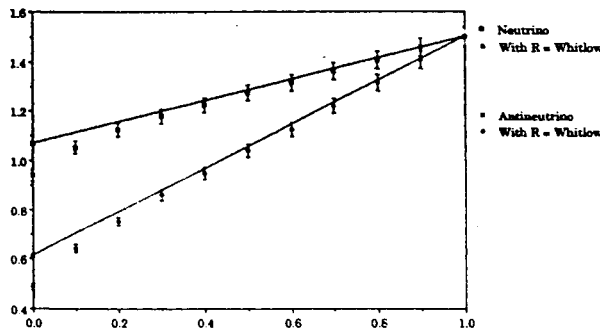
How (In Principle) to Extract all SIX Measured Structure Functions

CTEQ QCD Summer School
J. G. Morfin
June 2000

$$\frac{d\sigma^{vA}}{dx dQ^2} = \frac{G_F^2}{2\pi x} \left[\frac{1}{2} (F_2^{vA}(x, Q^2) + xF_3^{vA}(x, Q^2)) + \frac{(1-y)^2}{2} (F_2^{vA}(x, Q^2) - xF_3^{vA}(x, Q^2)) \right]$$

$$\frac{d\sigma^{\bar{v}A}}{dx dQ^2} = \frac{G_F^2}{2\pi x} \left[\frac{1}{2} (F_2^{\bar{v}A}(x, Q^2) - xF_3^{\bar{v}A}(x, Q^2)) + \frac{(1-y)^2}{2} (F_2^{\bar{v}A}(x, Q^2) + xF_3^{\bar{v}A}(x, Q^2)) \right]$$

$$\frac{\sigma(x, Q^2, (1-y)^2)}{G^2/2\pi x}$$



$$-2y^2 F_L(x, Q^2)$$

$$x = .10$$

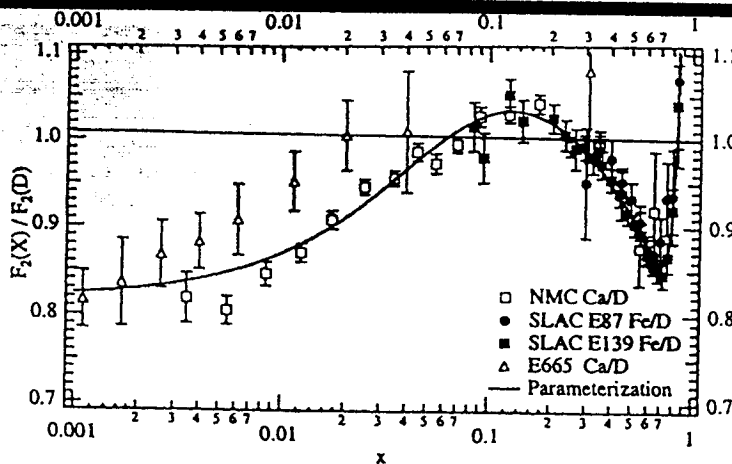
$$Q^2 = 3 \text{ GeV}^2$$

- Statistics: need sufficient events in each $x - Q^2$ bin to populate $(1-y)^2$ distribution for fit to R , slope and intercept. Goal: Order 1% or better.
- Beam systematics: minimize "wobble" in slope and intercept. Goal: Order 2% or better
- Measured quantities: need good calibration for confidence in $(1-y)^2$ and to understand migration/smearing between $x-Q^2$ bins. Goal: Better than CCFR / NuTeV.
- Particularly vicious complications involving treatment of the charm quark.



Effect of These Compromises: Heavy Nuclear Targets

CTEQ QCD
Summer School
J. G. Morfin
June 2000



- nuclear effects measured (with high statistics) in μ -A not in ν -A.
- from low-to-high x_{Bj} go through: shadowing, anti-shadowing, "EMC" effect, Fermi motion.

- In general, v/\bar{v} can isolate nuclear effects for valence quarks from the sum of all quarks.
- More specifically, recall

$$F_2^v - xF_3^v = 2(\bar{u} + \bar{d} + 2\bar{c}) = 2U + 4\bar{c}$$

$$F_2^{\bar{v}} - xF_3^{\bar{v}} = 2(\bar{u} + \bar{d} + 2\bar{s}) = 2U + 4\bar{s}$$

$$xF_3^v - xF_3^{\bar{v}} = 2[(s + \bar{s}) - (\bar{c} + c)] = 4\bar{s} - 4\bar{c}$$

- Could measure nuclear effects on individual partons...

