Lecture 15 Review

Fractal dimension of Sierpinski's gasket .Affine transformations and fractal examples.Root finding: N-R and bi-section algorithm.

HW aside: Newton-Raphson Beware

Software ain't magic...

Potential pathologies using N-R root finding technique



NR solution potentially sensitive to initial guess. <u>Beware</u>. (Graph, too.)

Particle in a Quantum Box



$$-\frac{\hbar^2}{2m}\frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x)$$
$$\mathcal{P} = |\psi(x)|^2 dx \qquad \qquad \int_{-\infty}^{\infty} dx \, |\psi(x)|^2 = 1$$

$$V(x) = \left\{ egin{array}{cc} -V_0 = -83\,{
m MeV}, & {
m for}\, |x| \leq a = 2\,{
m fm} \ 0, & {
m for}\, |x| > a = 2\,{
m fm} \end{array}
ight.$$

$$\frac{d^2\psi(x)}{dx^2} + \frac{2m}{\hbar^2}(E+V_0)\psi(x) = 0 \qquad \text{for } |\mathbf{x}| \le \mathbf{a}$$
$$\frac{d^2\psi(x)}{dx^2} + \frac{2m}{\hbar^2}E\psi(x) = 0 \qquad \text{for } |\mathbf{x}| > \mathbf{a}$$

$$\frac{2m}{\hbar^2} = \frac{2mc^2}{(\hbar c)^2} = \frac{2 \times 940 \,\mathrm{MeV}}{(197 \,\mathrm{Mev} - \mathrm{fm})^2} = 0.0483 \,\mathrm{MeV}^{-1} \mathrm{fm}^{-2}$$

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Particle in a Quantum Box

$$\psi(x) = \begin{cases} Ce^{\beta x}, & \text{for } -\infty < x < -a \\ B \cos \alpha x, & \text{for } -a < x < a \\ Ce^{-\beta x}, & \text{for } a < x < +\infty \end{cases} \qquad \alpha = \sqrt{\frac{2m(E+V_0)}{\hbar^2}} = \sqrt{0.0483(E+83)}$$

$$\beta = \sqrt{\frac{-2mE}{\hbar^2}} = \sqrt{-0.0483E}$$

$$B \cos \alpha a = Ce^{-\beta a} \quad \psi \text{ continuity}$$

$$-\alpha B \sin \alpha a = -\beta Ce^{-\beta a} \quad \psi' \text{ continuity}$$

$$\Rightarrow \alpha a \tan \alpha a - \beta a = 0$$

$$\sqrt{2m(E+V_0)} \tan \sqrt{\frac{2m(E+V_0)}{\hbar^2}} a - \sqrt{-2mE} = 0$$

$$\xi \tan \xi - \eta = 0 \quad \text{with} \quad \xi = \alpha a \quad \eta = \beta a$$

$$\xi^2 + \eta^2 = \frac{2mV_0a^2}{\hbar^2} = 16.08$$

$$f(E) = \xi \tan \xi - \eta = 0 \quad \longleftarrow \quad \text{Even parity solutions only.}$$

Particle in a Quantum Box (2)

$$\psi(x) = \begin{cases} Ce^{\beta x}, & \text{for } -\infty < x < -a \\ B \sin \alpha x, & \text{for } -a < x < a \\ Ce^{-\beta x}, & \text{for } a < x < +\infty \end{cases} \qquad \alpha = \sqrt{\frac{2m(E+V_0)}{\hbar^2}} = \sqrt{0.0483(E+83)}$$

$$B \sin \alpha a = Ce^{-\beta a} \quad \text{if } continuity$$

$$\alpha B \sin \alpha a = -\beta Ce^{-\beta a} \quad \text{if } continuity$$

$$\Rightarrow \alpha a \cot \alpha a + \beta a = 0$$

$$\sqrt{2m(E+V_0)} \cot \sqrt{\frac{2m(E+V_0)}{\hbar^2}} a + \sqrt{-2mE} = 0$$

$$\xi \cot \xi + \eta = 0 \quad \text{with} \quad \xi = \alpha a \quad \eta = \beta a$$

$$\xi^2 + \eta^2 = \frac{2mV_0a^2}{\hbar^2} = 16.08$$

$$f(E) = \xi \cot \xi + \eta = 0 \quad \longleftarrow \quad \text{Odd parity solutions only.}$$

Numerov Method for ODE Solution

Runge-Kutta (4th order) is our preferred solver for ODEs. Why not use it to solve Schrödinger equation? For <u>special case</u> of dy/dx missing (e.g., Schrödinger eqn) Use Numerov's method:

More accurate than RK-4th order AND w/ fewer steps.

