The Higgs: Above and Beyond

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Bag Ground Tra

Programme

- The Higgs what do we think we know?
- Above the Higgs what do we want to know?
- Beyond the Higgs what else should we know?

What We Think We Know

Photo credit: Horia Varlan

The Particle Checklist

- Spin – – – – – – – .



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"Signal Strength"







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ATLAS-CONF-2014-009



Couplings



K. Einsweiler, Moriond EW 2014

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Select events in a highly signal-like region...



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Phys. Lett. B 726 (2013), pp. 120-144

Construct a discriminant that separates signal from background, but also different signal hypotheses...



Each plot contains multiple Parity hypotheses, where appropriate.

Construct a discriminant that separates signal from background, but also different signal hypotheses...



Construct a discriminant that separates signal from background, but also different signal hypotheses...





A simple decision tree (DT). ATLAS used a Boosted DT ("BDT") for the spin-parity analysis. "Boosting" involves increasing the weight of mis-classified events and improving their classification. Stephen J. Sekula - SMU

Learn More About Statistical Data Analysis

- Narsky, I and Porter, F. "Statistical Analysis Techniques in Particle Physics: Fits, Density Estimation and Supervised Learning". Wiley-VCH; 1 edition (December 23, 2013).
- Bevan, A. "Statistical Data Analysis for the Physical Sciences". Cambridge University Press (June 28, 2013).

Construct a discriminant that separates signal from background, but also different signal hypotheses...



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Quantum Numbers

pp. 120-144



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Above the Higgs: What Do We Want To Know?

Photo credit: Tilemahos Efthimiadis

Wish List (a multi-decade program)

- The width of the Higgs boson
 - predicted in the Standard Model (~4 MeV)
- The decay modes and spin-parity of the Higgs
 - "Complete" the picture of couplings to bosons with $H \to Z \gamma$
 - Definitively measure bb and $\tau\tau$
 - Push the couplings and spin-parity uncertainties down
- Probe the Higgs Field directly
 - Try to get the Higgs self-coupling by looking at Higgs pair production
- Determine if this Higgs is "alone"
 - Search for additional Higgs Bosons in nature

Beyond the Higgs: Supersymmetry-Inspired Searches

in Preparation

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Higgs properties	Interpretation of ATLAS Higgs Measurements in Extensions of the Standard Model		Internal Review		

HSG6 Papers and Publications

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List continues off-screen...

HSG6 Webpage, Oct. 2013

HSG6 Data Notes:

- In preparation:
 - H+-> csbar CDS-INT CDS
 - Ratio method CDS-INT CDS SVN-INT SVN-CONF TWW
 - H+-> WZ CDS-INT CDS
 - a H+ -> a1W CDS-INT CDS
- 5/b results:
 - Title: Search for charged Higgs bosons decaying via H+ -> taunu in ttbar events using pp collision data at \sqrt(s) = 7 TeV with the ATLAS detector
 - Plots JHEP06(2012)039 arXiv:1204.2760 CERN-PH-EP-2012-083 ATLAS-CONF-2012-011 TWW
 - Supporting documents (internal notes):
 - Limit & combination supporting note: ATL-INT-PHYS-2012-044
 - Imp+jets supporting note: ATL-INT-PHYS-2012-045
 - tau+lep supporting note: ATL-INT-PHYS-2012-046
 - tau+jets supporting note: ATL-INT-PHYS-2012-047
- 1/b results:
 - Search for a charged Higgs boson decaying via H+ to tau(lep)+nu in ttbar events with one or two light leptons in the final state using 1.03/fb of pp collision data recorded at sqrt(s) = 7 TeV with the ATLAS detector. CDS-INT CDS-CONF TW/ki Plots
 - Search for Charged Higgs Bosons in the r*jets Final State in tt" Decays with 1.83 fb-1 of pp Collision Data Recorded at sqrt(s)=7 TeV with the ATLAS Experiment CDS-INT CDS-CONF TWiki Plots
 - A Search for a light charged Higgs boson decaying to cs' in pp collisions at vis = 7 TeV with the ATLAS detector CDS-INT CDS-CONF Plots
- 35/pb results;
 - Study of discriminative variables for charged Higgs boson searches in ttbar events with 35pb-1 of data from the ATLAS detector. CDS-INT CDS-CONF
 - Bata-driven estimation of the background to H+ searches with hadronic-tau final states. CDS-INT CDS-CONF TWIKE