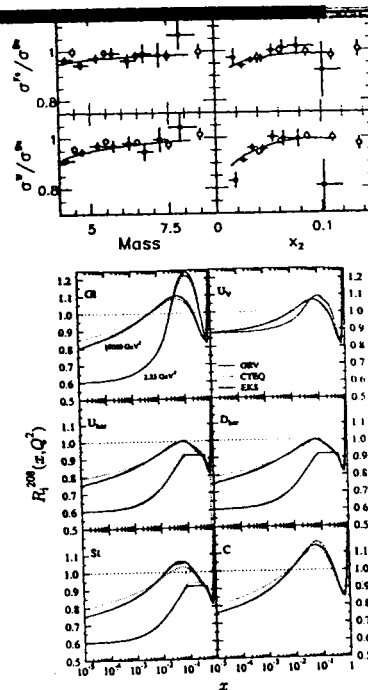


MINOS

Any Indication of a Difference in Nuclear Effects of Valence and Sea Quarks?

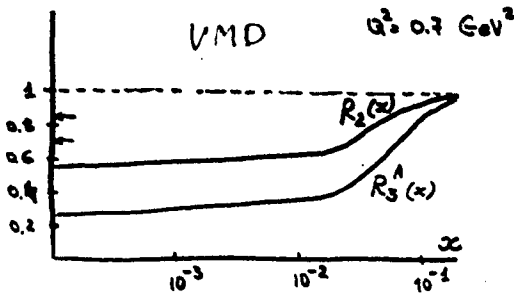
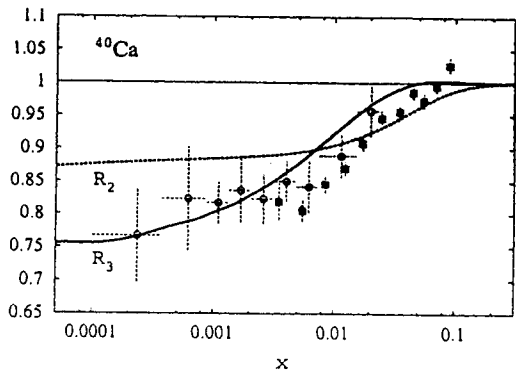
CTEQ QCD
Summer School
J. G. Morfin
June 2000

- Nuclear effects similar in Drell-Yan and DIS for $x < 0.1$. Then no "anti-shadowing" in D-Y^a while "anti-shadowing" seen in DIS (5-8% effect in NMC). Indication of difference in nuclear effects between valence & sea quarks?
- This quantified by K.J. Eskola^b et al within LO ($t = 2$) DGLAP using initial nuclear distributions from CTEQ4L and GRV-LO and assume scale evolution of nuclear parton densities is perturbative.



^a hep-ex/9906010 ^b hep-ph/9807297

- S.A.Kulagin has calculated shadowing for F_2 and xF_3 in ν -A interactions based on a non-perturbative parton model.
- At 5 GeV^2 , ratio of shadowing $xF_3 : F_2 \approx 0.5$ at $x = .02$ to ≈ 2 at $x = .0001$.
- Most recently calculated shadowing in VMD region (lower Q^2 : dominant for $x < .01$) and finds significantly stronger shadowing.



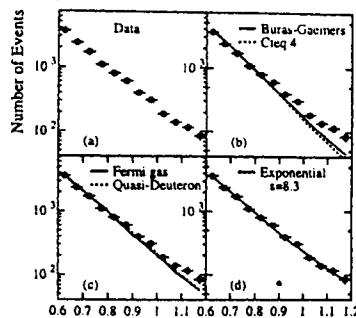
Another Open Nuclear Effects Question: Behavior of F_2 as $x \rightarrow 1.0$ in Nuclear Environment

CTEQ QCD
Summer School
J. G. Morfin
June 2000

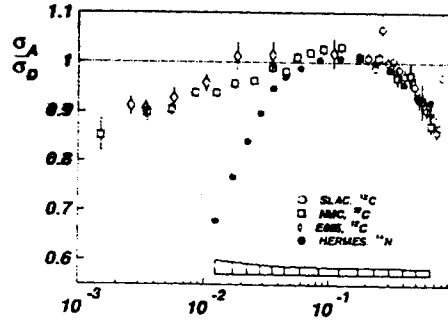
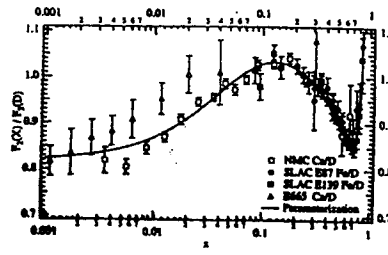
- Need to add more than Fermi gas model to simple nucleon model to reproduce behavior of F_2 at high x in nucleus.
- Few-nucleon-correlation and multi-quark models allow quarks to have higher momentum \rightarrow high x tail with $F_2 \propto e^{-ax}$.
- Analyses by SLAC, BCDMS, CEBAF and CCFR with values of a varying $7 < a < 17$ in various kinematical regions and targets.
- BCDMS and CCFR are in similar kinematical regions:

BCDMS ($\mu + C$): $a = 16.5 \pm 0.5$
CCFR ($\nu + Fe$): $a = 8.3 \pm 0.7 \pm 0.7(\text{syst.})$

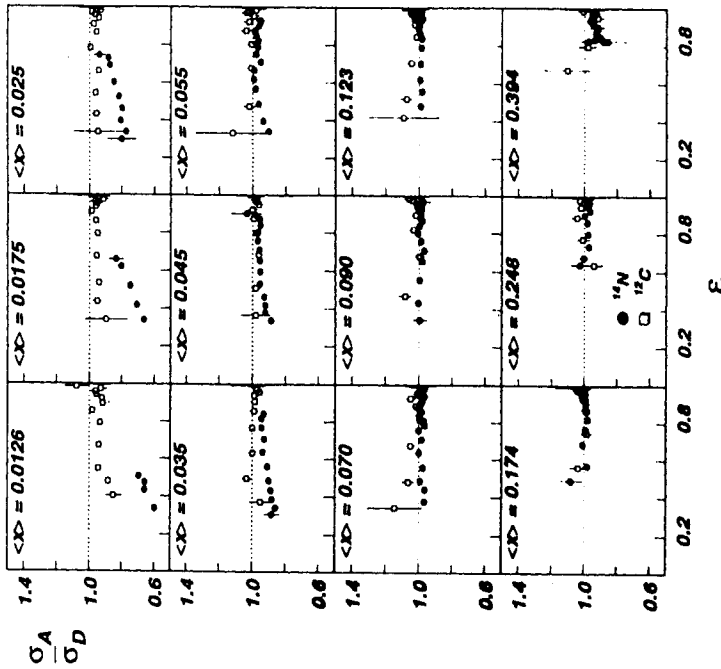
- Is a dependent on nucleus?
- Is a dependent on ν vs. μ ?



- NMC and E665 data for $Q^2 > 1.0$ and essentially no Q^2 dependence seen.
- HERMES data at $.01 < x < 0.1$ has $Q^2 < 1.0$.
- HERMES effect translates into an enhancement in F_L and a suppression in $2 \times F_1$ due to nuclear effects.
- Being tested with further data at HERMES and JLab.



More plots for HERMES effect



$$\bullet \frac{\sigma_A}{\sigma_D} = \frac{F_1^A(1+\epsilon R_A)(1+R_D)}{F_2^D(1+\epsilon R_D)(1+R_A)}$$

▷ So if $\epsilon \rightarrow 1$, $\frac{\sigma_A}{\sigma_D} = \frac{F_1^A}{F_2^D}$

$$\epsilon = \frac{2(1-y)}{1+(1-y)^2}$$

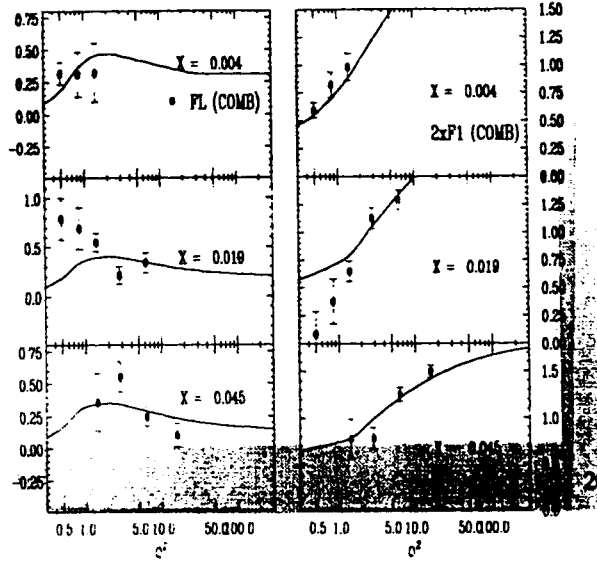
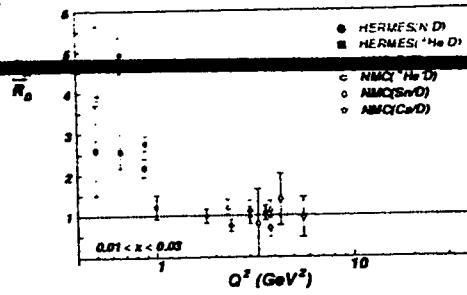
▷ A good agreement between NMC & HERMES N data

