

Phys 3344: Tuesday 13 Oct

Office Hours: Wed 5:00-6:00

Exam #2: this week Ch. 6-9 + Ch5. Fourier Transforms

Ch 9

Eigenvalues

Homework #8:

2020 FALL PHYS 3344					
#	DAY	LECTURE:	NOTES:	Chpt	TOPIC
1	TUE	08/25/20	First Class	1	Newton's Laws
2	THUR	08/27/20		2	Projectiles
3	TUE	09/01/20		3	Momentum & Angular Momentum
4	THUR	09/03/20		4	Energy
5	TUE	09/08/20		5	Oscillations
6	THUR	09/10/20			
7	TUE	09/15/20			
8	THUR	09/17/20		EXAM 1	
9	TUE	09/22/20		6	Calculus of Variations
10	THUR	09/24/20		7	Lagrange's Equation
11	TUE	09/29/20			
12	THUR	10/01/20		8	Two Body Problems
13	TUE	10/06/20			
14	THUR	10/08/20		9	Non-Inertial Frames
	TUE	10/13/20	Fall Break	10	Rotational Motion
15	THUR	10/15/20		EXAM 2	
16	TUE	10/20/20		10	Rotational Motion
17	THUR	10/22/20			
18	TUE	10/27/20		11	Coupled Oscillations
19	THUR	10/29/20			
20	TUE	11/03/20		13	Hamiltonian Mechanics
21	THUR	11/05/20	Drop Date		
22	TUE	11/10/20			
23	THUR	11/12/20		EXAM 3	
24	TUE	11/17/20		14	Collision Theory
25	THUR	11/19/20			
26	TUE	11/24/20		15	Special relativity
27	THUR	11/26/20	Thanksgiving	No Class	
28	TUE	12/01/20		No Class	
29	THUR	12/03/20	Last Class	Review	
	WED	Dec 16	FINAL EXAM	Wednesday Dec. 16, 2020, 11:30am - 2:30	
<i>Adjustments may be made depending on student interests/needs and unplanned events</i>					

$$S[q, q'] = \int L[q, q'] dq$$

$$L = T - U$$

$$\delta S[q, q'] = 0$$

$$\mathbf{F}_{\text{cor}} = 2m\dot{\mathbf{r}} \times \boldsymbol{\Omega} \quad \text{and} \quad \mathbf{F}_{\text{cf}} = m(\boldsymbol{\Omega} \times \mathbf{r}) \times \boldsymbol{\Omega}. \quad [\text{Eqs. (9.35) \& (9.36)}]$$

$$\frac{d}{dt} \frac{dL}{dq'} - \frac{dL}{dq} = \lambda \frac{df}{dq}$$

- the Coriolis force

$$\mathbf{F}_{\text{Coriolis}} = -2m\boldsymbol{\Omega} \times \mathbf{v}_r$$

- the centrifugal force

$$\mathbf{F}_{\text{centrifugal}} = -m\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r})$$

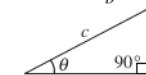
- and the Euler force

$$\mathbf{F}_{\text{Euler}} = -m \frac{d\boldsymbol{\Omega}}{dt} \times \mathbf{r}$$

f is constraint equation

$$\mu = \frac{m_1 m_2}{m_1 + m_2} \quad R_{CM} = \frac{r_1 m_1 + r_2 m_2}{m_1 + m_2}$$

$$r(\phi) = \frac{c}{1 + \epsilon \cos \phi} \quad \sin^2 \theta + \cos^2 \theta = 1$$

