

The Current Excitement in Neutrino Physics

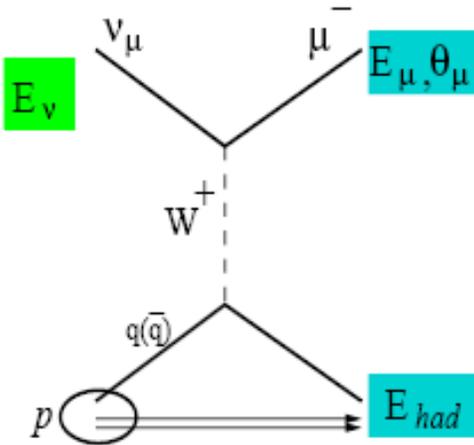
Rob Plunkett

Fermi National Accelerator Laboratory

*Lectures Presented to CTEQ 2007
Part 2*

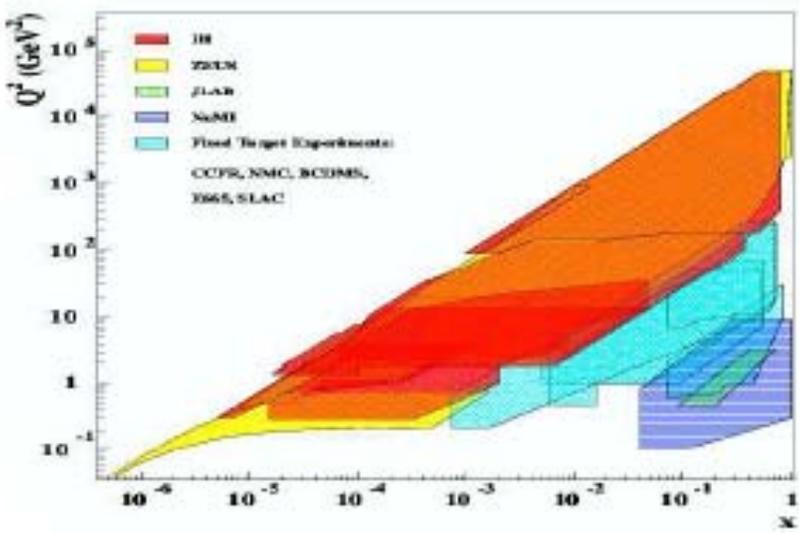
Deep Inelastic Scattering in Neutrino Physics

Classic, mature field
Still a lot to do!



Reconstruct $E_\nu = E_{HAD} + E_\mu$
 Shower energy resolution: $55\%/\sqrt{E}$
 Muon Momentum resolution: 6% range, 13% curvature

$Q^2 = 4E_\nu E_\mu \sin^2 \frac{\theta}{2}$,	Squared four momentum transfer
$x = \frac{Q^2}{2ME_{HAD}}$,	Fractional quark momentum
$y = \frac{E_{HAD}}{E_\nu}$,	Inelasticity
$W^2 = M^2 + 2ME_{HAD} - Q^2$,	Squared final state invariant mass



Coverage of x , Q^2 range for input to CTEQ!

More on Neutrino DIS

$$\frac{d^2\sigma^{\nu(\bar{\nu})}}{dxdy} = \frac{G_F^2}{\pi(1+Q^2/M_W^2)^2} \{Z(M, x, y, E, Q^2)F_2^{\nu(\bar{\nu})} \pm y(1-\frac{y}{2})xF_3^{\nu(\bar{\nu})}\}$$

Derived from basic structure of Lepton-hadron interactions.

F_2, F_3 are related via Quark-Parton model to PDF's

F_3 measures difference in quark-antiquark content of nucleon.

F_1 shows parton spin, not controversial to absorb into the others.

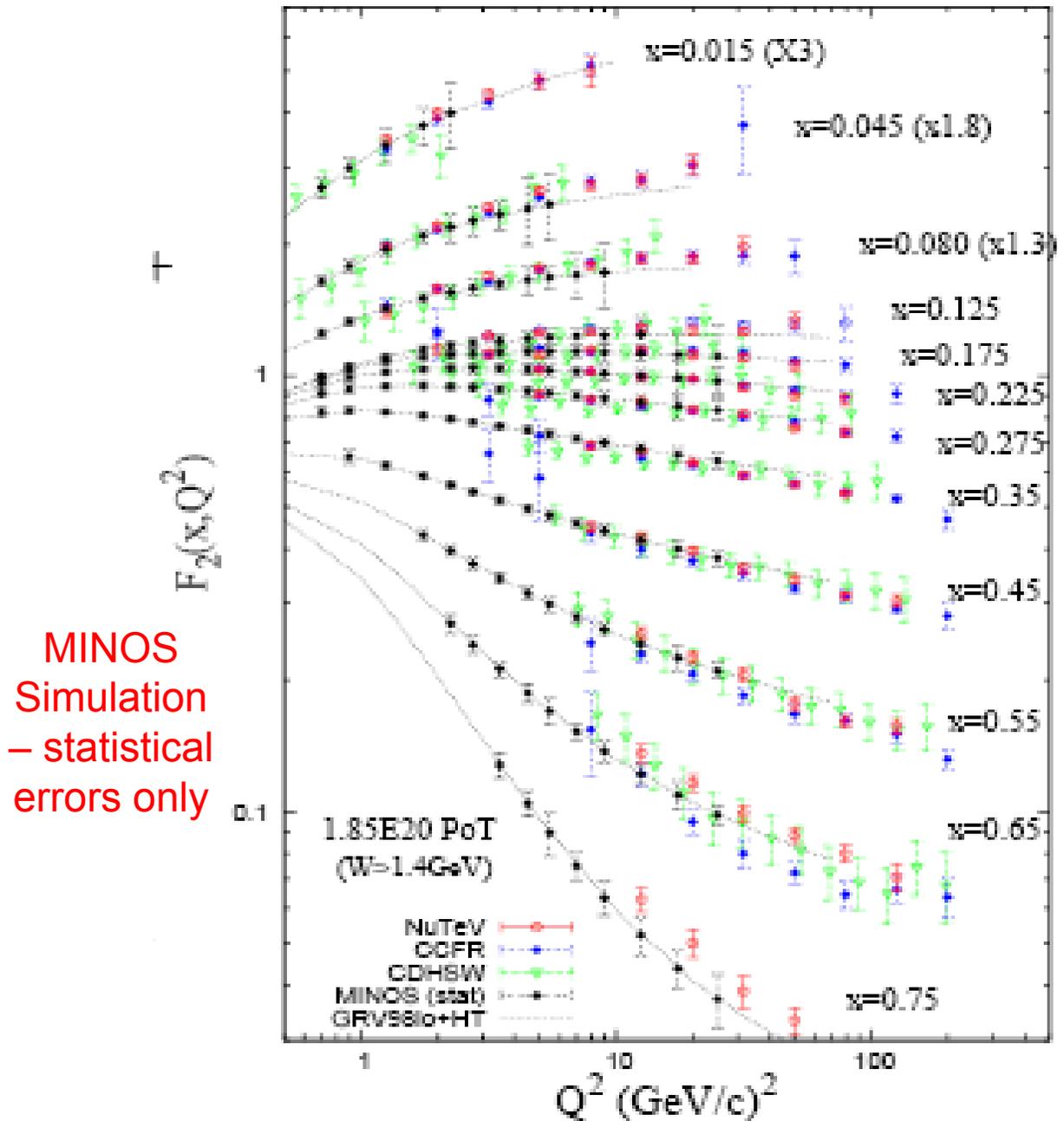
Extract F_2 from sum of neutrino, antineutrino cross-sections.

Only (anti)neutrinos can be used to get at F_3

Many world experts at this gathering!

Summary of Neutrino Experiment Data and Expected Extensions to Sensitivity

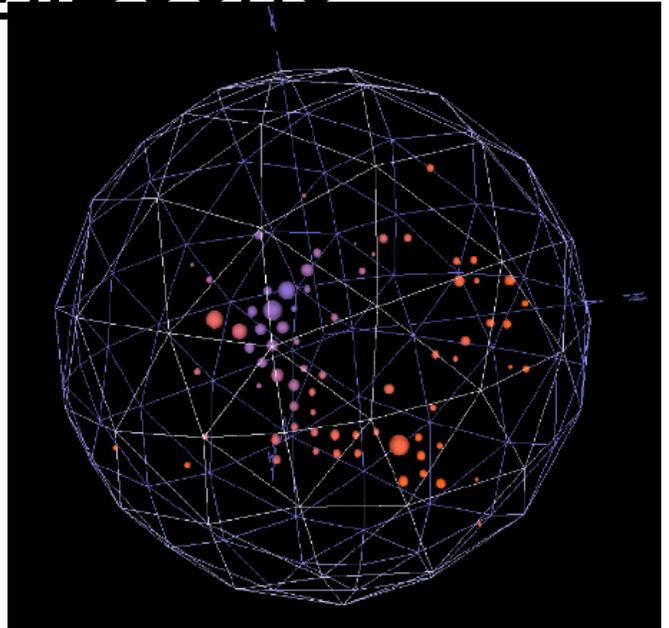
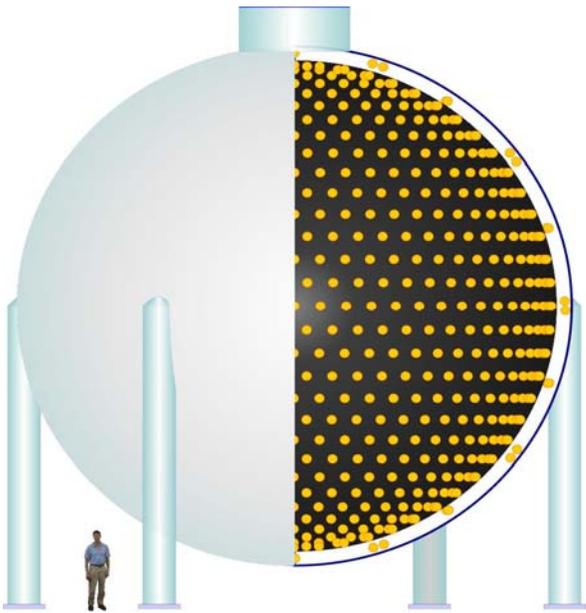
MINOS $F_2(x, Q^2)$ Sensitivity



Currently MINOS > 2.5 E20 Protons

An appearance experiment
using ν_μ at high values of
 Δm^2

- MiniBoone -



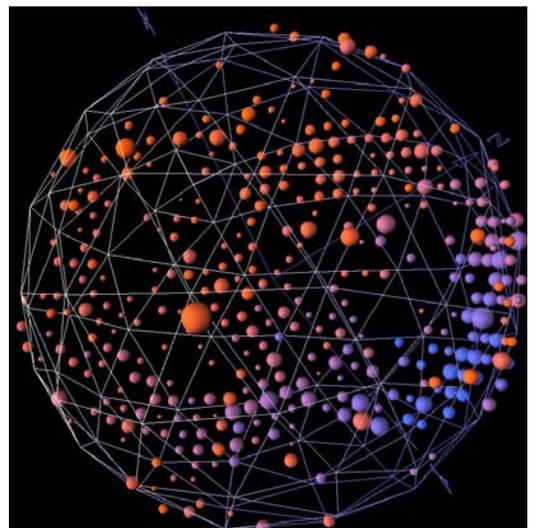
e from μ decay candidate.

Ragged outer edge of ring
from scattering, brems

12 m sphere, 950 K liters of oil.

1280 PMT's - 8" diameter

Cerenkov and Scintillation light



π^0 candidate –
overlapping rings,

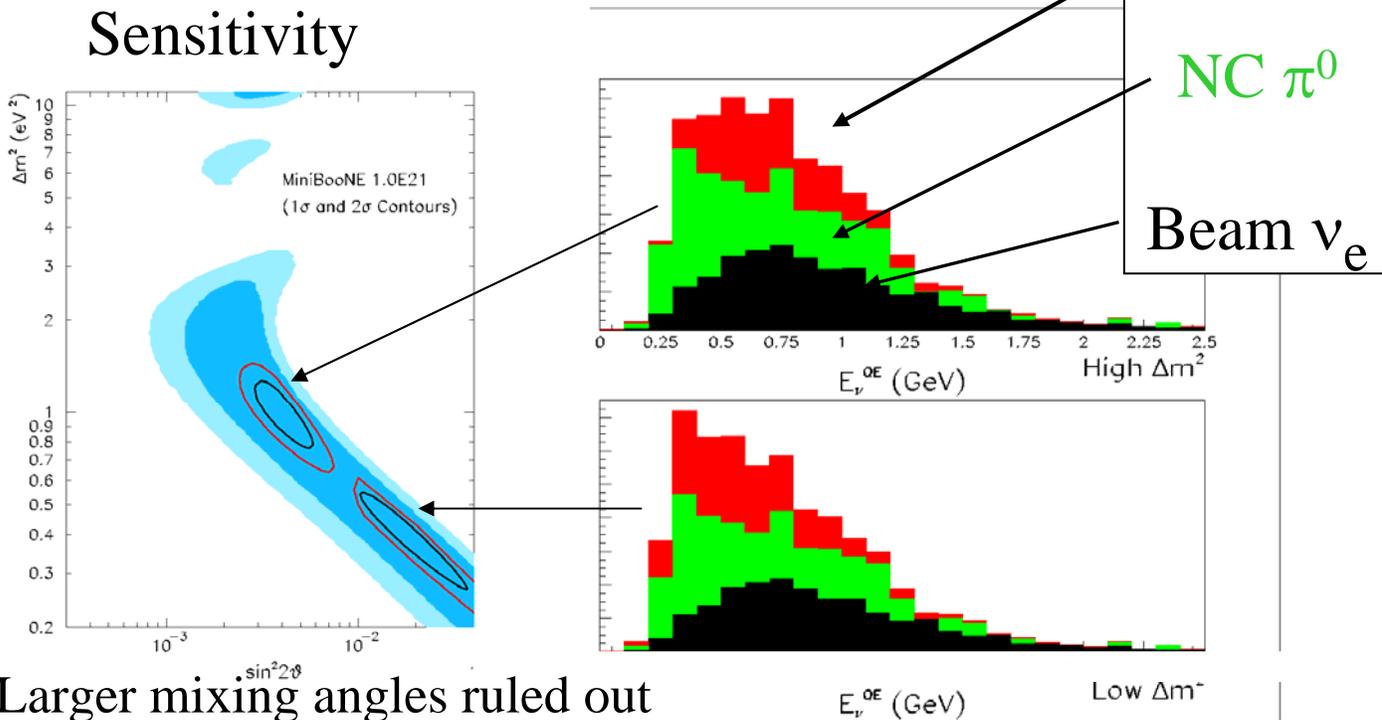
An appearance experiment using ν_μ at high values of Δm^2

Check/confirm LSND oscillation signal at Fermilab Booster

Different systematics from previous experiment

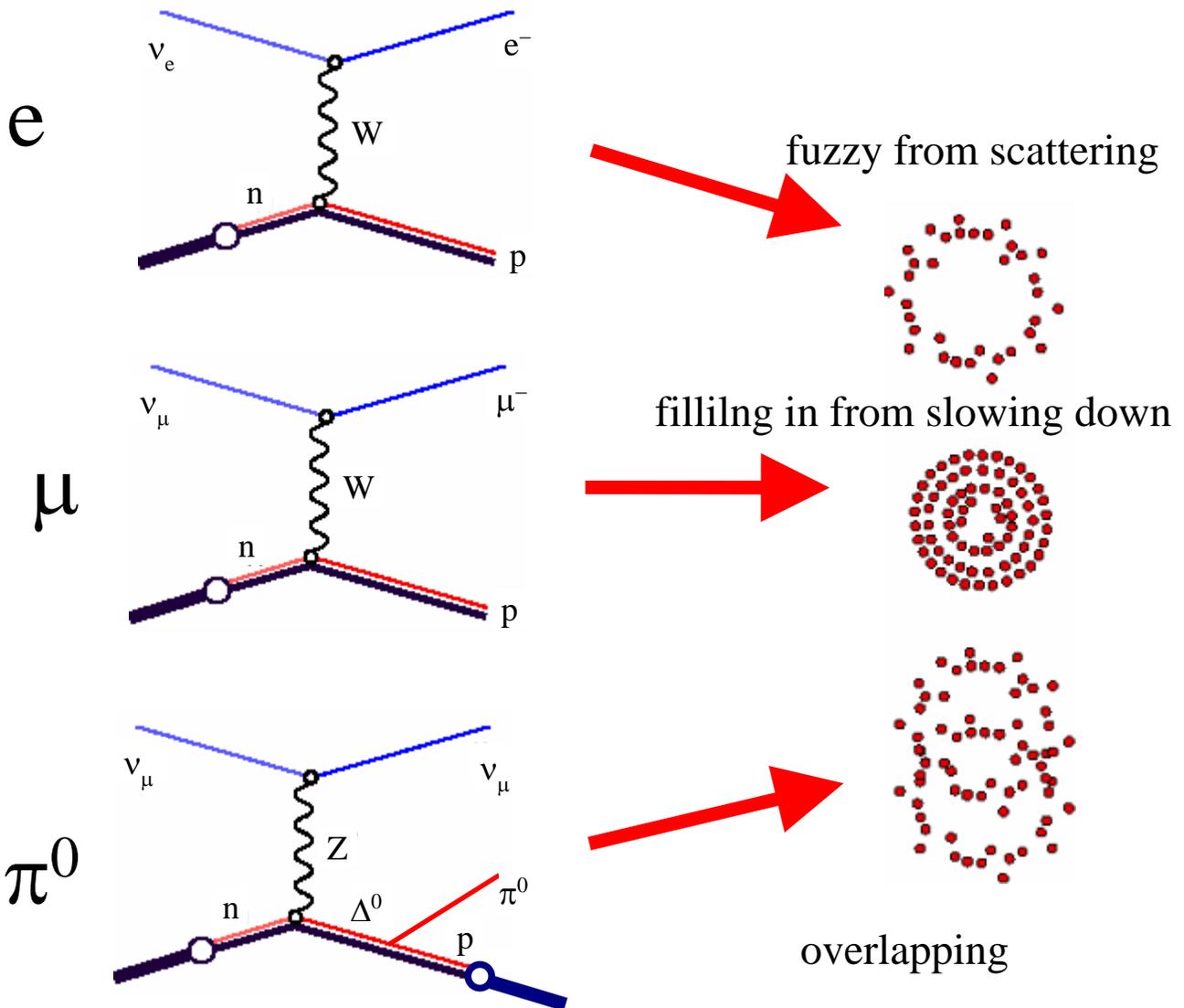
-L=540 m ~10x LSND

-E~500 MeV ~10x LSND



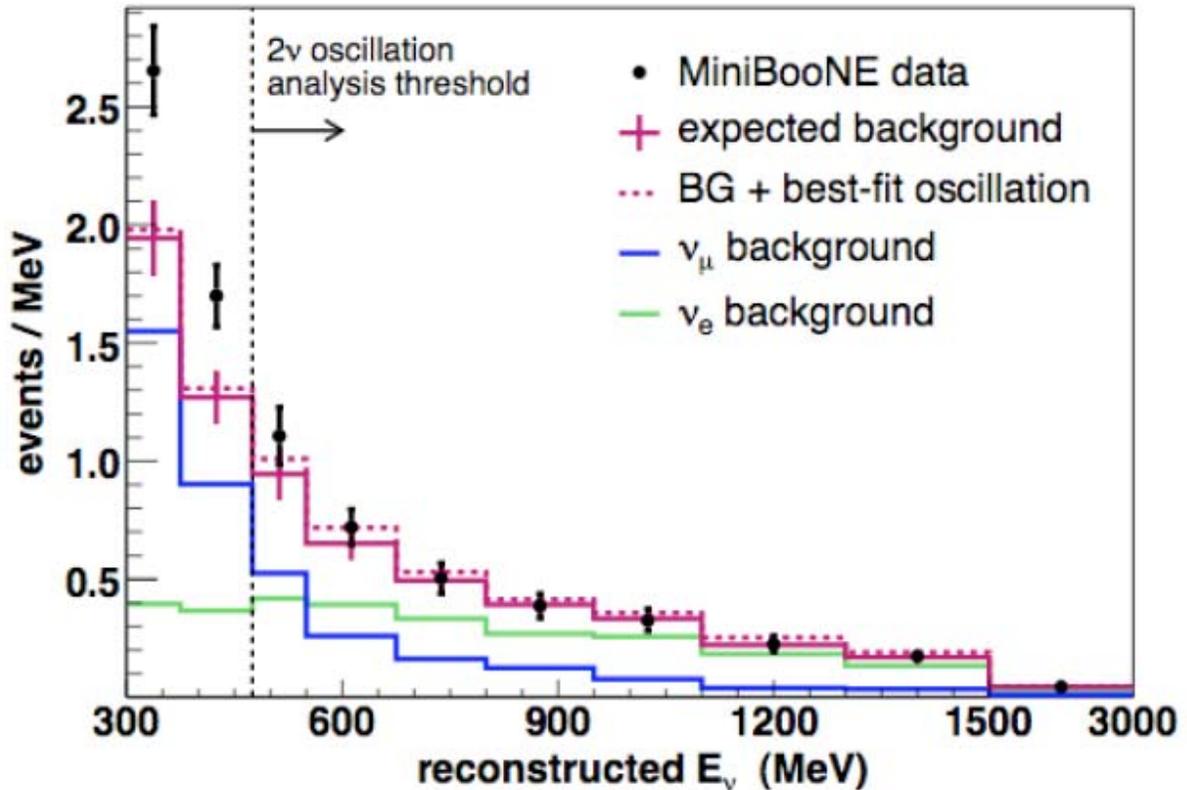
Larger mixing angles ruled out in the 80's

Event topologies in the MiniBoone detector



Graphic from S. Brice

MiniBoone Initial Result Blind Analysis (April 2007)



data: 380 events

expectation: 358 ± 19 (stat) ± 35 (sys) events

Inconsistent with ν_μ oscillation to ν_e as sole phenomenon (shape disagreement)

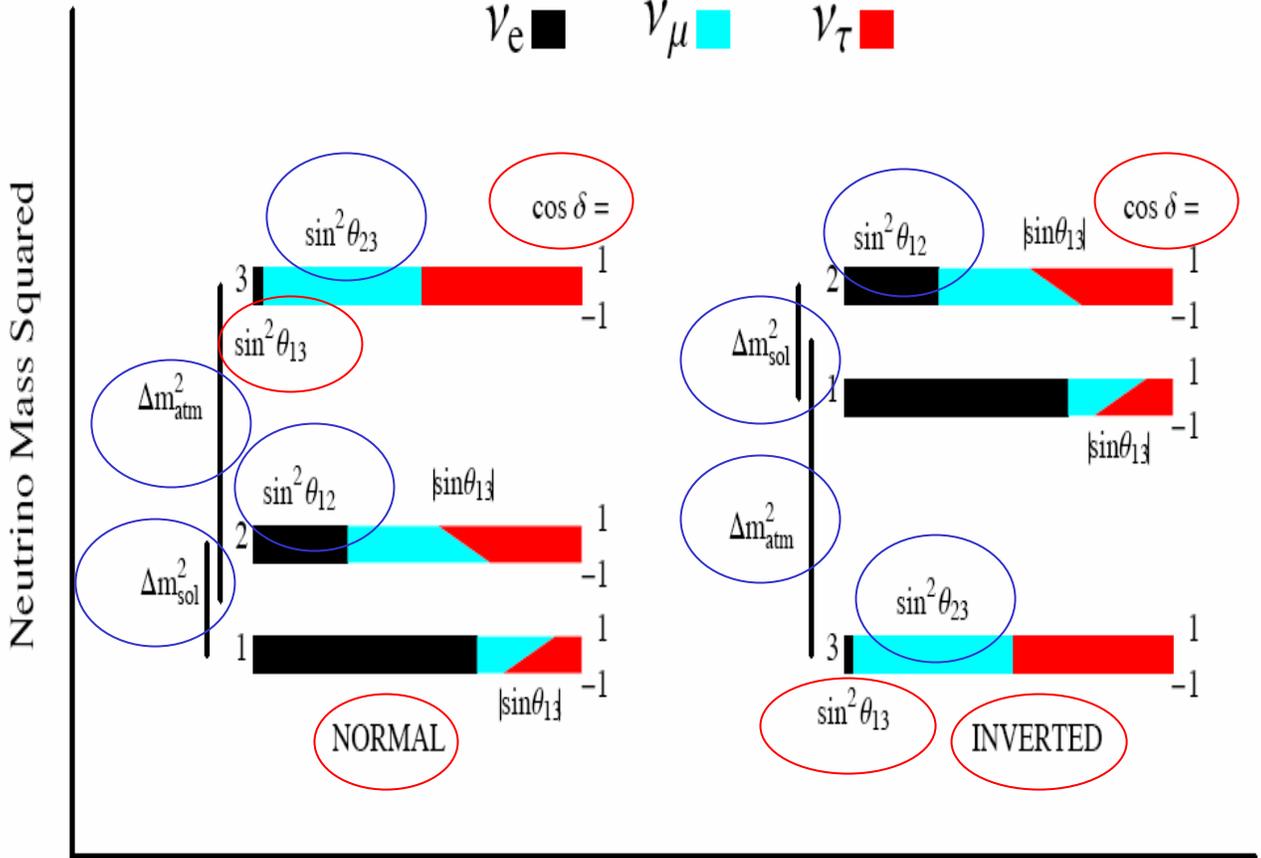
Disagreement at < 500 MeV still under investigation

Status of our Knowledge

○ Know to some extent

○ Don't know

ν_e ■ ν_μ ■ ν_τ ■



Fractional Flavor Content varying $\cos \delta$

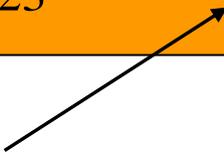
O. Mena and S. Parke, hep-ph/0312131

ν_e Appearance and the Importance of θ_{13}

In our formalism, with $\Delta_{12} \cong 0$ (solar),

$$P_{\nu_{\mu} \rightarrow \nu_e} = 4 \cos^2 \theta_{13} \sin^2 \theta_{13} \sin^2 \theta_{23}$$

or finally to leading order and using $\cos \theta \sin \theta = \frac{\sin 2\theta}{2}$

$$P_{\nu_{\mu} \rightarrow \nu_e} = \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 (\Delta_{13})$$


This small number will be seen again and again.

Currently $\sin^2 2\theta_{13} \leq 0.14$

Reactor Experiments - sooner and later

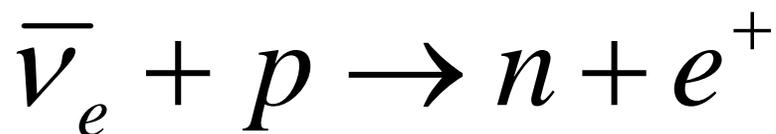
Reactor experiments provide an alternative route to measurement of $\sin^2 2\theta_{13}$.

$P(\nu_e \rightarrow \nu_e) = 1 - \sin^2 2\theta_{13} \sin^2(\Delta m_{\text{atm}}^2 L/4E) + \mathcal{O}(\text{solar})$
Free from CP, hierarchy asymmetries.

Disappearance experiments at modest baselines

Ballpark: $\langle E\nu \rangle \cong 3 \text{ MeV at } 1 \text{ km gives } L/E = .003$
comparable to MINOS

Reaction

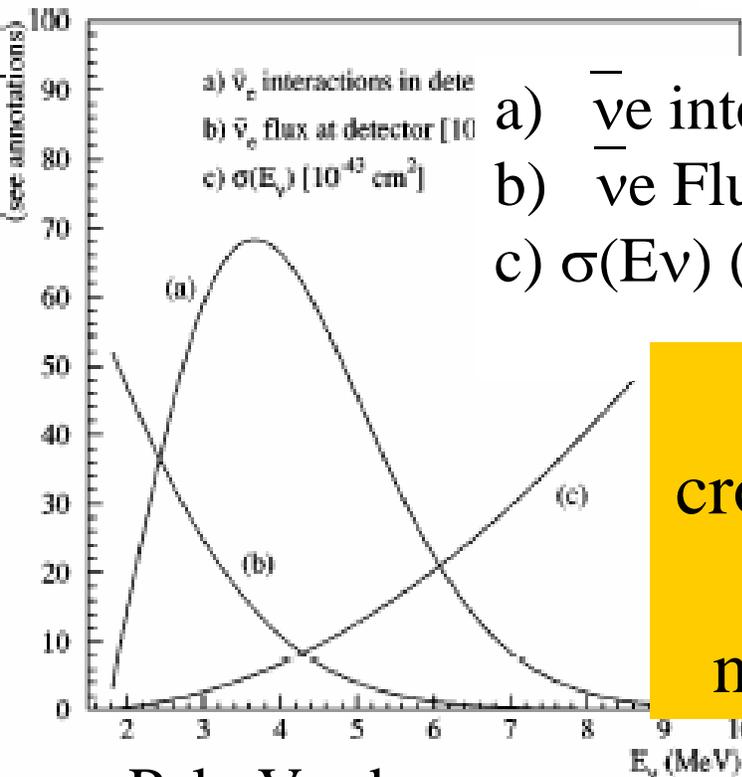
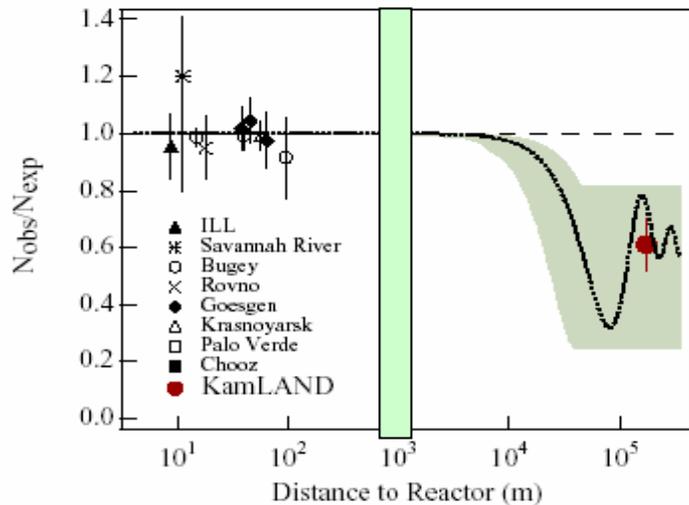


Detect via annihilation γ , followed by delayed γ from n-capture on nucleus

Reactor Experimental Landscape

KamLAND sees a 40% deficit/shape at 200km related to Δm^2_{12}

Search for a 1-5% deficit/shape at ~ 1 km related to Δm^2_{13}



Palo Verde

Courtesy KamLAND, Palo Verde, Reyna

≠ from Palo Verde

- a) $\bar{\nu}_e$ interactions/MeV-day
- b) $\bar{\nu}_e$ Flux ($10^8/\text{MeV}\cdot\text{s}\cdot\text{cm}^2$)
- c) $\sigma(E\nu)$ (10^{-43} cm^2)

High flux, low cross-section, results in acceptable numbers of events

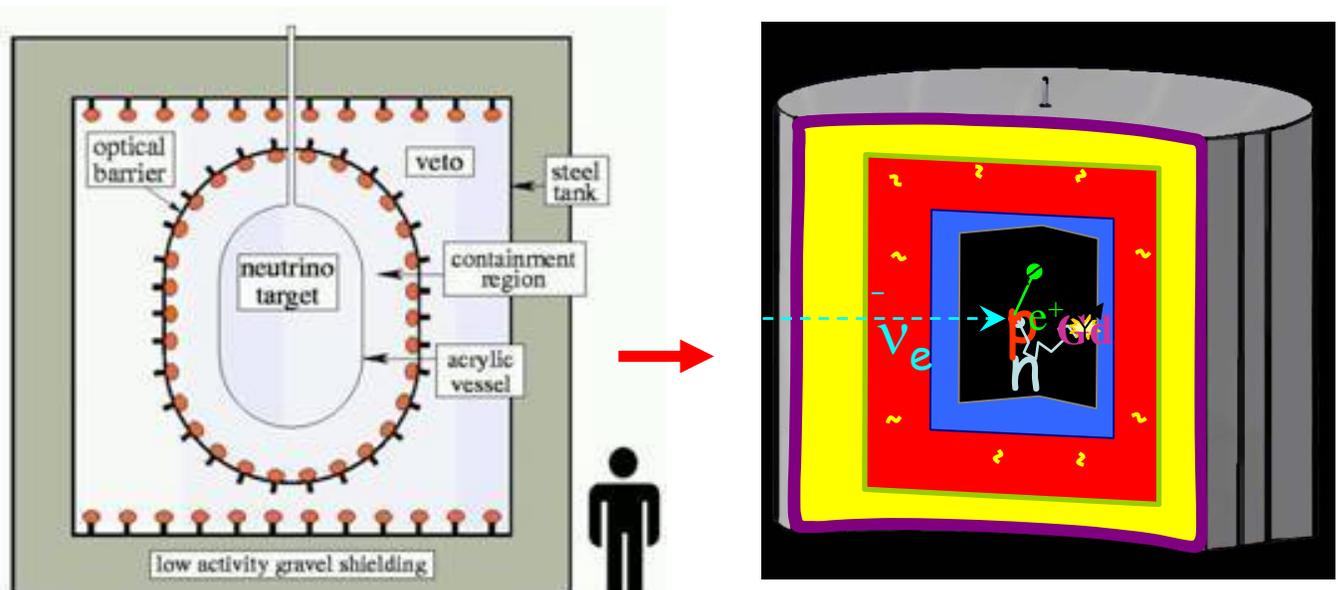
Near term for reactor

θ_{13} Double Chooz

Currently world's best limit $\sin^2 2\theta_{13} < 0.14$

Improve by adding near detector, increasing mass of detector(s) 5T \rightarrow ~ 10 T

Reactor Power substantially increased.