▷ Strong R dependence

In $e/p-N$ scattering

$$\frac{d\sigma}{d\Omega} \sim 2 \times F_1(1+\epsilon R)$$

- Use CCFR F_2 data and R_{world} (with no HERMES effect) to determine F_L and $2xF_1$.
- Observe an interesting blip at $x = 0.019$ in the "right" Q^2 region.
- Why no effect in the $x = .004$ bin?
- Any reason to expect a difference in nuclear dependence of R: ν vs. e/μ ?



Goals in Study of Nuclear Effects with $\nu/\bar{\nu}$ scattering

- Overall Goal: Measure nuclear effects across full x_{Bj} -range in $\nu/\bar{\nu}$ scattering off a variety of targets.
- Goal: Measure nuclear effects separately for F_2 and xF_3 . What are the nuclear effects for valence quarks alone?
- Goal: Determine nuclear dependence of R at low x and Q^2
- Long-term Goal: High statistics $\nu/\bar{\nu}$ scattering experiment on H_2 and D_2 as well as heavy nuclei.

■ NuMI

- Uses high intensity Main Injector with 120 GeV protons and horn focussed π & K beam for ν .
- MINOS (Oscillation Expt.) using low-energy ν beam on schedule for 2003.
- 2nd generation high energy DIS expt. in 2006?

■ Neutrino Factory

- Uses muon storage ring with production of ν_μ and ν_e via decay of μ^\pm
- Possible in 2010 - 2015 time frame



NuMI $\nu/\bar{\nu}$ events: Intensity and Kinematic Domain

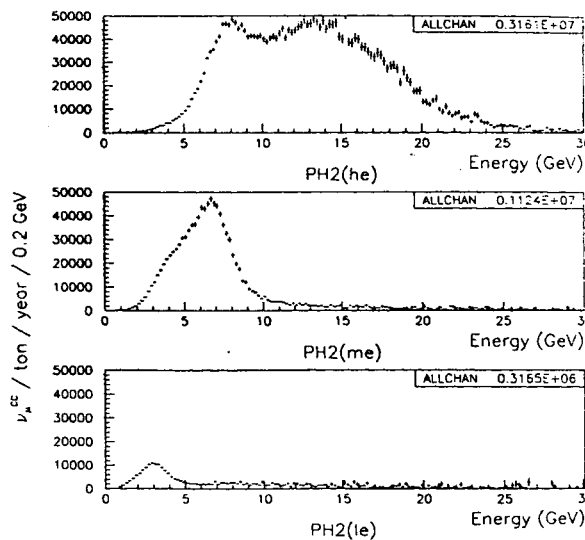
CTEQ QCD
Summer School
J. G. Morfin
June 2000

3.7×10^{20} protons/year
on the NuMI target
will yield

2.5×10^6 CC ν
 0.6×10^6 CC $\bar{\nu}$
events / year-ton

high energy-
configuration

Event Rates at MINOS Near Detector



	$W^2 > 5 \text{ GeV}^2$ (26 M events)						In units of 10^5 events/year	
$x = 0$	$Q^2 = 0$	2	4	8	16	$>30 \text{ GeV}^2$		
.03	22.6	.1	.04					
.05	15.6	.6	.1	.05				
.075	17.6	2.4	.5	.06	.01			
.10	14.9	5.6	.6	.2	.02			
.125	11.5	6.5	1.7	.25	.02			
.15	8.9	7.1	3.1	.2	.03			
.20	13.1	11.9	8.5	1.3	.3			
.30	12.1	18.3	17.8	6.5	.8	.01		
.45	2.2	13.6	17.1	11.9	2.0	.03		
.65	.9	1.7	6.9	6.9	2.8	.2		
1.0	.1	.5	1.4	.8	.5	.01		

(Expect $\approx 100,000$ events $x > 0.75$)



Modified Near Detector (MIDIS) for a High Energy Beam Run

CTEQ QCD
Summer School
J. G. Morfin
June 2000

- Standard MINOS near detector: Solid scintillator planes with fiber readout between 2.5 cm thick Fe plates (0.40 ton fiducial/plate) yielding 16 tons in 40 plane target region.
- This translates into 40 M ν events/yr. and 10 M $\bar{\nu}$ / yr. in the fiducial volume.