MECHANICS		ELECTRICITY	
$v_x = v_{x0} + a_x t$	a = acceleration	$ F_E = k \frac{ q_1 q_2 }{r^2}$	A = area
$x = x_0 + v_{x0} t + \frac{1}{2} a_x t^2$	A = amplitude	F = force	I = current
$v_x^2 = v_{x0}^2 + 2 a_x (x - x_0)$	d = distance	ℓ = length	P = power
$\ddot{a} = \frac{\sum \vec{F}}{m} = \frac{\vec{F}_{\text{net}}}{m}$	E = energy	q = charge	R = resistance
$ \vec{F}_f \leq \mu \vec{F}_n $	F = force	r = separation	t = time
$a_c = \frac{v^2}{r}$	I = rotational inertia	V = electric potential	ρ = resistivity
$\vec{p} = m \vec{v}$	K = kinetic energy		
$\Delta \vec{p} = \vec{F} \Delta t$	k = spring constant		
$K = \frac{1}{2} m v^2$	L = angular momentum		
$\Delta E = W = F_{ } d = F d \cos \theta$	ℓ = length		
$P = \frac{\Delta E}{\Delta t}$	m = mass		
$\theta = \theta_0 + \omega_0 t + \frac{1}{2} \alpha t^2$	P = power		
$\omega = \omega_0 + \alpha t$	p = momentum		
$x = A \cos(2\pi f t)$	r = radius or separation		
$\ddot{a} = \frac{\sum \vec{F}}{I} = \frac{\vec{F}_{\text{net}}}{I}$	T = period		
$\tau = r \perp F = r F \sin \theta$	t = time		
$L = I \omega$	U = potential energy		
$\Delta L = \tau \Delta t$	V = volume		
$K = \frac{1}{2} I \omega^2$	v = speed		
$ \vec{F}_s = k \vec{x} $	W = work done on a system		
$U_s = \frac{1}{2} k x^2$	x = position		
$\rho = \frac{m}{V}$	y = height		
	α = angular acceleration		
	μ = coefficient of friction		
	θ = angle		
	ρ = density		
	τ = torque		
	ω = angular speed		
	$\Delta U_g = mg \Delta y$		
	$T_s = 2\pi \sqrt{\frac{m}{k}}$		
	$T_p = 2\pi \sqrt{\frac{\ell}{g}}$		
	$ \vec{F}_g = G \frac{m_1 m_2}{r^2}$		
	$\vec{g} = \frac{\vec{F}_g}{m}$		
	$U_G = -\frac{G m_1 m_2}{r}$		
WAVES		GEOMETRY AND TRIGONOMETRY	
		f = frequency	
		$\lambda = \frac{v}{f}$	v = speed
			λ = wavelength
GEOMETRY AND TRIGONOMETRY			
		Rectangle	A = area
		$A = bh$	C = circumference
		Triangle	V = volume
		$A = \frac{1}{2}bh$	S = surface area
		Circle	b = base
			h = height
			ℓ = length
			w = width
			r = radius
		Rectangular solid	$c^2 = a^2 + b^2$
		$V = \ell wh$	
		Cylinder	$\sin \theta = \frac{a}{c}$
		$V = \pi r^2 \ell$	$\cos \theta = \frac{b}{c}$
		$S = 2\pi r \ell + 2\pi r^2$	
		Sphere	$\tan \theta = \frac{a}{b}$
		$V = \frac{4}{3} \pi r^3$	
		$S = 4\pi r^2$	
			

CHAPTER 6 Calculus of Variations 215

-
- 6.1 Two Examples 216
 - 6.2 The Euler-Lagrange Equation 218
 - 6.3 Applications of the Euler-Lagrange Equation 221
 - 6.4 More than Two Variables 226
 - Principal Definitions and Equations of Chapter 6 230
 - Problems for Chapter 6 230

CHAPTER 7 Lagrange's Equations 237

-
- 7.1 Lagrange's Equations for Unconstrained Motion 238
 - 7.2 Constrained Systems; an Example 245
 - 7.3 Constrained Systems in General 247
 - 7.4 Proof of Lagrange's Equations with Constraints 250
 - 7.5 Examples of Lagrange's Equations 254
 - 7.6 Generalized Momenta and Ignorable Coordinates 266
 - 7.7 Conclusion 267
 - 7.8 More about Conservation Laws* 268
 - 7.9 Lagrange's Equations for Magnetic Forces* 272
 - 7.10 Lagrange Multipliers and Constraint Forces* 275
 - Principal Definitions and Equations of Chapter 7 280
 - Problems for Chapter 7 281

CHAPTER 8 Two-Body Central-Force Problems 293

-
- 8.1 The Problem 293
 - 8.2 CM and Relative Coordinates; Reduced Mass 295
 - 8.3 The Equations of Motion 297
 - 8.4 The Equivalent One-Dimensional Problem 300
 - 8.5 The Equation of the Orbit 305
 - 8.6 The Kepler Orbits 308
 - 8.7 The Unbounded Kepler Orbits 313
 - 8.8 Changes of Orbit 315
 - Principal Definitions and Equations of Chapter 8 319
 - Problems for Chapter 8 320

CHAPTER 9 Mechanics in Noninertial Frames 327

-
- 9.1 Acceleration without Rotation 327
 - 9.2 The Tides 330
 - 9.3 The Angular Velocity Vector 336
 - 9.4 Time Derivatives in a Rotating Frame 339