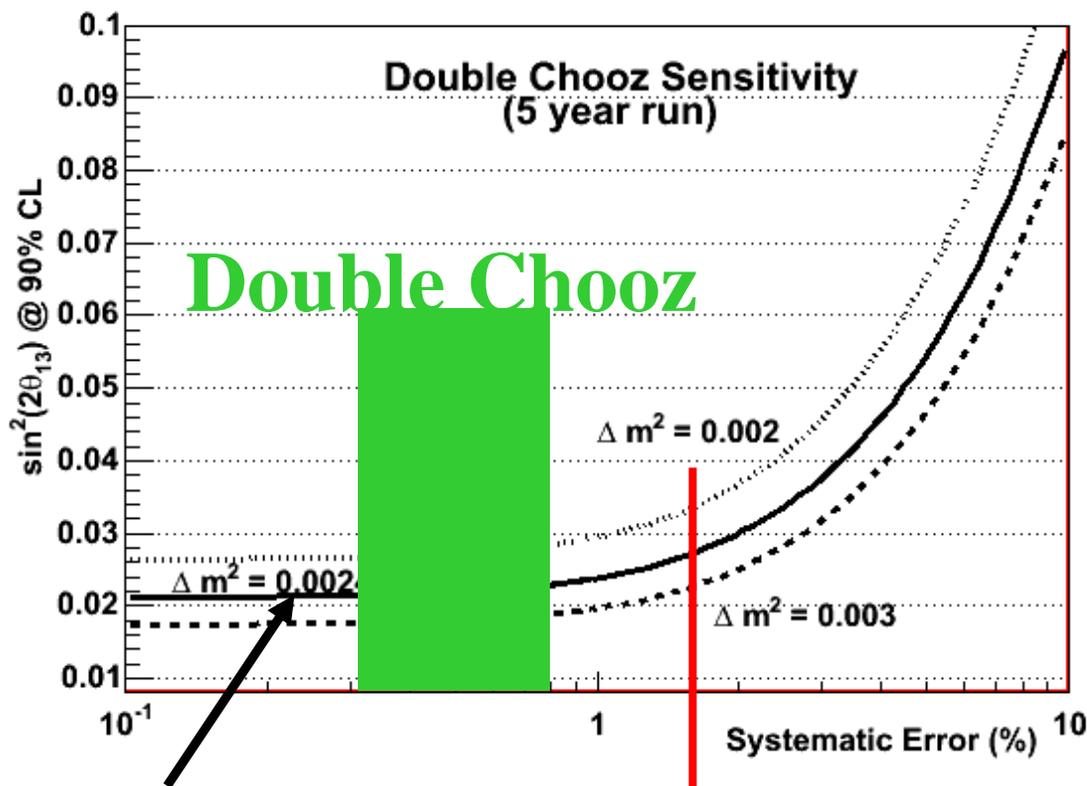
Distances : Near 100-200m Far: ~ 1 km

Depths: Near 30-40m Far: ~ 100 m

Exposure: 12 GW-T-year \Rightarrow 200-300 GW-T-year

Events: 2700 \Rightarrow 40K

Double Chooz Sensitivity

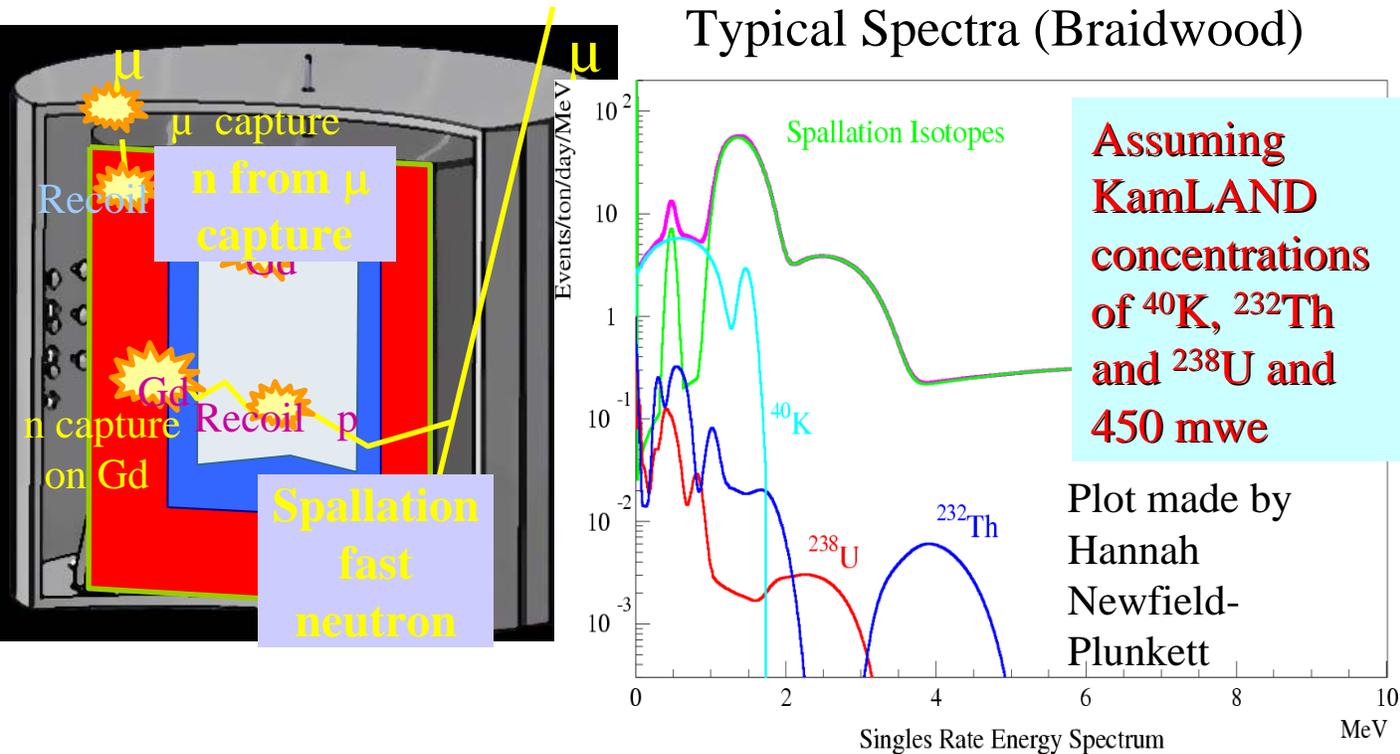


statistical limit

Original Chooz
Detector Error

	Chooz	Double-Chooz
Reactor cross section	1.9 %	—
Number of protons	0.8 %	0.2 %
Detector efficiency	1.5 %	0.5 %
Reactor power	0.7 %	—
Energy per fission	0.6 %	—

Background reduction in DoubleChooz



Background coming from radioactive materials, cosmic sources.

Spallation neutrons from μ , direct β -emitter creation

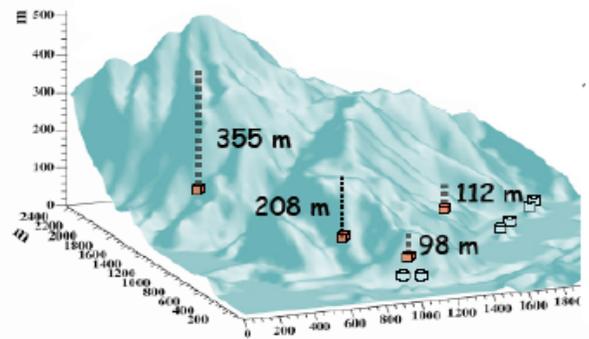
Control via detector design, reactor-off comparison with simulation

Expect total background subtraction systematic $\sim 1\%$

A Next Generation Reactor Experiment – Daya Bay (China)

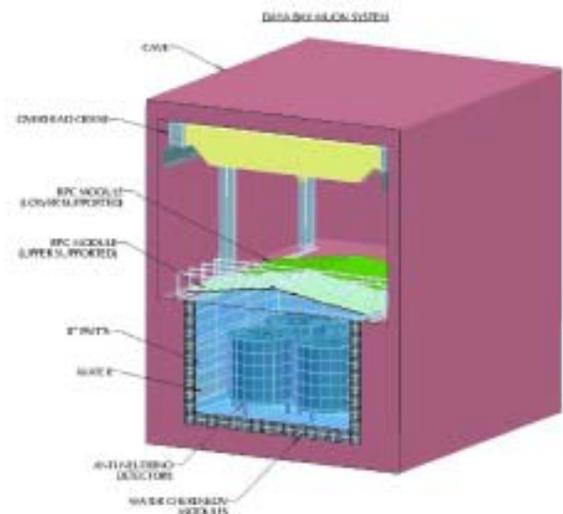
Main avenues of improvement

- Increased Mass
- Depth

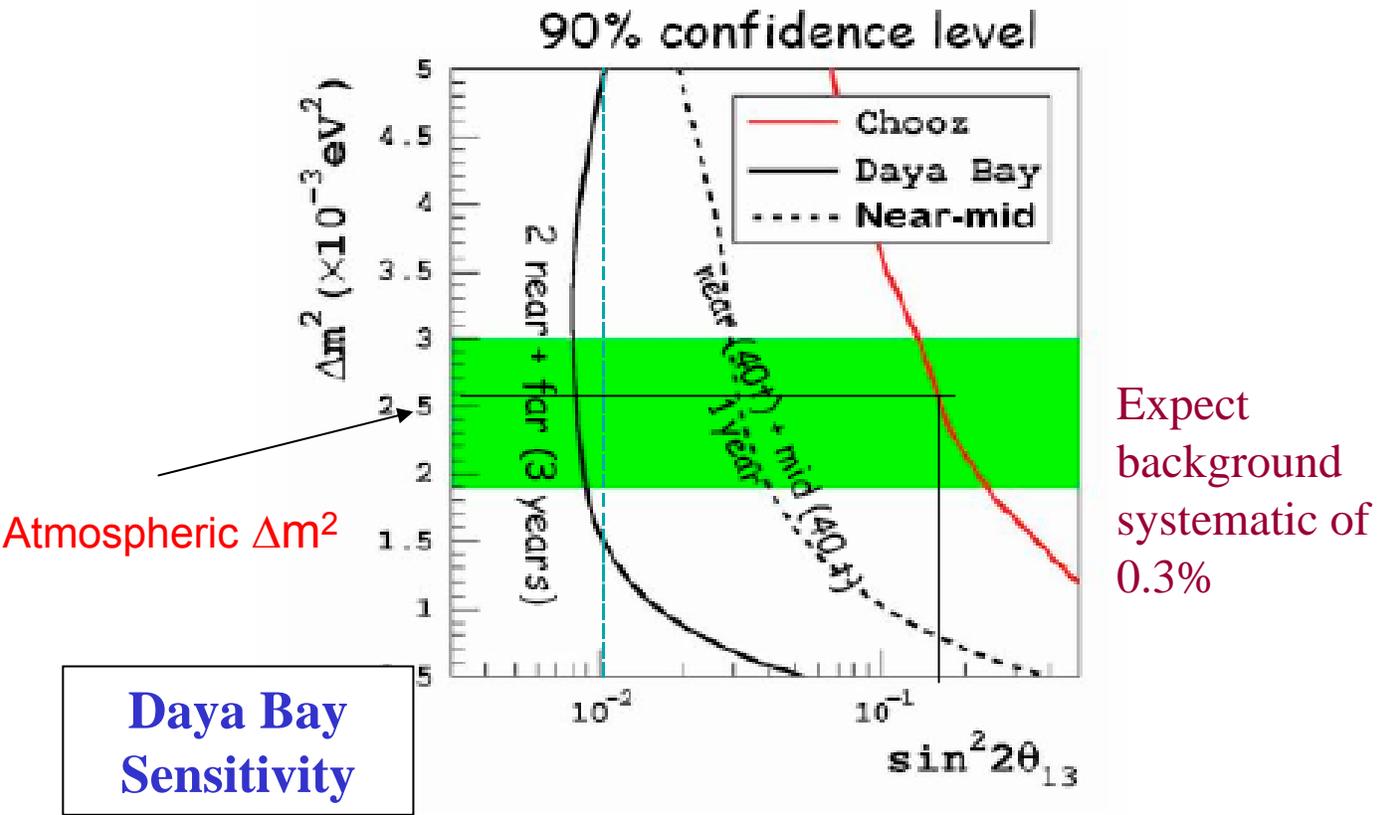


4 “Modules” at far position, 2 at each near.

Daya Bay, China
8 x 20 T detectors
(fiducial)
1000 mwe (far)



Sensitivity and Systematics



- Use rate and spectral shape
- input relative detector systematic error of 0.38%

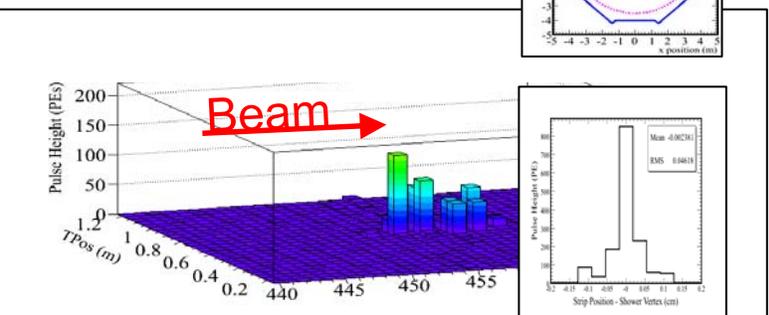
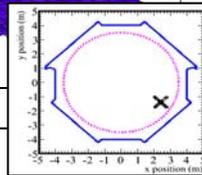
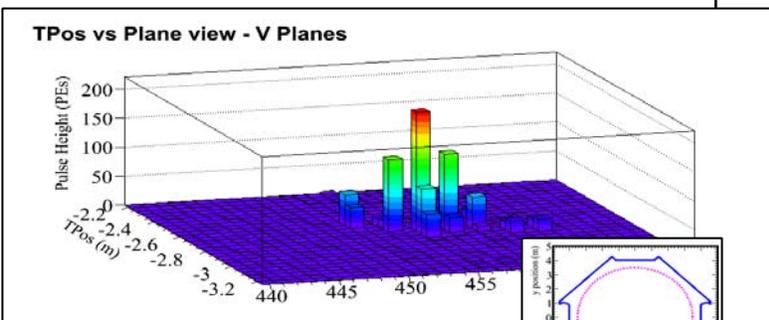
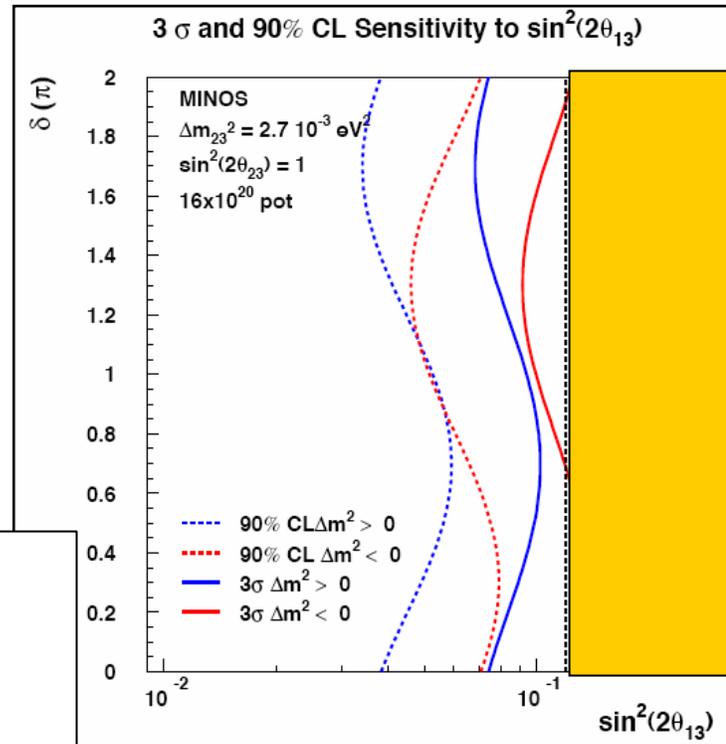
Source of uncertainty		Chooz (absolute)	Daya Bay (relative)		
			Baseline	Goal	Goal w/Swapping
# protons	H/C ratio	0.8	0.2	0.1	0
	Mass	-	0.2	0.02	
Detector Efficiency	Energy cuts	0.8	0.2	0.1	0.1
	Position cuts	0.32	0.0	0.0	0.0
	Time cuts	0.4	0.1	0.03	0.03
	H/Gd ratio	1.0	0.1	0.1	0.0
	n multiplicity	0.5	0.05	0.05	0.05
	Trigger	0	0.01	0.01	0.01
	Live time	0	< 0.01	< 0.01	< 0.01
Total detector-related uncertainty		1.7%	0.38%	0.18%	0.12%

Overall detector + background about 0.5% systematics

MINOS electron sensitivity

Challenging because of detector granularity – typical electron is 8 planes long, 4 strips wide.

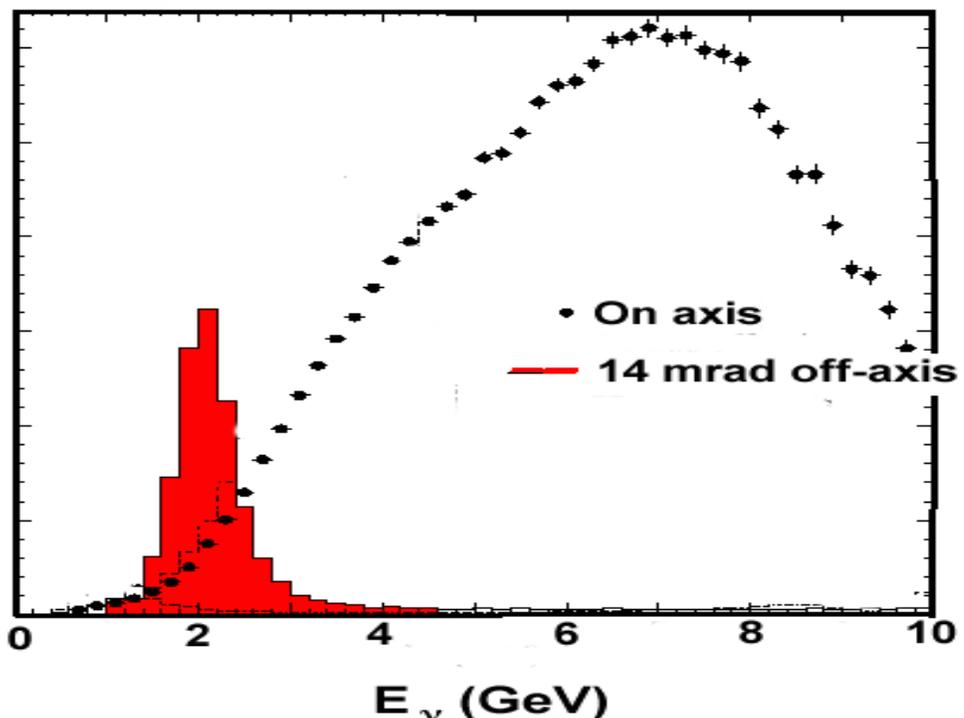
Background high – especially misidentified NC.



Multi-year run

Off-axis beams and NOvA

- Off-axis neutrino beams provide narrow-band kinematics
 - Reduces backgrounds
 - mis-id NC
 - ν_e 's from K decay (wrong kinematics)
- Increases flux at oscillation maximum.
- This provides a good setting for ν_e appearance experiments
 - Will be focus of Japanese T2K effort.
- NuMI beam already exists, can be exploited by construction of new detector.
 - NOvA proposal addresses this need.



Understand Off-axis Mathematically

Model beam with $\theta^* = \pi/2$

$$P_T = P^* \sin \theta^*$$

$$P_L = \gamma P^* (1 + \cos \theta^*) \cong E_\nu$$

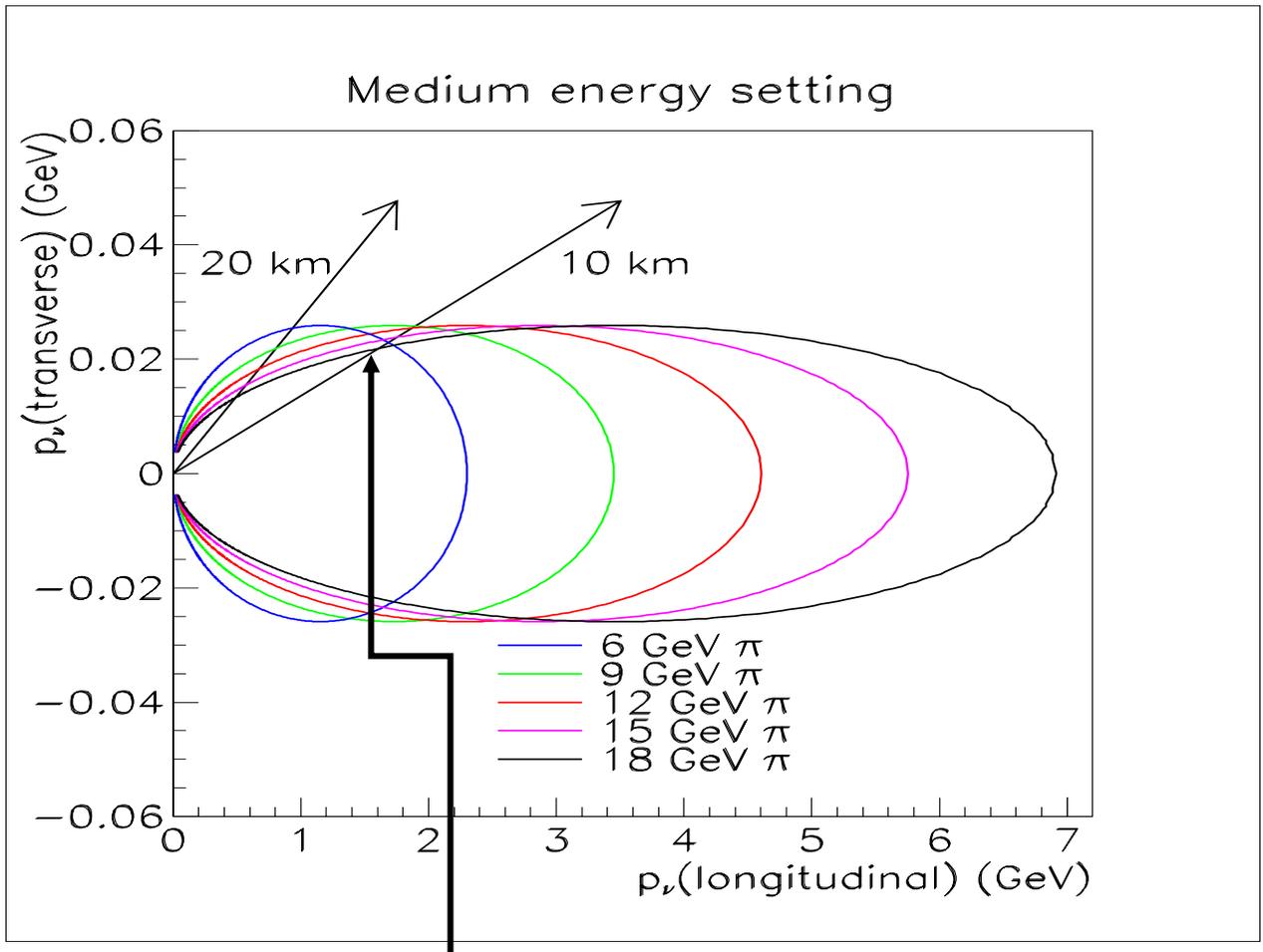
We fix $\theta_{lab} \equiv P_T / P_L$

At $\pi/2$ $\Delta P_T \cong 0$ as vary θ^*

$$\Delta P_L = \frac{\Delta P_T}{\theta_{lab}} \cong 0$$

Constraint on angle means γ compensates θ^* to fix P_L (at fixed θ_{lab})

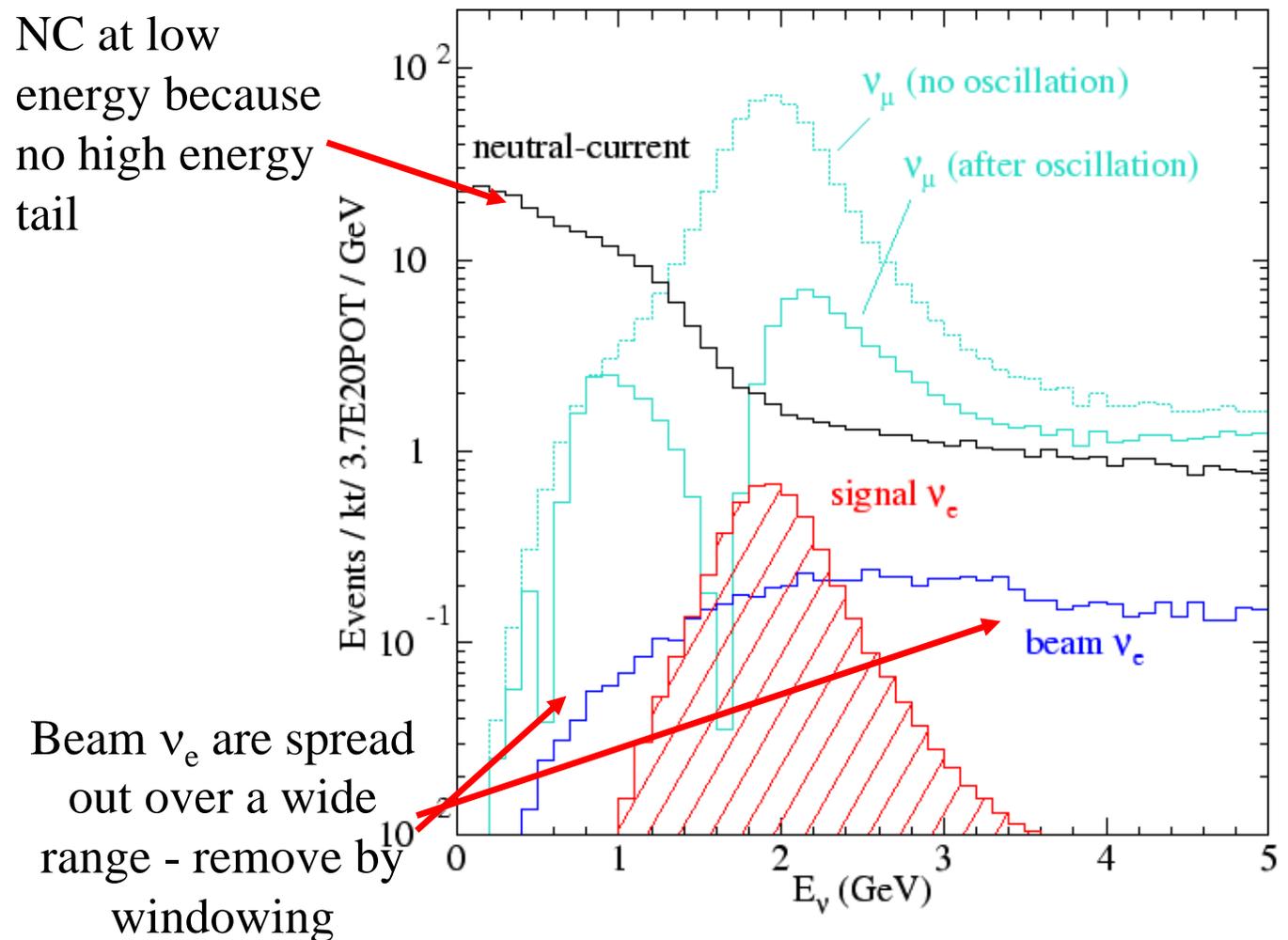
Off Axis Beam Graphically



Length of this vector very
insensitive to pion γ

$$E(\theta^* = \pi/2) = 1/2 E(\theta^* = 0)$$

Combination of off-axis “narrow-band” beam and good energy resolution.

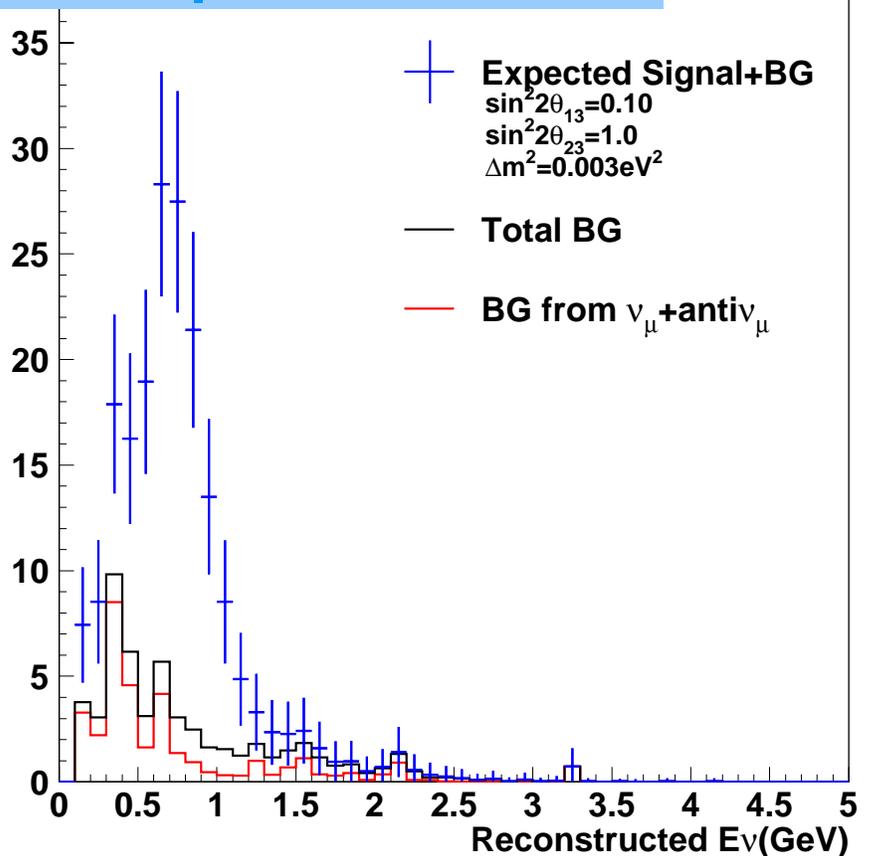


Backgrounds show a different energy spectrum and can be removed during analysis.

Sidebar - T2K Project at KEK



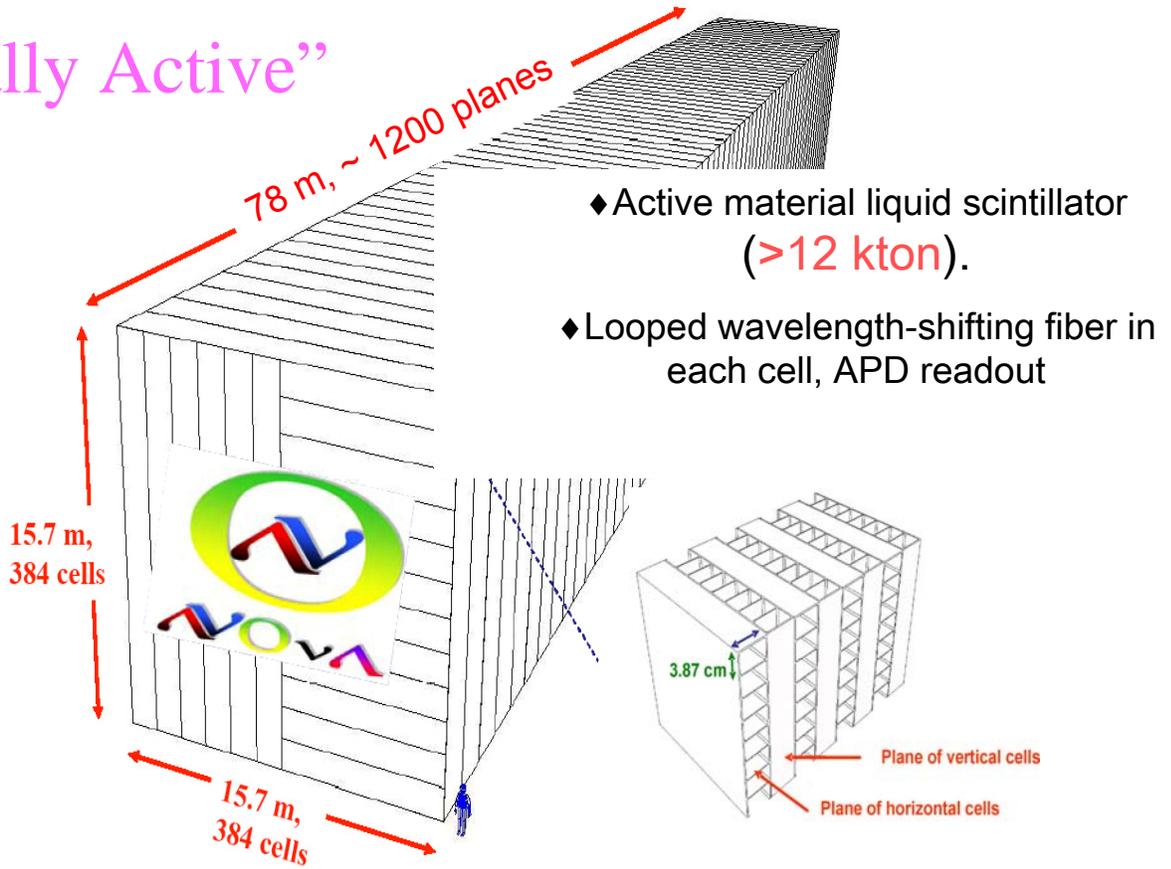
Hope for 2009
startup



Results of 5 years running

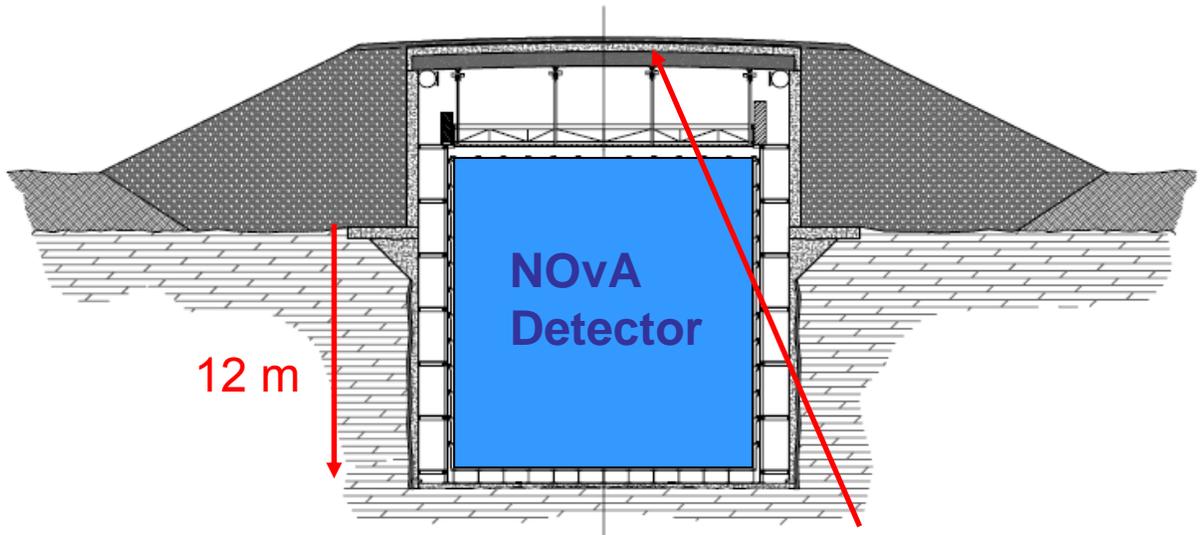
NOvA, a large electron appearance experiment

“Totally Active”

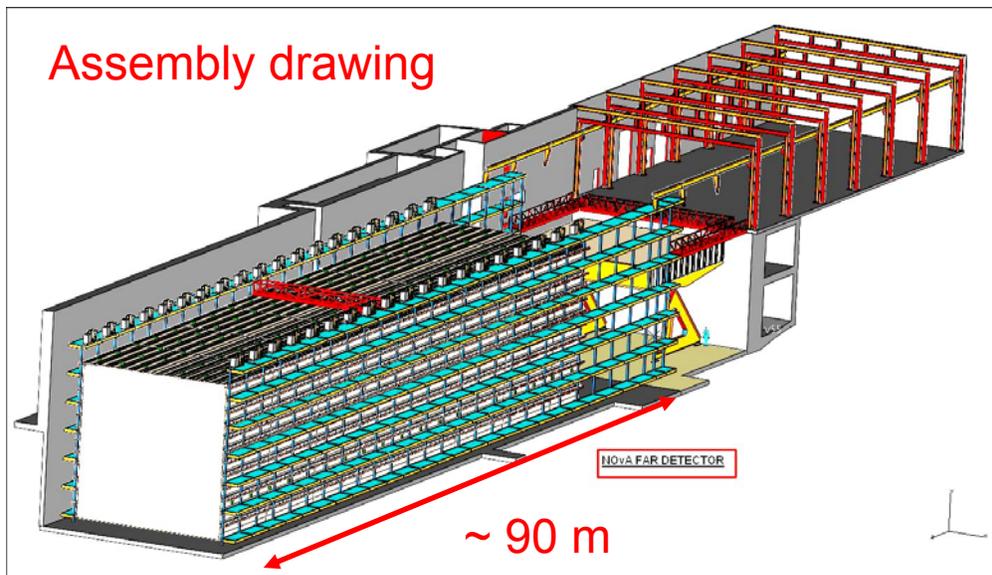


- 18 kton detector
- Baseline ~810 km
- ~1400 ν_{μ} CC events per $7e20$ POT ($\Delta m^2 = 2.5 \times 10^{-3}$)
- Electron ID efficiency 24%
- For $\sin^2 2\theta_{13} \sim 0.1$ would see ~100 ν_e interactions in 5 years