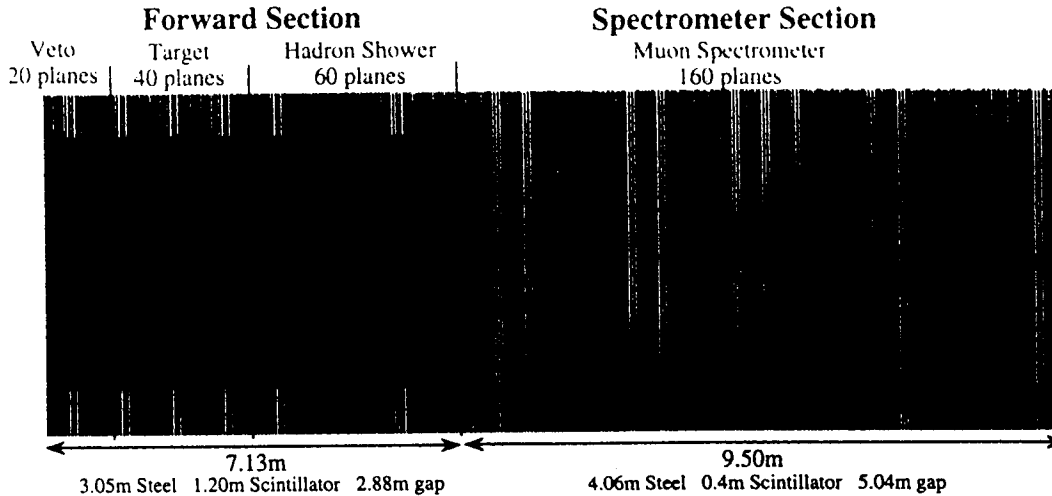
MIDIS Near Detector

- Retain 11 fiducial tons of Fe with 26 planes.
- Assume 12 (instrumented) planes of C (≈ 4 M fiducial events / yr.), 2 plane of W (5 M events) and 2 year ν and 2 year $\bar{\nu}$ run.
- Ratios: beam syst. cancel.
- Assume C-Fe-W relative target syst. same as Tevatron Muon Expt. O (1 %) added in quadrature.

Solid scintillator planes with fiber readout sandwiched between 2.5 cm thick Fe plates.

Instrumented out to $r \approx 1\text{m}$: 0.6 ton/plate

0.40 ton fiducial: 1.00 M ν events/plate - yr: $\bar{\nu} = \nu / 4.0$



13

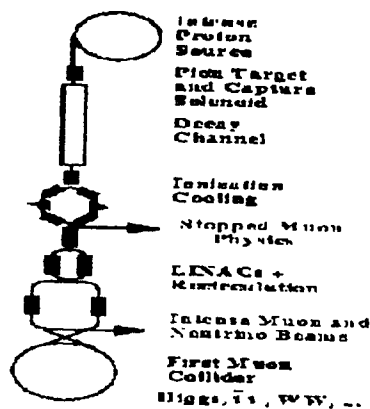


Further in the Future: DIS with Neutrinos from a Muon Storage-Ring/Collider

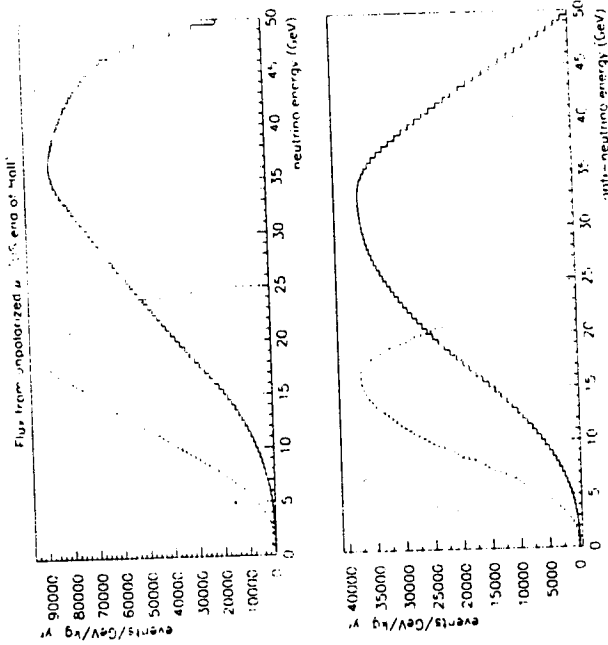
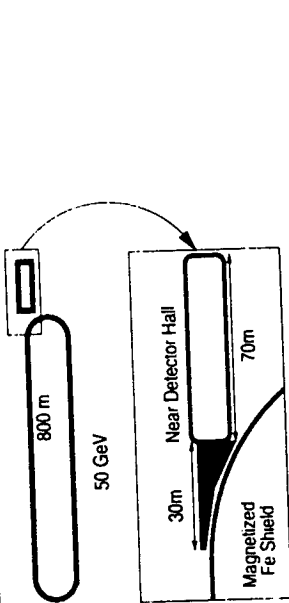
CTEQ QCD
Summer School
J. G. Morfin
June 2000

How does a muon collider work?