HSG6 Simulation Notes in CDS

- Embedding Technique for the Itbar Background Estimation in Charged Higgs Boson Searches
- Data-driven measurement of the fake tau contribution from light jets and application for the tibar background estimation in charged Higgs Searches
- Light Charged Higgs Boson Searches for H+ to tau+ nu and H+ to c shar in Early LHC Data at the ATLAS Experiment
- An ATLAS Search for the NMSSM Ideal Higgs Scenario
- Update on Light Charged Higgs Boson Searches with Early LHC Data at 7 TeV in the ATLAS Experiment

Other Public Results

- Expected Sensitivity in Light Charged Higgs Boson Searches for H+ to tau+nu and H+ to c+sbar...
- Public Plots from ATL-PHYS-PUB-2010-006
- Computing and Software Commissioning (CSC) Book in CDS

HSG6 Webpage, ~Oct. 1, 2012

NaN

(not quite fair – BSM Higgs analysis efforts were distributed among multiple parent HSGx groups. Nonetheless, 2011-2013 was quite a fast ride!)

HSG6 TWiki, ~Oct. 1, 2011

MSSM: Important Features

$$h^0, \ H^0, \ A^0, \ H^{\pm}$$

$$M_{H^{\pm}}^{2} = M_{A}^{2} + M_{W^{\pm}}^{2}$$

$$M_A$$
, tan(β), X_t ,
 M_2 , μ , M_{SUSY}

Coupling Mixing Angle Mass Dependence Dependence Huu $sin(\alpha)/sin(\beta)$ m $\cos(\alpha)/\cos(\beta)$ Hdd m_d Auu $\cot(\beta)$ m Add $tan(\beta)$ m $m_d \tan\beta (1 + \gamma_5) + m_u \cot\beta (1 - \gamma_5)$ H[±]ud

Five physical Higgs bosons (2 CP-even, one CP-odd, and 2 electrically charged)

Tree-level mass relationship

Free parameters

Embedded in the MSSM is a Type-II Two-Higgs Doublet Model (2HDM)

c.f. Physics Reports, Volume 459, Issues 1-6, 2008, Pages 1-241

2HDMs

	type-I	type-II
ξ_h^u	$\sin(\beta - \alpha) + \cos(\beta - \alpha)/\tan\beta$	$\sin(\beta - \alpha) + \cos(\beta - \alpha)/\tan\beta$
ξ_h^d	$\sin(\beta - \alpha) + \cos(\beta - \alpha)/\tan\beta$	$\sin(\beta - \alpha) - \cos(\beta - \alpha) \cdot \tan\beta$
ξ_h^l	$\sin(\beta - \alpha) + \cos(\beta - \alpha) / \tan\beta$	$\sin(\beta - \alpha) - \cos(\beta - \alpha) \cdot \tan\beta$
ξ_H^u	$\cos(\beta - \alpha) - \sin(\beta - \alpha) / \tan\beta$	$\cos(\beta - \alpha) - \sin(\beta - \alpha) / \tan\beta$
ξ_{H}^{d}	$\cos(\beta - \alpha) - \sin(\beta - \alpha) / \tan\beta$	$\cos(\beta - \alpha) + \sin(\beta - \alpha) \cdot \tan\beta$
ξ_{H}^{l}	$\cos(\beta - \alpha) - \sin(\beta - \alpha) / \tan\beta$	$\cos(\beta - \alpha) + \sin(\beta - \alpha) \cdot \tan\beta$
ξ^u_A	$1/\tan\beta$	$1/\tan\beta$
ξ^d_A	$-1/\tan\beta$	$\tan\beta$
$\xi_A^{\hat{l}}$	$-1/\tan\beta$	$\tan\beta$

Table 1: Yukawa coupling coefficients of the neutral bosons of the type-I and type-II 2HDMs for up-type quarks (*u*), down-type quarks (*d*) and charged leptons (*l*). These coefficients are defined such that the Yukawa Lagrangian terms are $-(m_f/v)\bar{f}f\phi$ and $i(m_f/v)\bar{f}\gamma_5 fA$ where f = u, d, l and $\phi = h, H$.

A Comment on Notation

- h
 - the lightest (lowest-mass) neutral CP-even Higgs in the theory
- H
 - the next-lightest (heavier) neutral CP-even Higgs
- A priori, we don't know which one we've discovered

MSSM: m_h-max scenario

$$\begin{split} m_t &= 174.3 \; \text{GeV}, \quad M_{SUSY} = 1 \; \text{TeV}, \quad \mu = 200 \; \text{GeV}, \quad M_2 = 200 \; \text{GeV}, \\ X_t^{\text{OS}} &= 2 \; M_{SUSY} \; (\text{FD calculation}), \quad X_t^{\overline{\text{MS}}} = \sqrt{6} \; M_{SUSY} \; (\text{RG calculation}) \\ A_b &= A_t, \quad m_{\tilde{g}} = 0.8 \; M_{SUSY} \; . \end{split}$$

Eur.Phys.J.C26:601-607,2003

Designed to maximize the SM-like Higgs (h^{0}) mass ($m_{h} \sim 135$ GeV).