Large semi-underground building



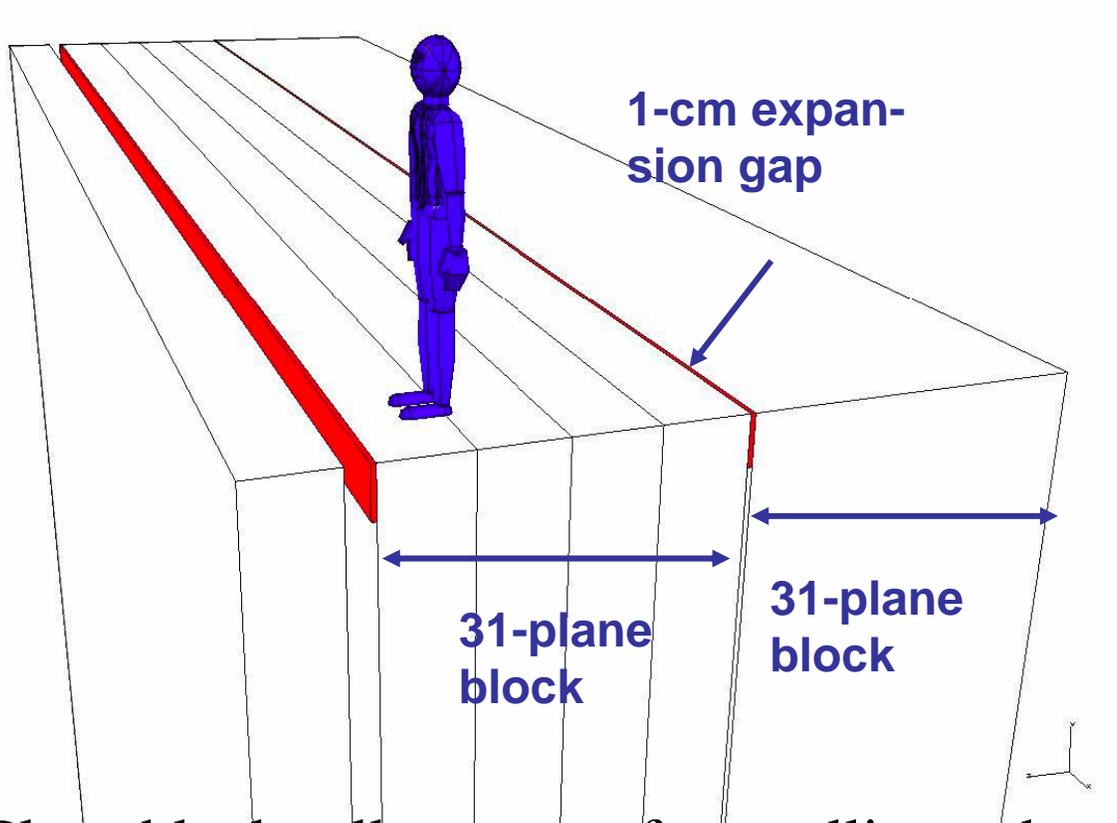
Barite cover reduces background from cosmic γ
(can fake electron)



Probable location at Ash River, MN
14 mrad off axis

Trigger on timing \rightarrow surface location OK

*For structural reasons,
detector is built in 31 plane
blocks*



31 Plane blocks allow space for swelling at bottom of tubes.

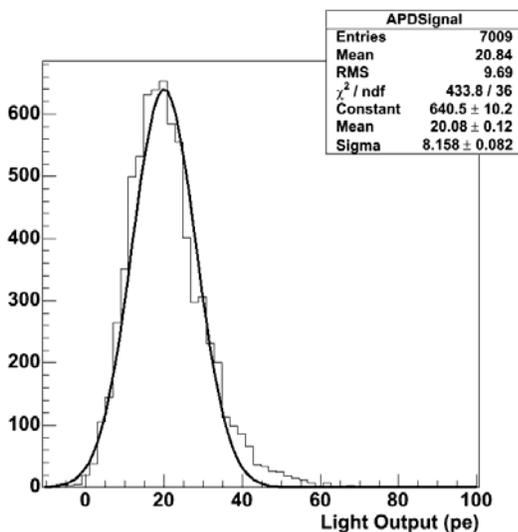
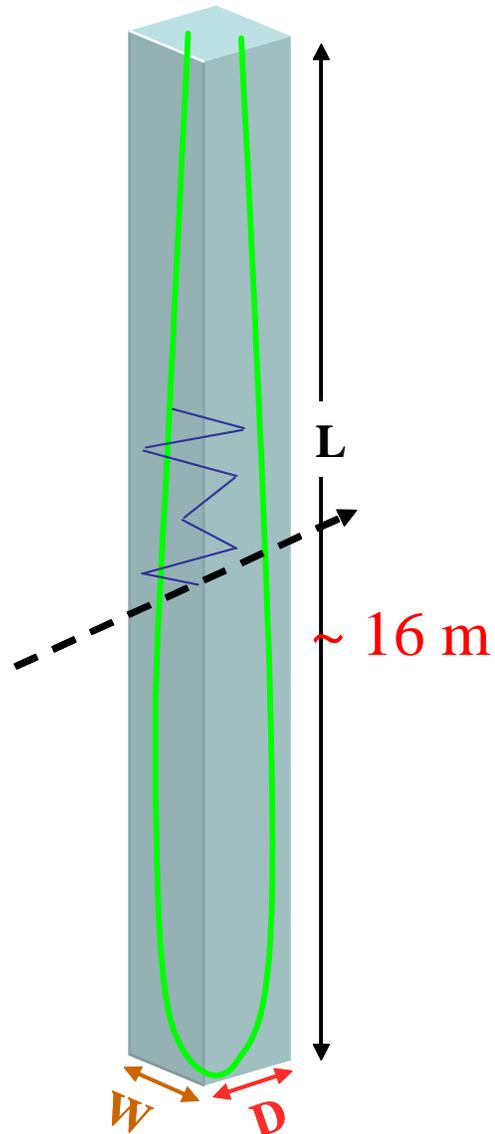
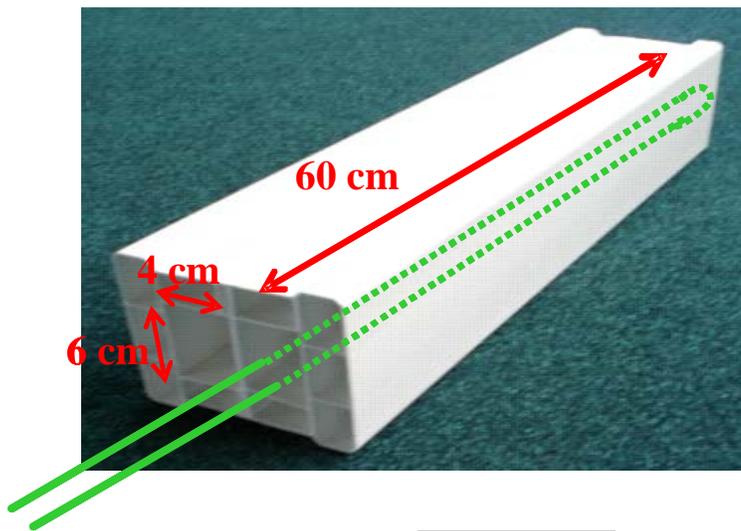
Structure erected block-by-block

Weight of each block (empty): ~125 tons

Extruded PVC tubes filled with liquid scintillator

Prototype with full-length fiber
Readout cosmic muons

To 1 APD pixel



Prototype APD readout showing light
output

Review of Oscillations in Matter

Recall $\nu_e + e$ scattering adds a potential

$$H_{matter} = H_{vac} + V_W \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

Leading after some manipulation to

$$P_{\nu_\mu \rightarrow \nu_e} = \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\Delta_{13}^2}{(\Delta_{13} \mp aL)^2} \sin^2(\Delta_{13} \mp aL)$$

where a is

$$a = \frac{G_F N_e}{\sqrt{2}} \approx (3700 \text{ km})^{-1} \left(\frac{\rho}{2.8 \text{ g cm}^{-3}} \right)$$

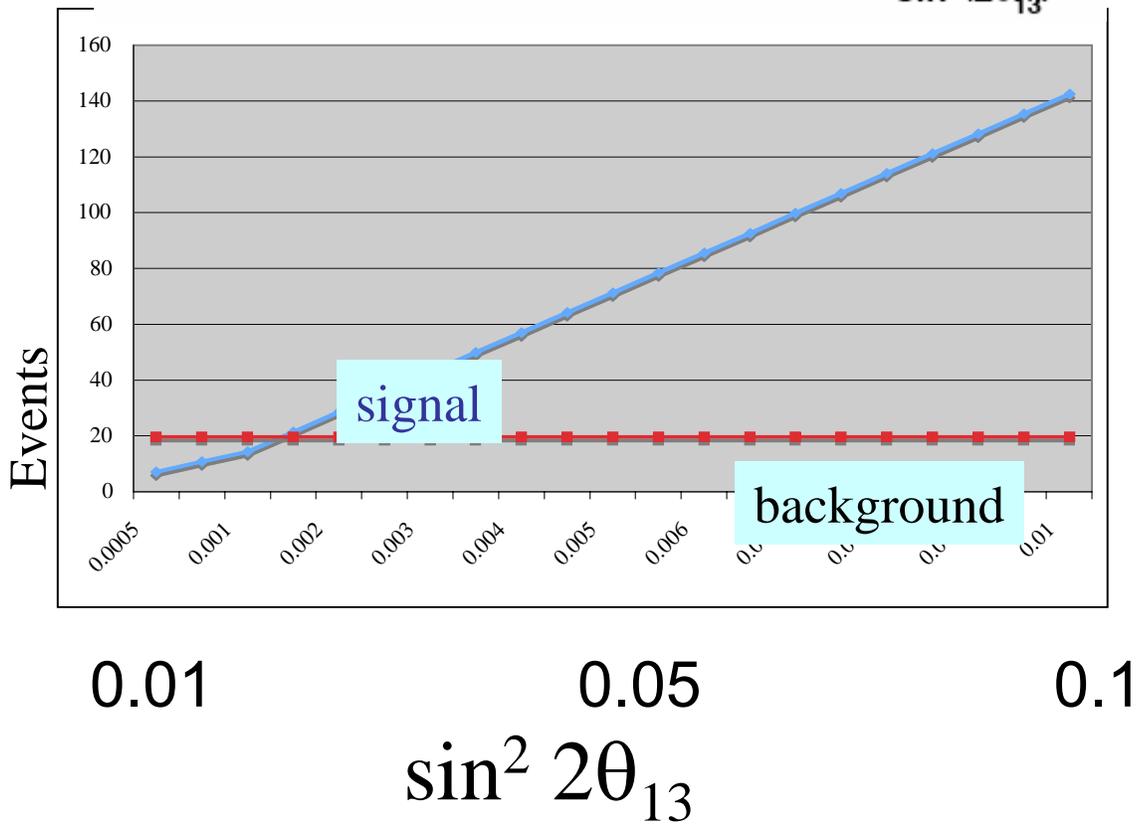
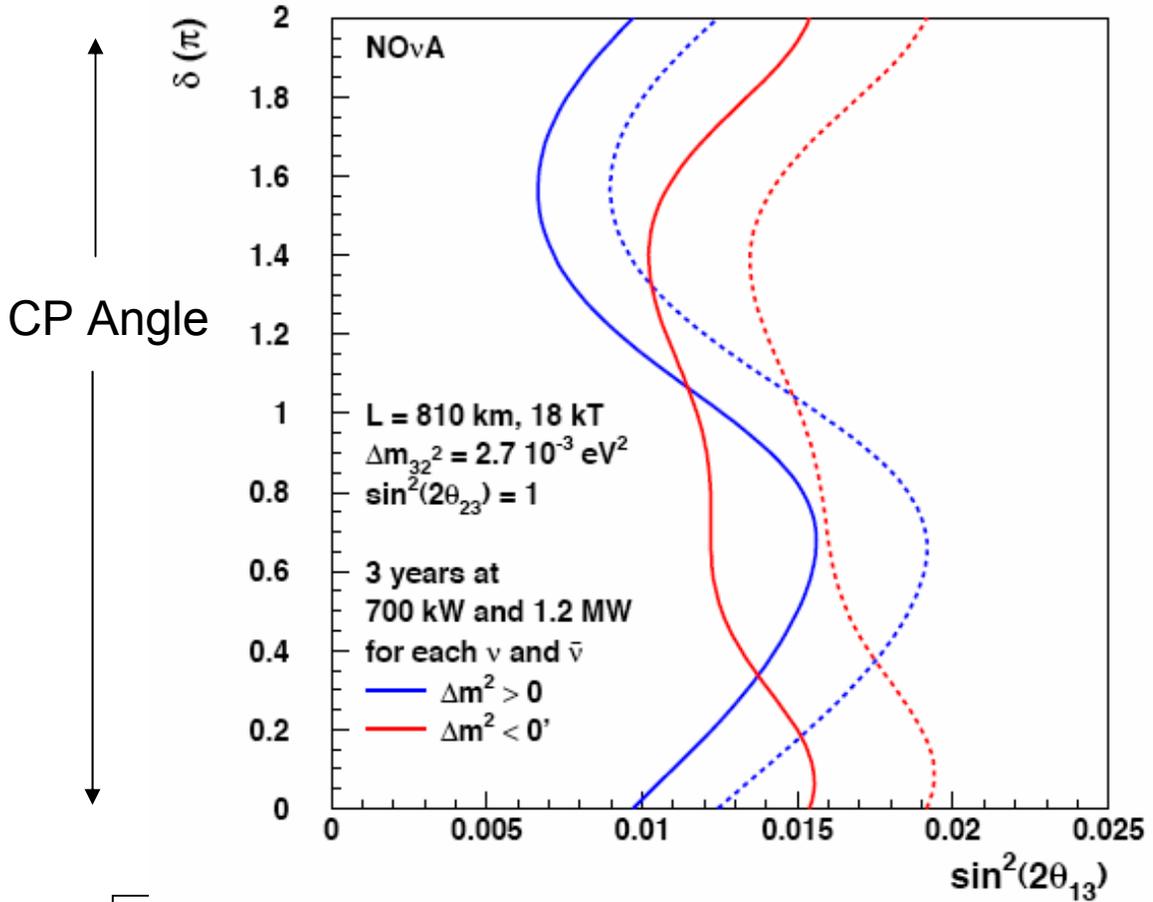
new mixing
strength

new effective
mass

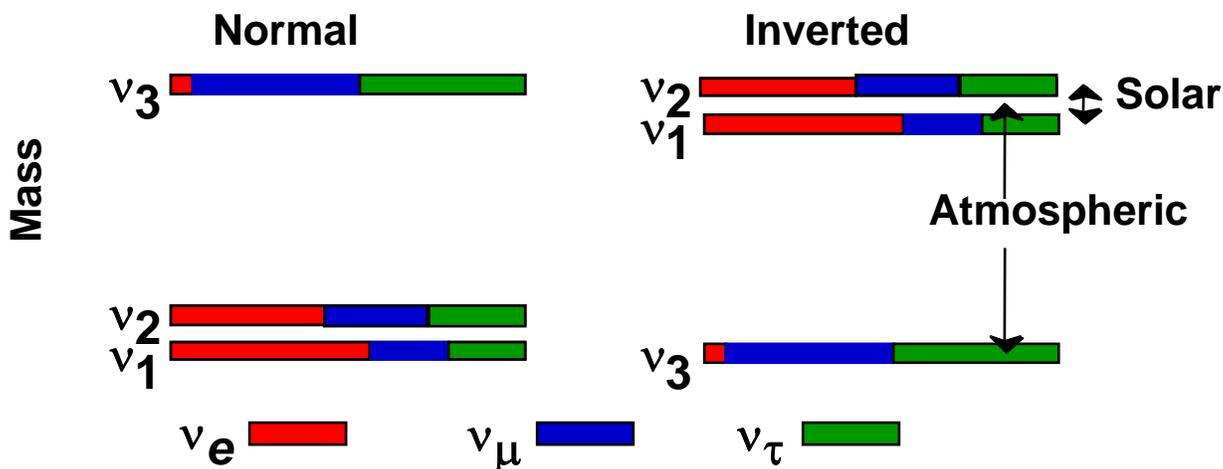
and the sign changes between neutrinos and antineutrinos

3- ν 's: Change is OK for all orders n
of $\sin^n \Delta_{12}$ and $\sin^n \Delta_{13}$

NOvA Sensitivity to θ_{13}



Mass Hierarchy and NOvA



Major unsolved problem

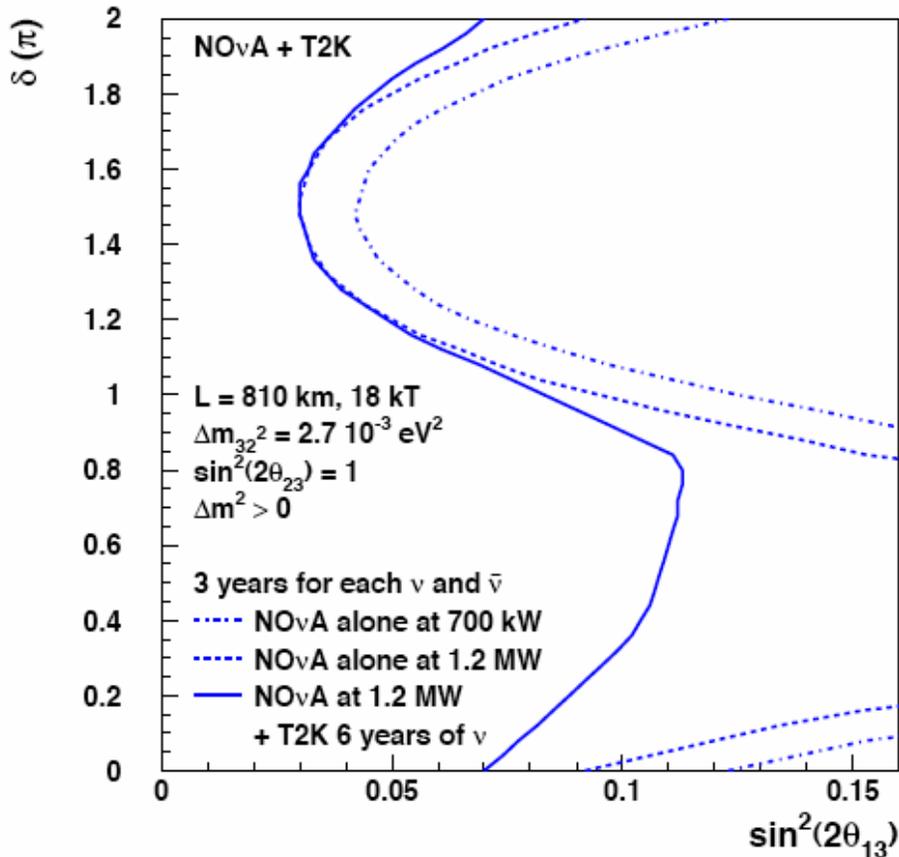
Requires matter effects to observe

Hence a real “target of opportunity” for a long baseline.

To leading order and zero CP violation, acts as if neutrinos and antineutrinos were interchanged.

CP violation makes it more complicated

NO ν A Coverage range for mass ordering



Upgraded Fermilab complex could produce more beam, on order 1-2 MW

Best results from worldwide combinations

Next Order and CP Violation

- Do not neglect Δ_{12}
- Add back CP phase δ
- Terms group as

$$\left| e^{-i(\delta + \Delta_{32})} A_{atm} + A_{solar} \right|^2$$

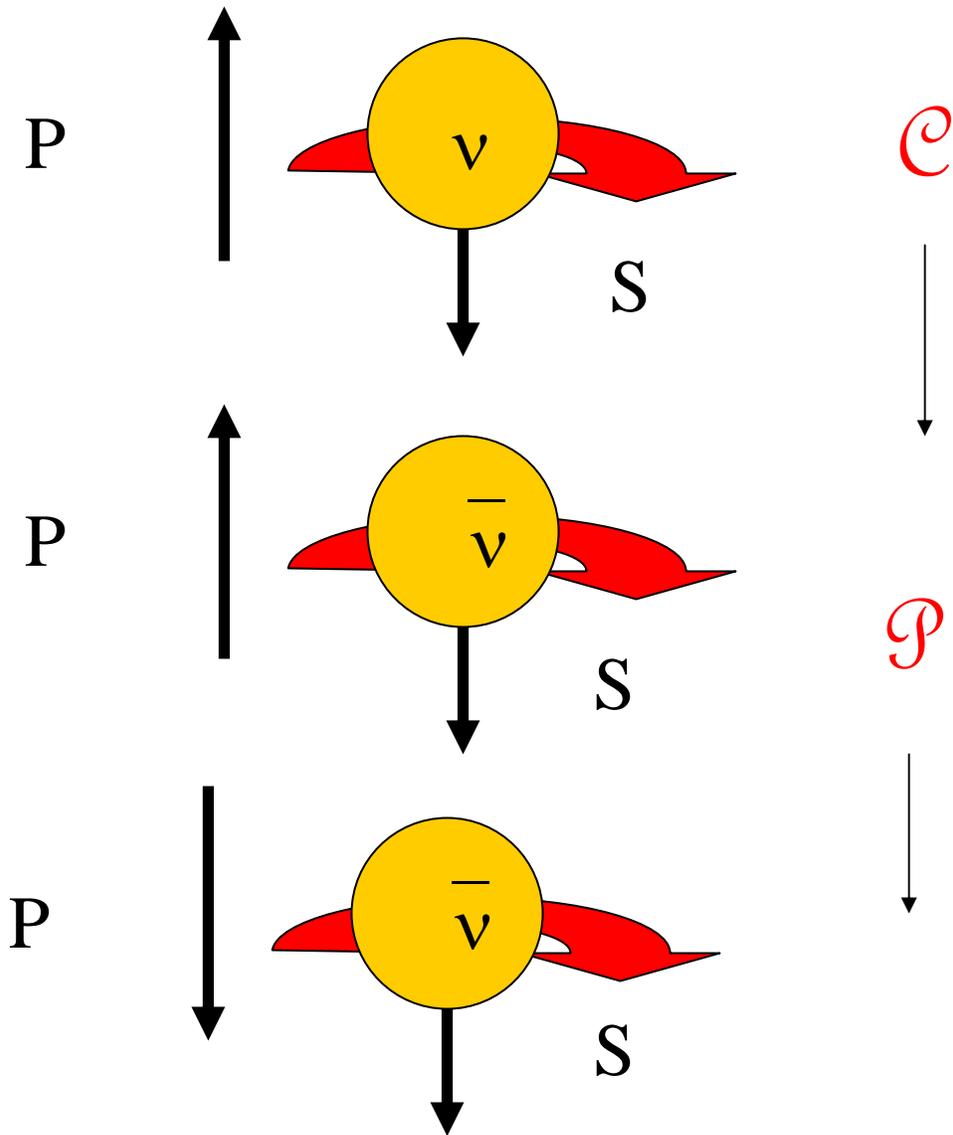
(Before we looked only at A_{atm}^2)

Giving

$$\begin{aligned}
 P(v_\mu \rightarrow v_e) = & \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \Delta_{31} \quad \text{Atmospheric} \\
 & + \cos^2 \theta_{13} \cos^2 \theta_{23} \sin^2 2\theta_{12} \sin^2 \Delta_{21} \quad \text{Solar} \\
 & + J_r \sin \Delta_{21} \sin \Delta_{31} \quad \text{Interference} \\
 & (\cos \Delta_{32} \cos \delta - \sin \Delta_{32} \sin \delta)
 \end{aligned}$$

Note J_r a combination of angles $\cong 0.9 \sin^2 \theta_{13}$

CP and neutrinos



CP gives right-handed $\bar{\nu}$, Dirac partner of left-handed ν

These are the states created by the weak interaction

CP and Oscillation Experiments

$$P(\nu_\alpha \rightarrow \nu_\beta)$$

CP

$$P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$$

T

$$P(\bar{\nu}_\beta \rightarrow \bar{\nu}_\alpha)$$

CPT conservation gives

$$\Delta P_{\alpha\alpha}^T = 0$$

$$\Rightarrow \Delta P_{\alpha\alpha}^{CP} = 0$$

Disappearance experiments can't see CP violation, only appearance experiments!

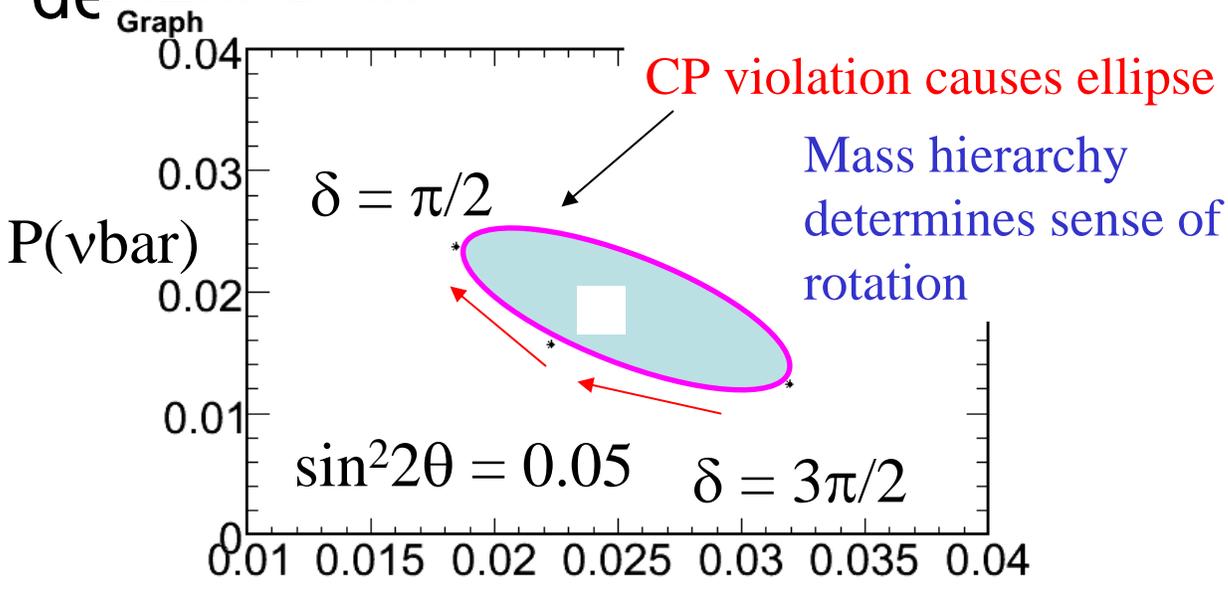
CP and Anti-Neutrinos

- CP δ changes sign between ν and $\bar{\nu}$.

$$\frac{P_{\bar{\nu}_\mu \rightarrow \bar{\nu}_e}}{P_{\nu_\mu \rightarrow \nu_e}}$$

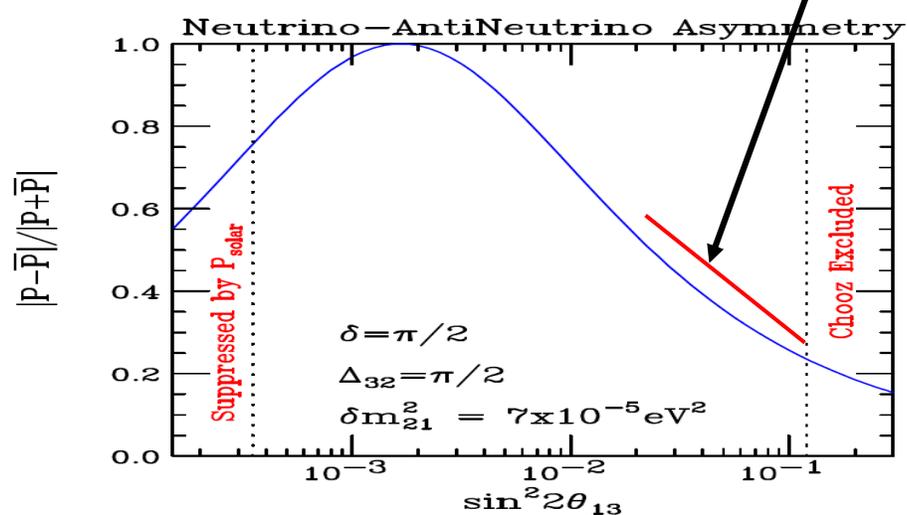
- As a result depends on δ

$$P_{\nu_\mu \rightarrow \nu_e}$$



Relevant for near-term experiments

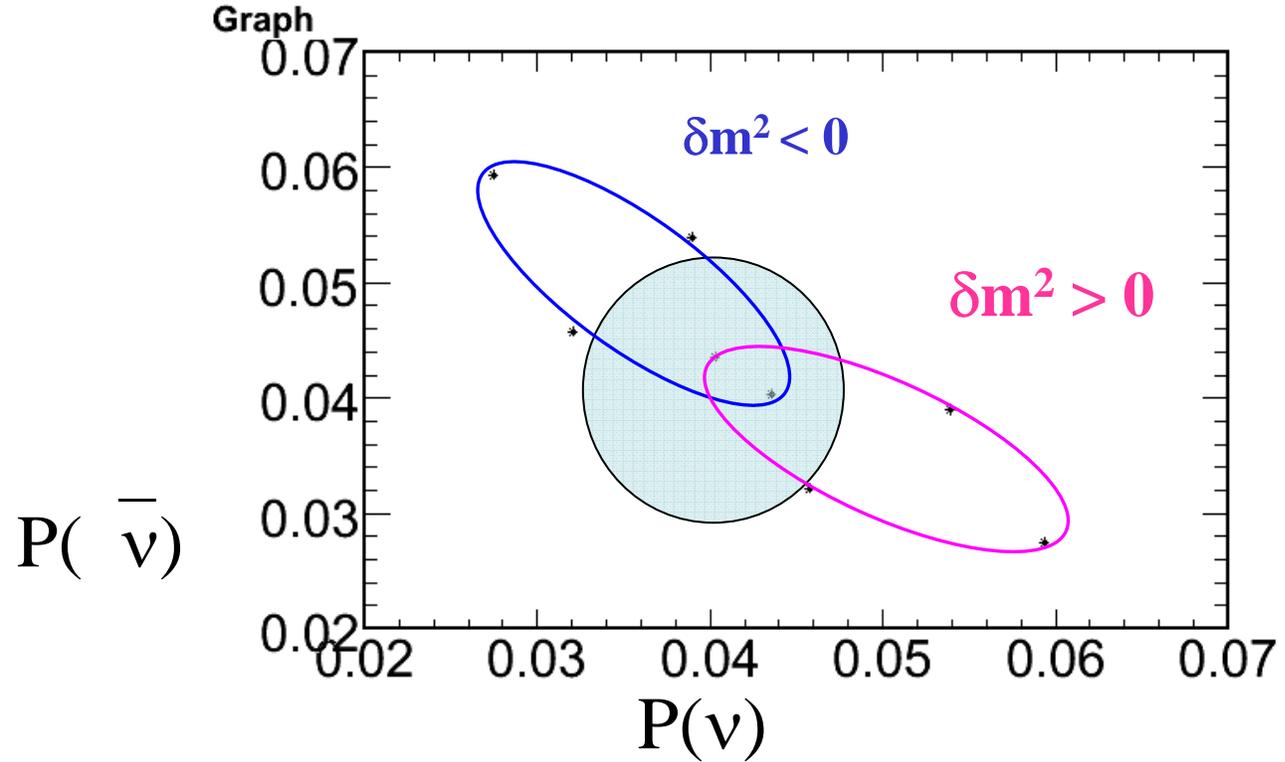
Ef



all θ_{13}
 $A_{\text{solar}} = A_{\text{atm}}$

Matter Effects
Separate the CP
Ellipses

$\sin^2 2\theta_{13} = 0.10$



Typical data for multiple years of running

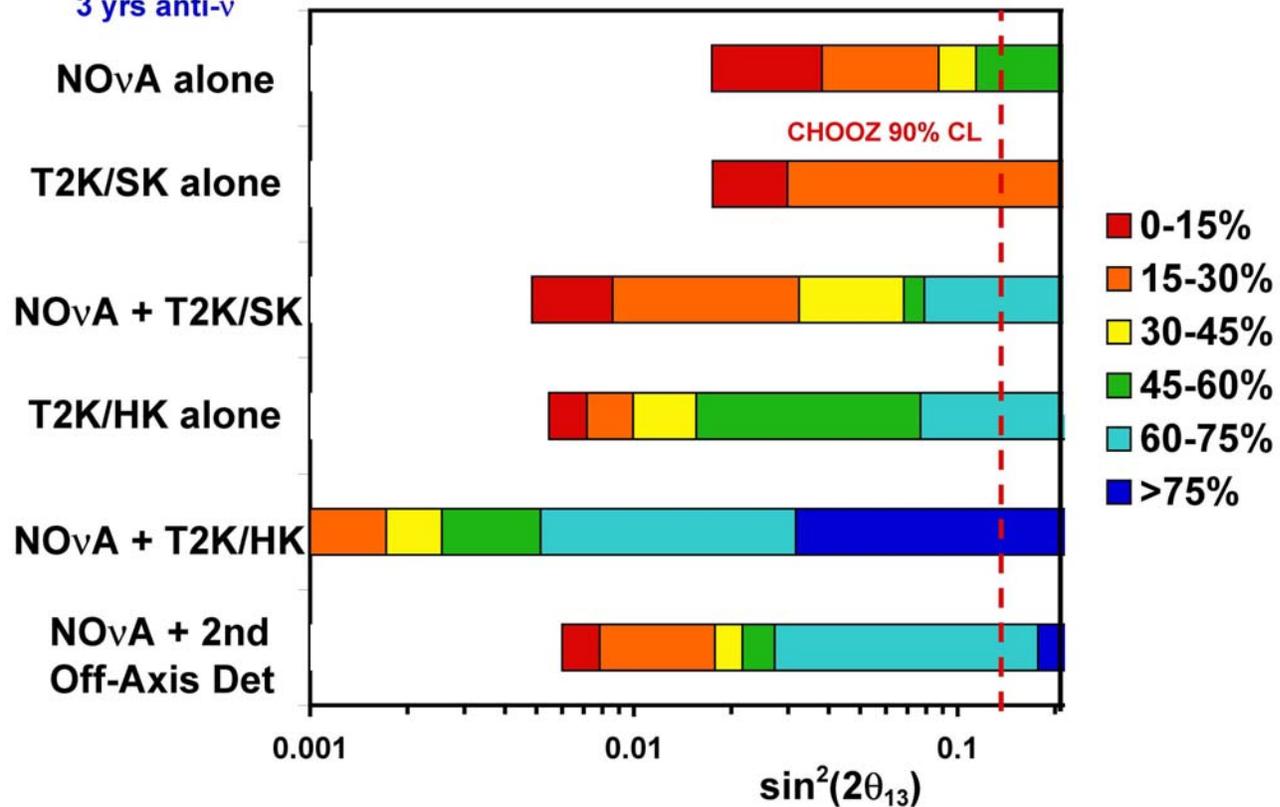
With adequate data, can distinguish mass hierarchy for some range of δ_{CP} , depending on strength of oscillation

To make progress on CP violation requires all the handles.

3 σ Determination of CP Violation

3 yrs ν and
3 yrs anti- ν

In all cases NO ν A with PD and T2K with 4 MW



Need lots of protons ==> upgrades of beams

Great help from combining experiments!

The Next Stages

- The angle θ_{13} will be measured by the next generation of experiments
 - NoVA and T2K
 - Reactor Experiments
- For many values of δ , mass hierarchy will be resolved
 - For θ_{13} bigger than about 0.04
- Clues about CP violation
- Further progress will require additional detectors either
 - Bigger
 - Strategically placed.
- Large effort, potential big payoff

The scientific stakes could not be higher!

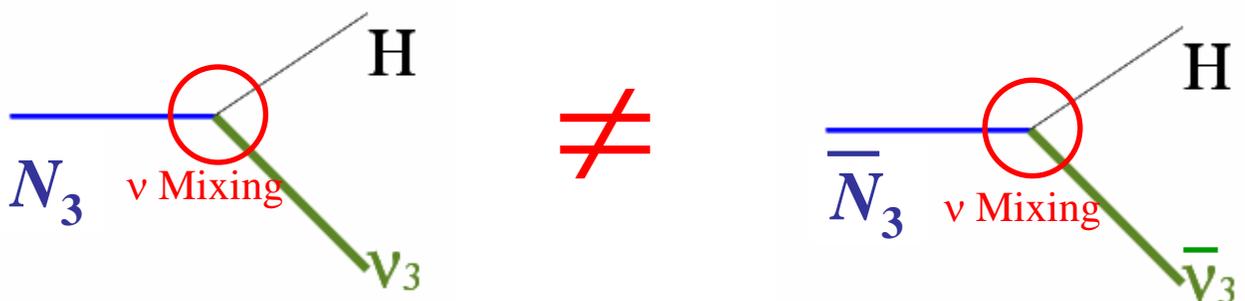
Leptogenesis and Cosmology

Universe has essentially no antimatter
BUT

Cosmology requires more CP-violation than the standard model provides to make this happen.