- Intense Proton Source (8 - 16 GeV, 10^{21-22} /year)
- Pion Target and Capture Solenoid
- Decay Channel (yield 10^{20-21} muons /year)
- Ionization Cooling Section
- LINAC + Storage Ring
- Muon Collider
- DIS from μP collider and intense ν beams



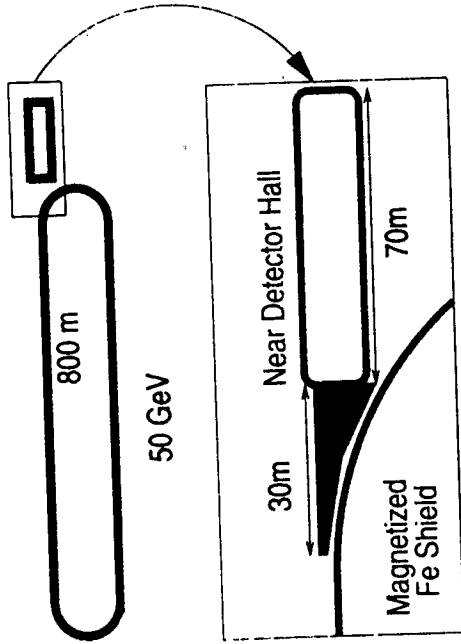
Fluxes



50, 25, 10 GeV μ beams shown
 If lengths scale with $E_\mu \implies$ short-baseline $N_\nu \propto E_\nu$

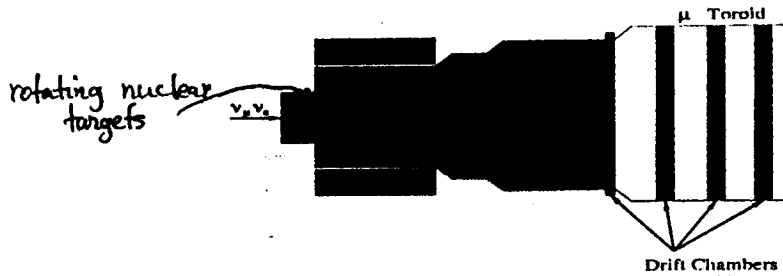
The Opportunity

10^{20} /yr μ decays in the green straight section \implies
 5-8% of all interactions in detector 10 cm in radius
 40-50% of all interactions in detector 50 cm in radius
 $\approx 1.5-3 \times 10^6 \times \frac{E_\mu}{50 \text{ GeV}}/\text{kg/yr}$ at beam center
 (FNAL Neutrino Factory Study Group Parameters)



c.f. *Journal of High Energy Physics*

Beam	$\langle E_\nu \rangle$ [GeV]	ν per
NuTeV/CCFR (Fermilab)	100	
CHORUS/NOMAD (CERN)	30	
MINOS Near (Fermilab)	15	
Neutrino Factory	30	$5 \sim 10$



- Liquid deuterium cryogenic target followed by rotating high A targets.
- All targets ($r = 10$ cm) acquire $10^7 \nu_\mu$ events and $4 \times 10^6 \bar{\nu}_e$ /year.
- D_2 @ 50 cm long, C @ 3.7 cm and W @ 0.45 cm



