

Preliminaries:Take $n, m \in \mathbb{Z}$ integers

$$\begin{aligned}
 I &= \int_0^{2\pi} dx \, e^{inx} \cdot \left(e^{imx} \right)^* = \int_0^{2\pi} dx \, e^{i(n-m)x} \\
 &= \frac{e^{i(n-m)x}}{i(n-m)} \Big|_0^{2\pi} = \frac{e^{i(n-m)2\pi} - e^0}{i(n-m)}
 \end{aligned}$$

$$= 0 \quad \text{if } n \neq m$$

If $n = m$:

$$I = \int_0^{2\pi} dx \cdot 1 = 2\pi$$

$$\circ \circ \int_0^{2\pi} dx \, e^{inx} \left(e^{imx} \right)^* = \begin{cases} 0 & \text{if } n \neq m \\ 2\pi & \text{if } n = m \end{cases}$$

$$\equiv 2\pi \delta_{n,m}$$

$$\int_0^{2\pi} dx \, e^{inx} \left(e^{imx} \right)^* = 2\pi \delta_{n,m}$$

Orthogonality!!!

Likewise:

$$\int_0^{2\pi} dx \sin(nx) \sin(mx) = \begin{cases} \pi \delta_{n,m} \\ 0 \quad \text{For } n=m=0 \end{cases}$$

$$\int_0^{2\pi} dx \cos(nx) \cos(mx) = \begin{cases} \pi \delta_{n,m} \\ 2\pi \quad \text{For } n=m=0 \end{cases}$$

$$\int_0^{2\pi} dx \sin(nx) \cos(mx) = 0$$

Short cut:

$$\sin^2 + \cos^2 = 1$$

Suggests

$$\overline{\sin^2} = \overline{\cos^2} = \frac{1}{2}$$

$$\int_0^{2\pi} dx \sin^2 = \int_0^{2\pi} dx \cdot \frac{1}{2} = \frac{1}{2} \cdot 2\pi = \pi$$

Likewise:

$$\int_0^{2\pi} dx \cos^2 = \pi$$

Fourier Series :

$$F(x) = \sum_{n=-\infty}^{\infty} C_n e^{c'n x}$$

Goal: Find C_n :

$$\int_0^{2\pi} dx e^{-c'm x} F(x) = \int_0^{2\pi} dx \sum_{n=-\infty}^{\infty} C_n e^{c'(n-m)x}$$

$$= \sum_n C_n \underbrace{\int_0^{2\pi} dx e^{c'(n-m)x}}_{2\pi \delta_{n,m}} = 2\pi C_m$$

$$C_m = \frac{1}{2\pi} \int_0^{2\pi} dx e^{-c'm x} F(x)$$

Similar Derivation :

$$F(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(nx) + \sum_{n=1}^{\infty} b_n \sin(nx)$$

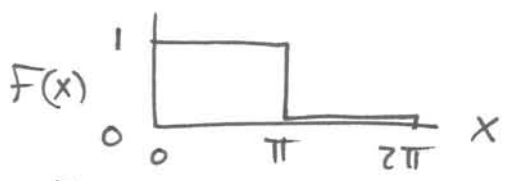
$$a_n = \frac{1}{\pi} \int_0^{2\pi} dx \cos(nx) F(x)$$

$$b_n = \frac{1}{\pi} \int_0^{2\pi} dx \sin(nx) F(x)$$

Note: $C_n = \frac{1}{2} (a_n - i b_n) \equiv \frac{1}{2\pi} \int_0^{2\pi} dx e^{-c'n x} F(x)$