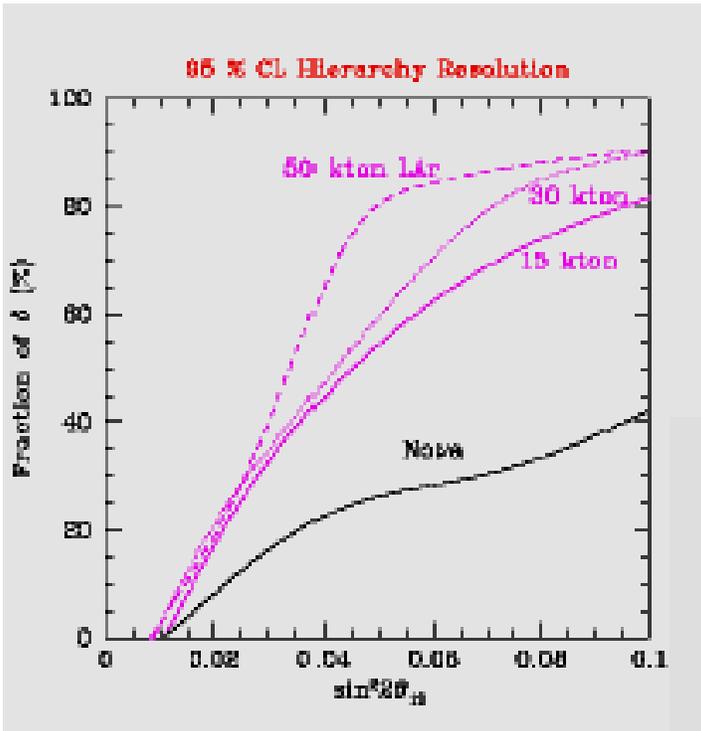
QCD can convert a lepton asymmetry to a quark one if there is enough CP violation in the neutrinos!

A heavy-mass partner of the neutrino is the most likely contributor. Recall see-saw.

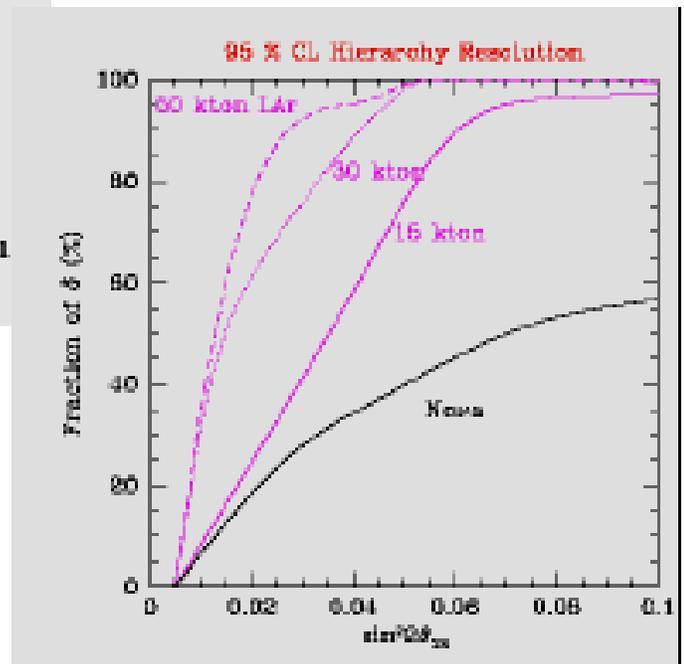


After Inflation at $T \sim 10^{12}$ GeV

Advantages of a second detector



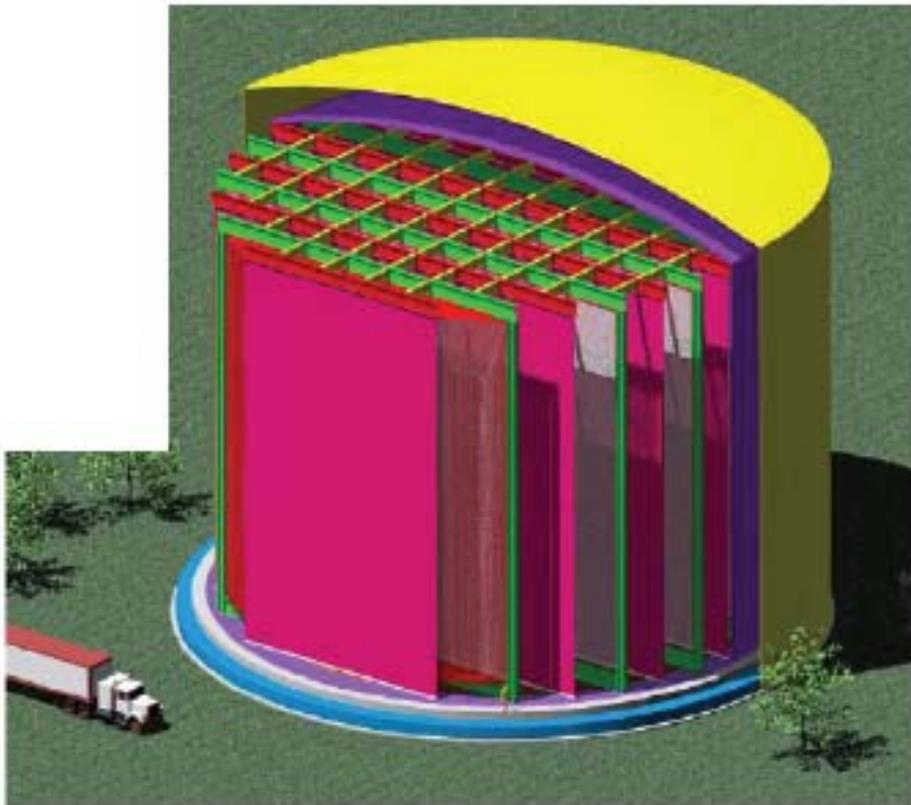
Upgraded Intensity
($25 \times 10^{20}/\text{year}$)
6 years one detector +
4 years both



Nominal Intensity ($6.5 \times 10^{20}/\text{year}$)
6 years one detector + 8
years both

One scenario with a second detector located at
200 km, with 0.7 GeV beam!

*“For best results”, 2nd
detector should be very
efficient.*



One potential technology is large liquid Argon (LAr) TPC

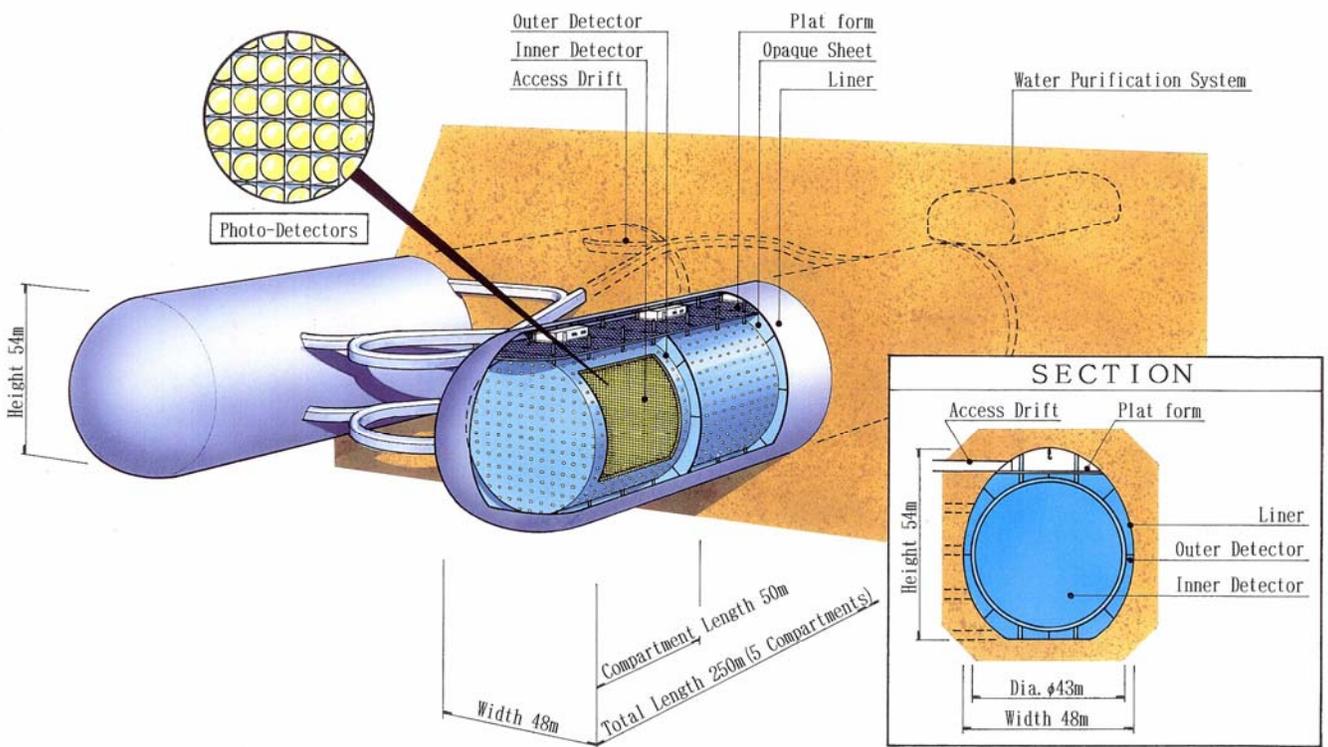
- Industrial size tanks
- Challenges from purity, noise, long drifts, cost
- **Extremely high efficiency (claim ~90%)**

Sidebar - Hyper-K megaton detector

1 MTON Water Cherenkov (like Super-K)

0.54 MTON fiducial volume

200 K PMT's



Tunnel shape cavity helps excavation and optimizes detector performance

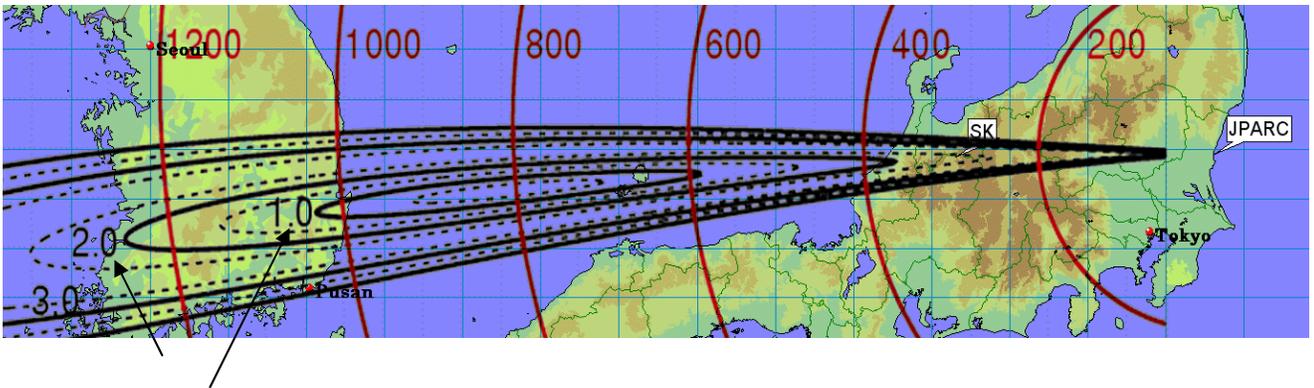
< 50 m deep for PMT's

Light path < 100 m

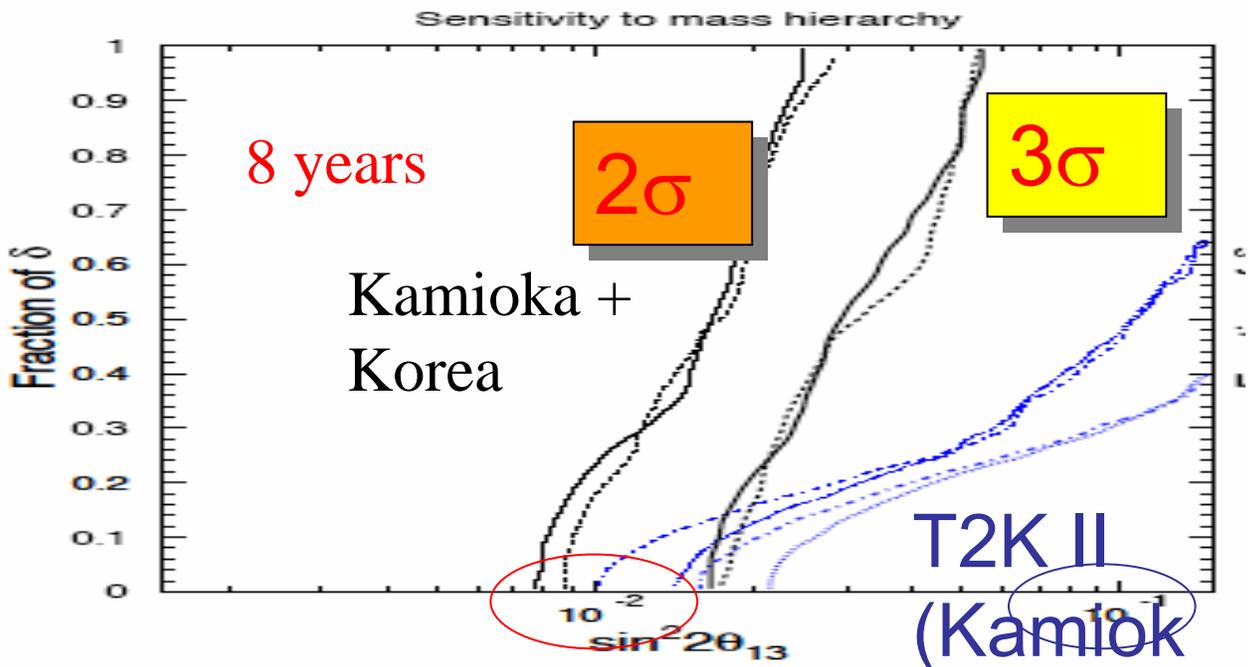
Dual cavities allows for staged construction, maintenance

Potential for Korean 2nd Detector

Split fiducial mass in 2 pieces (0.27 Mton each)



Off-axis angle (degrees) in 3 dimensions

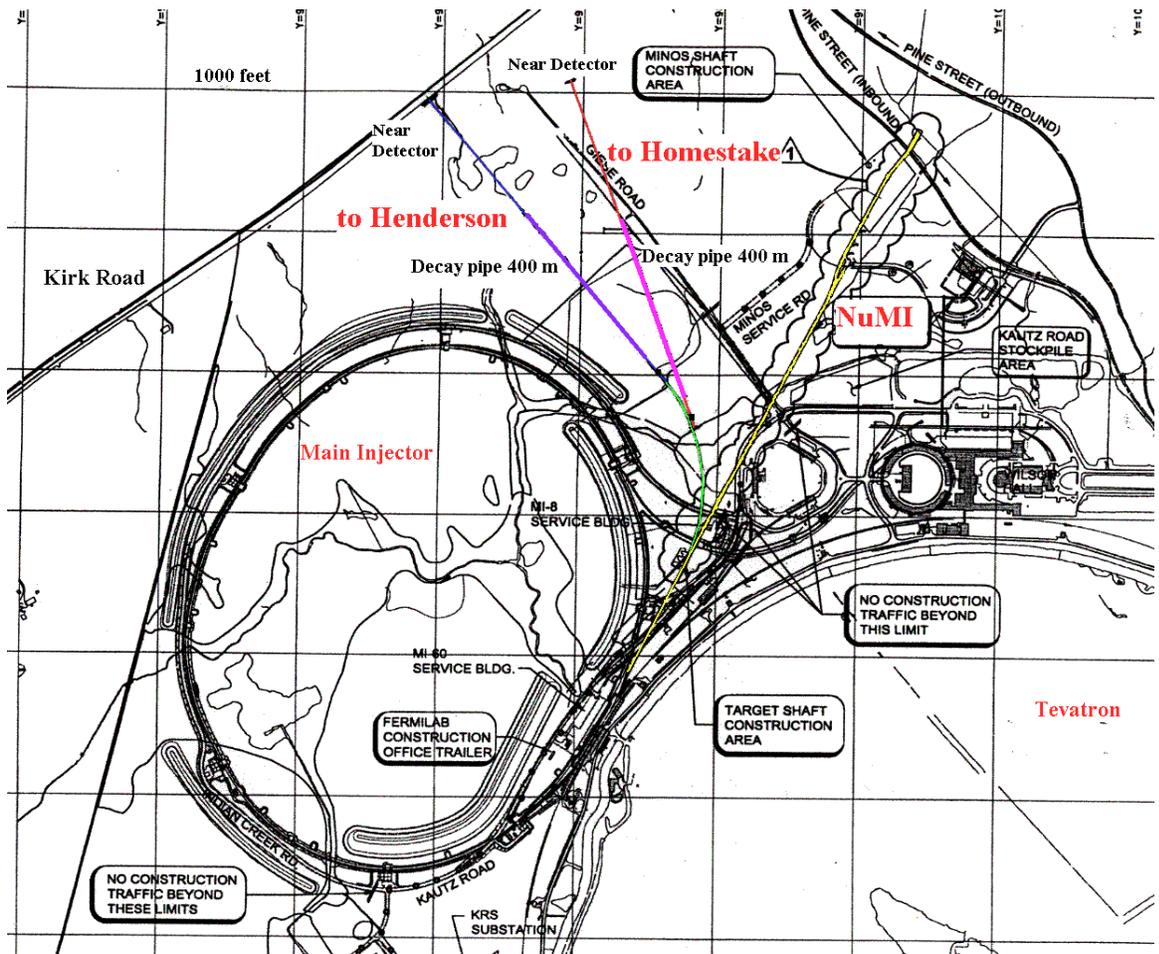


Gives matter effect so that mass hierarchy can be studied

T2K II
(Kamiok
a)

hep-ph/0504026

Can a beamline to a large DUSEL detector fit at Fermilab?

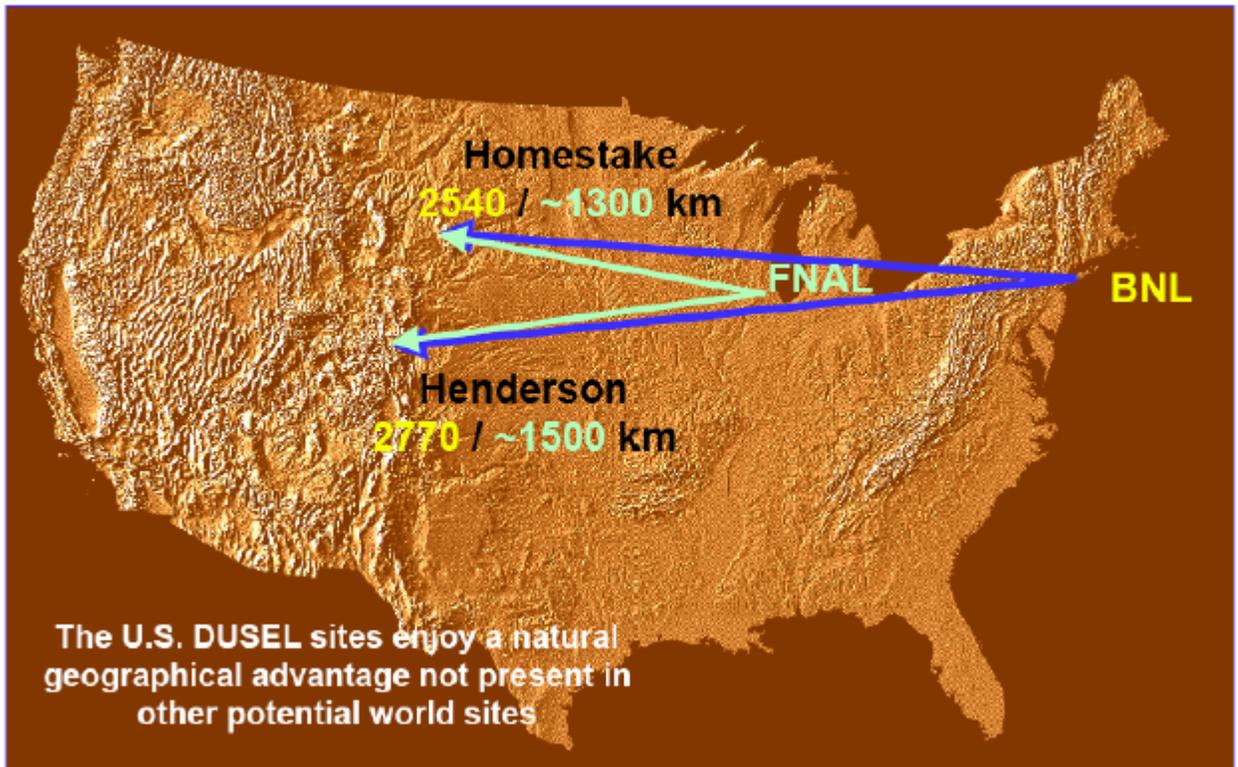


400 m decay pipe for use with low energy beam

Homestake (SD) 1289 km, ~ -6 degrees

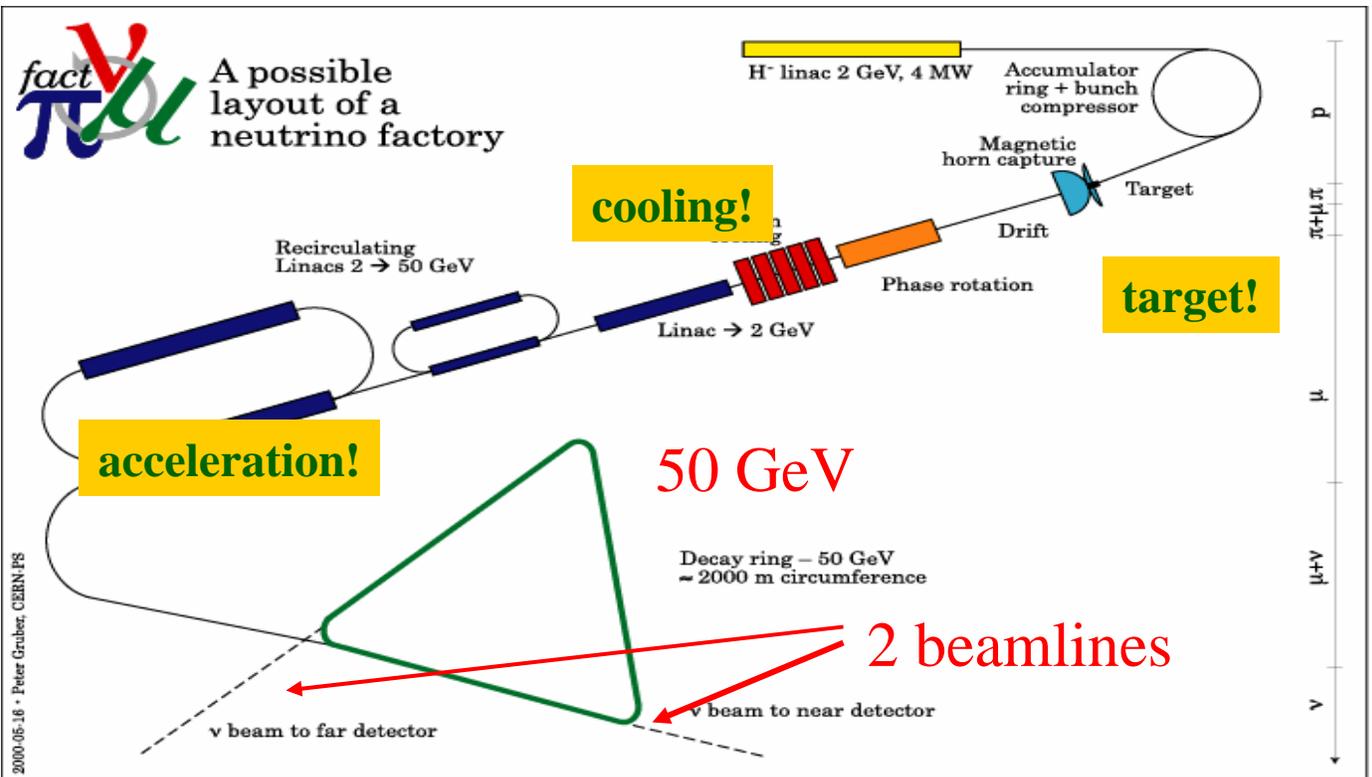
Henderson (CO) 1495 km, ~ -7 degrees

Preferred DUSEL sites in the United States



Ideal for long-baseline experiments
Focus now is on Fermilab as host
Lots of challenges, but lots of interest.

Advanced concepts- Neutrino Factory



Based on decays of stored muons.

After Blondel

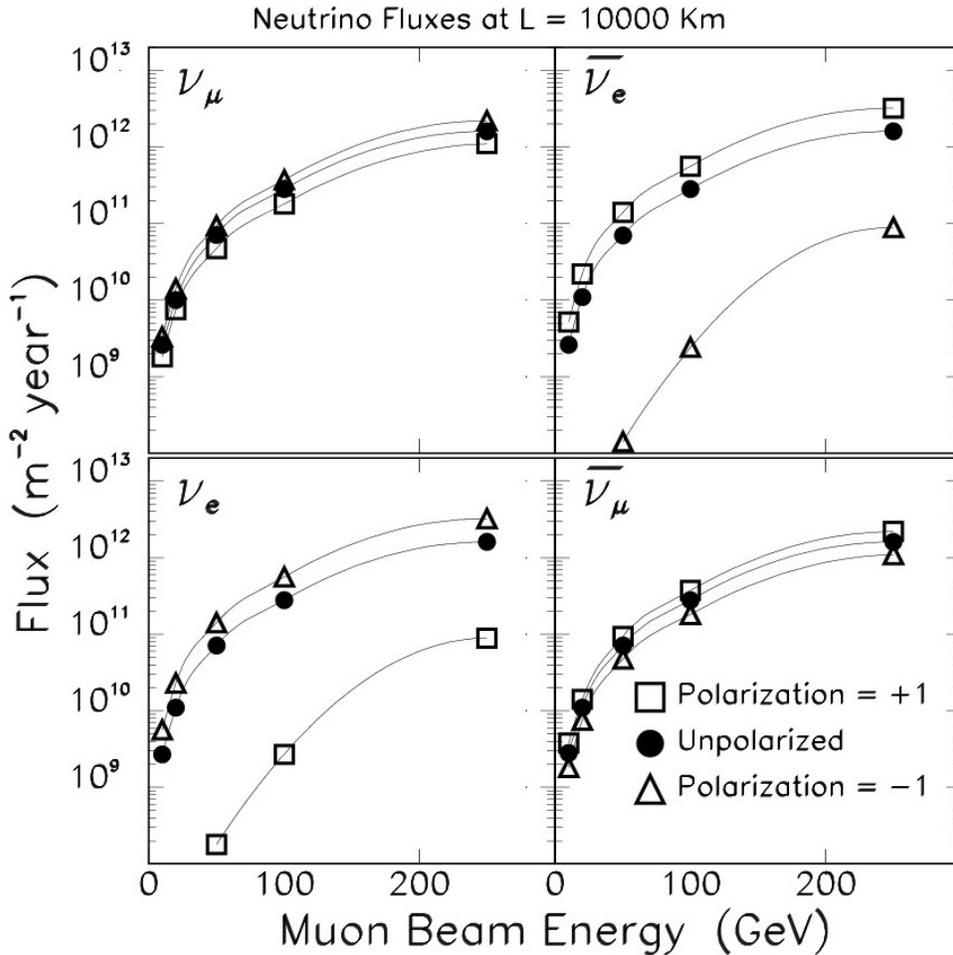
$$\mu^- \rightarrow e^- + \nu_\mu + \bar{\nu}_e$$

$$\mu^+ \rightarrow e^+ + \bar{\nu}_\mu + \nu_e$$

Muon cooling schemes adopted from μ -collider designs.

Designs produce $1-5 \times 10^{20}$ μ decays per year

Neutrino Factory Fluxes



In forward direction

$$F_{\nu_{\mu}}(x) \propto E_{\mu}^2 x^2 \left[(3 - 2x) + P_{\mu} (1 - 2x) \right]$$

$$F_{\bar{\nu}_e}(x) \propto E_{\mu}^2 x^2 \left[(1 - x) + P_{\mu} (1 - x) \right]$$

where $x = \frac{E_{\nu}}{E_{\mu}}$ and P is the μ polarization

Statistical power of Neutrino Factory

5 Years data taking

$$\sin^2 2\theta_{13} = 0.1 \quad \sin^2 2\theta_{13} = 0.01$$

Calculations of W. Winter

Expt	Signal	Bkg	Signal	Bkg
MINOS	49.1	108	6.7	109
ICARUS	31.8	69.1	4.5	70.3
OPERA	11.2	28.3	1.6	28.6
T2K	132	22.7	16.9	23.5
NOvA	186	19.7	23.0	20.7
NOvA+FPD	716	75.6	88.6	79.5
NuFact nu	29752	44.9	4071	44.9
NuFact nubar	7737	82.0	1116	82.0

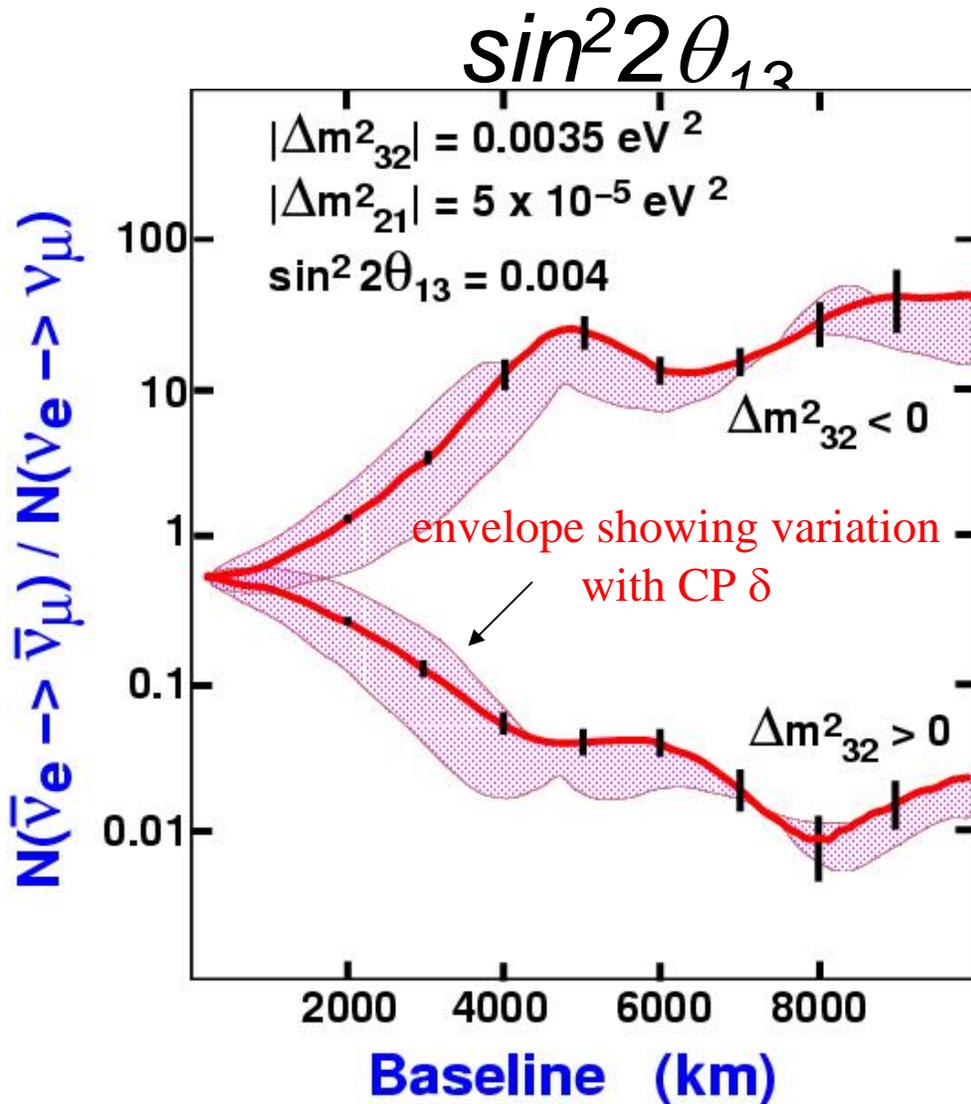
Normal Hierarchy, $\delta = 0$

NUFACT: Beam = 3×10^{20} decays/yr,

E = 50 GeV, $M_{\text{det}} = 100$ kt, baseline 7300 km

Note this is a large magnetized detector
because signal is “wrong-sign” μ

Neutrino Factory is sensitive to very low

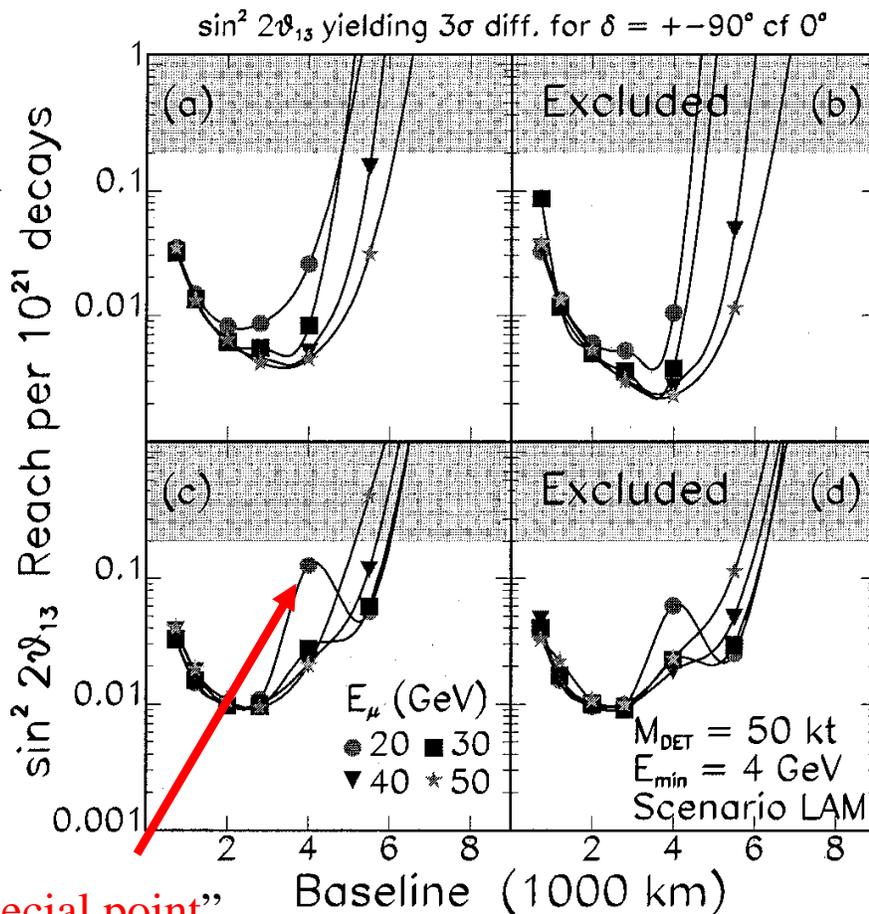


Ratio of oscillation rates from μ^-
and μ^+ vs. baseline

Conditions: 10^{20} μ from 20 GeV
 ν -factory, 50 kT detector

Interplay of mixing strength and baseline for CP determination at a ν -factory

$\sin^2 2\theta_{13}$ reach vs. baseline
for CP discrimination



$0, \pi/2$
distinct

$0, -\pi/2$
distinct

1st “special point”

normal

inverted

Conditions: 10^{21} μ from 20 GeV
 ν -factory, 50 kT detector

Practicalities - cooling in two directions

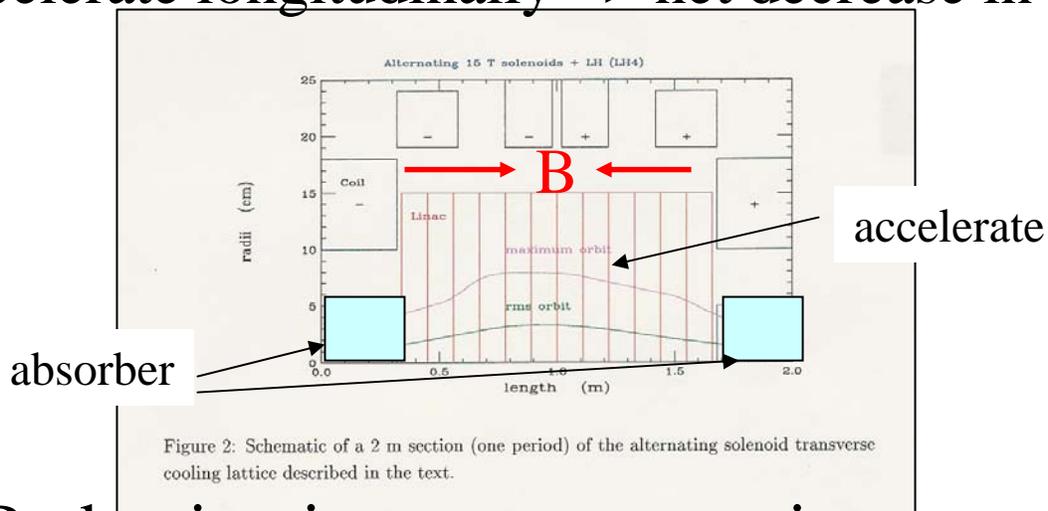
μ^+ and μ^- from decays naturally require cooling both in P_T and P_L

Current thinking centers on *ionization cooling*.