Compare Event Rates for Two Future Experiments

CTEQ QCD
Summer School
J. G. Morfin
June 2000

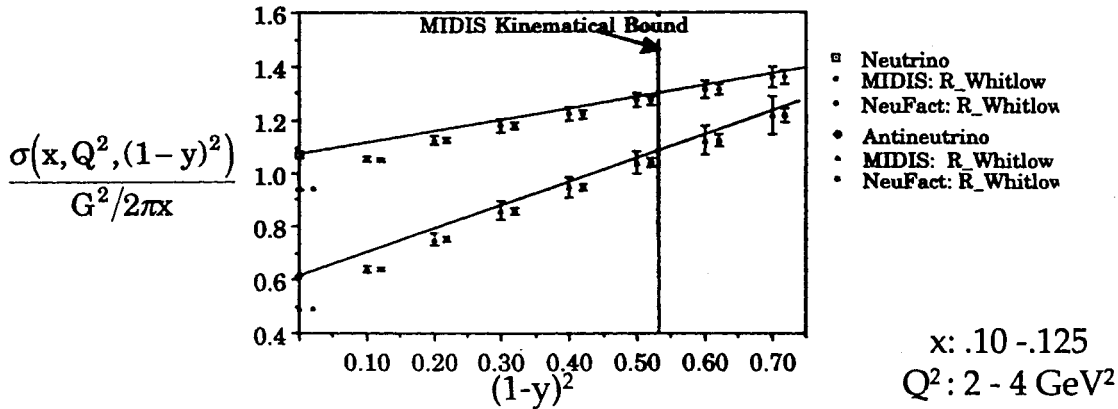
MIDIS in the NuMI Beam

- 3.7×10^{20} protons/year on the NuMI target.
- 2.5×10^6 CC ν or 0.6×10^6 CC $\bar{\nu}$ events /ton-year (in high energy configuration)
- Know $\phi(E_\nu)$ and $\int \phi(E_\nu)$ to $< 2\%$
- First Beam 2003.
HE beam run > 2006

Near Detector/50 GeV NuFact (50 m away from 800 m straight section)

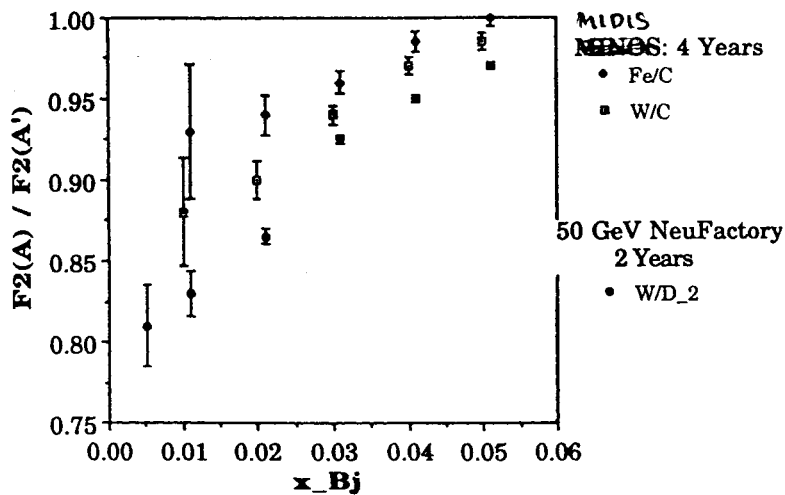
- $10^{20} \mu$ decays/year
- 4.0×10^6 CC ν_μ and 1.7×10^6 CC ν_e events/kg-year (fid. vol. $r < 10$ cm)
- Know $\phi(E_\nu)$ and $\int \phi(E_\nu)$ to $< 1.0 \%$
- First Beam a bit after MINOS runs are finished

- MIDIS: 26 M v (Fe) evts (6.5 M $\bar{\nu}$ evts) /year in fid. vol.
- Analysis based on four year run
- NuFact: Analysis based on 2 year run
- Statistical errors **negligible** at small values of $(1-y)^2$
- R value now dominated by systematic errors

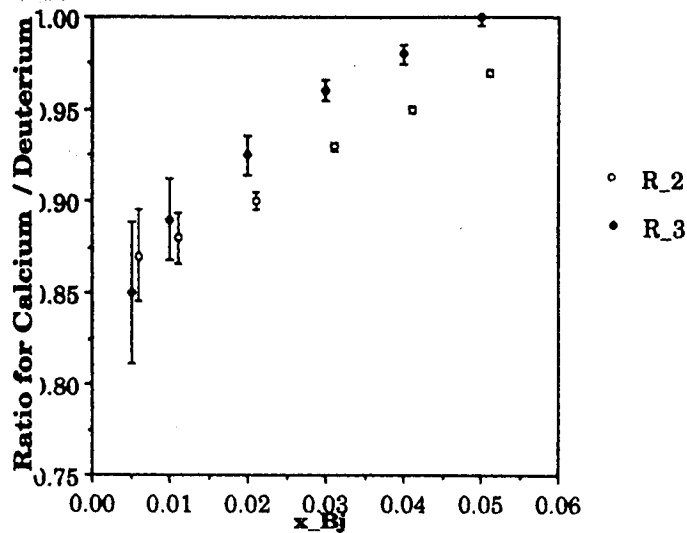


MIDIS / NuFact v Nuclear Shadowing Measurement

CTEQ QCD
 Summer School
 J. G. Morfin
 June 2000



Kulagin's Non Perturbative Parton Model



SUMMARY

CTEQ QCD
Summer School
J. G. Morfin
June 2000

- NuMI/MIDIS will offer an intense, lower energy ν beam to explore with impressive statistics and reasonable systematics among other topics:
 - a first look at precision measurements of separate $\nu / \bar{\nu}$ F_2 , xF_3 & R .
 - study of $\nu / \bar{\nu}$ nuclear effects, particularly HERMES effect!