When applied to potential wells, Numerov's method is an example of solving ODE w/ <u>boundary conditions</u> (as opposed to ICs).

Numerov Method for ODE Solution (2)

$$-rac{\hbar^2}{2m}rac{d^2\psi(x)}{dx^2}+V(x)\psi(x)=E\psi(x)$$

$$rac{d^2\psi(x)}{dx^2}+k^2(x)\psi(x)=0$$

$$k^2(x) \equiv (rac{2m}{\hbar^2}) \left\{ egin{array}{cc} E - V_0 & |x| < a \ E & |x| > a \end{array}
ight.$$

Recall that E < 0 for bound states

$$(rac{2m}{\hbar^2})=0.04829\,{
m Mev^{-1}\,fm^{-2}}$$



Numerov Method for ODE Solution (3)

$$\begin{split} \psi(x+h) &= \psi(x) + h\psi' + \frac{h^2}{2}\psi'' + \frac{h^3}{3!}\psi''' + \frac{h^4}{4!}\psi^{\text{IV}} + \dots \\ \psi(x-h) &= \psi(x) - h\psi' + \frac{h^2}{2}\psi'' - \frac{h^3}{3!}\psi''' + \frac{h^4}{4!}\psi^{\text{IV}} + \dots \\ &\Rightarrow \psi(x+h) + \psi(x-h) \simeq 2\psi(x) + h^2\psi'' + \frac{h^4}{12}\psi^{\text{IV}} + \mathcal{O}(h^6) \\ \psi''(x) &\simeq \frac{\psi(x+h) - \psi(x-h) - 2\psi(x)}{h^2} - \frac{h^2}{12}\psi^{\text{IV}}(x) + \mathcal{O}(h^6) \\ \hline \text{Trick:} \quad \text{Apply} \ 1 + \frac{h^2}{12}\frac{d^2}{dx^2} \quad \text{to TISE} \\ &\left(1 + \frac{h^2}{12}\frac{d^2}{dx^2}\right) \left(\frac{d^2\psi}{dx^2} + k^2(x)\psi\right) = 0 \end{split}$$

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Numerov Method for ODE Solution (5)

$$\psi(x+h) \simeq rac{2[1-rac{5}{12}h^2k^2(x)]\psi(x)-[1+rac{h^2}{12}k^2(x-h)]\psi(x-h)}{1+rac{h^2}{12}k^2(x+h)}$$

$$\psi_{i+1} \simeq rac{2(1-rac{5}{12}h^2k_i^2)\psi_i - (1+rac{h^2}{12}k_{i-1}^2)\psi_{i-1}}{1+rac{h^2}{12}k_{i+1}^2}$$

See numerov.cc

numerov.cc discussion

$$\psi_{i+1} \simeq \frac{2(1 - \frac{5}{12}h^2k_i^2)\psi_i - (1 + \frac{h^2}{12}k_{i-1}^2)\psi_{i-1}}{1 + \frac{h^2}{12}k_{i+1}^2}$$

Basic algorithm. Start w/ ψ_0 and ψ_1 specified.

-a

a

 $-V_0$

Symmetric QM wells have ψ 's of definite <u>parity</u>. Recall, $\psi(-x) = \psi(x)$, "even" parity — $\psi(-x) = -\psi(x)$, "odd" parity

 $|\psi(x)|^2$ and $|\psi(-x)|^2$ have physical significance. ψ and ψ ' must be continuous.

numerov.cc lets you pick parity. (By selecting sign of ψ_1 .) numerov.cc as coded does not check for continuity of ψ' . X

Numerov Technique Lab Exercise

Modify numerov.cc

Limit number of possible iterations to, say, 1000. Modification should tell you if you reached this limit.

> Alter potential from a square well to a V-shaped well.



Find energy eigenvalues and eigenstates

(i.e., find permissible energy levels & corresponding wavefunctions.)

Summary

Application of root finding to particle in a quantum box Numerov technique for ODE solution. Numerical solution of square well potential.