-
- 9.5 Newton's Second Law in a Rotating Frame 342
 - 9.6 The Centrifugal Force 344
 - 9.7 The Coriolis Force 348
 - 9.8 Free Fall and the Coriolis Force 351
 - 9.9 The Foucault Pendulum 354
 - 9.10 Coriolis Force and Coriolis Acceleration 358
 - Principal Definitions and Equations of Chapter 9 359
 - Problems for Chapter 9 360

CHAPTER 10 Rotational Motion of Rigid Bodies 367

-
- 10.1 Properties of the Center of Mass 367
 - 10.2 Rotation about a Fixed Axis 372
 - 10.3 Rotation about Any Axis; the Inertia Tensor 378
 - 10.4 Principal Axes of Inertia 387
 - 10.5 Finding the Principal Axes; Eigenvalue Equations 389
 - 10.6 Precession of a Top due to a Weak Torque 392
 - 10.7 Euler's Equations 394
 - 10.8 Euler's Equations with Zero Torque 397
 - 10.9 Euler Angles* 401
 - 10.10 Motion of a Spinning Top* 403
 - Principal Definitions and Equations of Chapter 10 407
 - Problems for Chapter 10 408

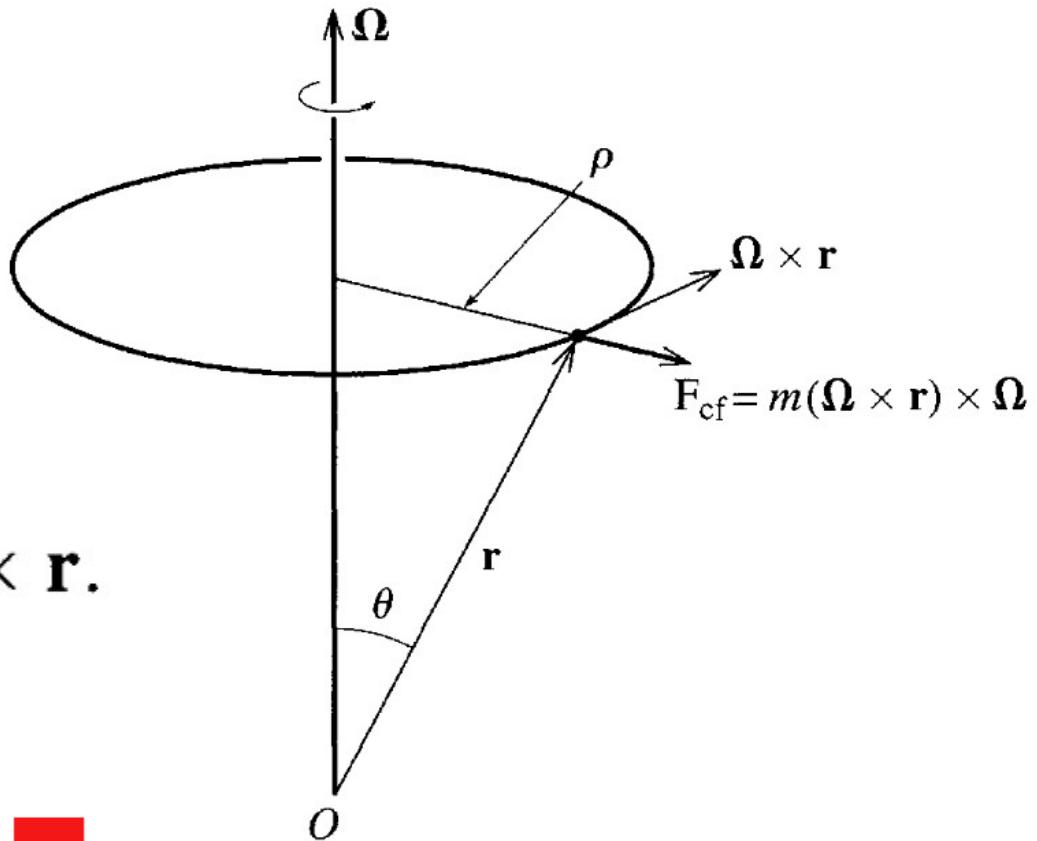
CHAPTER 11 Coupled Oscillators and Normal Modes 417

-
- 11.1 Two Masses and Three Springs 417
 - 11.2 Identical Springs and Equal Masses 421
 - 11.3 Two Weakly Coupled Oscillators 426
 - 11.4 Lagrangian Approach: The Double Pendulum 430
 - 11.5 The General Case 436
 - 11.6 Three Coupled Pendulums 441
 - 11.7 Normal Coordinates* 444
 - Principal Definitions and Equations of Chapter 11 447
 - Problems for Chapter 11 448

