

Lagrangian

$$L(q, \dot{q})$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}} \right) - \frac{\partial L}{\partial q} = 0$$

$$\frac{\partial L}{\partial \dot{q}} = \overset{\circ}{p}$$

$$\frac{\partial L}{\partial q} = -\overset{\circ}{p}$$

Hamiltonian

$$H(q, p)$$

$$H(q, p) = \overset{\circ}{q} p - L(q, \dot{q})$$

$$\frac{\partial H}{\partial p} = \overset{\circ}{q}$$

$$\frac{\partial H}{\partial q} = -\overset{\circ}{p}$$

SHO Lagrangian

$$L = T - V = \frac{m \dot{q}^2}{2} - \frac{k q^2}{2}$$

$$E_{eqs} \Rightarrow m \ddot{q} + k q = 0$$

$$q = e^{i\omega t} \Rightarrow \omega = \sqrt{\frac{k}{m}}$$

One 2nd-order DEQ.

SHO Hamiltonian

$$H = T + V = \frac{p^2}{2m} + \frac{k q^2}{2}$$

$$\frac{\partial H}{\partial p} = \frac{p}{m} = \overset{\circ}{q}$$

$$\frac{\partial H}{\partial q} = k q = -\overset{\circ}{p}$$

$$p = m \dot{q}$$

$$-k q = \overset{\circ}{p}$$

$$p = m v$$

$$-k x = F$$

Hooke's Law

Goldstein ch 9.3

Canonical TX

1

$$H = \frac{P^2}{2m} + \frac{Kq^2}{2} \equiv \frac{1}{2m} (P^2 + m^2 \omega^2 q^2)$$

FF

$$P = F(P) \cos Q$$
$$q = F(P) \sin Q$$

$$K = H = \frac{F^2(P)}{2m}$$

Cyclic in Q

Generating Function $F_1(q, Q) = \frac{m\omega q^2}{2} \cot Q$

For $F_1(q, Q)$ $P = \frac{\partial F_1}{\partial q} = m\omega q \cot Q$ (1)

$$Q = -\frac{\partial F_1}{\partial P} = \frac{m\omega q^2}{2 \sin^2 Q}$$
 (2)

$$K = H + \frac{\partial F_1}{\partial t} \equiv H$$

(2) $\Rightarrow q^2 = \frac{2P}{m\omega} \sin^2 Q$

(1) $\Rightarrow P^2 = m\omega^2 q^2 \frac{2P}{m\omega} \sin^2 Q \frac{\cos^2 Q}{\sin^2 Q} = 2m\omega P \cos^2 Q$

$$q = \sqrt{\frac{2P}{m\omega}} \sin Q$$

$$H = K = \omega P$$

$$P = \sqrt{2P m \omega} \cos Q$$

$$F(P) = \sqrt{2P m \omega}$$

Goldstein
TYPO

$$K(Q, P) = \omega P \equiv H(Q, P)$$

2

$$\frac{\partial H}{\partial P} = \omega = +\dot{Q} \Rightarrow Q = \omega t + \alpha$$

$$\frac{\partial H}{\partial Q} = 0 = -\dot{P} \Rightarrow P = \text{const} = \frac{E}{\omega}$$

Having solved $\{Q, P\} \Rightarrow \{q, p\}$

$$q = \sqrt{\frac{2P}{m\omega}} \sin Q = \sqrt{\frac{2E}{m\omega^2}} \sin(\omega t + \alpha) \quad \leftarrow \text{out of phase}$$