TRANSVERSE COOLING

Reduce momentum uniformly in all coordinates by ionization in a liquid medium (liquid H_2).

Re-accelerate longitudinally \rightarrow net decrease in P_T/P_L



Reduction in transverse emittance ϵ_T

Note $\epsilon_T \sim$ Area of phase ellipse in $x, \text{ angle space}$

Longitudinal Cooling

Solenoids

Absorbers

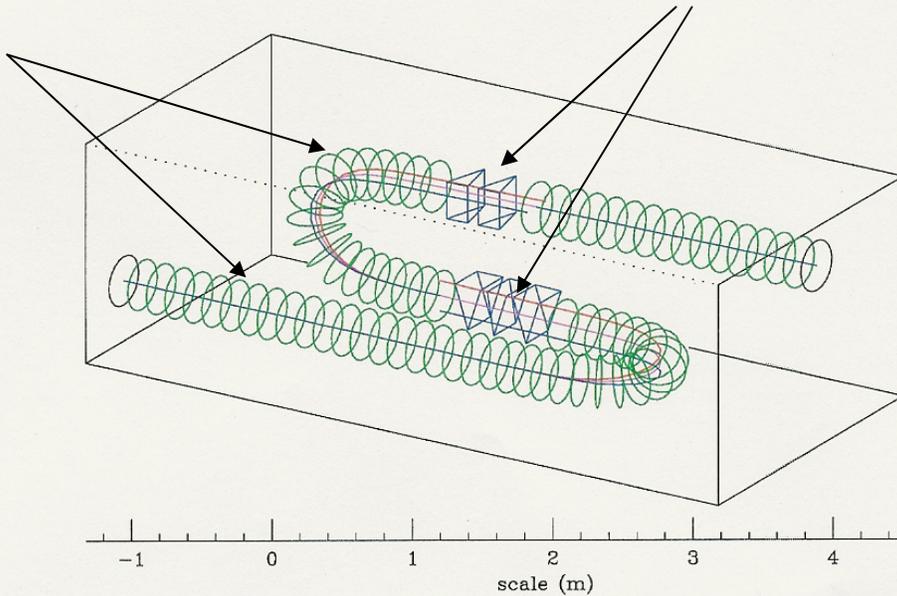
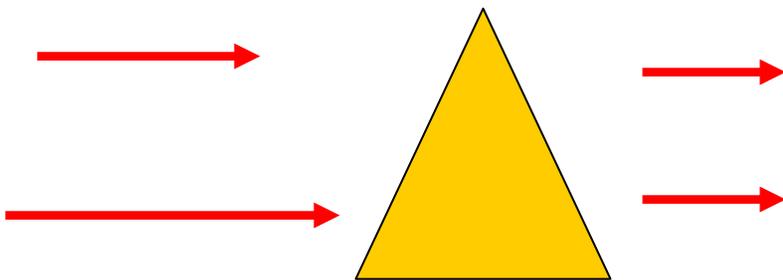


Figure 3: Schematic of the bent solenoid longitudinal emittance exchange section.

Curved solenoids introduce dispersion, e.g.

$$y = f(E)$$



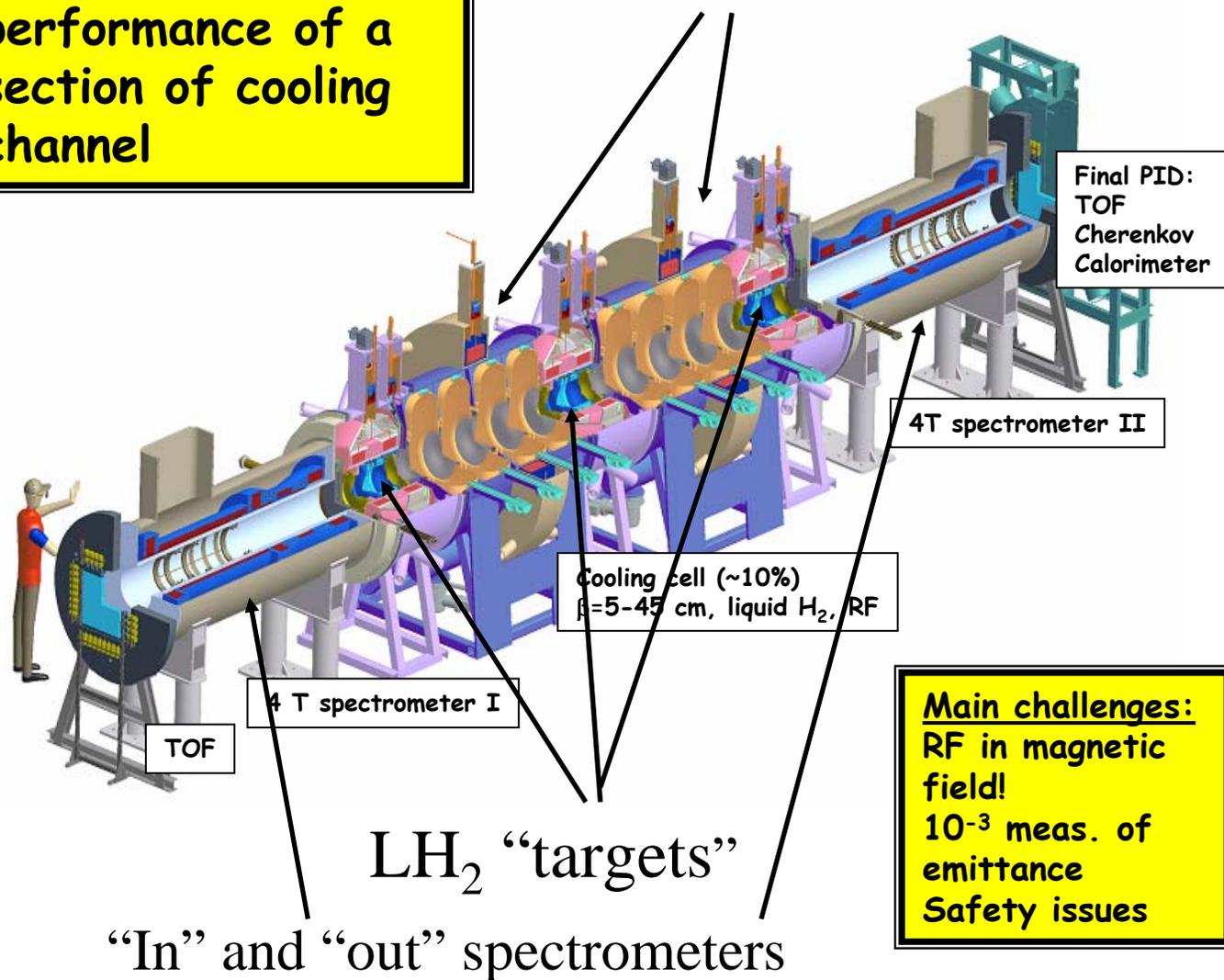
High energy sees more material than low
=> spread reduced

MICE Experiment at RAL



Aims: demonstrate feasibility & performance of a section of cooling channel

RF Cavities



Main challenges:
RF in magnetic field!
 10^{-3} meas. of emittance
Safety issues

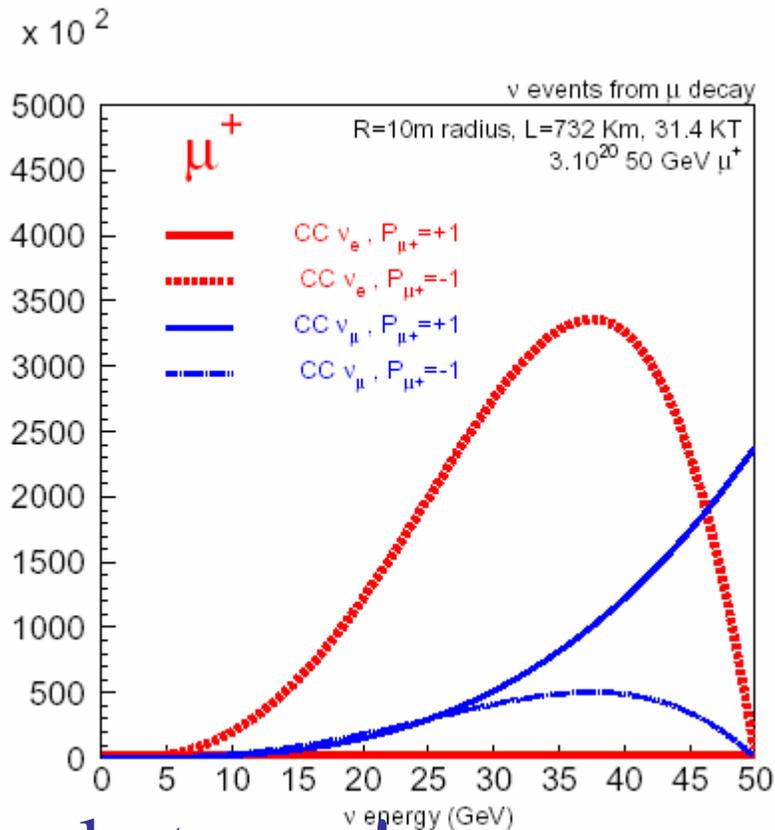
Build a prototype cooling channel

Cool 200 MeV beam by 10%

Useful or just for fun?

A.Blondel, “Muon Polarisation in the Neutrino Factory”

By switching from negative muons with 50% negative polarisation to positive muons with 50% negative polarisation, one can change the ratio of CC ν_e to CC ν_μ by a factor 20, this ratio is only 5 in absence of polarisation. This must be useful.



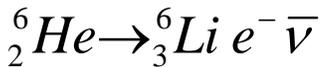
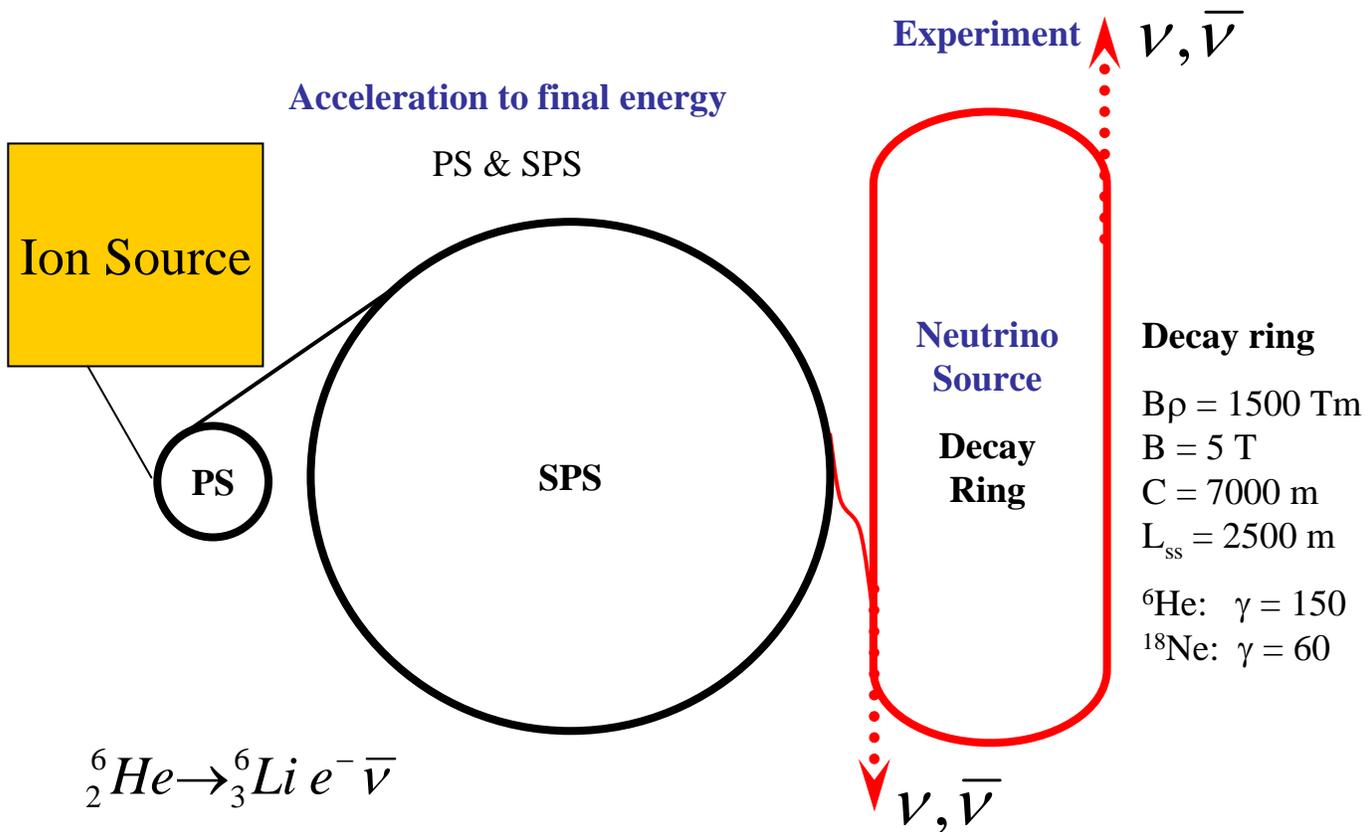
Too early to say!

Beta-beams - an alternative source of clean

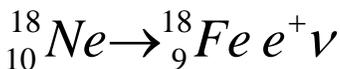
ν_e and $\bar{\nu}_e$ beams

Acceleration

Neutrino source



$$\text{Average } E_{cms} = 1.937 \text{ MeV}$$



$$\text{Average } E_{cms} = 1.86 \text{ MeV}$$

Need high γ because energy low.

CERN concept
(for Frejus, 130 km)

(from Lindross, NUFACT 05)

About 10^{18} decays/year

Any relevance of this β -beam concept for Fermilab?

Extraordinarily difficult if it can be done at all

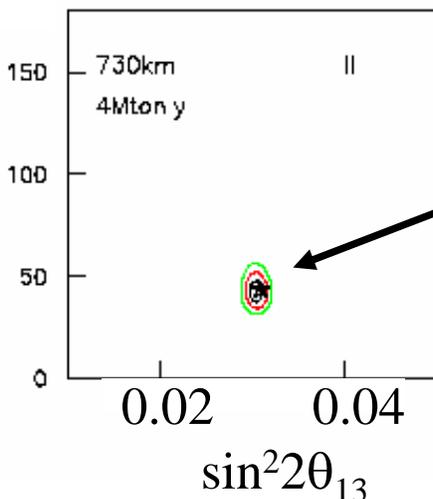
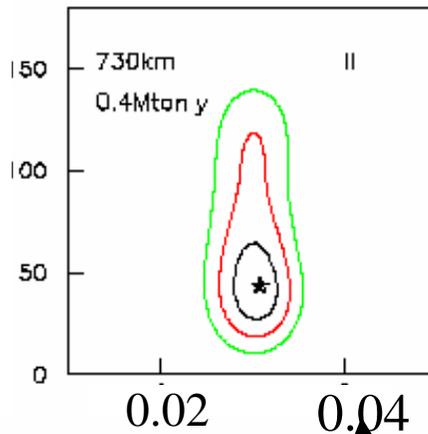
Build ion source

Rigorous control of losses (quenching)

Tunnel activation

Discussed in literature, conferences

Distinguish δ from $(0, \pi)$



$\sin^2 2\theta_{13}$

Soudan distance with 2 different (large) exposures

High γ gives an advantage that improves competitiveness with combination of conventional “superbeam” experiments

Conclusions – Neutrinos are Critical Physics

- Since about 1998, the neutrino has been demanding extension of the Standard Model.
- Many experiments feeding on the new phenomena already.
- Many more in the pipeline.
- Lots of basic physics, lengthy hard work required, geological distances

- All this must be telling us something important!

BACKUP SLIDES

Final Sidebar - Polarization at a neutrino factory

Polarization is defined as $\max \langle \sigma \cdot e \rangle = P$

(e is direction of polarization vector)

In the reaction $\pi^+ \rightarrow \mu^+ \nu_\mu$

the muon polarization is $-v/c = -P_\mu^*/E_\mu^* = -0.27$

Spin rotation in the magnetic and electric fields of an accelerator has been shown to decrease this to $\sim 18\%$, unless special steps are taken.

Why interesting? Recall that

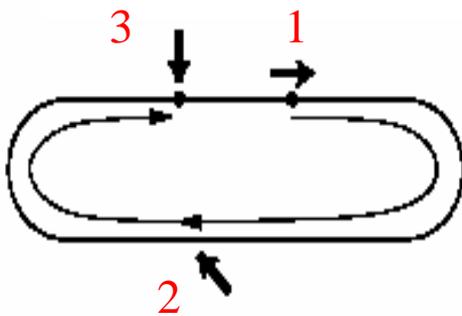
$$F_{\bar{\nu}_e}(x) \propto E_\mu^2 x^2 \left[(1-x) + P_\mu (1-x) \right]$$

If $P_\mu = 1$, the flux of $\bar{\nu}_e$ vanishes!

Preserving polarization at a storage ring

Spin tune depends on magic energy (principle of g-2 experiment)

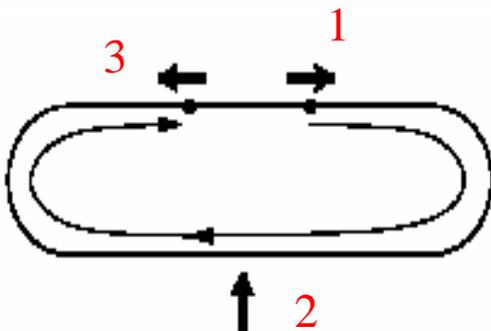
$$\nu = a_{\mu} \gamma = \frac{g_{\mu} - 2}{2} \frac{E_{\text{beam}}}{m_{\mu}} = \frac{E_{\text{beam}} (\text{GeV})}{90.6223(6)}$$



E = 22.656 GeV

“Normal”

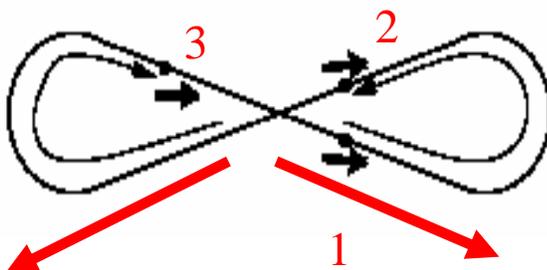
$\nu = 1/4$



E = 45.311 GeV

“Reversing”

$\nu = 1/2$

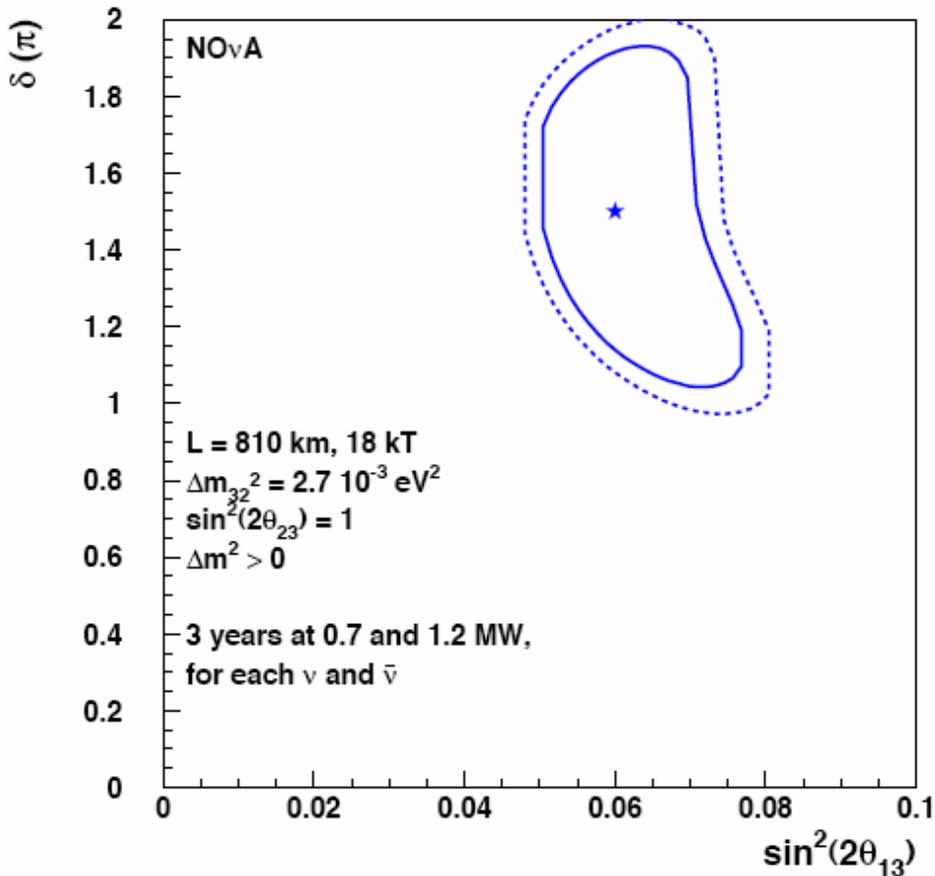


Any Energy

“Bowtie” = 2
beams

$\nu = 0$

NOvA Coverage range for mass ordering



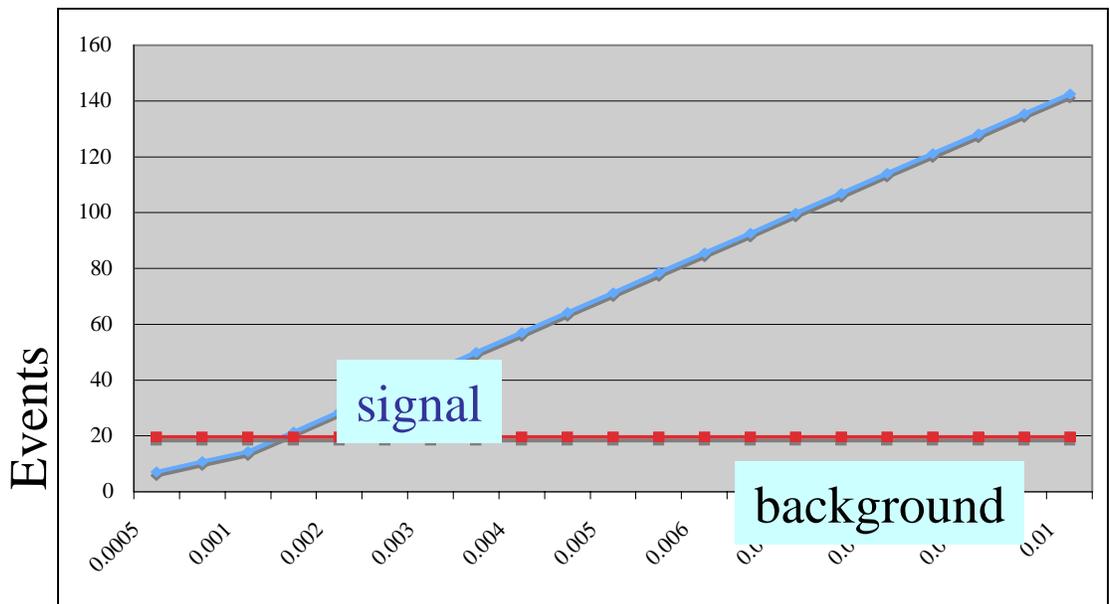
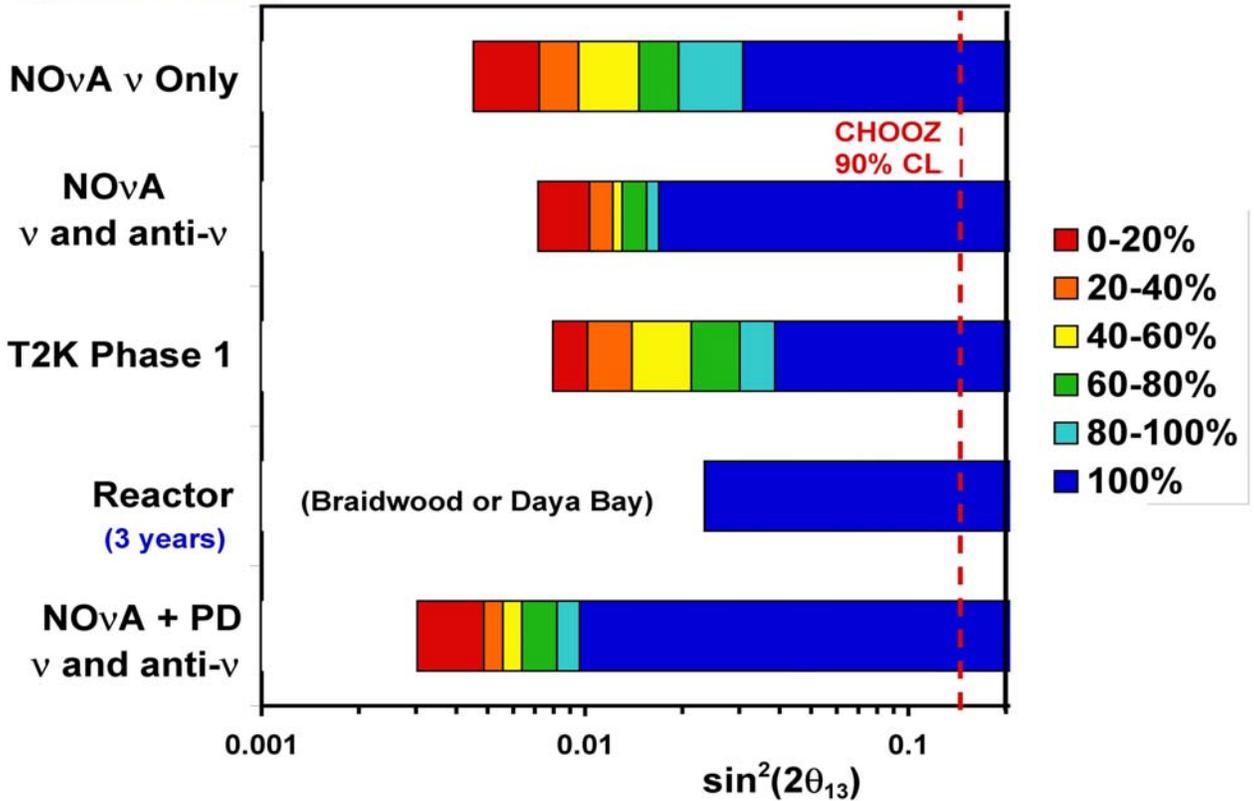
~~Fermilab~~ Upgraded Fermilab complex could produce more beam, on order 1-2 MW

Best results from worldwide combinations

NOvA Sensitivity to θ_{13}

5 years of running

3 σ Discovery Limits for $\theta_{13} \neq 0$



0.01

0.05

0.1

$\sin^2 2\theta_{13}$

Sidebar - magic baseline
removes degeneracies
(that second “special” point)

$$\begin{aligned}
 P_{\text{app}} &\simeq \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2[(1 - \hat{A})\Delta]}{(1 - \hat{A})^2} \\
 &\pm \alpha \sin 2\theta_{13} \xi \sin \delta_{\text{CP}} \sin(\Delta) \frac{\sin(\hat{A}\Delta) \sin[(1 - \hat{A})\Delta]}{\hat{A} (1 - \hat{A})} \\
 &+ \alpha \sin 2\theta_{13} \xi \cos \delta_{\text{CP}} \cos(\Delta) \frac{\sin(\hat{A}\Delta) \sin[(1 - \hat{A})\Delta]}{\hat{A} (1 - \hat{A})} \\
 &+ \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2},
 \end{aligned}$$

$$\Leftrightarrow \sin(\hat{A}\Delta) = 0$$

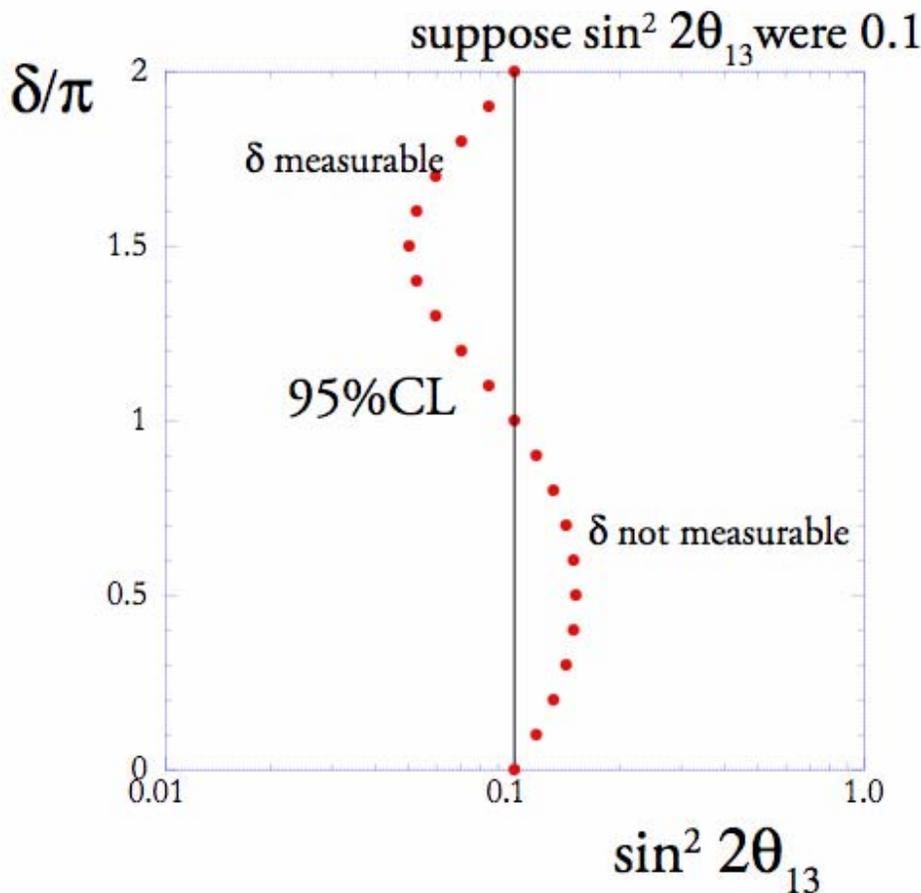
$$\Leftrightarrow \sqrt{2}G_F n_e(L)L = 2\pi$$

$$\Leftrightarrow L \sim 7500 \text{ km}$$

Removes CP dependence at this baseline, regardless of conditions.

Best for $\sin^2 2\theta_{13} < 0.01$, where CP degeneracies are largest

Can discuss measurements of
quantities in terms of what
fraction of CP-space they can
see



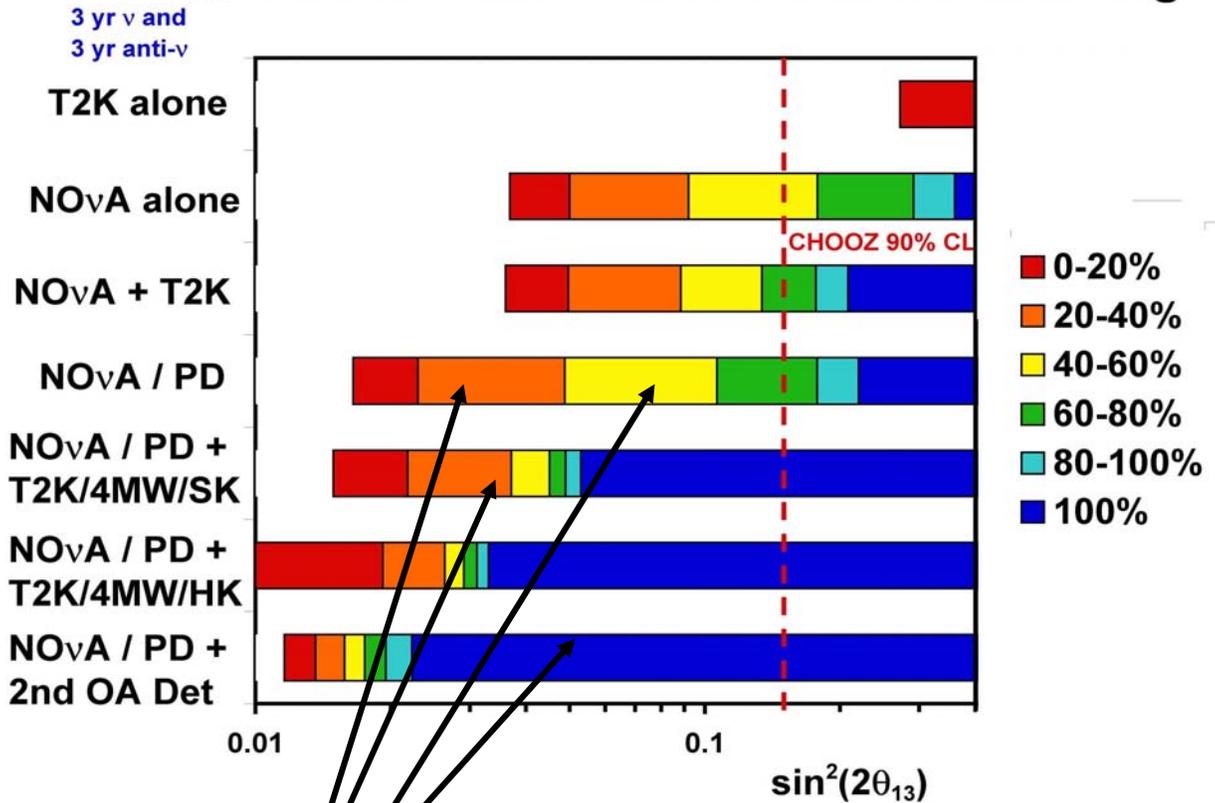
This is just a guide-the-eye example to illustrate the following plots.

In reality, overlaps, degeneracies can make situation quite complicated.

Graphic - after R. Bernstein

NOvA Coverage range for mass ordering

95% CL Determination of the Mass Ordering



Upgraded Fermilab complex could produce more beam, on order 1-2 MW

Best results from worldwide combinations

Flux Comparison between Off-Axis and Forward Beams

$$P_T(\pi/2) = P^* \sin(\pi/2) = P^*$$

$$P_L(\pi/2) = \gamma P^* = \frac{1}{2} P_L(0)$$

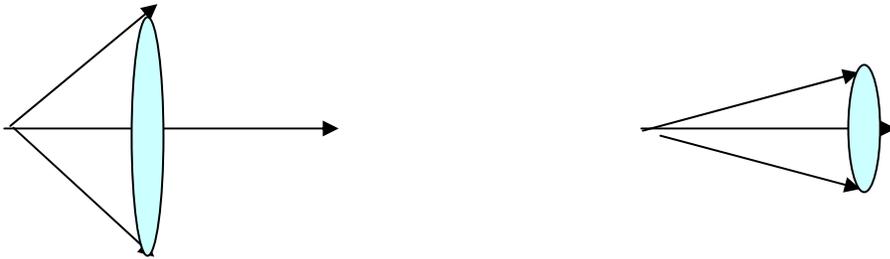
$$\frac{P_L}{P_T} = \gamma \Rightarrow \gamma_{\pi/2} \theta_{lab} = 1$$

Compare solid angle between fully forward beam ($\theta = 0$) and model off-axis beam.

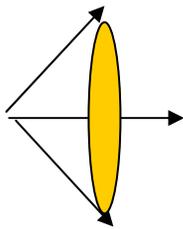
$$\Delta\Omega_{\pi/2} \propto \left(\frac{2\gamma_{\pi/2}}{1 + \gamma_{\pi/2}^2 \theta_{lab}^2} \right)^2 = (2\gamma_0)^2 = \Delta\Omega_0$$

Hence flux can stay constant or even increase in region of interest.