PART II Further Topics 455**CHAPTER 12** Nonlinear Mechanics and Chaos 457

-
- 12.1 Linearity and Nonlinearity 458
 - 12.2 The Driven Damped Pendulum DDP 462
 - 12.3 Some Expected Features of the DDP 463

$$\left(\frac{d\mathbf{Q}}{dt}\right)_{S_0} = \left(\frac{d\mathbf{Q}}{dt}\right)_S + \boldsymbol{\Omega} \times \mathbf{Q}.$$



$$\left(\frac{d\mathbf{r}}{dt}\right)_{S_0} = \left(\frac{d\mathbf{r}}{dt}\right)_S + \boldsymbol{\Omega} \times \mathbf{r}.$$

$$\left(\frac{d\textcolor{red}{\square}}{dt}\right)_{S_0} = \left(\frac{d\textcolor{red}{\square}}{dt}\right)_S + \boldsymbol{\Omega} \times \textcolor{red}{\square}$$

The vector $\boldsymbol{\Omega} \times \mathbf{r}$ is the velocity of an object as it is dragged eastward with speed $\Omega\rho$ by the earth's rotation. Therefore, the centrifugal force, $m(\boldsymbol{\Omega} \times \mathbf{r}) \times \boldsymbol{\Omega}$, points radially outward from the axis and has magnitude $m\Omega^2\rho$.

Relation between velocities in the two frames [edit]

A velocity of an object is the time-derivative of the object's position, or

$$\mathbf{v} \stackrel{\text{def}}{=} \frac{d\mathbf{r}}{dt}$$

The time derivative of a position $\mathbf{r}(t)$ in a rotating reference frame has two components, one from the explicit time dependence due to motion of the particle itself, and another from the frame's own rotation. Applying the result of the previous subsection to the displacement $\mathbf{r}(t)$, the [velocities](#) in the two reference frames are related by the equation

$$\mathbf{v}_i \stackrel{\text{def}}{=} \frac{d\mathbf{r}}{dt} = \left(\frac{d\mathbf{r}}{dt} \right)_r + \boldsymbol{\Omega} \times \mathbf{r} = \mathbf{v}_r + \boldsymbol{\Omega} \times \mathbf{r},$$

where subscript i means the inertial frame of reference, and r means the rotating frame of reference.

Relation between accelerations in the two frames [edit]

Acceleration is the second time derivative of position, or the first time derivative of velocity

$$\mathbf{a}_i \stackrel{\text{def}}{=} \left(\frac{d^2\mathbf{r}}{dt^2} \right)_i = \left(\frac{d\mathbf{v}}{dt} \right)_i = \left[\left(\frac{d}{dt} \right)_r + \boldsymbol{\Omega} \times \right] \left[\left(\frac{d\mathbf{r}}{dt} \right)_r + \boldsymbol{\Omega} \times \mathbf{r} \right],$$

where subscript i means the inertial frame of reference. Carrying out the [differentiations](#) and re-arranging some terms yields the acceleration *relative to the rotating* reference frame, \mathbf{a}_r

$$\mathbf{a}_r = \mathbf{a}_i - 2\boldsymbol{\Omega} \times \mathbf{v}_r - \boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r}) - \frac{d\boldsymbol{\Omega}}{dt} \times \mathbf{r}$$

where $\mathbf{a}_r \stackrel{\text{def}}{=} \left(\frac{d^2\mathbf{r}}{dt^2} \right)_r$ is the apparent acceleration in the rotating reference frame, the term $-\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r})$ represents [centrifugal acceleration](#), and

the term $-2\boldsymbol{\Omega} \times \mathbf{v}_r$ is the [Coriolis acceleration](#). The last term $(-\frac{d\boldsymbol{\Omega}}{dt} \times \mathbf{r})$ is the [Euler acceleration](#) and is zero in uniformly rotating frames.

Newton's second law in the two frames [edit]

When the expression for acceleration is multiplied by the mass of the particle, the three extra terms on the right-hand side result in fictitious forces in the rotating reference frame, that is, apparent forces that result from being in a [non-inertial reference frame](#), rather than from any physical interaction between bodies.