- NuFact, a "few years" after MIDIS, with H_2 , D_2 , and nuclear targets as well as polarized targets:
 - improve statistical and systematic errors of MIDIS measurements.
 - Complete measurements of the partonic spin structure of nucleon.
 - High statistics charm production with "clean" final state.
 - Precision electroweak measurements; ν -e scattering.

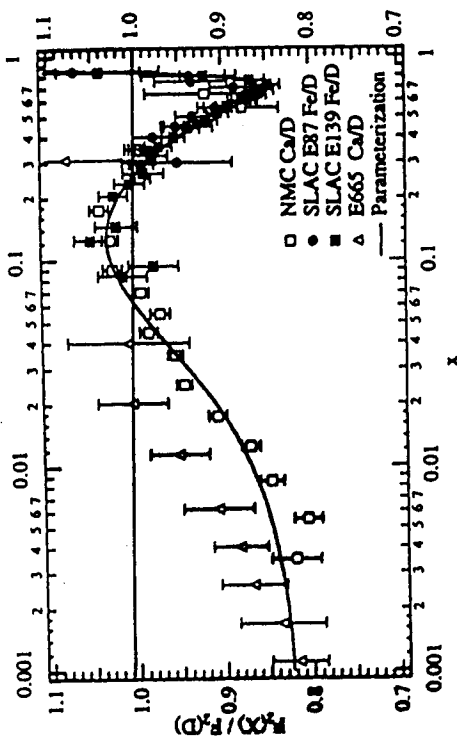
Nuclear Effects in D.I.S

ν -DIS experiments needed nuclear targets to increase luminosity ($R_{Fe}/R_{H_2} \approx 10^2$)

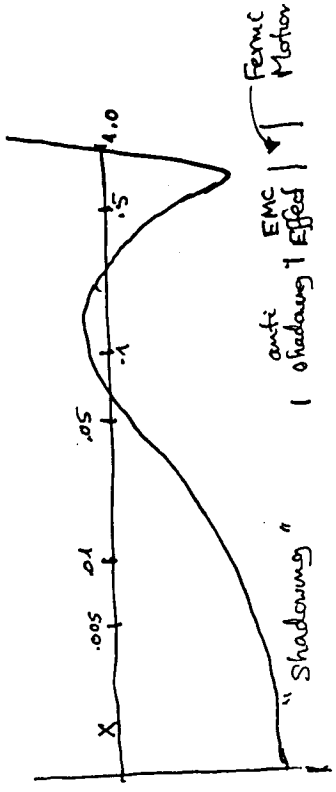
In 1983 EMC compared results of μ -Fe with μ -D₂ and discovered $F_2^{D_2} \neq F_2^{Fe}$

SLAC experimenters went back and reanalysed old data and confirmed "EMC Effects" Then ran a dedicated experiment

Further experimental results: BCDMS, NMC, E665 extended range in x_Bj \leftarrow α^2 increased statistics and minimized systematic errors



$$\frac{F_2^A}{F_2^{D_2}}$$



- 1) Shadowing
 - a) Vector Meson Dominance (primarily low x)



- b) Parton Fusion Models
 - Uncertainty Principle $\Delta z \sim \frac{1}{x}$

For low x Δz greater than nucleon size \Rightarrow partons from neighboring nucleons or quarks at higher x created by $g \bar{g}$ pairs at k or $g \bar{g} \rightarrow g$ at higher x

- 2) Anti-Shadowing
 - a) Momentum Sum Rule
 - b) Regge models
 - c) Contributions from real part of soft amplitudes

EMC Effect

- a) Conventional Nuclear Physics - Bound nuclear mass \neq free nuclear mass $x = Q^2/2M^2$
Can reproduce data but internally inconsistent
- b) Cluster Models - multiquark (6, 9, 12, ...) clusters in the nucleus or presence of a π or a Δ in the nucleus ...
- c) QCD models - reproduce the data

Fermi Motion

Smeared measured F_2