Example of simple detector analysis - μ/e discrimination



Muon - long range (late hits), cone fills in as muon slows.

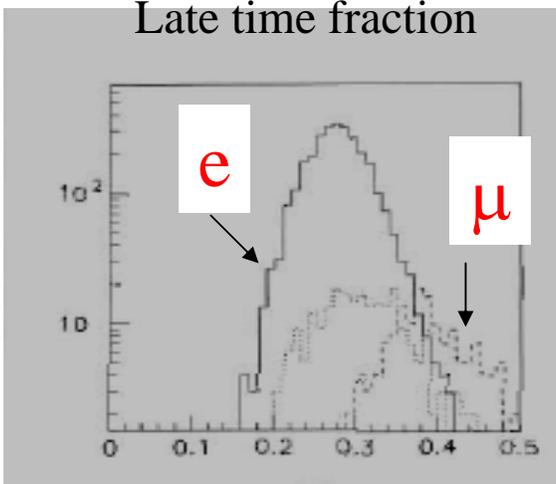


recall Cherenkov angle

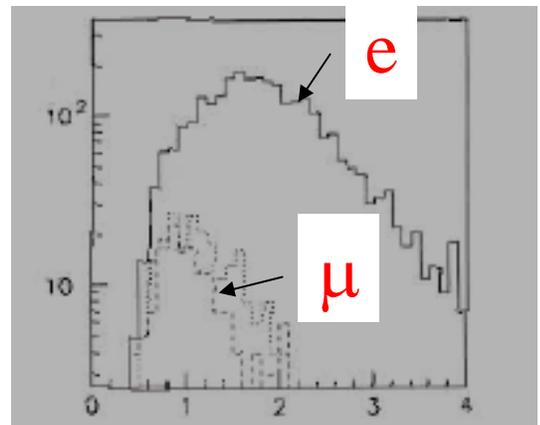
$$\cos \theta = \frac{1}{\beta n}$$

Electron - short range (prompt), most light in outer part

Late time fraction



Large angle light fraction



Note: this is a very old illustration of how this can be done. MiniBoone now uses likelihood analysis techniques



Neutrinos North! The MINOS Project at Fermilab

Presented by Rob Plunkett

Fermi National Accelerator Laboratory

- What's so special about neutrinos?
- Ways and means of small masses
- Neutrinos from a mix
- A glimpse at the data
- A real stretch - the MINOS long baseline experiment
- The truth is out there...soon.

The International Neutrino World

A short summary of the neutrino history

Large History of the neutrino

1898

Discovery of the radioactivity
The first car's races (70 km/h maximum!)

1926

Problem with beta radioactivity
Between two wars, people dance the Charleston and the Boston

1930

Pauli invents the neutrino particle
Crisis of 1929...

1933

Fermi baptizes the neutrino and builds his theory of weak interaction
Hitler gets power in Germany

1956

First discovery of the neutrino by an experiment ν_e
Riots in Budapest. Indochina. Cold war. Atmospheric tests of thermonuclear bombs.

1962

Discovery of an other type of neutrino: the ν_μ
Missiles of Cuba, not far from the End!

1974

Discovery of neutral currents thanks to the neutrinos
Energy crisis... first oil problem

1977

Discovery of the tau (co-particle of ν_τ)
Energy crisis... second oil problem

1985-1989

Pseudo-discovery of a neutrino at 17 keV
The man on the moon 20 years ago. Mutations in European East countries

1990

The mystery of dark matter comes again
Berlin wall finally falls!

1991

LEP experiments show that there are only three light neutrinos
Elsine is the new President of Russia. Iraq invades Koweit

1994

First proclamation of possible neutrinos oscillations seen by LSND experiment

1995

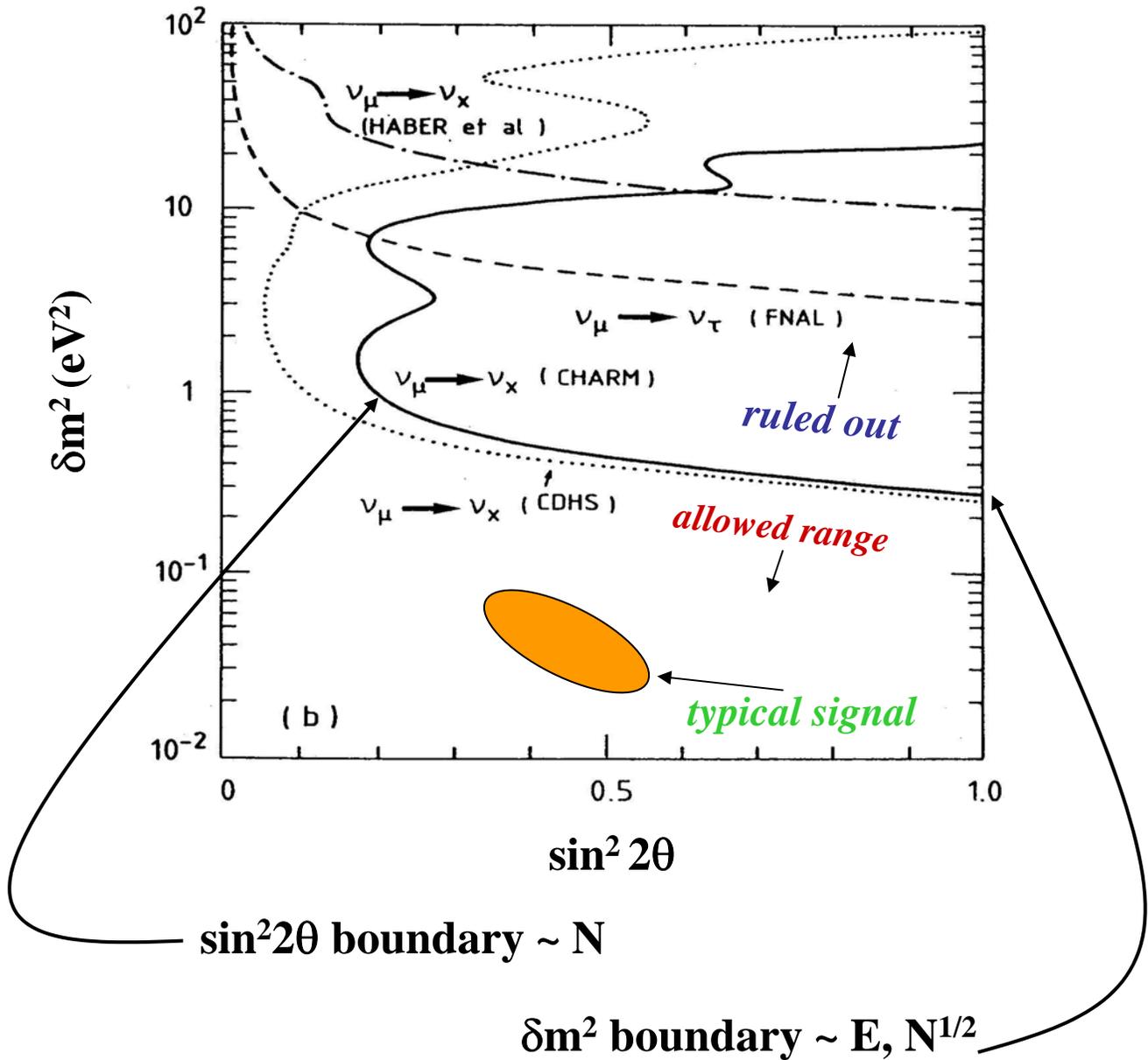
Missing solar neutrinos confirmed by GALLEX
Internet and its WWW explode... and LAPP is 19 years old

1998

May be neutrinos oscillations seen in LSND and Super-Kamiokande
Human rights declaration is 50 years old, 1968 is 30 years old



How to Read an Oscillations Plot



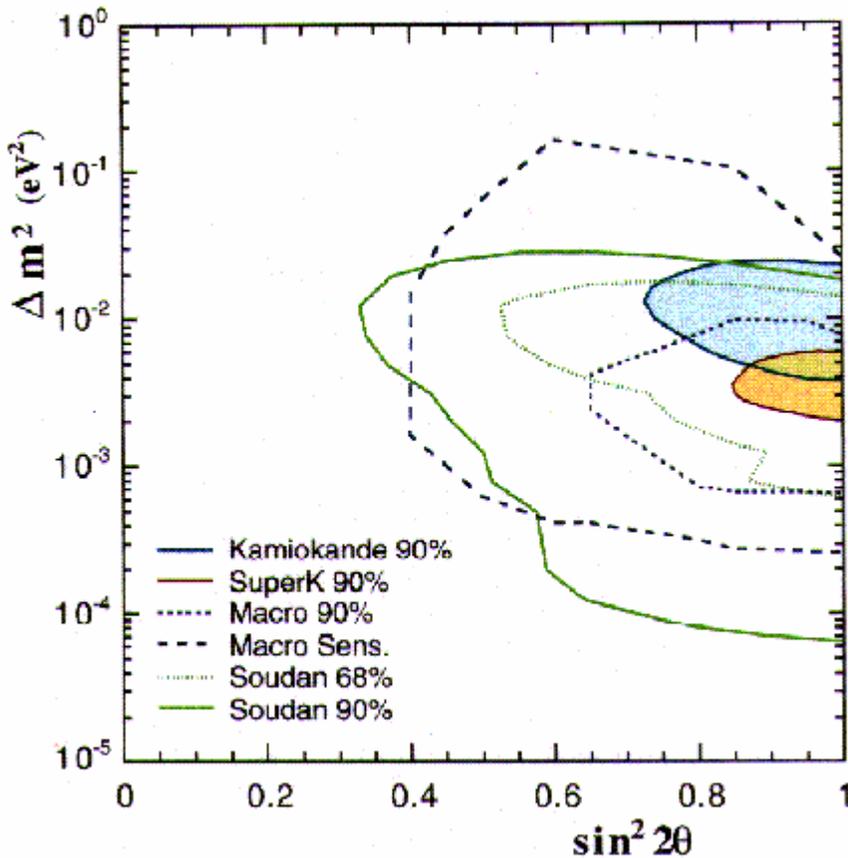
Fits to Oscillation Hypothesis from Atmospheric Experiments

SuperK Best Fit:

$$\Delta m^2 = 3.5 \times 10^{-3} \text{ eV}^2$$

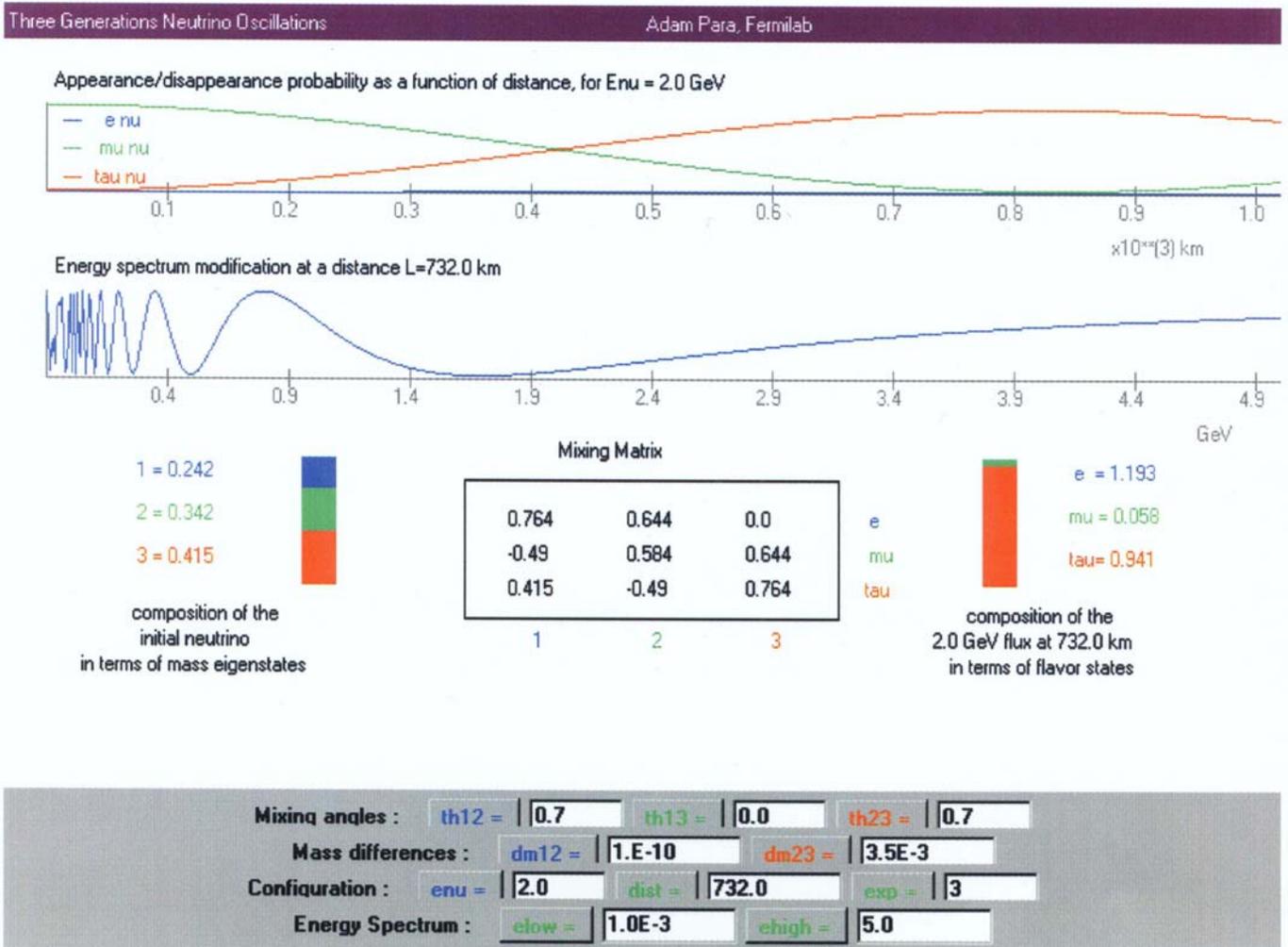
$$\sin^2 2\theta = 1.0$$

$$\alpha = 0.06$$



3 Families of Neutrinos Mixing

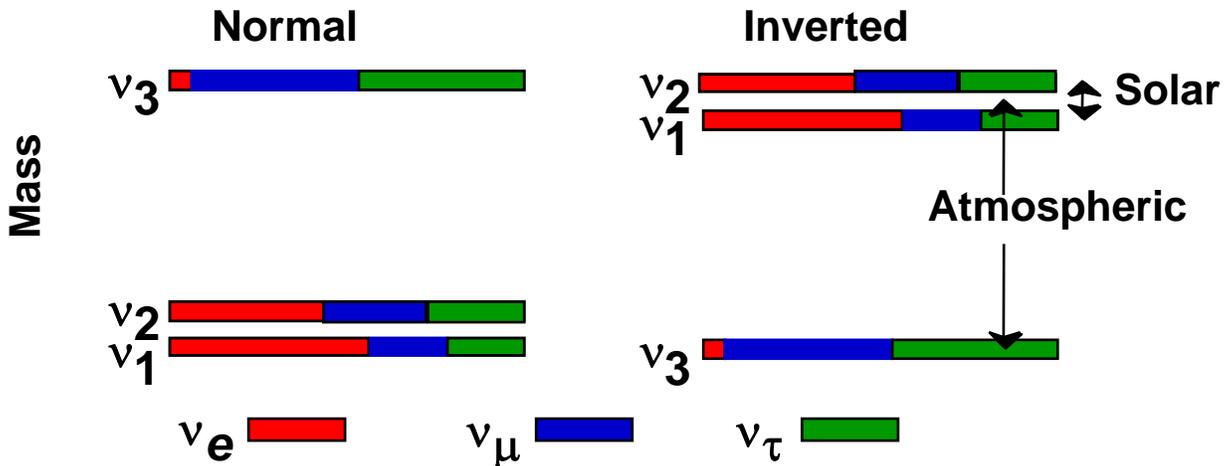
Web-based tool... (courtesy A. Para)



Implementing general model...

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Mass Hierarchy and NOvA



Major unsolved problem

Requires matter effects to observe

Hence a real “target of opportunity” for a long baseline.

To leading order and zero CP violation, acts as if neutrinos and antineutrinos were interchanged.

CP violation makes it more complicated

Review of Oscillations in Matter

Recall $\nu_e + e$ scattering adds a potential

$$H_{matter} = H_{vac} + V_W \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

Leading after some manipulation to

$$P_{\nu_\mu \rightarrow \nu_e} = \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\Delta_{13}^2}{(\Delta_{13} \mp aL)^2} \sin^2(\Delta_{13} \mp aL)$$

where a is

$$a = \frac{G_F N_e}{\sqrt{2}} \approx (3700 \text{ km})^{-1} \left(\frac{\rho}{2.8 \text{ g cm}^{-3}} \right)$$

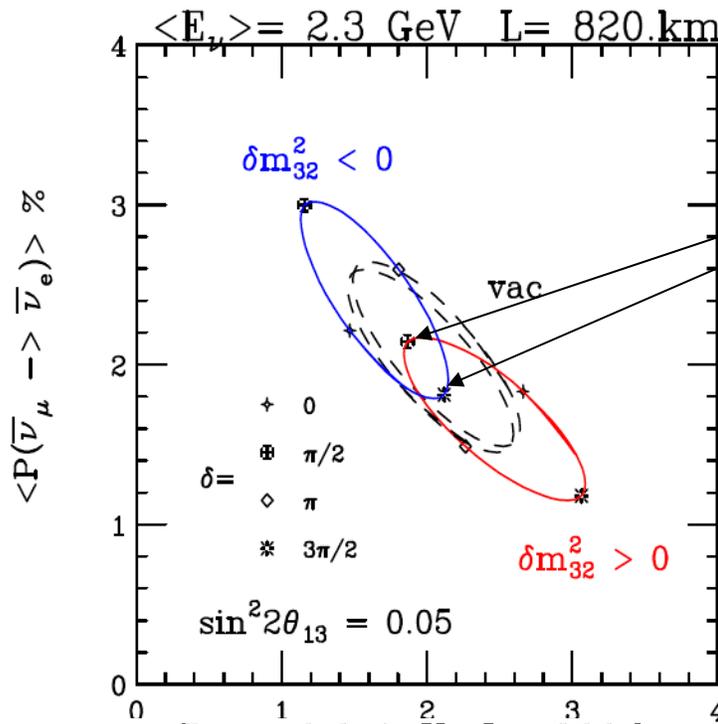
new mixing
strength

new effective
mass

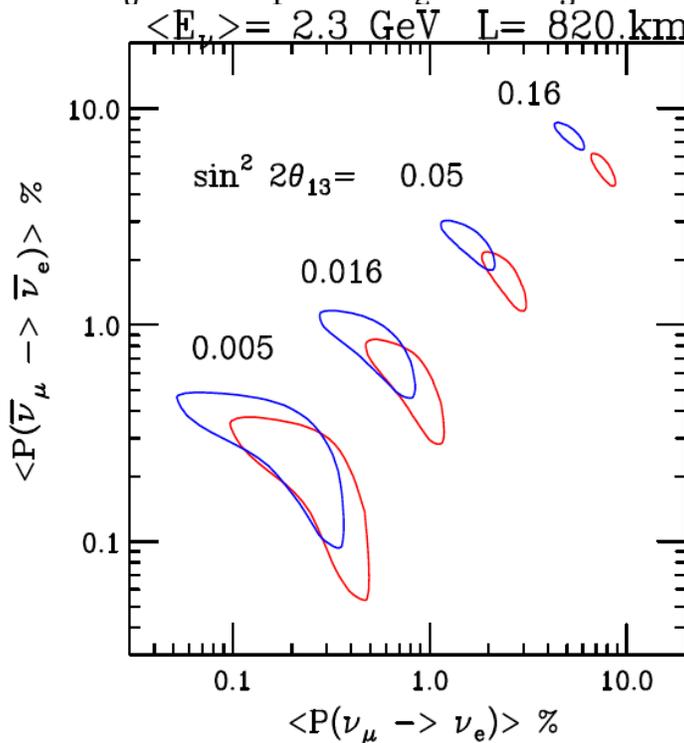
and the sign changes between neutrinos and antineutrinos

3- ν 's: Change is OK for all orders n
of $\sin^n \Delta_{12}$ and $\sin^n \Delta_{13}$

Matter effects split the CP ellipses



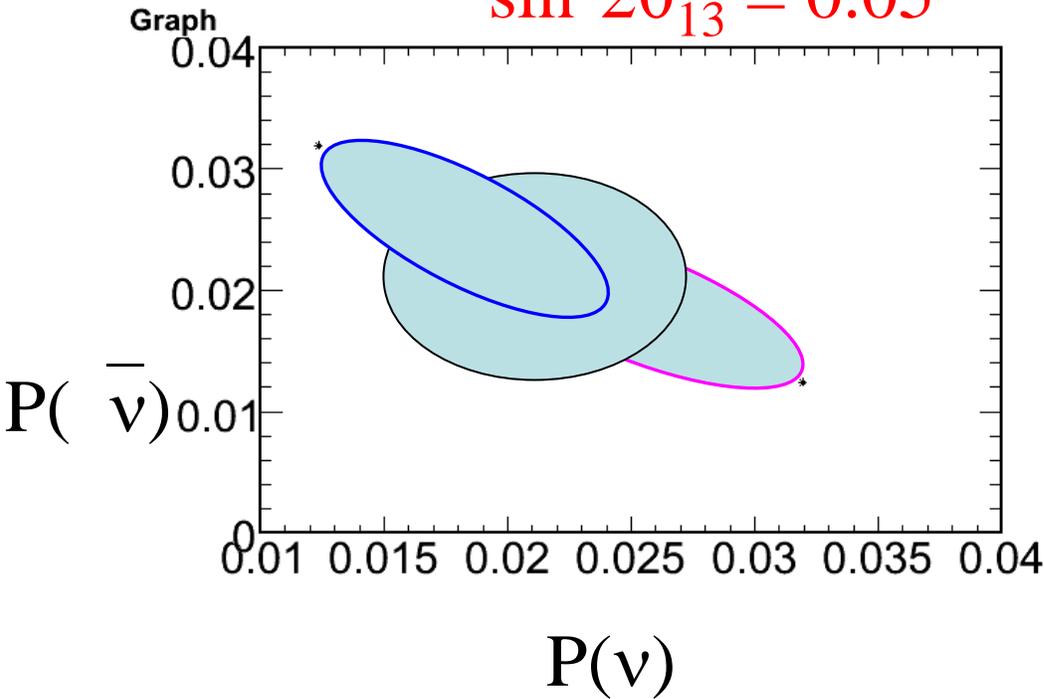
Opposite rotation senses can increase overlap



Ellipse separation is leading order effect, grows with $\sin^2 2\theta$

2- σ data estimate
superimposed on CP-
ellipses

$\sin^2 2\theta_{13} = 0.05$



Scale from NOvA proposal

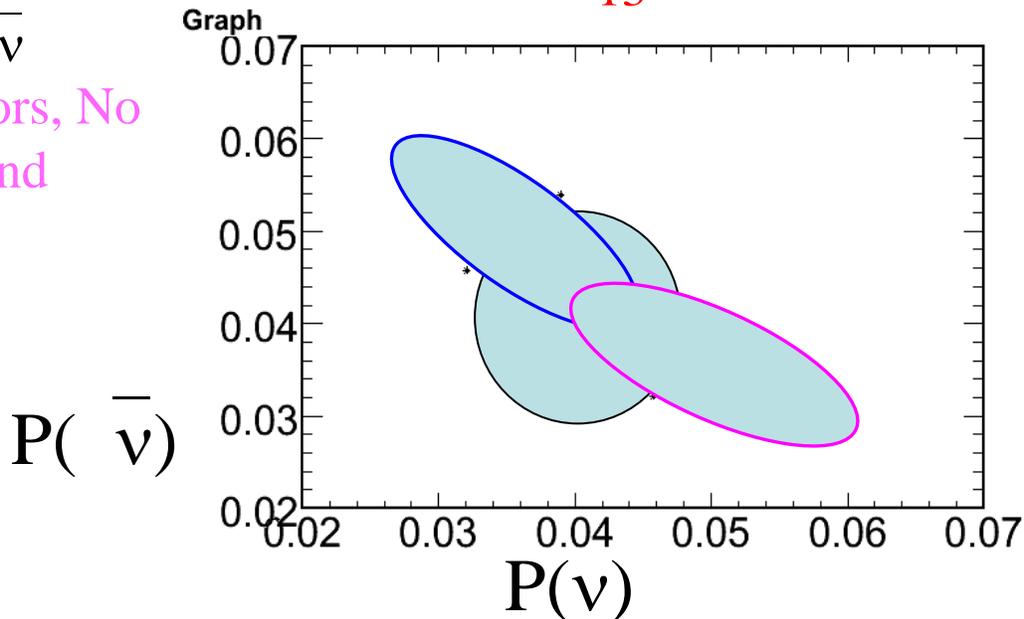
for $\sin^2 2\theta_{13}$

3 years nominal NuMI for

both ν , $\bar{\nu}$

Statistical errors, No
background

$\sin^2 2\theta_{13} = 0.10$



MINOS Beam Event Characteristics

Simple event selections for both detectors.

Far Detector

50 μ s window around beam spill

Reconstructed track within fiducial volume (70% for CC)

Track angle along beam direction.

Data and beam quality cuts (96%)

Only an unknown fraction of the far detector data is used for checks and testing.

Near Detector

Fiducial cuts using track or event vertex for candidate neutral currents

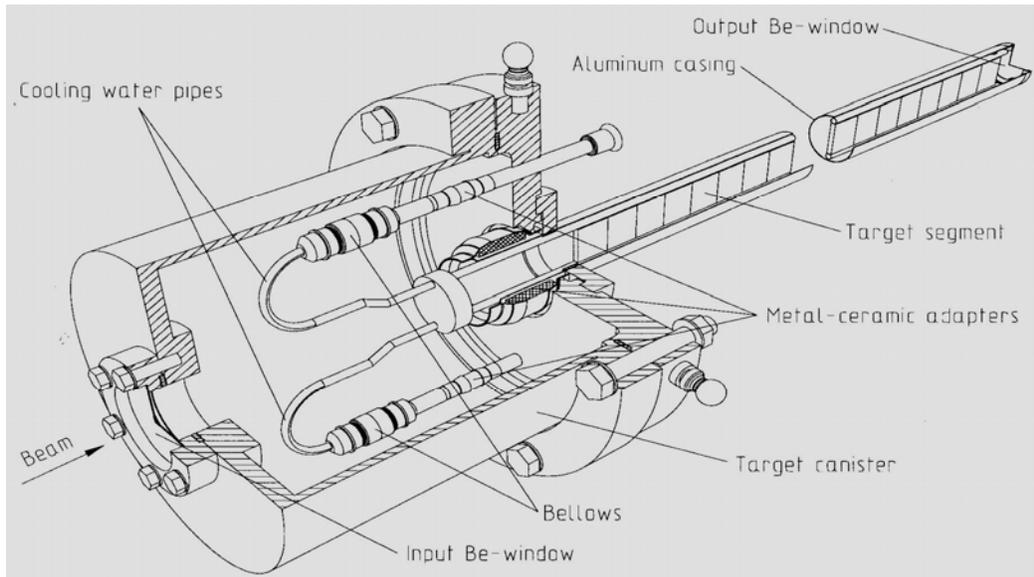
Track quality cuts for events with tracks

Beam quality cuts

A Tanalizing but Scattered Outlook

- Plausible evidence from 3 sources for the oscillation phenomenon
 - No show-stoppers uncovered in any of the channels
- Unfortunately, all point in different directions!
 - $\sum \delta m_i^2 = 0$ by definition
 - Can't accomodate atmospheric, solar, and LSND
- Non-standard physics solution would be required
 - Like a 4th, “sterile” neutrino, ν_s , with no standard model interactions.
- Definitive resolution will require a new generation of accelerator experiments ==> Fermilab

NuMI Target - where the decaying particles start



60 cm graphite fins - water cooled
Be windows at either end

Absorbs 40 KW of power at design intensities.

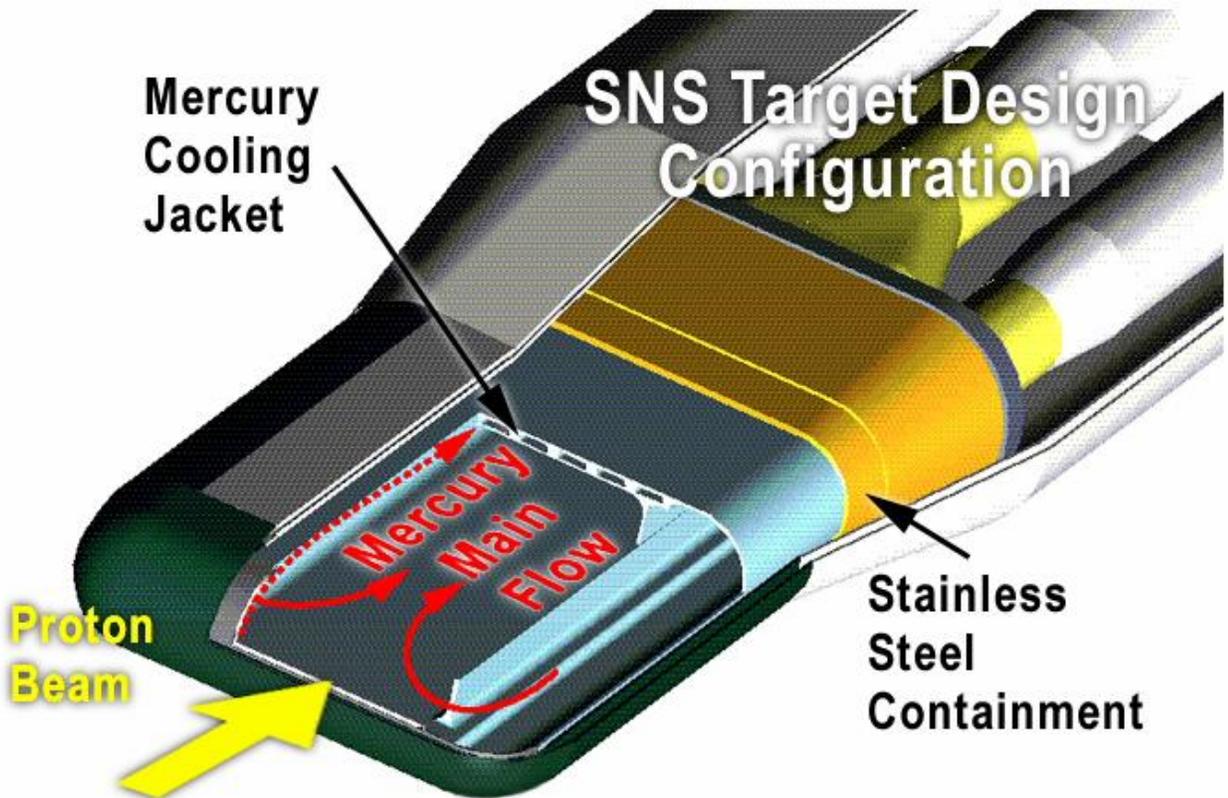
Careful modeling because distribution of particles coming out affects ν flux

Sidebar- advanced targeting

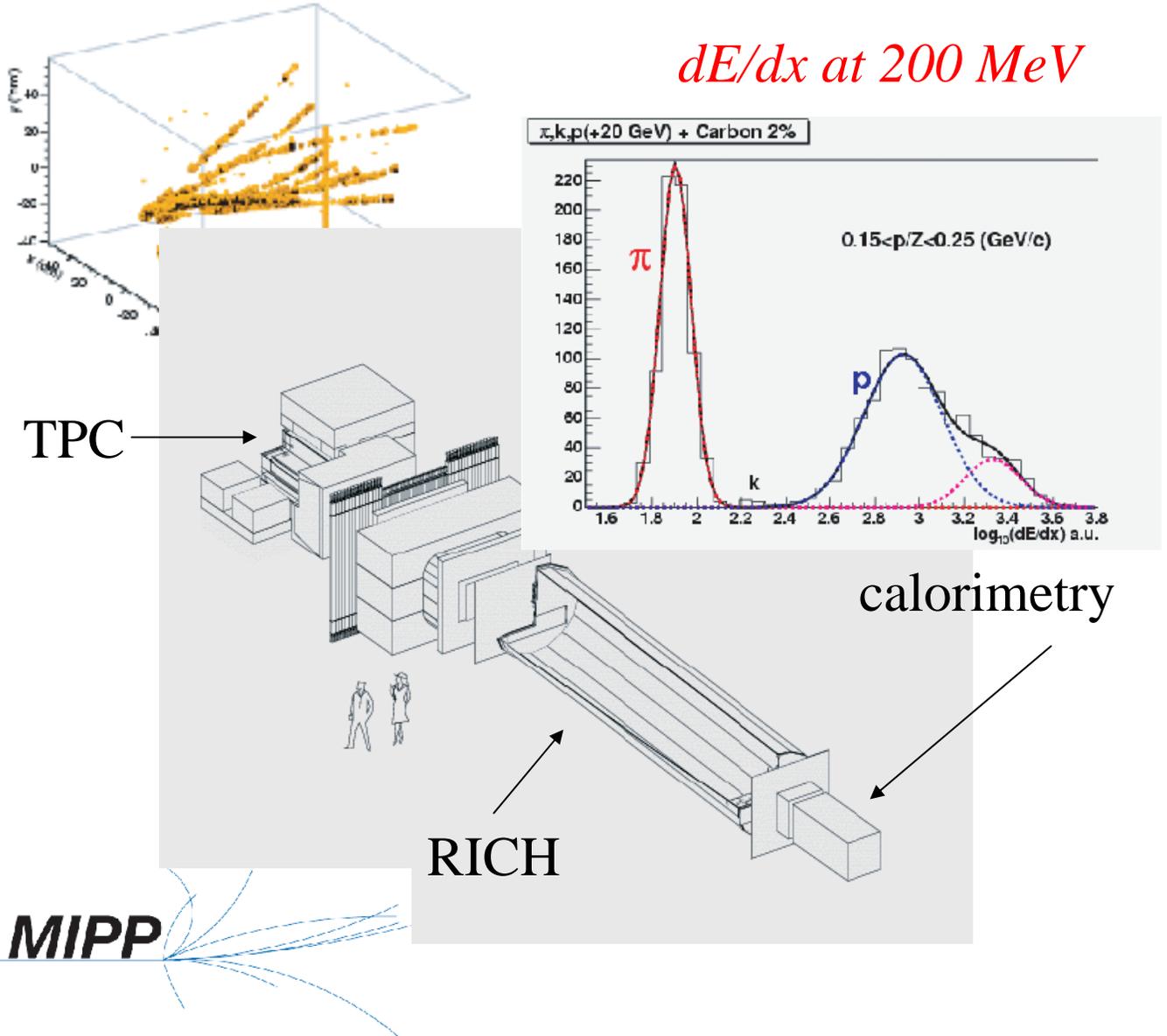
At very high beam powers, targeting becomes more difficult.

Radiation, heat transfer, shock issues.

SNS flowing liquid mercury target absorbs 2 MW!