Using [Newton's second law of motion](#) $\mathbf{F} = m\mathbf{a}$, we obtain:^{[1][2][3][5][6]}

- the [Coriolis force](#)

$$\mathbf{F}_{\text{Coriolis}} = -2m\boldsymbol{\Omega} \times \mathbf{v}_r$$

- the [centrifugal force](#)

$$\mathbf{F}_{\text{centrifugal}} = -m\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r})$$

- and the [Euler force](#)

$$\mathbf{F}_{\text{Euler}} = -m \frac{d\boldsymbol{\Omega}}{dt} \times \mathbf{r}$$

where m is the mass of the object being acted upon by these fictitious forces. Notice that all three forces vanish when the frame is not rotating, that is, when $\boldsymbol{\Omega} = 0$.

For completeness, the inertial acceleration \mathbf{a}_i due to impressed external forces \mathbf{F}_{imp} can be determined from the total physical force in the inertial (non-rotating) frame (for example, force from physical interactions such as [electromagnetic forces](#)) using [Newton's second law](#) in the inertial frame:

$$\mathbf{F}_{\text{imp}} = m\mathbf{a}_i$$

Newton's law in the rotating frame then becomes

$$\mathbf{F}_r = \mathbf{F}_{\text{imp}} + \mathbf{F}_{\text{centrifugal}} + \mathbf{F}_{\text{Coriolis}} + \mathbf{F}_{\text{Euler}} = m\mathbf{a}_r .$$

In other words, to handle the laws of motion in a rotating reference frame:^{[6][7][8]}

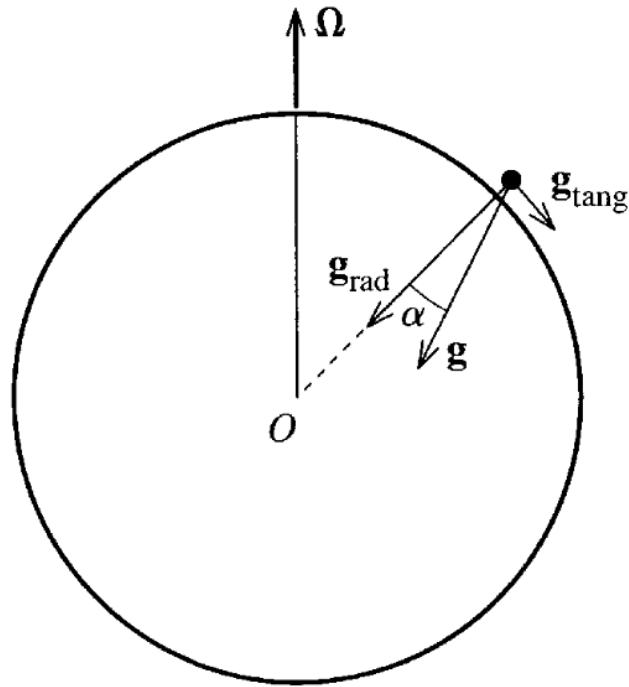


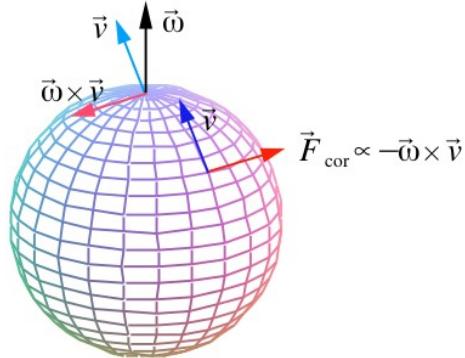
Figure 9.11 Because of the centrifugal force, the free-fall acceleration \mathbf{g} has a nonzero tangential component (greatly exaggerated here) and \mathbf{g} deviates from the radial direction by the small angle α .

Coriolis: about 3:00min

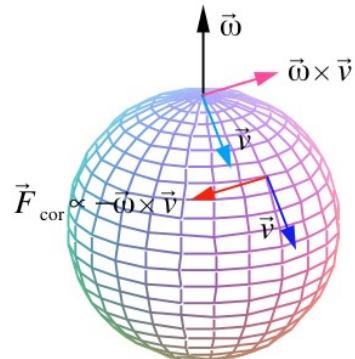
<https://youtu.be/6L5UD240mCQ>

<https://youtu.be/f8lwL2ZtDTc>

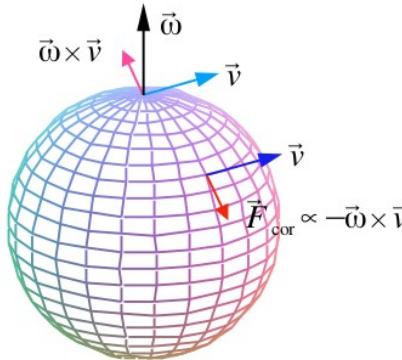
Body moving north (v_θ negative, other components zero); coriolis force eastward:



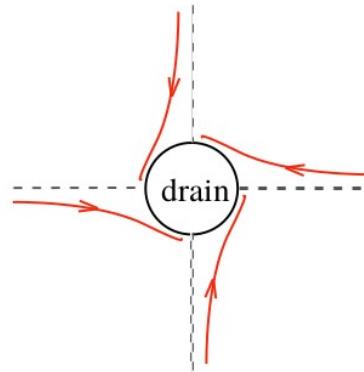
Body moving south (v_θ positive, other components zero); coriolis force westward:



Body moving east (v_ϕ positive, other components zero); coriolis force has component in southward direction:

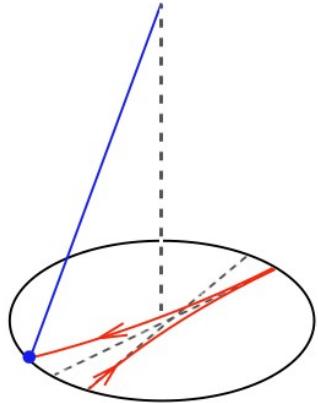


We see that for any body moving on the surface of the earth in the northern hemisphere, the coriolis force deflects it to the *right*. This is responsible for the counterclockwise rotation of the bath water as it drains from your tub, as viewed from above:



The same applies to the direction of air flow around an area of low atmospheric pressure.

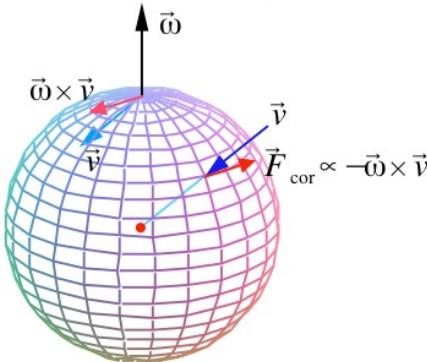
Another case where the coriolis force is important is the *Foucault pendulum* (pronounced “Foo-ko”). This is a very large pendulum which you sometimes see in the lobbies of big important buildings. The plane in which the pendulum swings back and forth *precesses*, or turns, slowly around in a clockwise direction as viewed from above. The following diagram shows why:



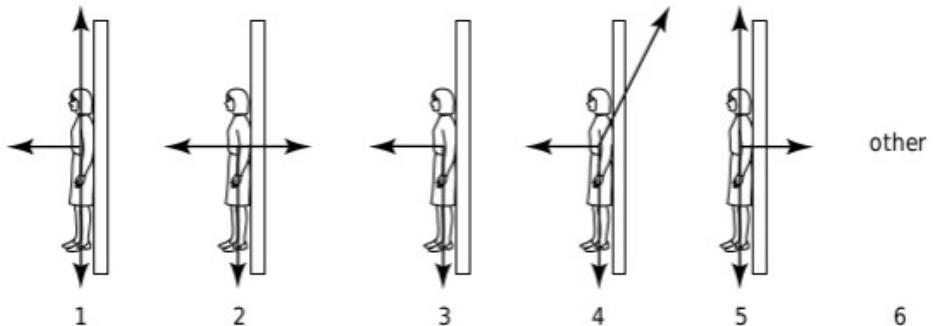
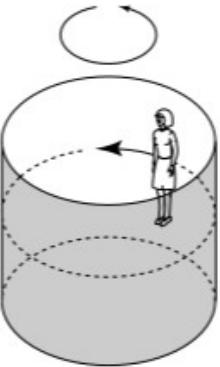
As the pendulum swings across, the coriolis force pushes it to the right. On the way back it is also pushed to the right, and this just rotates the plane of the pendulum further in the clockwise sense. The Foucault pendulum is a rather striking demonstration of the rotation of the earth. Here is a [movie](#) showing the Foucault pendulum from above.