Example #1-A

C.F., P. 819



$$F(x) = \begin{cases} 1; & x \in [0, \pi] \\ 0; & x \in [\pi, 2\pi] \end{cases}$$

$$a_0 = \frac{1}{\pi} \int_0^{\pi} dx = 1$$

$$a_n = \frac{1}{\pi} \int_0^{2\pi} dx \cos(nx) F(x) = \frac{1}{\pi} \int_0^{\pi} dx \cos(nx) \cdot 1 = 0$$

$$b_n = \frac{1}{\pi} \int_0^{2\pi} dx \sin(nx) F(x) = \frac{1}{\pi} \int_0^{\pi} dx \sin(nx) \cdot 1$$

$$= \frac{1}{\pi} \left[\frac{-\cos(nx)}{n} \right]_0^{\pi} = \frac{1}{\pi} \left[\frac{-(-1)^n + 1}{n} \right] = \begin{cases} \frac{2}{n\pi} & n = \text{odd} \\ 0 & n = \text{even} \end{cases}$$

∞

$$F(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} b_n \sin(nx)$$

$$= \frac{a_0}{2} + \sum_{n=1,3,5,\dots} \frac{2}{n\pi} \sin(nx)$$

$$= \frac{1}{2} + \frac{2}{\pi} \left[\frac{\sin(x)}{1} + \frac{\sin(3x)}{3} + \frac{\sin(5x)}{5} + \dots \right]$$

Example #1-B

$$C_n = \frac{1}{2\pi} \int_0^{2\pi} dx \left(e^{c'nx} \right)^* F(x) = \frac{1}{2\pi} \int_0^{2\pi} dx e^{-c'nx} F(x)$$

$$= \frac{1}{2\pi} \int_0^{\pi} dx e^{-c'nx} \cdot 1 = \frac{1}{2\pi} \left[\frac{e^{-c'nx}}{-c'n} \right]_0^{\pi}$$

$$= \frac{1}{2\pi} \left[\frac{e^{-c'n\pi} - 1}{-c'n} \right] = \frac{1}{2\pi} \left[\frac{(-1)^n - 1}{-c'n} \right] =$$

$$C_n = \begin{cases} \frac{-i}{n\pi} & n = \text{odd} \\ 0 & n = \text{even} \end{cases}$$

$$C_0 = \frac{1}{2\pi} \int_0^{2\pi} dx \cdot 1 \cdot F(x) = \frac{1}{2}$$

$$F(x) = \sum_{n=-\infty}^{\infty} C_n e^{c'nx} = \frac{1}{2} + \sum_{\substack{n=1,3,5 \\ -1,-3,-5}} \left(\frac{-i}{n\pi} \right) e^{c'nx}$$

$$= \frac{1}{2} + \sum_{n=1,3,5} \frac{-i \cos(nx) + \sin(nx)}{n\pi}$$

$$+ \sum_{+n=1,3,5} \frac{-i \cos(-nx) + \sin(-nx)}{-n\pi}$$

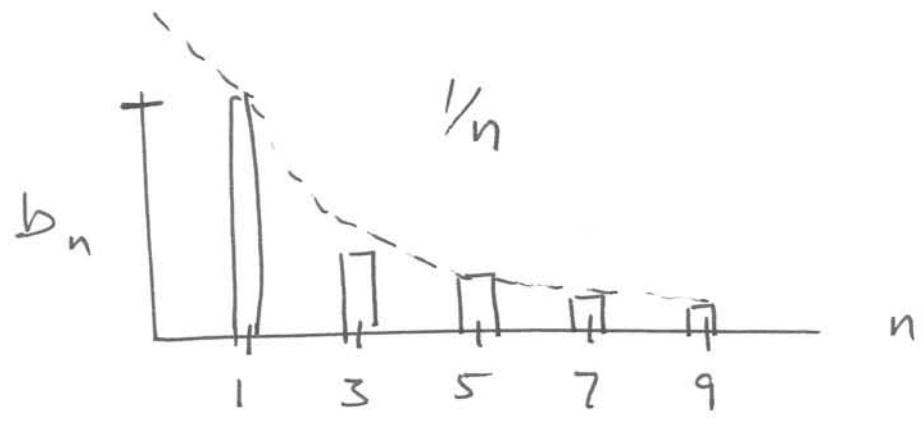
$$= \frac{1}{2} + \sum_{n=1,3,5} \frac{\sin(nx)}{n\pi} + \sum_{+n=1,3,5} \frac{-\sin(nx)}{-n}$$