$$p = \sqrt{2m\omega P} \cos Q = \sqrt{2mE} \cos(\omega t + \alpha)$$

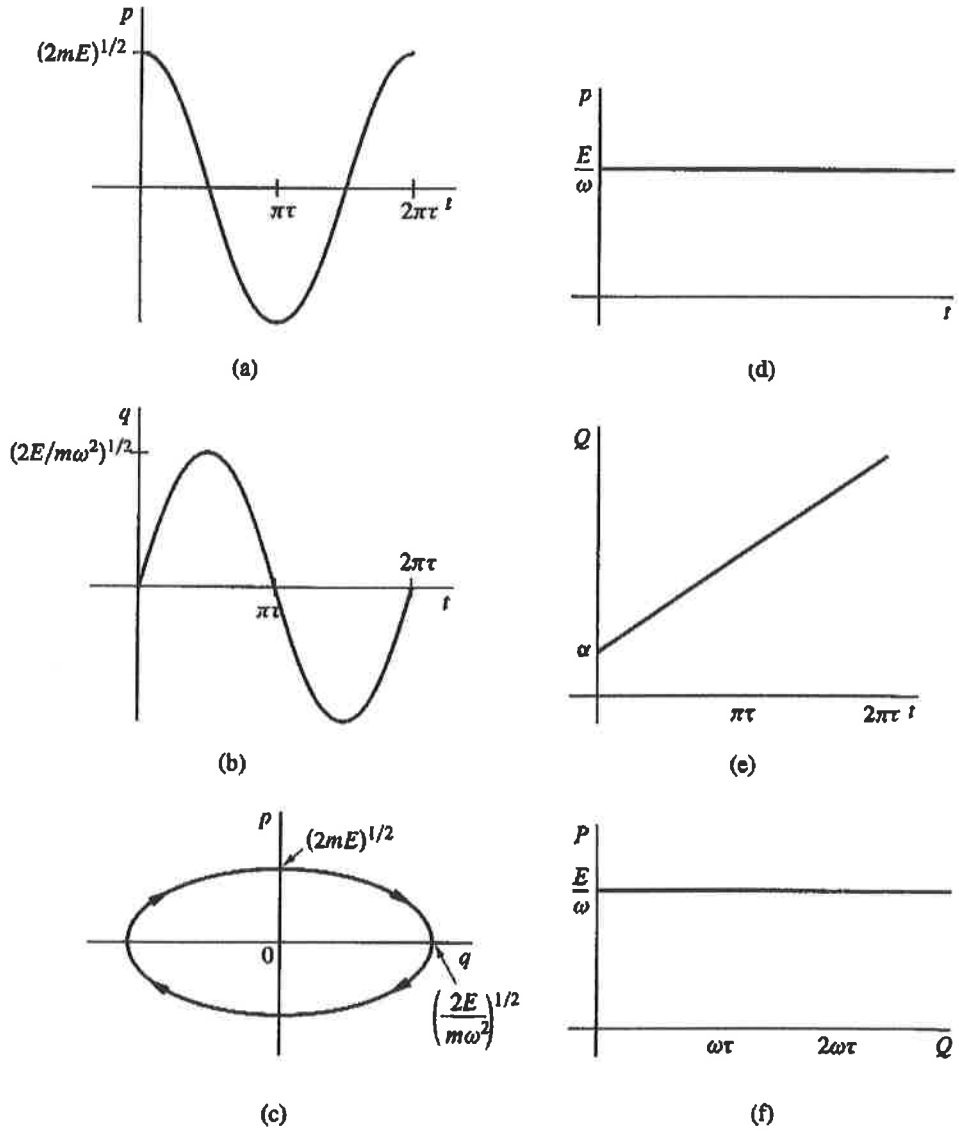


FIGURE 9.1 The harmonic oscillator in two canonical coordinate systems. Drawings (a)–(c) show the q, p system and (d)–(f) show the P, Q system.

When we invoke quantum mechanics, we write $E = \hbar\omega$, where $\hbar = h/2\pi$, and h is Planck's constant. The coordinate and momentum q and p can be normalized as

$$q' = \sqrt{\frac{m\omega^2}{2E}} q \quad \text{and} \quad p' = \frac{p}{\sqrt{2mE}}.$$

to make the phase space plot of p' versus q' a circle of area π . This normalized form will be useful in Section 11.1 on chaos.