Last but not least, there is also an eastward force on a body falling vertically:

Body falling vertically (v_r negative, other components zero); coriolis force eastward:

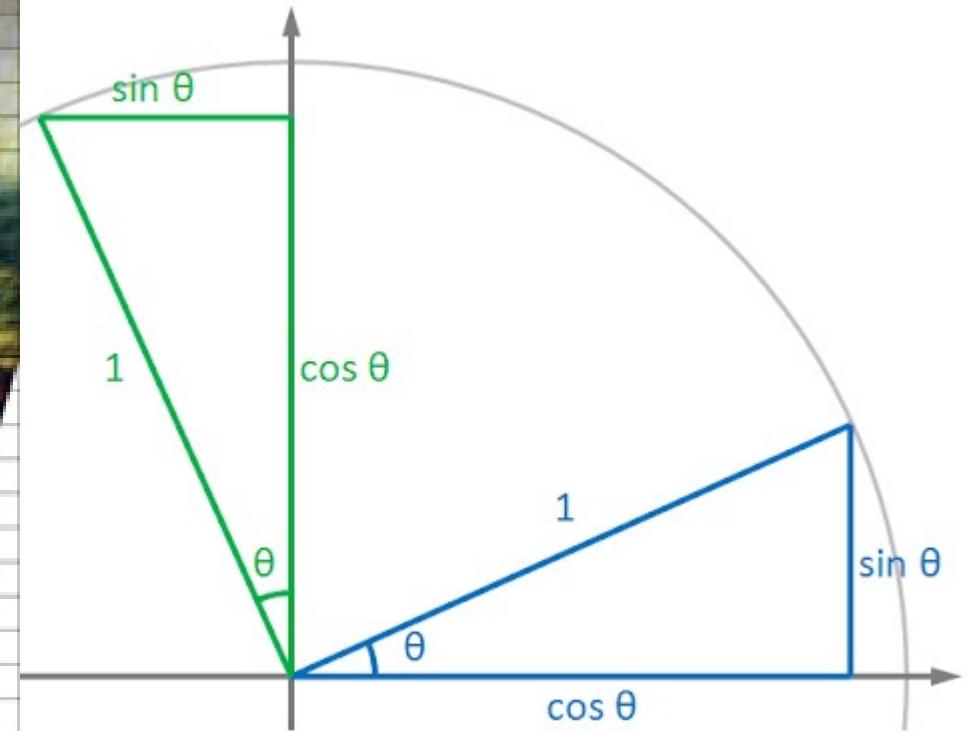
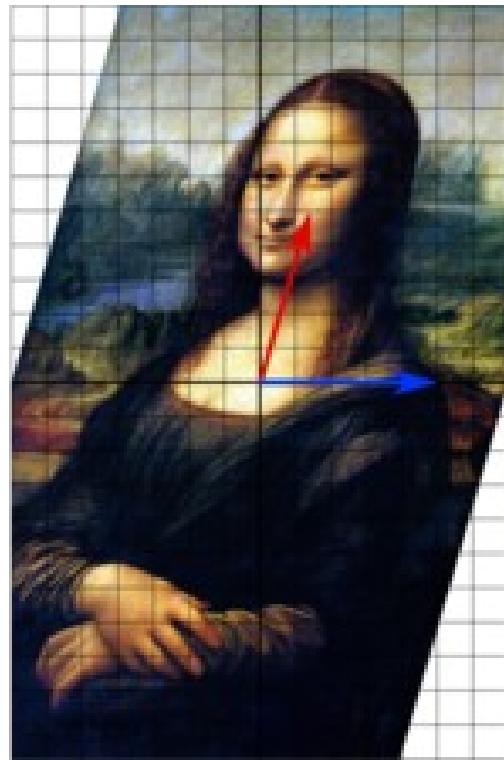
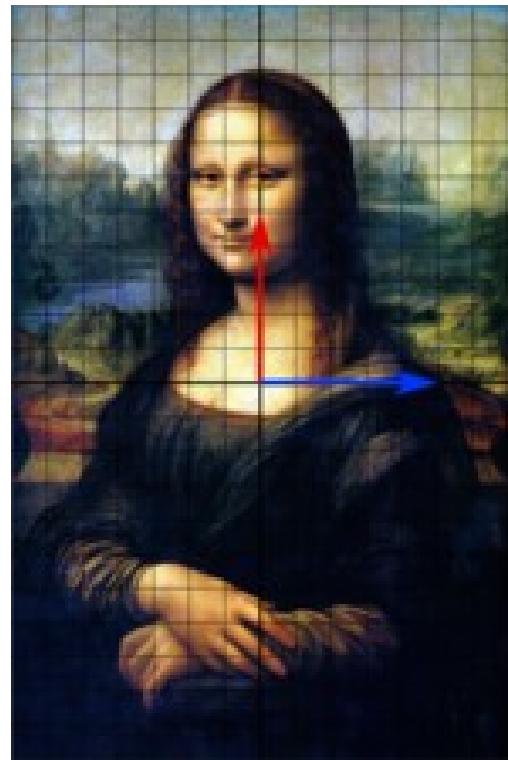


A rider in a “barrel of fun” finds herself stuck with her back to the wall. Which diagram correctly shows the forces acting on her?