$$= \frac{1}{2} + \sum_{n=1,3,5} \left(\frac{2}{n\pi} \right) \sin(nx)$$

$$= \frac{1}{2} + \frac{2}{\pi} \left[\frac{\sin(x)}{1} + \frac{\sin(3x)}{3} + \frac{\sin(5x)}{5} + \dots \right]$$

Same as Example #1-A

$$b_n = \frac{2}{\pi} \frac{1}{n} \text{ for } n = \text{odd}$$



Continuous Fourier Transform :

$$F(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dw \, g(w) e^{-i\omega x}$$

$$g(w) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dx \, F(x) e^{+i\omega x}$$

Check :

$$F(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dw \left\{ \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dx' F(x') e^{i\omega x'} \right\} e^{-i\omega x}$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} dx' F(x') \int_{-\infty}^{\infty} dw e^{i\omega(x'-x)}$$

$\underbrace{\int_{-\infty}^{\infty} dw e^{i\omega(x'-x)}}_{2\pi \delta(x'-x)}$

$$= 1 \cdot \int_{-\infty}^{\infty} dx' F(x') \delta(x'-x)$$

$$F(x) = F(x)$$

Note :

$$2\pi \delta(x-x') = \int dw e^{i\omega(x-x')}$$

Green's Functions :

Hard

Problem : $\mathbb{D} F(x) = \phi(x)$

Easier

Problem : $\mathbb{D} G(x) = \delta(x)$

or $\mathbb{D} G(x-x') = \delta(x-x')$

If we solve easier problem, we can solve hard problem, because :

$$F(x) = \int_{-\infty}^{\infty} G(x-x') \phi(x') dx'$$

Check :

$$\mathbb{D} F(x) = \int_{-\infty}^{\infty} \underbrace{\mathbb{D} G(x-x')}_{\delta(x-x')} \phi(x') dx'$$

$$= \int_{-\infty}^{\infty} dx' \delta(x-x') \phi(x')$$

$$\mathbb{D} F(x) = \phi(x)$$

Q.E.D.

Now, how to solve easy problems:

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Use Fourier Transforms.

Turns differential eq \Rightarrow algebraic eq.

Ex. $(a D^2 + b D + c) F(x) = \delta(x)$

Use: $F(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} d\omega g(\omega) e^{-i\omega x}$

$\delta(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{\pm i\omega x}$ } either sign works

∞

$(a D^2 + b D + c) \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} d\omega g(\omega) e^{-i\omega x}$

$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} d\omega g(\omega) [-\omega^2 a - i\omega b + c] e^{-i\omega x}$

$\equiv \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{-i\omega x}$

∞ $\Rightarrow g(\omega) [-\omega^2 a - i\omega b + c] = \frac{1}{\sqrt{2\pi}}$

$g(\omega) = \frac{1}{\sqrt{2\pi}} \frac{1}{[-\omega^2 a - i\omega b + c]}$

Full solution:

$$F(x) = \frac{1}{\sqrt{2\pi}} \int dw g(w) e^{-i\omega x}$$

$$= \frac{1}{\sqrt{2\pi}} \int dw \frac{e^{-i\omega x}}{[-\omega^2 a - i\omega b + c]}$$

Example

Schrodinger Equation:

Guess answer is wave: $\psi(x) = e^{+iPx/\hbar}$

$$\nabla \psi = \frac{iP}{\hbar} \psi \Rightarrow P = -i\hbar \nabla$$

$$\nabla^2 \psi = -\frac{P^2}{\hbar^2} \psi \Rightarrow P^2 = -\hbar^2 \nabla^2$$

$$E = \frac{1}{2} m v^2 = \frac{P^2}{2m} \quad \text{with } P = mv$$

∞ $E = \frac{P^2}{2m} = \frac{-\hbar^2 \nabla^2}{2m}$

or $E \psi = \frac{-\hbar^2 \nabla^2}{2m} \psi$

check: $\psi = e^{+iPx/\hbar}$

$$\nabla^2 \psi = -\frac{P^2}{\hbar^2} \psi$$

∞ $E \psi = \frac{-\hbar^2}{2m} \nabla^2 \psi = \frac{-\hbar^2}{2m} \left(-\frac{P^2}{\hbar^2} \psi \right) = \frac{+P^2}{2m} \psi$

$\Rightarrow E \psi = \frac{P^2}{2m} \psi$ check.

Momentum Representation

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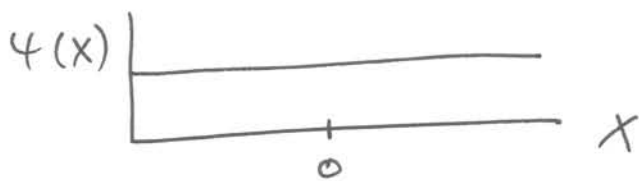
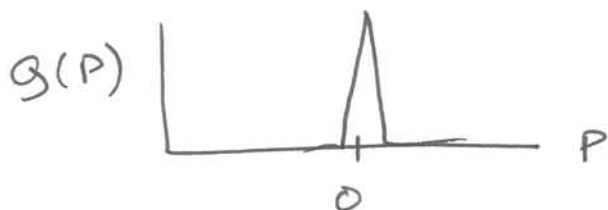
$$\psi(x) = \frac{1}{\sqrt{2\pi\hbar}} \int dp g(p) e^{+iPx/\hbar}$$

$$g(p) = \frac{1}{\sqrt{2\pi\hbar}} \int dx \psi(x) e^{-iPx/\hbar}$$

Example: Let $g(p) = \delta(p)$

$$\psi(x) = \frac{1}{\sqrt{2\pi\hbar}} \int dp \delta(p) e^{iPx/\hbar}$$

$$= \frac{1}{\sqrt{2\pi\hbar}} \cdot 1 = \text{const.}$$

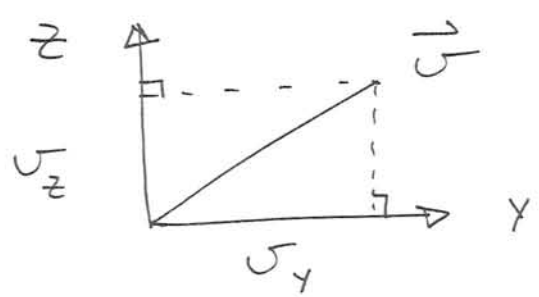


What do you expect for:

a) $g(p) = e^{-p^2 x^2 / \hbar^2}$



Fourier Transforms: (An Analogy)



$$u_y = \vec{y} \cdot \vec{u} = \langle y | u \rangle$$

$$u_z = \vec{z} \cdot \vec{u} = \langle z | u \rangle$$

$$\vec{u} = \hat{y} u_y + \hat{z} u_z$$

$$|u\rangle = |y\rangle \langle y | u \rangle + |z\rangle \langle z | u \rangle$$

"divide" by $|u\rangle$

$$\hat{1} = |y\rangle \langle y| + |z\rangle \langle z| \equiv \sum_i |x_i\rangle \langle x_i|$$