However, we now know the mass of the h^{0} (~126 GeV), so m_{h} -max is a bit too aggressive.

$$\begin{split} m_t &= 173.2 \; {\rm GeV}, \\ M_{\rm SUSY} &= 1000 \; {\rm GeV}, \\ \mu &= 200 \; {\rm GeV}, \\ M_2 &= 200 \; {\rm GeV}, \\ X_t^{\rm OS} &= 2 \; M_{\rm SUSY} \; ({\rm FD \; calculation}), \\ X_t^{\rm \overline{MS}} &= \sqrt{6} \; M_{\rm SUSY} \; ({\rm RG \; calculation}), \\ A_b &= A_\tau = A_t, \\ m_{\tilde{g}} &= 1500 \; {\rm GeV}, \\ M_{\tilde{l}_3} &= 1000 \; {\rm GeV} \; . \end{split}$$

arXiv:1307.1347

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Comments

- Many results I show today employ m_h-max
 - this is largely for backward-compatibility of interpretations over the last 1-2 years.
 - LHC Higgs Cross-Section Working Group has looked at the effect on exclusions in tan β vs. m_{H+} using m_h-max or m_h-mod+
 - no significant differences seen for channels like A/H $\rightarrow \tau \tau$ or H⁺ $\rightarrow \tau \nu$
 - in those channels, exclusions in m_h -max are basically transferable to m_h -mod+

Above and Beyond the Higgs: The Search for H[±]

H⁺ Production

H⁺ Production

Dominant for masses below $m_t - m_b \approx 169$ GeV

Dominant for larger masses

arXiv:1307.1347

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arXiv:1307.1347

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 $H^+ \rightarrow \tau^+ \nu (m_{H^+} < m_t - m_b)$

Different experimental approaches are

 H^+

ŧ

W

- hadronic tau decays more prone to QCD
- leptonic tau/W decays yield more significant MET and present reconstruction/resolution challenges

W

 \rightarrow hadrons

 $\tau + lep$

channel

 ν_{τ}

 $\bar{\nu}_{\ell}$

hadronic tau lepton $\tau + je$ (looks like a narrow jet with 1 or 3 charged particles) (BR = 65%). Missing Transver

 $p_{T} > 40 \text{ GeV}$

τ + jets channel ^{ith}

Missing Transverse Energy (MET) from neutrinos (from the H⁺ and tau decays. MET > 65 GeV (light search) or > 80 GeV (heavy search)

Jet (p_T > 25 GeV), consistent with b-quark fragmentation (e.g. displaced secondary vertex, lepton activity in the jet, etc.)

A pair of jets $(p_{T} > 25 \text{ GeV})$ from W-boson decay (BR = 67.6%)

Cartoon of production and decay in the plane transverse to the proton beams... hadronic tau lepton (looks like a narrow jet with 1 or 3 charged particles) (BR = 65%). Missing Tr

(BR = 65%). $p_{\tau} > 40 \text{ GeV}$

τ + jets channel

Missing Transverse Energy (MET) from neutrinos (from the H⁺ and tau decays. MET > 65 GeV (light search) or > 80 GeV (heavy search)

Jet (p_T > 25 GeV), consistent with b-quark fragmentation (e.g. displaced secondary vertex, lepton activity in the jet, etc.)

Major Backgrounds:

- Fake tau leptons (e.g. from jets or electrons mis-identified as tau leptons)
- Real tau leptons (e.g. from W+jets or tt events with taus in the final state)

A pair of jets $(p_{T} > 25 \text{ GeV})$ from W-boson decay (BR = 67.6%)

Cartoon of production and decay in the plane transverse to the proton beams...

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Trigger Efficiency Corrections

Technique for this analysis was developed by A. Randle-Conde and S.S. Hadronic top from t → Wb → (τ ν) b

"Tag and probe" approach:

- Trigger on the muon ("tag")
- Reconstruct the event using signal analysis selection
- Apply tau+MET trigger at the end ("probe") and compare efficiency in data and MC simulation

Muon from t \rightarrow Wb \rightarrow ($\mu \nu$) b

b-tagged jet

MET

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Trigger Systematics

	$p_{\rm T}^{\tau} = (40, 70) {\rm GeV}$	$p_{\rm T}^{\rm T} = (70, 500) {\rm GeV}$
$E_{\rm T}^{\rm miss} = (65, 80) {\rm GeV}$	$0.94 \pm 0.16 \pm 0.18$	$0.89 \pm 0.14 \pm 0.26$
$E_{\rm T}^{\rm miss} = (80, 100) {\rm GeV}$	$1 \pm 0.1 \pm 0.4$	$1.1 \pm 0.15 \pm 0.24$
$E_{\rm T}^{\rm miss} = (100, 500) {\rm GeV}$	$0.93 \pm 0.09 \pm 0.18$	$0.94 \pm 0.10 \pm 0.31$

The biggest source of systematic uncertainty is due to the effect of kinematic bias – there is more MET, typically, in the control sample than in real signal events. We've learned to deal with this and for the publication we expect to reduce greatly these systematics.