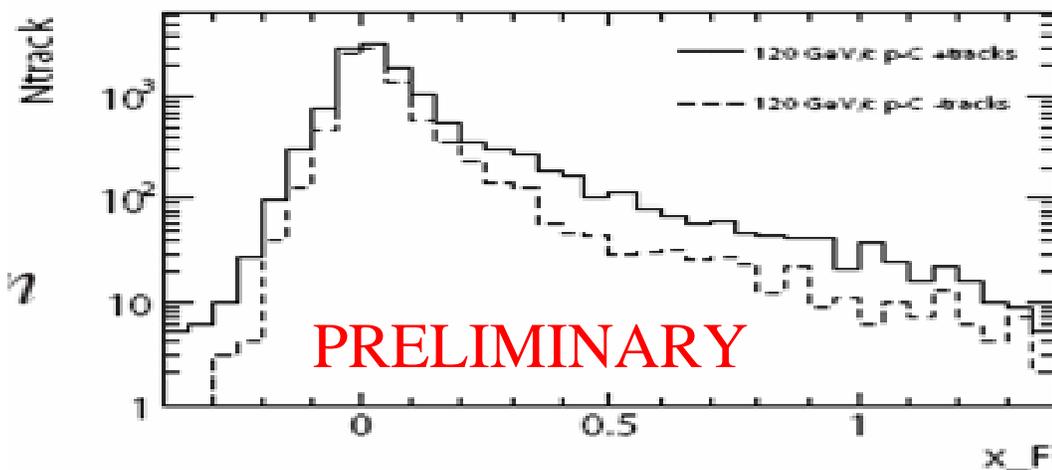
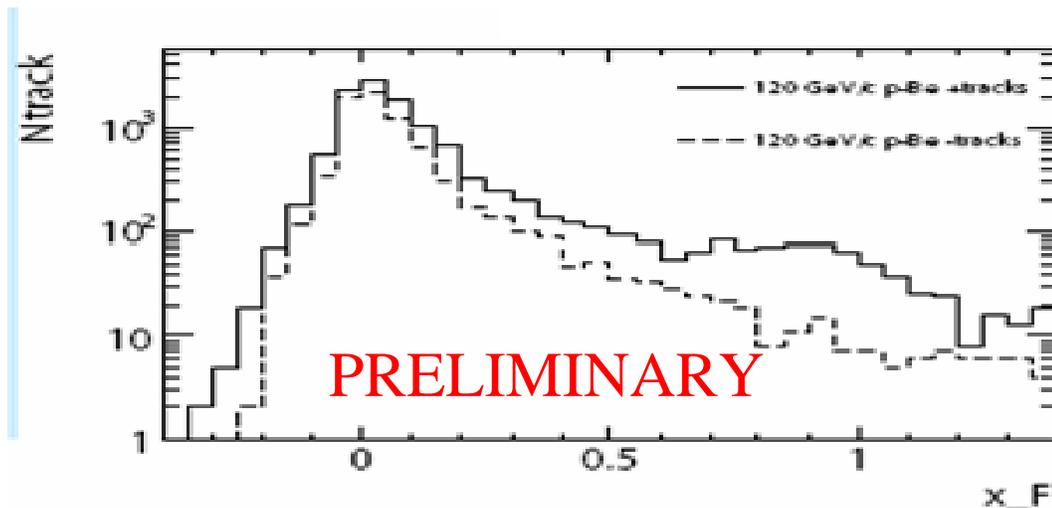


Example of experimental program- MIPP experiment at Fermilab



Study p, K + A interactions from 5-85 GeV
Study p + A interactions at 120 GeV
Measurements on H, Be, C, various metals

Typical MIPP Preliminary Data



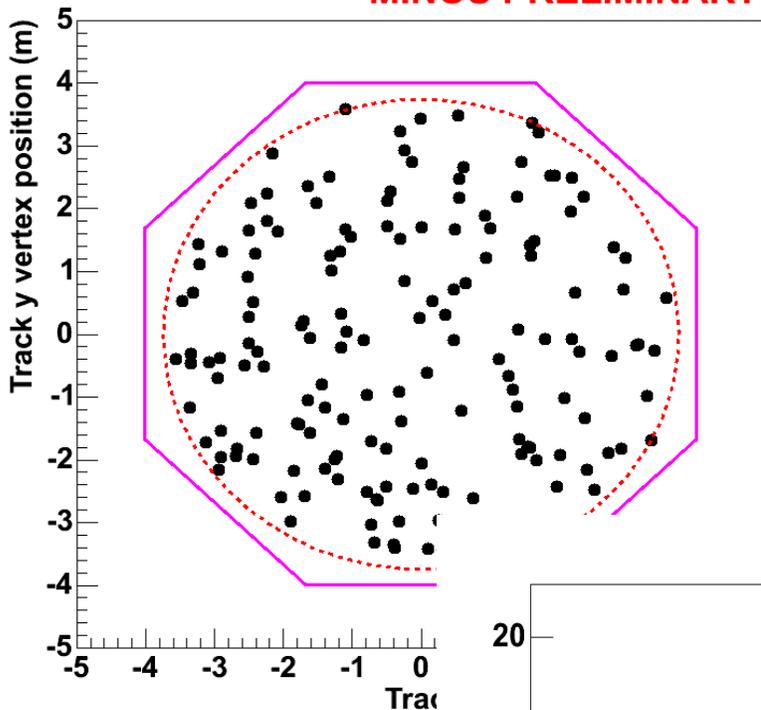
Inclusive x_F distributions for Be (top) and C (bottom) targets

Working through target list including NuMI target

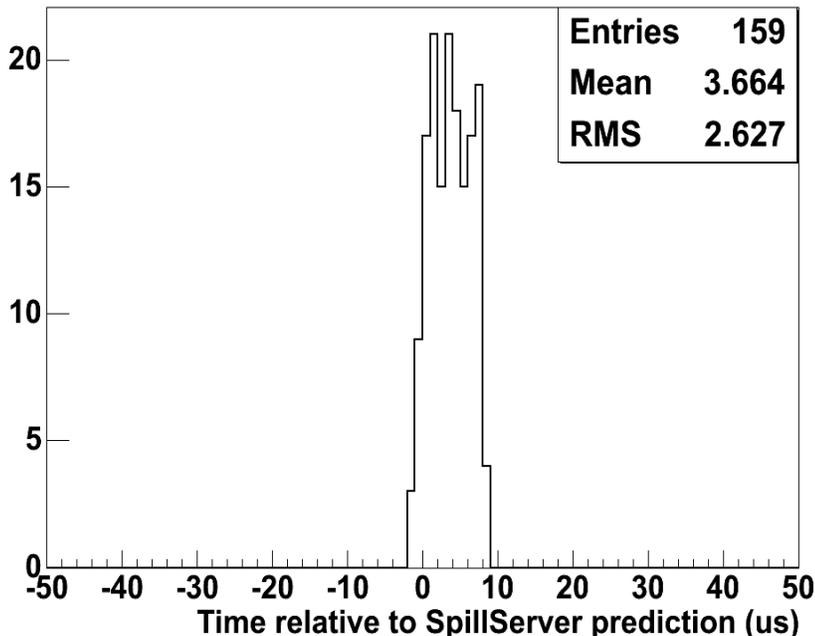
Courtesy Messier, NOVE-06

Vertex and Timing Distribution of Far Detector Events

MINOS PRELIMINARY



MINOS PRELIMINARY

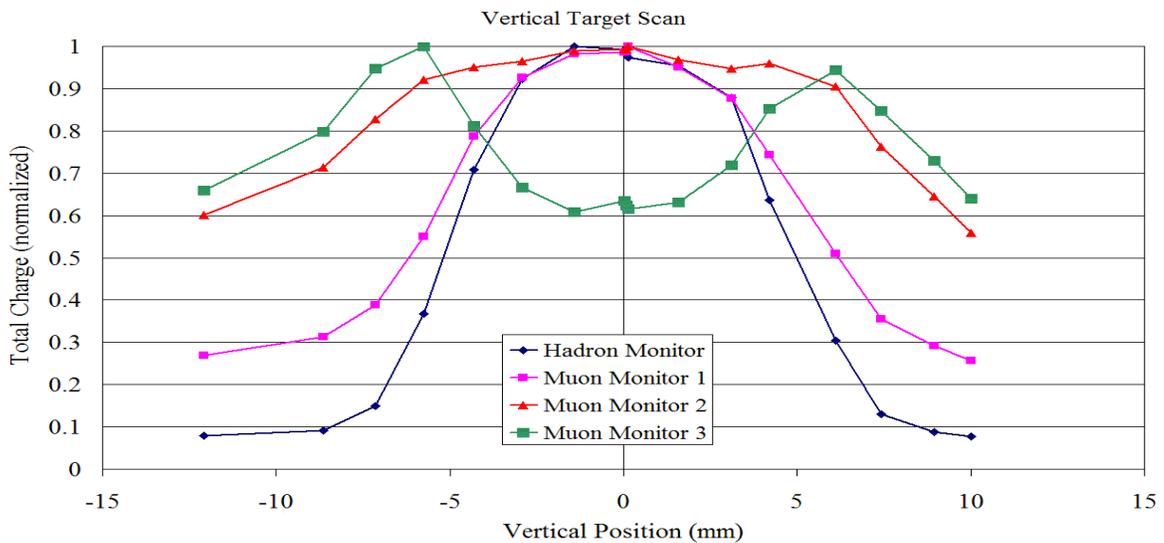
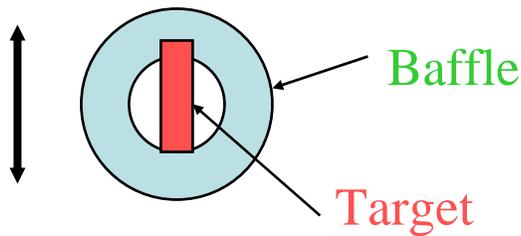
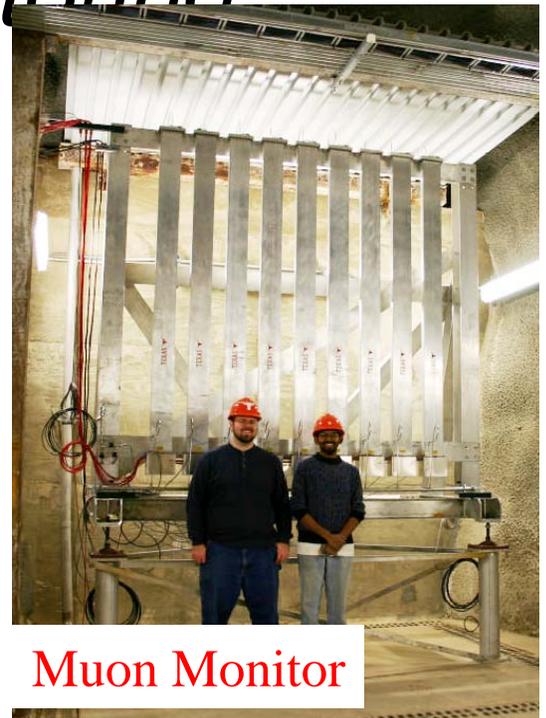
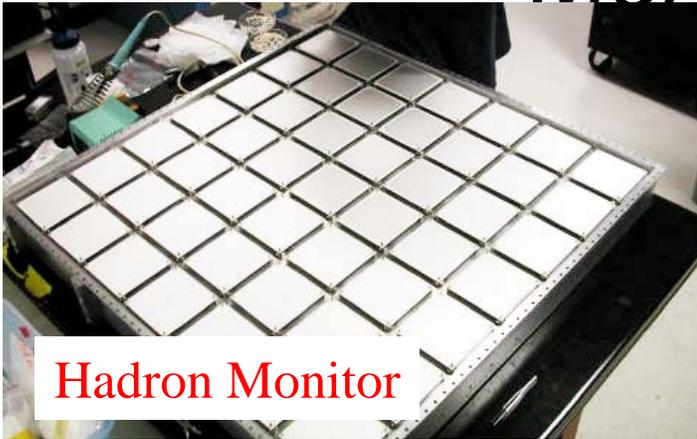


The cuts described result in 159 neutrino events.

Protons used for this work 9.3×10^{19}

Caution: This is MINOS open sample only

Secondary Beam Monitoring



Scan of beam across target shows edges, increased production in outer baffle - real detail!

Neutrinos from π, K Decay - Spectrum

Two-body decay $\pi \rightarrow \mu\nu$ of CM energy E^*

Lorenz Transformation of P_L, E

$$M = \gamma(1 + \beta \cos \theta^*)$$
$$M \cos \theta = \gamma(\cos \theta^* + \beta)$$

using

$$M \equiv E_{lab} / E^*$$

$$M = \frac{1}{\gamma(1 - \beta \cos \theta)}$$
$$\cong \frac{1}{\gamma(1 - \cos \theta)} \cong \frac{2\gamma}{1 + \gamma^2 \theta^2}$$

Spectrum depends on angle

Sharper for increased π energy

Neutrinos from π , K Decay - Flux

Lorenz transform P_T this time

$$E \sin \theta = E^* \sin \theta^*$$

↓

$$M \sin \theta = \sin \theta^*$$

for small angles

$$\sin \theta \cong \theta \Rightarrow M \theta = \theta^*$$

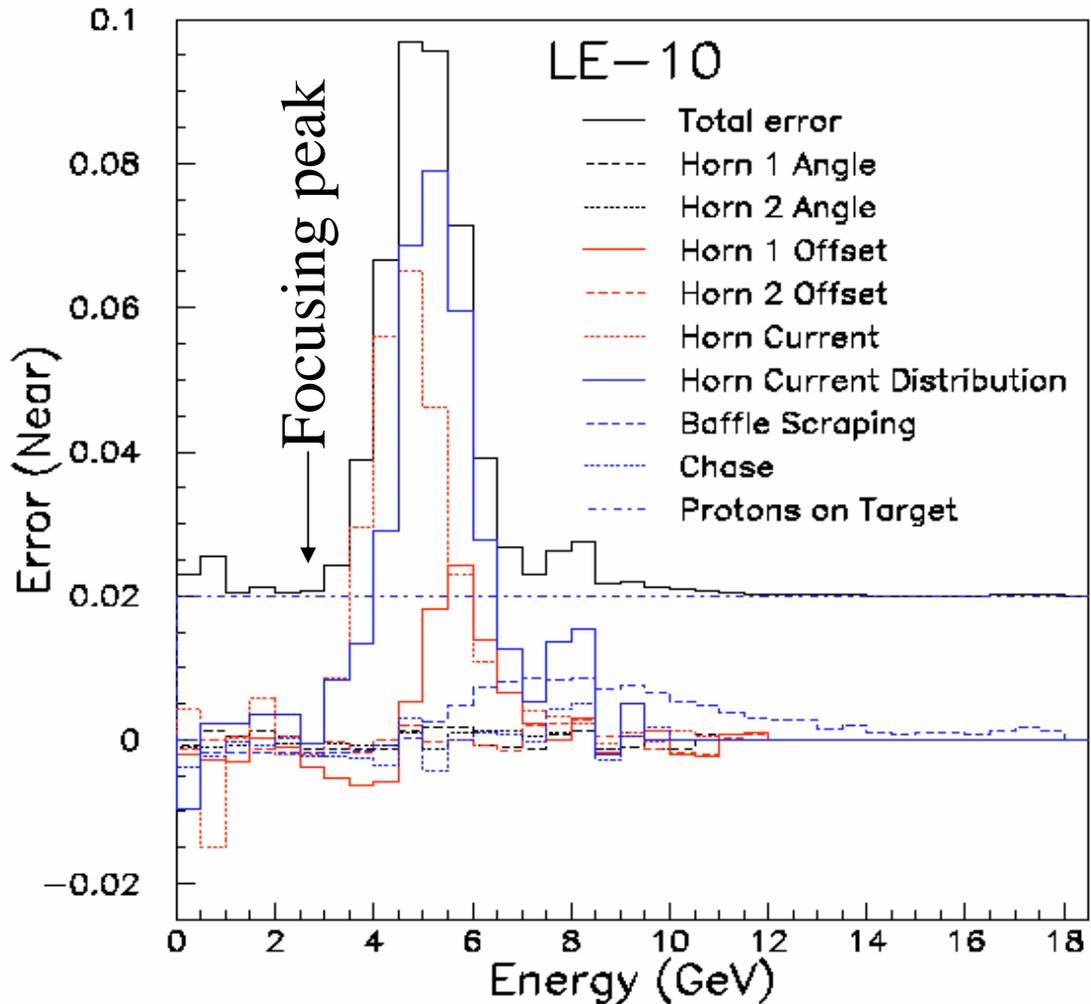
$$\Rightarrow \frac{d\theta^*}{d\theta} = M$$

*actually, this is valid for all values of θ^**

giving flux compression factor

$$\frac{d\Omega^*}{d\Omega} = \frac{\sin \theta^* d\theta^*}{\sin \theta d\theta} = M^2 = \frac{4\gamma^2}{(1 + \gamma^2 \theta^2)^2}$$

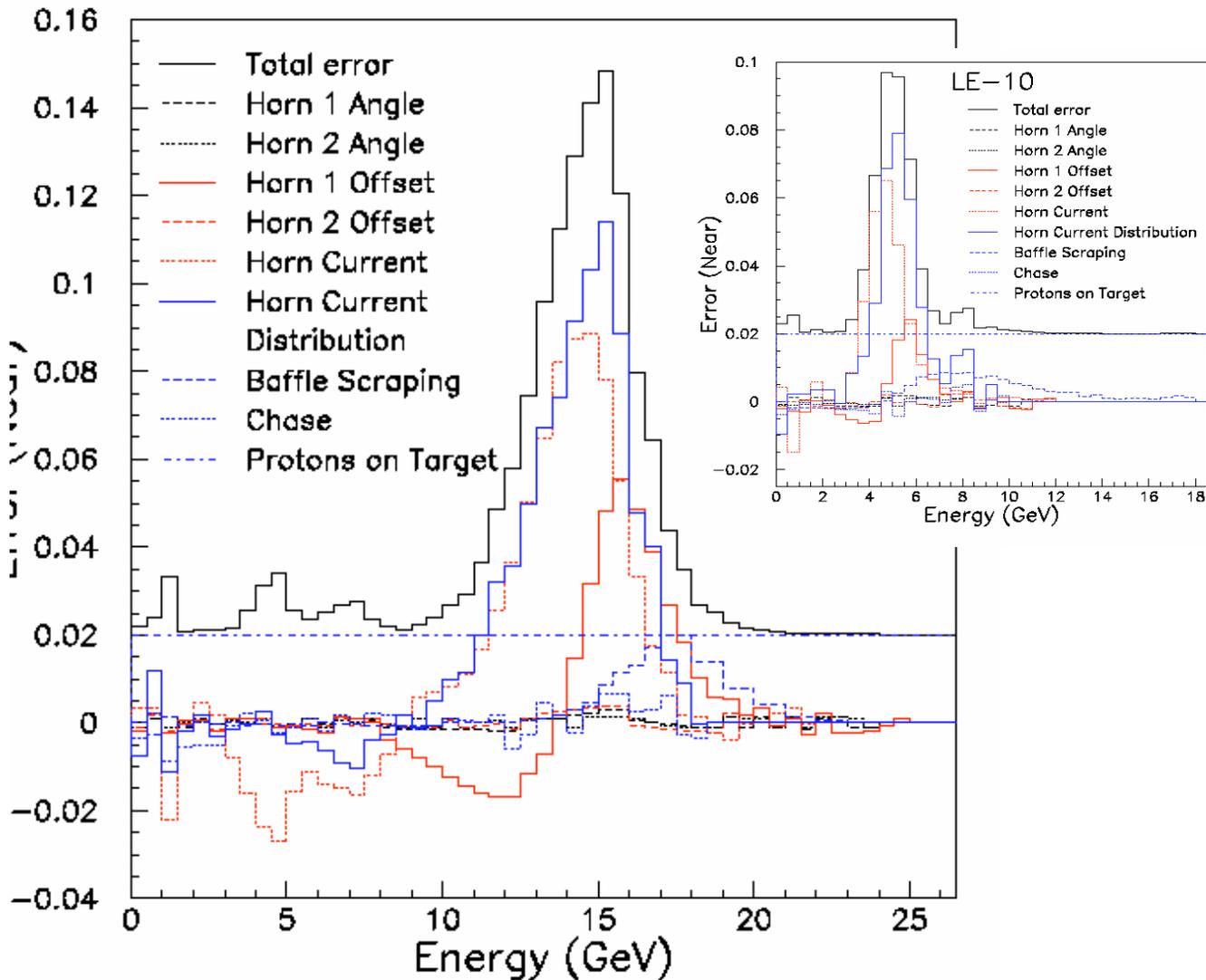
Beamline Flux Fractional Uncertainties in NuMI Beam



Kopp, et. al

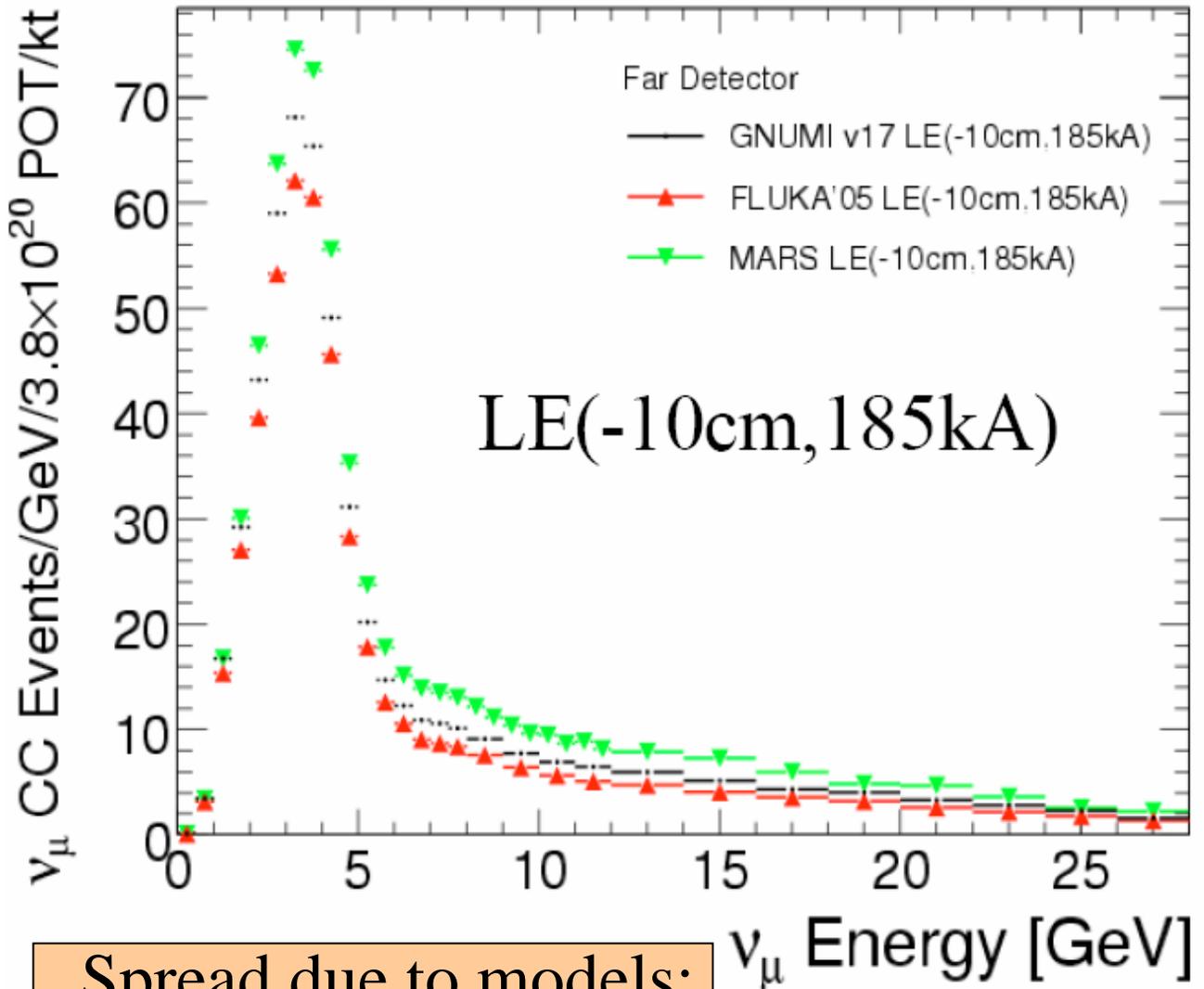
Does not include errors from
hadron production model

Study typical uncertainties with beam energy evolution



NuMI high energy beam at near location
(inset repeats the low energy plot)

Spread of a neutrino spectrum due to parent hadron uncertainties



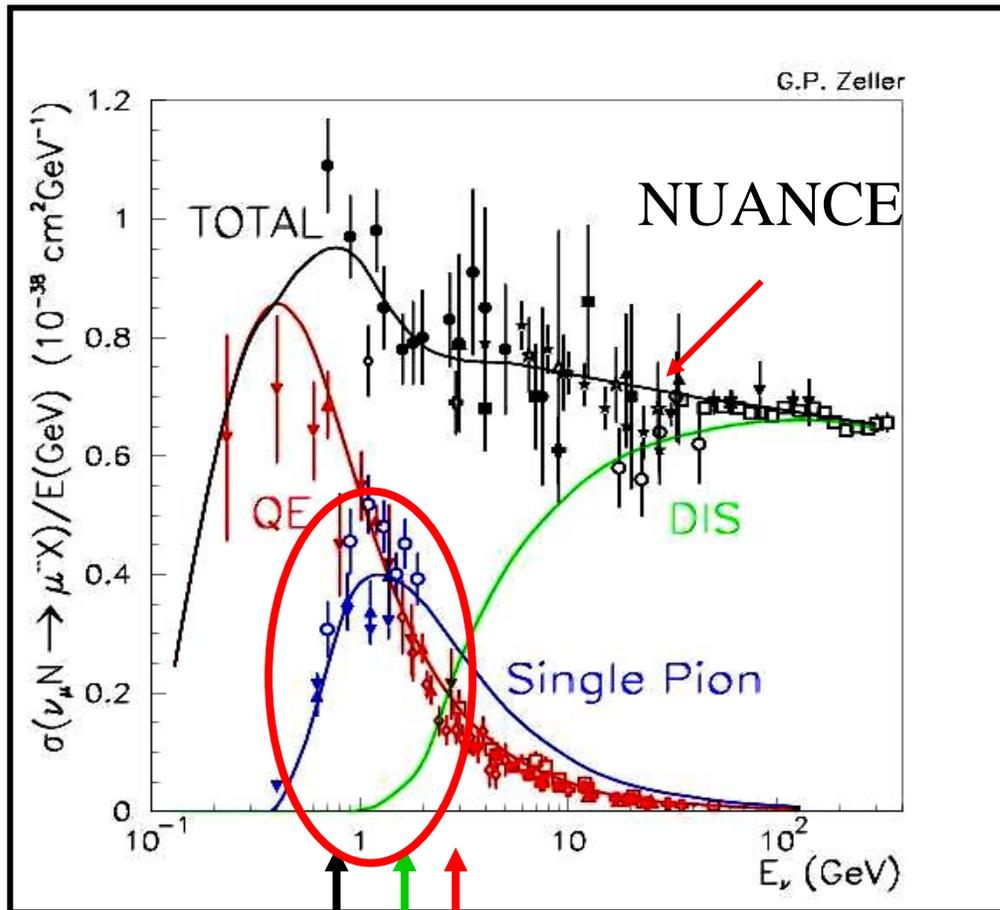
Spread due to models:

8% (peak)

15% (tail)

Kopp, et. al.

Interlude - the physics of 0.5-5 GeV Neutrino Interactions



↑ **MINOS, NuMI**
↑ **K2K,**
↑ **MiniBooNE, T2K**
← **NOVA**
← **Super-K atmospheric ν**

Complex physics modeled as a combination of low-multiplicity processes

Pions are produced via intermediate resonant states

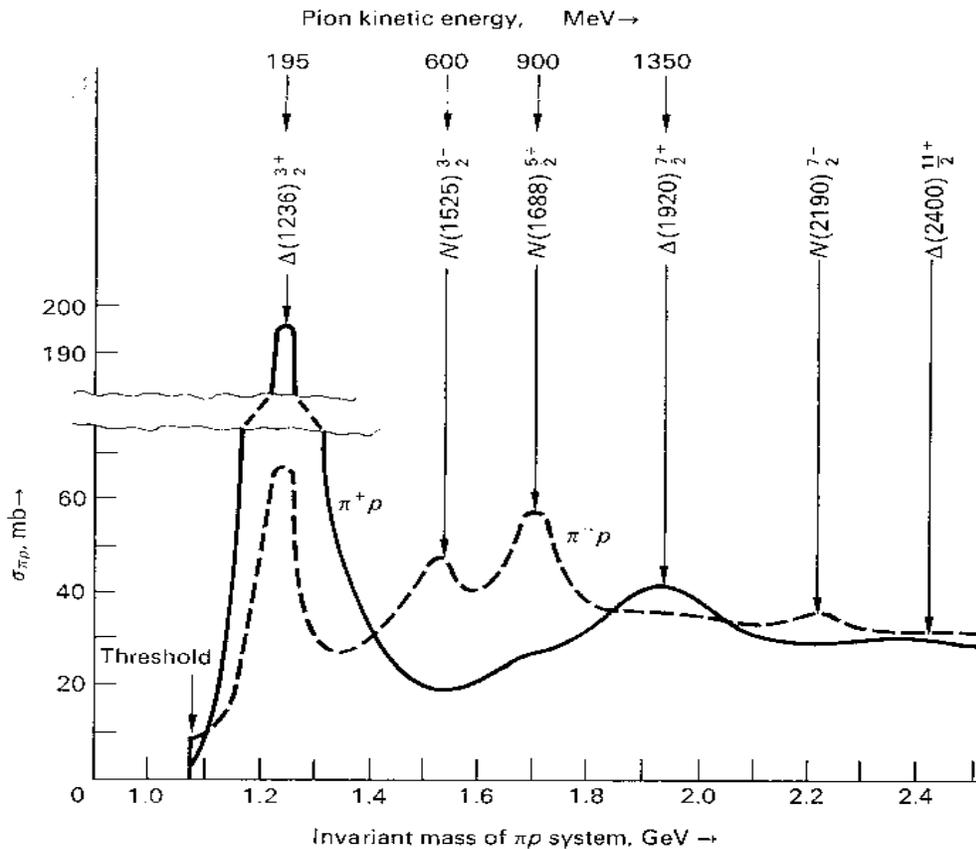
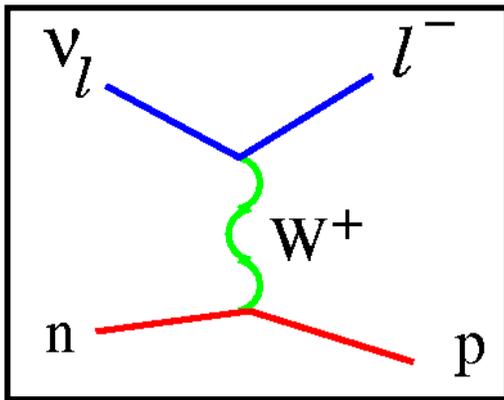


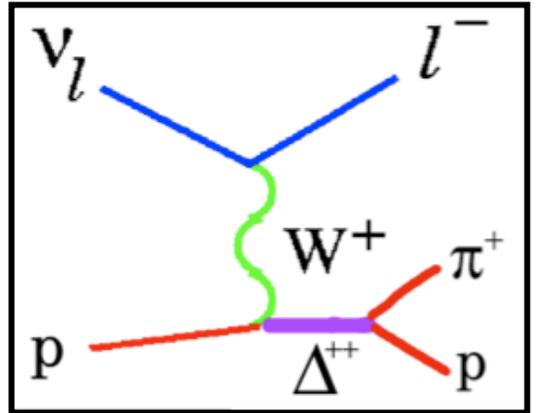
Fig. 4.6 Variation of total cross section for π^+ and π^- mesons on protons, with incident pion energy. The symbol Δ refers to resonances of $l = \frac{3}{2}$; N refers to $l = \frac{1}{2}$. The positions of only a few of the known states, together with their spin-parity assignments, are given.

Resonances as observed in πn scattering experiments
(Figure from Perkins, Introduction to High Energy Physics)

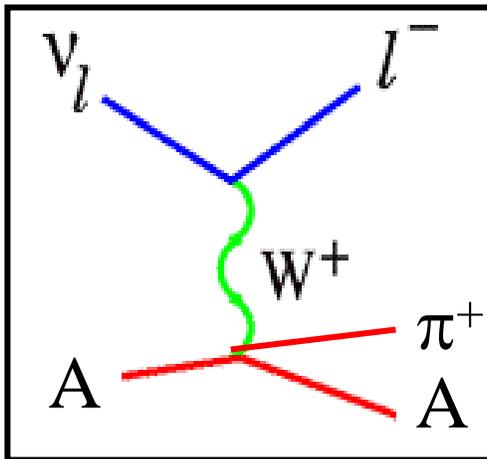
Different models have underlying physics in common



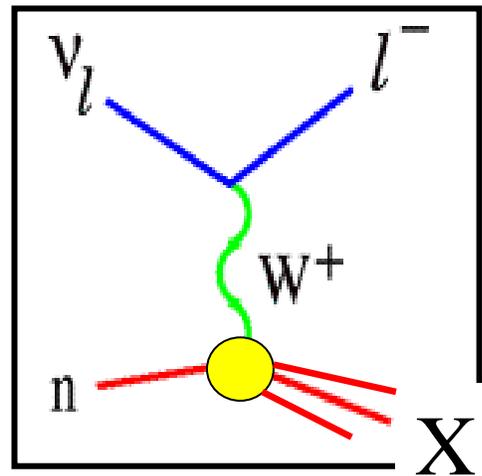
Quasi-elastic



Resonant π production



Coherent production



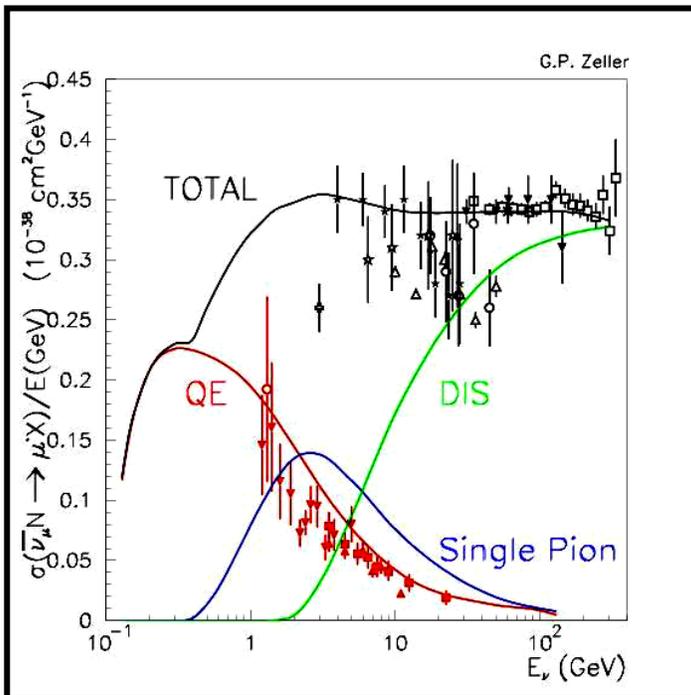
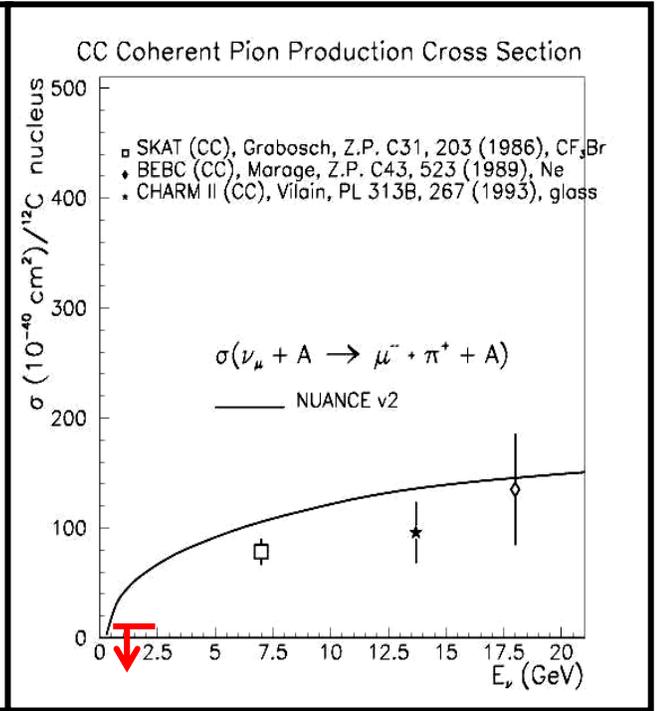
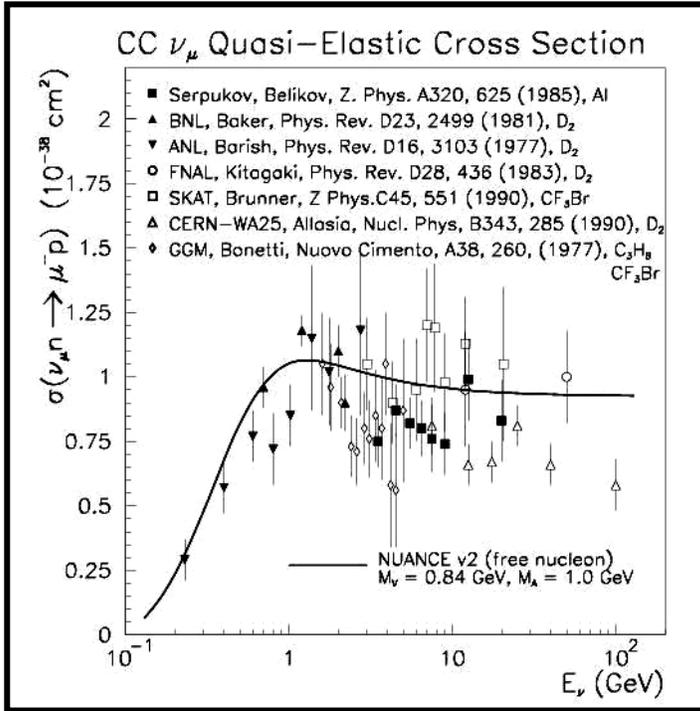
DIS

Different models combine channels differently.

e.g. NUANCE - coherent addition of resonances

NEUGEN - incoherent addition

Existing data not strongly constraining

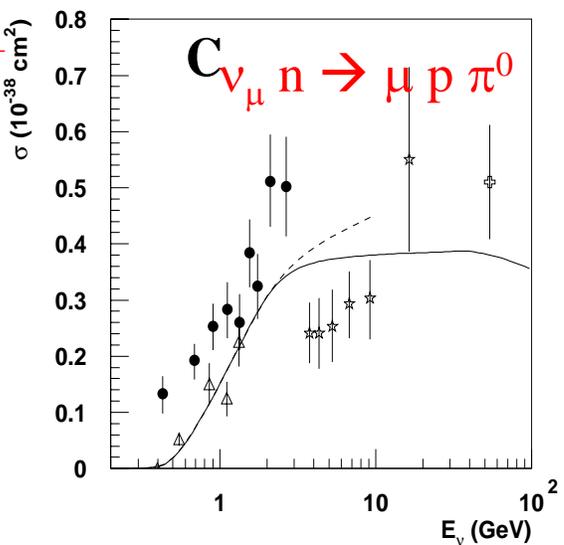
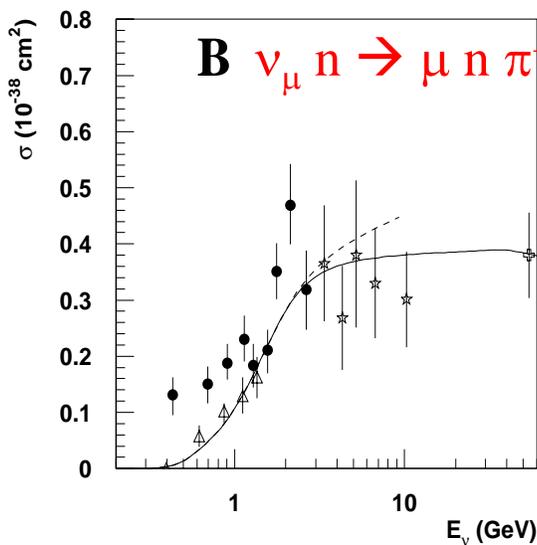
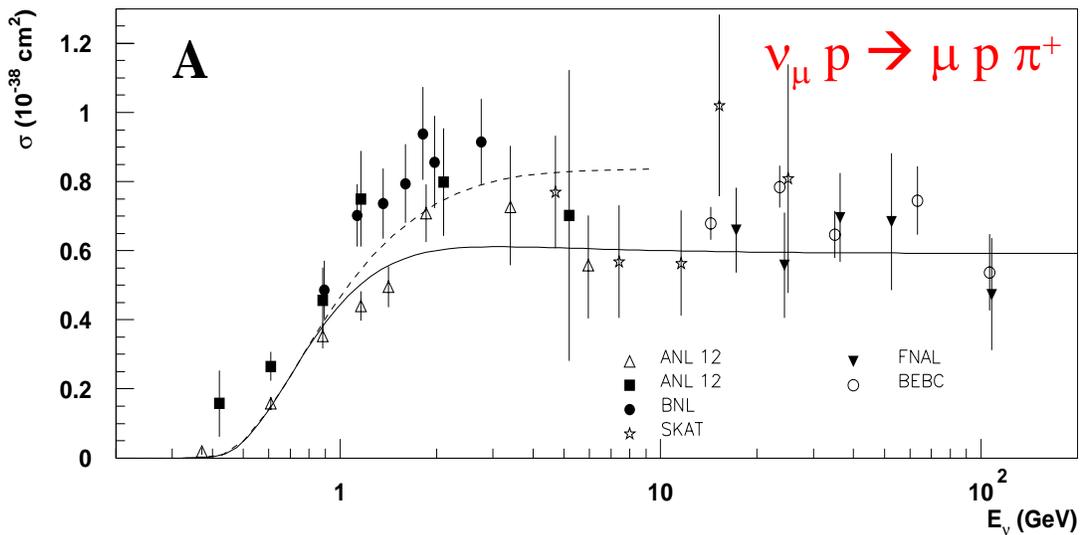


Coherent, ν -bar data basically nonexistent in region of interest.