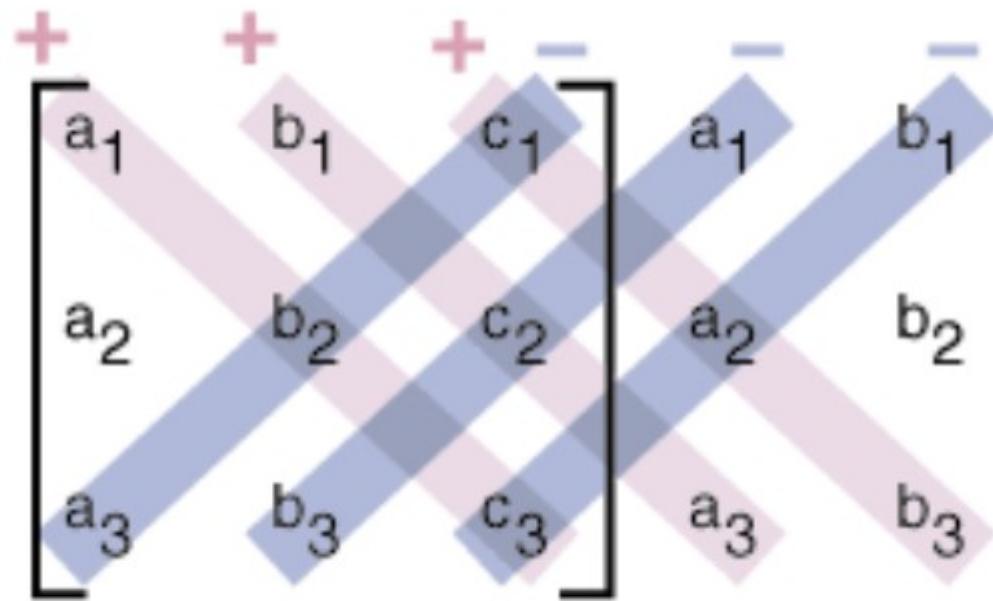


Consider two people on opposite sides of a rotating merry-go-round. One of them throws a ball toward the other. In which frame of reference is the path of the ball straight when viewed from above: (a) the frame of the merry-go-round or (b) that of Earth?

1. (a) only
2. (a) and (b)—although the paths appear to curve
3. (b) only
4. neither; because it's thrown while in circular motion, the ball travels along a curved path.

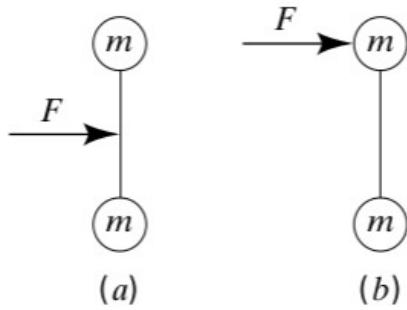


$$\det \begin{vmatrix} a & b \\ c & d \end{vmatrix} = ad - bc$$



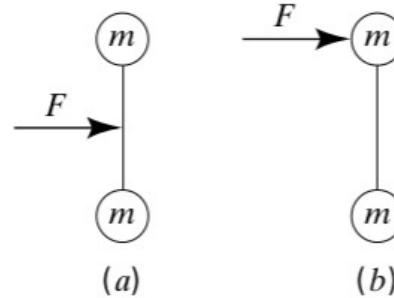
$$\det A = (a_1 b_2 c_3 + b_1 c_2 a_3 + c_1 a_2 b_3) - (a_3 b_2 c_1 + b_3 c_2 a_1 + c_3 a_2 b_1)$$

A force F is applied to a dumbbell for a time interval Δt , first as in (a) and then as in (b). In which case does the dumbbell acquire the greater center-of-mass speed?



1. (a)
2. (b)
3. no difference
4. The answer depends on the rotational inertia of the dumbbell.

A force F is applied to a dumbbell for a time interval Δt , first as in (a) and then as in (b). In which case does the dumbbell acquire the greater energy?



1. (a)
2. (b)
3. no difference
4. The answer depends on the rotational inertia of the dumbbell.