\vec{x}_i vector \uparrow
 Projection on \vec{x}_i \uparrow

Works for any vector

$$|w\rangle = \left(\sum_i |x_i\rangle \langle x_i| \right) |w\rangle$$

$$= \sum_i |x_i\rangle \langle x_i | w \rangle$$

basis vector \uparrow projection coefficient a number \uparrow

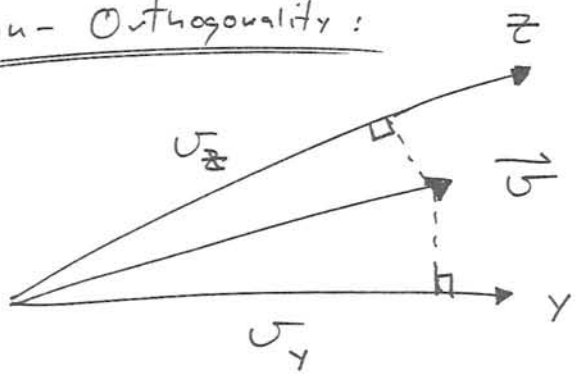
Orthogonality :



$$x_i \cdot x_j = \delta_{ij}$$

$$\langle x_i | x_j \rangle = \delta_{ij}$$

Non-Orthogonality :



$$\sigma_y = y \cdot \sigma = \langle y | \sigma \rangle$$

$$\sigma_z = z \cdot \sigma = \langle z | \sigma \rangle$$

Note:

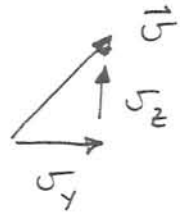
$$\vec{\sigma} \neq \sigma_y \hat{y} + \sigma_z \hat{z}$$

$$\neq |y\rangle\langle y| + |z\rangle\langle z|$$

$$\neq \sum_i |x_i\rangle\langle x_i|$$

Completeness :

$$\hat{1} = |y\rangle\langle y| + |z\rangle\langle z| = \sum_i |x_i\rangle\langle x_i|$$



$$\vec{\sigma} = \hat{y} \sigma_y + \hat{z} \sigma_z$$

$$|\sigma\rangle = |y\rangle\langle y|\sigma\rangle + |z\rangle\langle z|\sigma\rangle$$

This would not work if we left out - say - y

$$\hat{1} \neq 0 + |z\rangle\langle z| \quad \vec{\sigma} \neq 0 + \hat{z} \sigma_z$$

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Change from vectors $\vec{v} = |v\rangle$ to functions $F(x) = |F\rangle$

$$F(x) = |F\rangle = 2 \sin(2x) + 3 \sin(3x)$$

$$|F\rangle = 2 |X_2\rangle + 3 |X_3\rangle$$

$$|X_2\rangle = \sin(2x)$$

$$|X_3\rangle = \sin(3x)$$

$$\langle a | b \rangle = \frac{1}{\pi} \int_0^{2\pi} dx \ a^* b$$

$$\circ\circ \quad |F\rangle = (|X_2\rangle\langle X_2| + |X_3\rangle\langle X_3|) |F\rangle$$

$$= |X_2\rangle \langle X_2 | F \rangle + |X_3\rangle \langle X_3 | F \rangle$$

$$\langle X_2 | F \rangle = \frac{1}{\pi} \int_0^{2\pi} dx \ \sin(2x) [2 \sin(2x) + 3 \sin(3x)]$$

$$= \frac{1}{\pi} \int_0^{2\pi} dx \ \left[2 \underbrace{\sin^2(2x)}_{\frac{1}{2}} + 3 \underbrace{\sin(2x) \sin(3x)}_0 \right]$$

$$= \frac{1}{\pi} \cdot 2\pi \cdot 2 \cdot \frac{1}{2} = 2$$

Likewise

$$\langle X_3 | F \rangle = \frac{1}{\pi} \cdot 2\pi \cdot 3 \cdot \frac{1}{2} = 3$$

$$\circ\circ \quad |F\rangle = |X_2\rangle \langle X_2 | F \rangle + |X_3\rangle \langle X_3 | F \rangle$$

$$= \sin(2x) \cdot 2 + \sin(3x) \cdot 3$$

A perfect match.
(designed that way)