Jet $\rightarrow \tau$ Fake Rate Systematics

Variation	Shift (±%)	Shift (±%)		
	Light H^+ event selection	Heavy H^+ event selection		
True τ contamination	3	3		
Jet composition	10	10		
Statistical uncertainty on p_m	16	14		
Statistical uncertainty on p_r	7	8		
$\tau_{\rm had-vis}$ e-veto uncertainty	4	4		
$\tau_{had-vis}$ identification uncertainty	9	11		

Simulation Systematics

Variation	Shift up (%)	Shift down (%)	Shift up (%)	Shift down (%)	
	Light H^+ event selection		Light H^+ event selection Heavy H^+ event selection		event selection
b jet (mis-)tag efficiency uncertainty	3.1	-3.4	2.9	-3.2	
Jet energy scale uncertainties	3.7	-4.8	7.1	-6.8	
JVF uncertainty	2.2	-1.9	2.2	-2.1	
$E_{\rm T}^{\rm miss}$ uncertainties	0.4	0.3	-0.6	-0.2	
$\tau_{had-vis}$ e-veto uncertainty	0.02	-0.02	0.01	-0.01	
$\tau_{had-vis}$ energy scale uncertainty	3.6	-3.8	3.6	-3.8	
$\tau_{had-vis}$ identification uncertainty	3.8	-3.8	3.7	-3.7	
Pile-up uncertainties	0.9	-1.5	2.6	-2.1	

Theory Systematics

Source of uncertainty	Normalization uncertainty
Light H^{\pm}	
Generator model $(b\bar{b}W^{\pm}H^{\pm})$	9%
Generator model $(b\bar{b}W^+W^-)$	9%
$t\bar{t}$ cross section uncertainty	10%
Jet production rate (SM and H^{\pm})	11%
Heavy H^{\pm}	
Generator model (H^{\pm})	2 – 9%
Generator model (SM)	8%
$t\bar{t}$ cross section uncertainty	10%
Jet production rate (H^{\pm})	4.4 - 11%
Jet production rate (SM)	11%

 $H^+ \rightarrow \tau^+ \nu (\tau_{had} + jets)$

m_T [GeV] 51

No evidence for a signal, so we set upper limits at the 95% C.L. on . . .

- The branching fraction B(t \rightarrow H⁺b) (assuming H⁺ \rightarrow τ v is 100%) for the low-mass search. Limits range between 0.2-2.1%
- The product $\sigma(pp \rightarrow H^+ t (b)) \times B(H^+ \rightarrow \tau \nu)$ for the high-mass search. Limits range between 0.01-0.9 pb.

ATLAS-CONF-2013-090

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Model-Dependent Results

Current: 2012 Data

Fade in

space for

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Fade in

space for

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Fade in

for the heavy H⁺ - more channels in 8 TeV and a lot more data and creativity at 14 TeV will push this effort along rapidly! Stephen J. Sekula - SMU

Conclusions and Outlook

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The Road Ahead

- Run 2 2015-2017
 - expect to integrate $\sim 100/fb$
- Long-Shutdown 2: ~2018-2019
 - "Phase I upgrade"
 - prepare for running at 2x10³⁴/cm²/s
- Running 2019-2021
- Long-Shutdown 3: ~2022-2023
 - "Phase II upgrade"
 - prepare for running at 5x10³⁴/cm²/s
- Running ~2024-...

- Above the Higgs
 - There is much work for Run 2:
 - improved statistics and upgraded detector will allow us to push the statistical uncertainties down on Higgs couplings and quantum numbers, as well as improve access to the fermion (and other) final states.
 - likely that the revisiting of reconstruction and identification efforts will greatly improve systematic uncertainties
- Beyond the Higgs
 - So far, in Run 1, there are no clear direct signs of additional Higgs bosons
 - The MSSM is all but ruled out if SUSY is real, it's eluded us in Run 1 and if we don't see it in Run 2 well . . . then . . .
 - Let's just be open for surprises . . . !