Parameters in the models are tuned to exclusive channel cross sections.

$$\sigma_{tot} = \sigma_{QE} + \sum_i (\sigma_{res}^i + f_i \sigma_{DIS}^i)$$

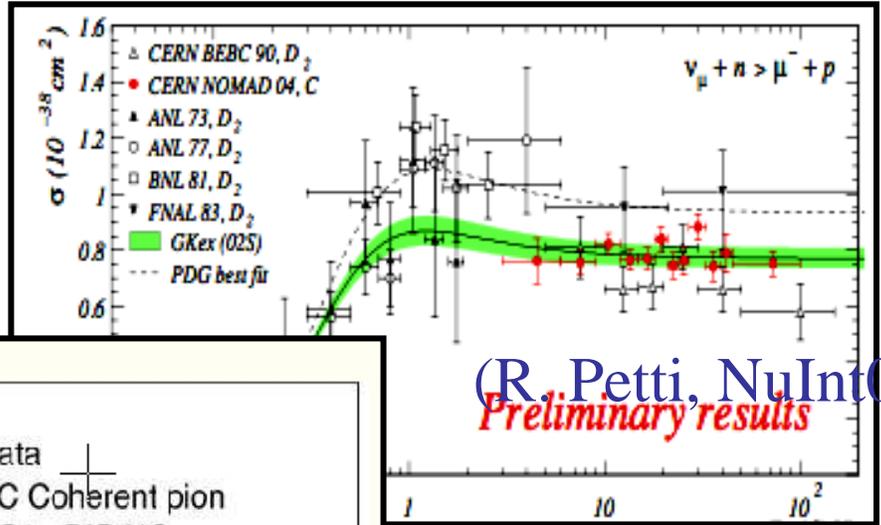
Charged Current Single Pion Production



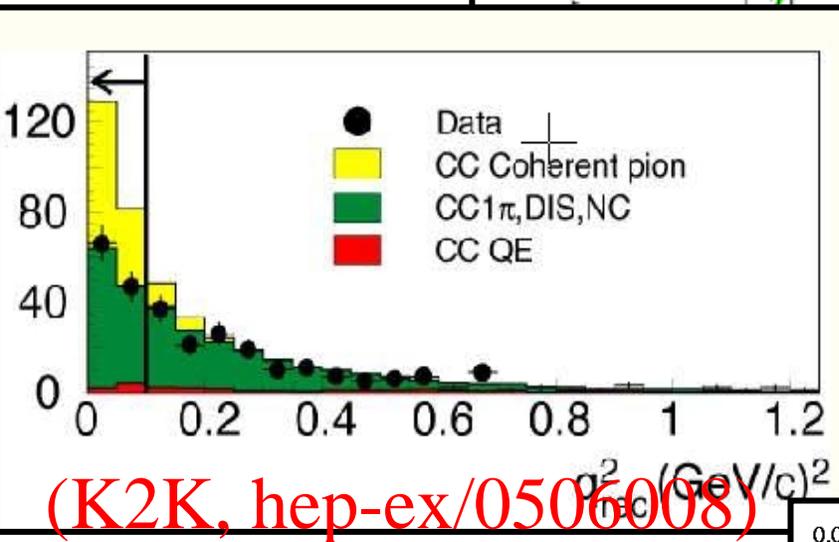
Examples from NEUGEN courtesy H. Gallagher

Worldwide effort to improve knowledge

•NOMAD, ^{12}C



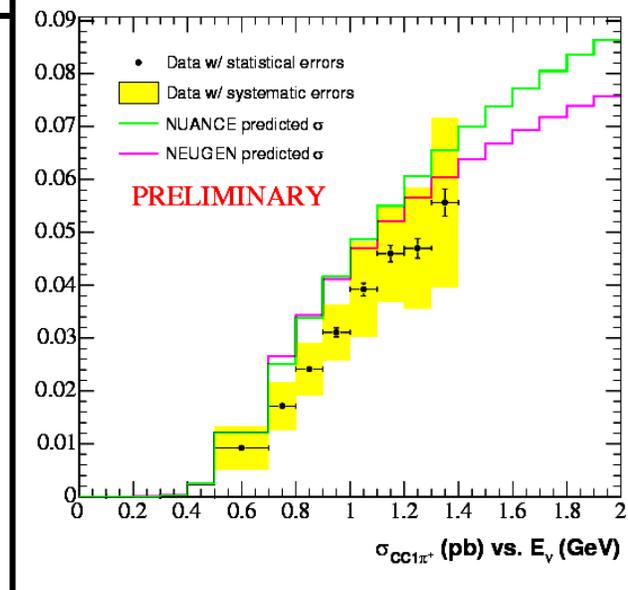
(R. Petti, NuInt05)
Preliminary results



← K2K,
 coherent π^+

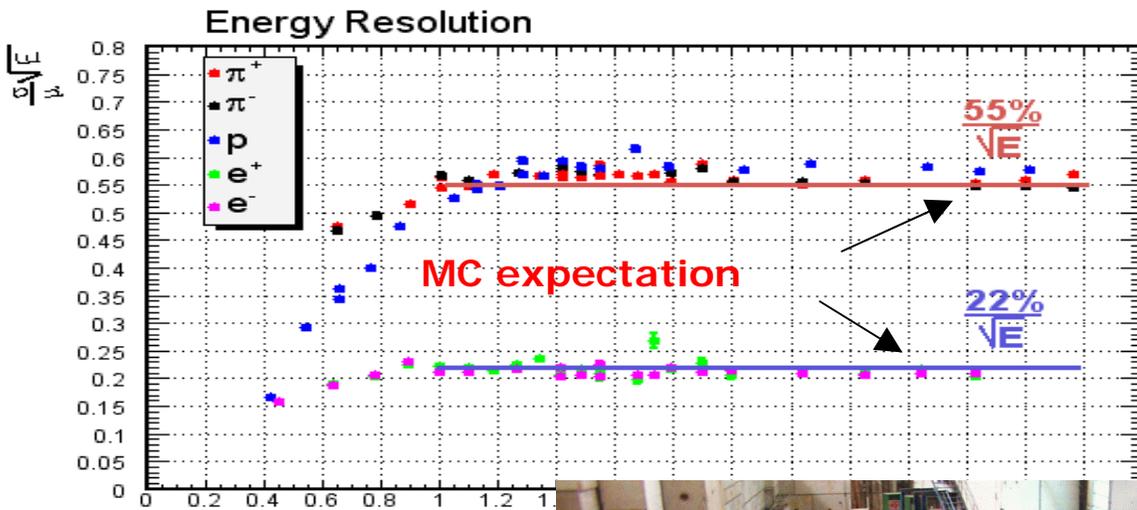
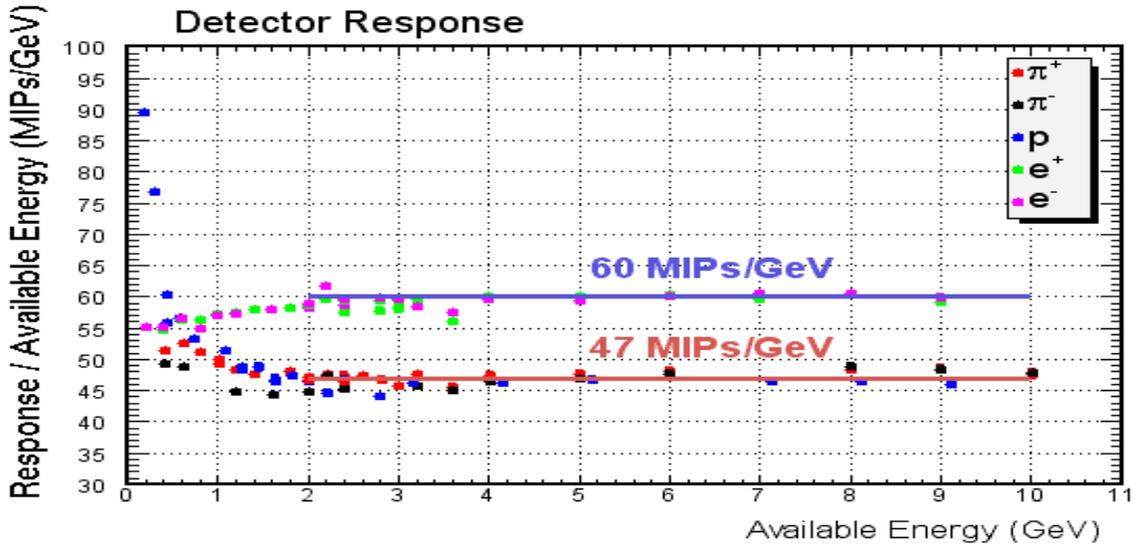
(K2K, hep-ex/0506008)

MiniBoone single π
 (Monroe, Wasco)



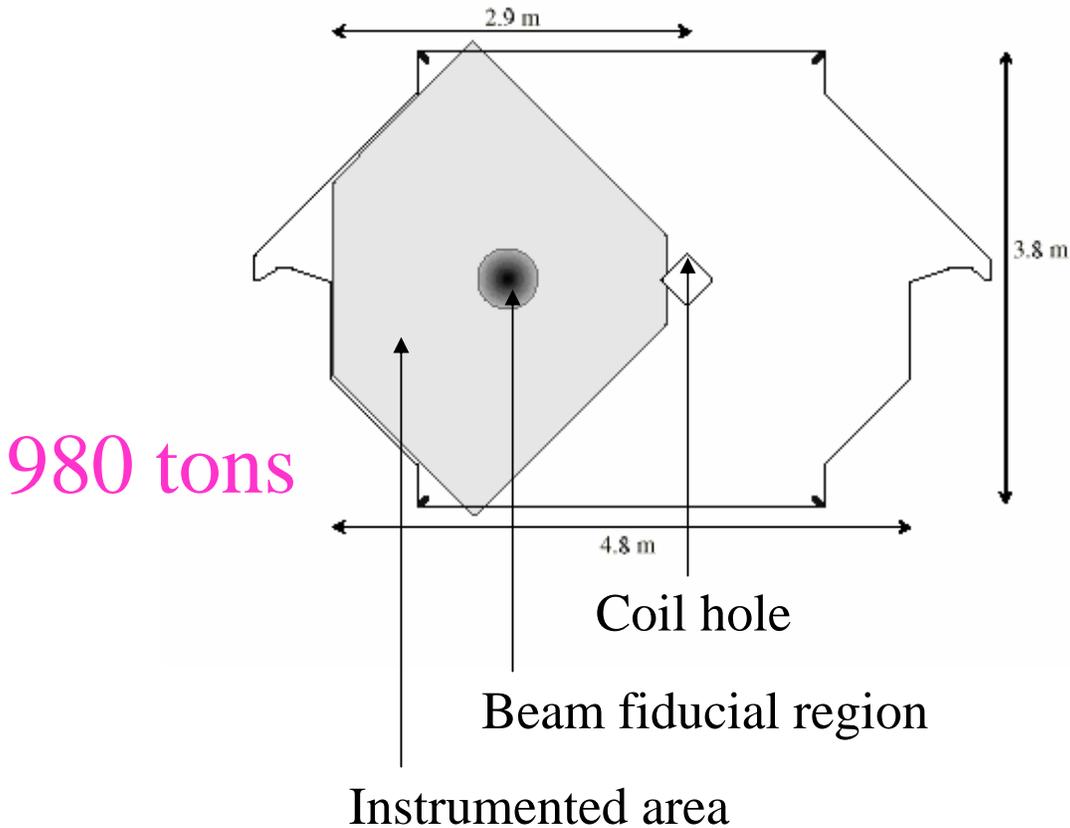
Plots courtesy G. Zeller,
 NOvE-06

MINOS Calibration Detector at CERN



60-plane 'micro - MINOS'
Also checked near/far electronics

Minos Near Detector Architecture

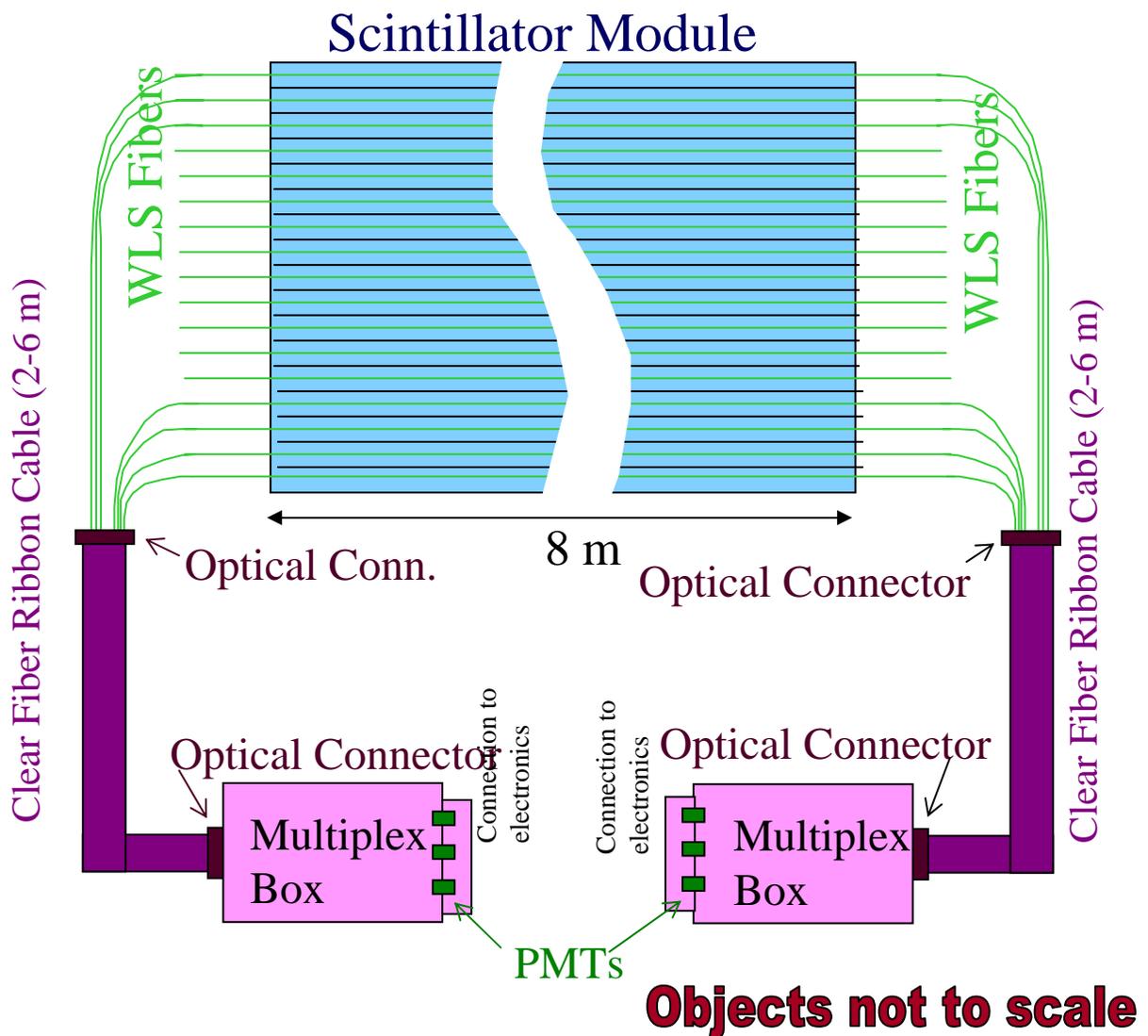


Single steel plates (280) , shorter modules

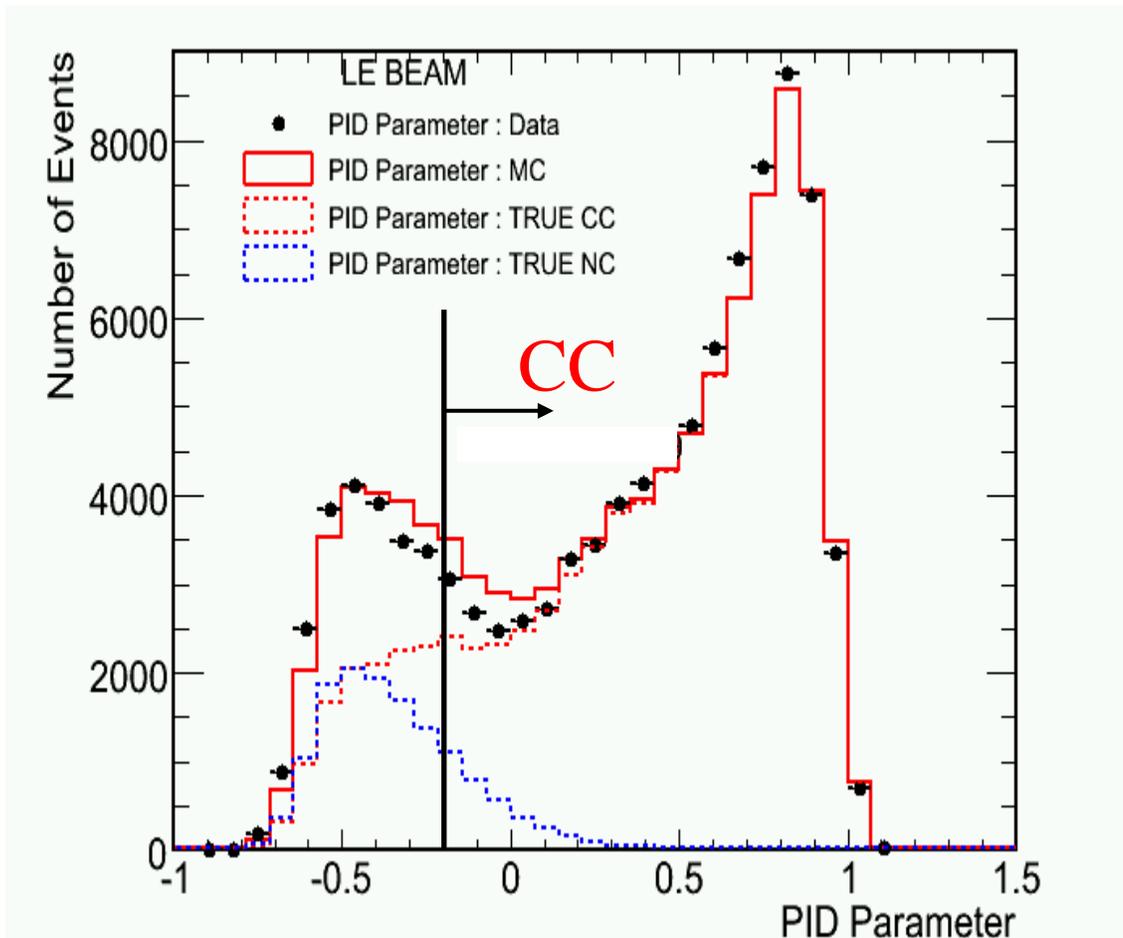
Calorimeter (1st 3/7) is partially instrumented
except for 1/5 of planes with full coverage

Muon Spectrometer section has only every 5th
plane instrumented

MINOS Detector Technology



Performance of Charged Current Selection Algorithm

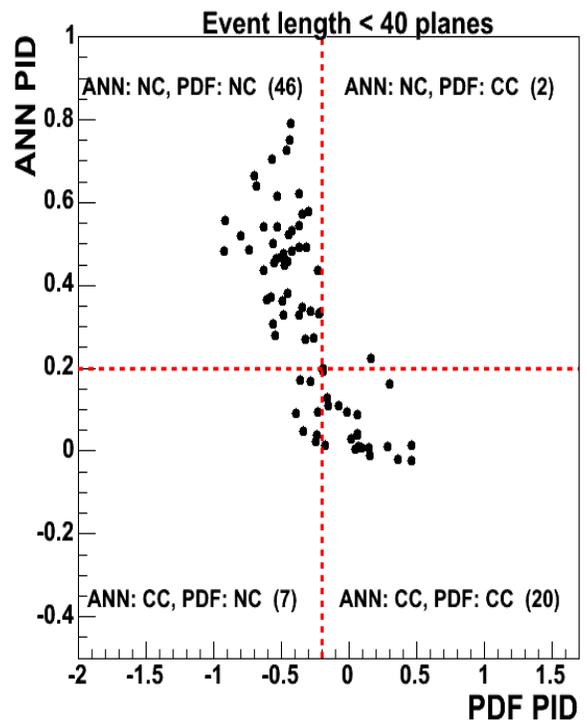
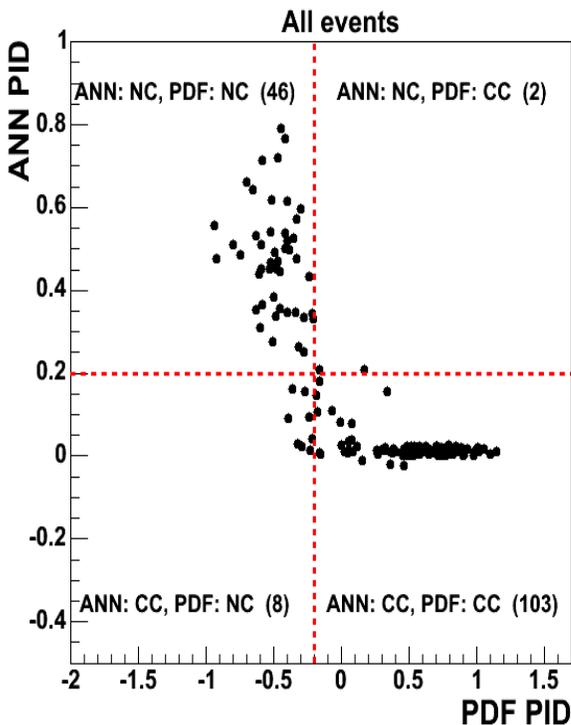


With cut at -0.2, MC estimate of efficiency is 87%, with purity 98%

Stability of CC Selection Algorithms

MINOS PRELIMINARY

Correlations between PID parameters



Likelihood-Based

Likelihood-Based

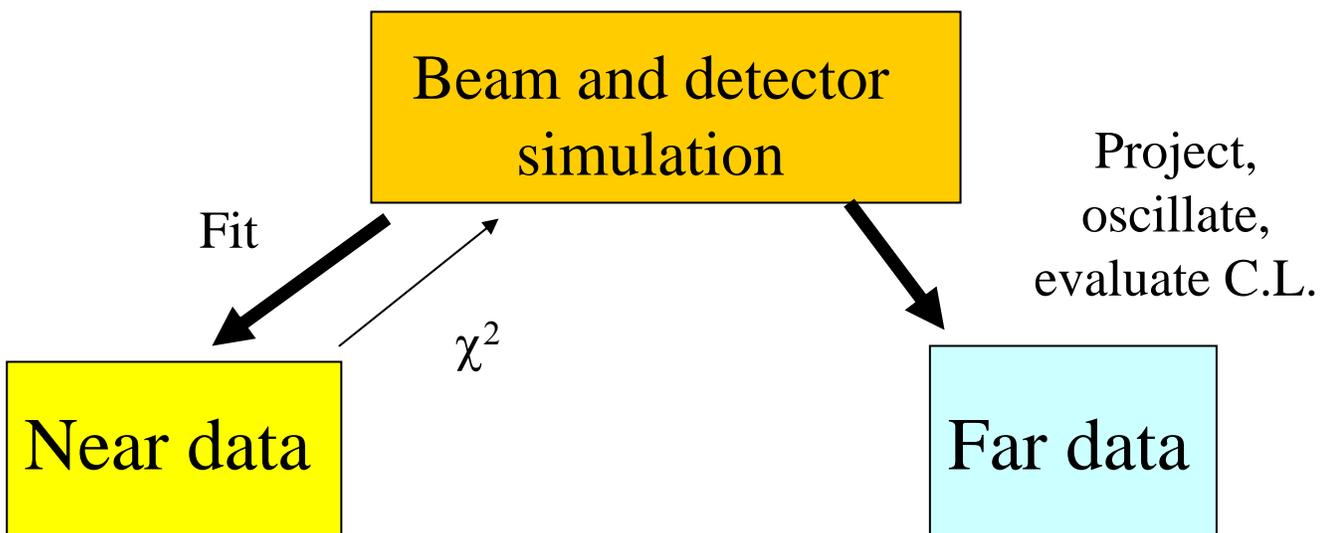
Excellent overlap between algorithms
for charged-current selection.

Extrapolations and use of near detector data

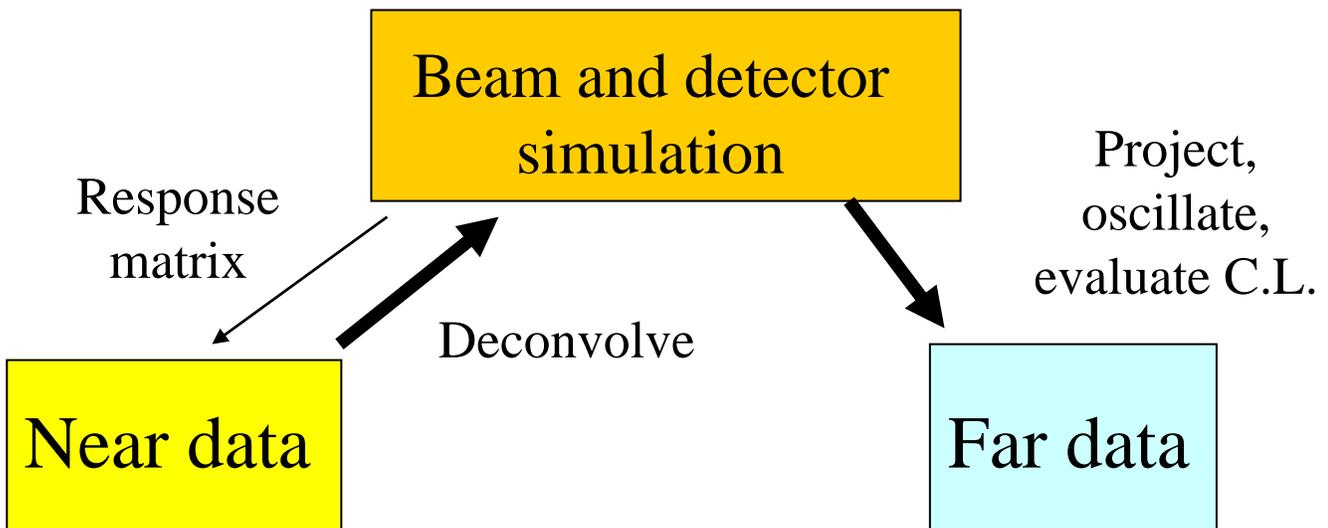
Two representative techniques

many common inputs

I) Fitting the near data



II) Deconvoluting the near data

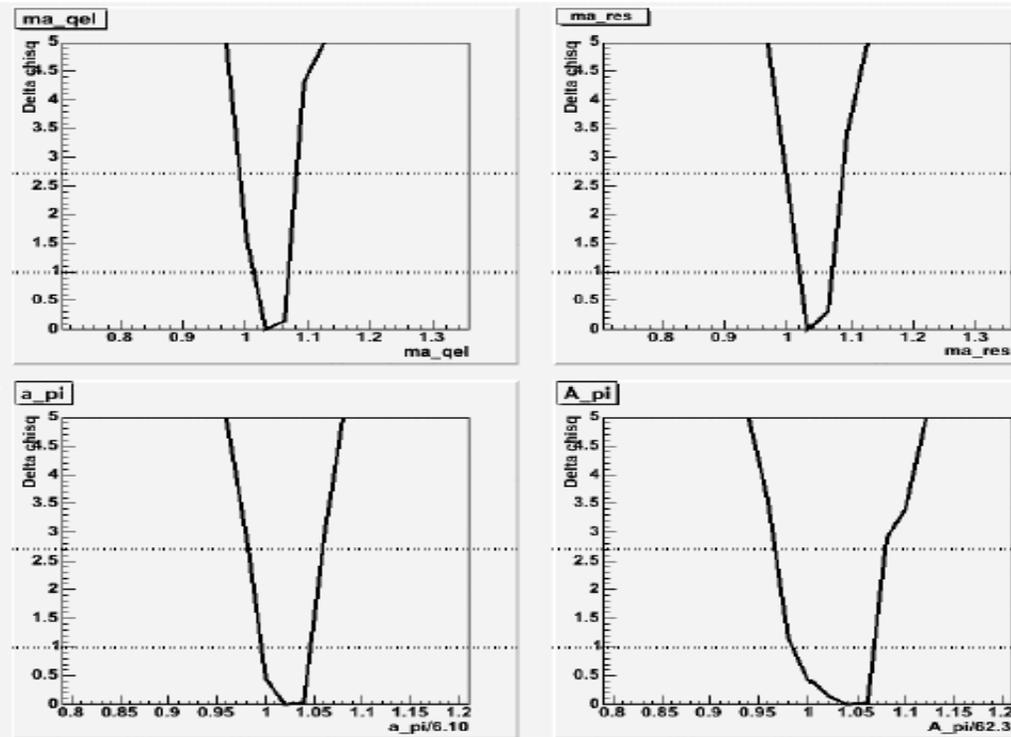


Example of typical fitting procedure - MINOS Mock

Data Challenge

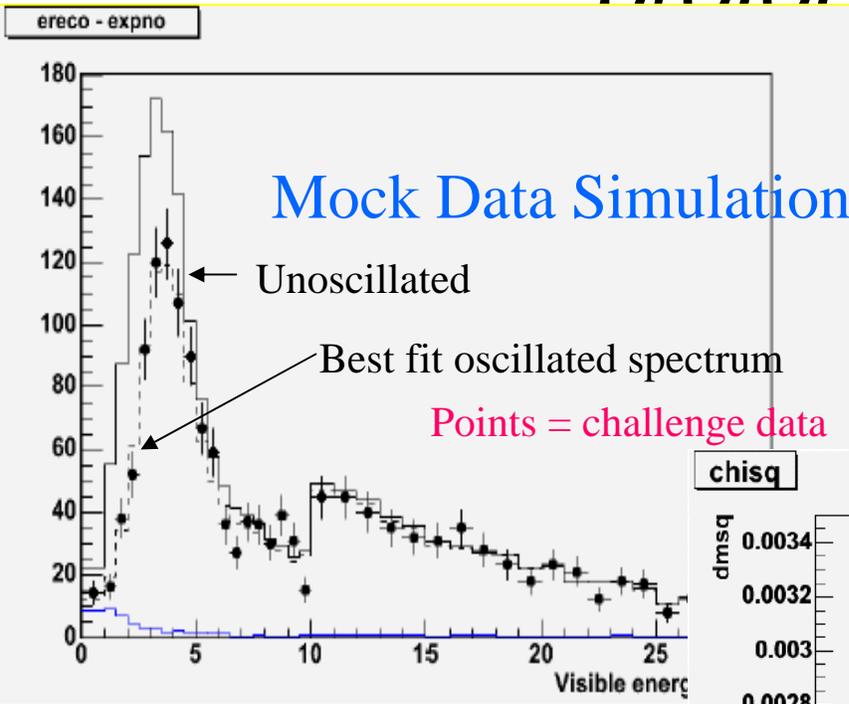
$$\chi^2(\vec{\alpha}) = \sum_{i=1}^{nbins} \frac{[n_i - m_i(\vec{\alpha})]^2}{\sigma_i^2} + \sum_{j=1}^{nsyst} \frac{\varepsilon_j^2}{\sigma_j^2}$$

where $\vec{\alpha}$ is a vector of parameters and ε is a fractional (uncorrelated) systematic error.



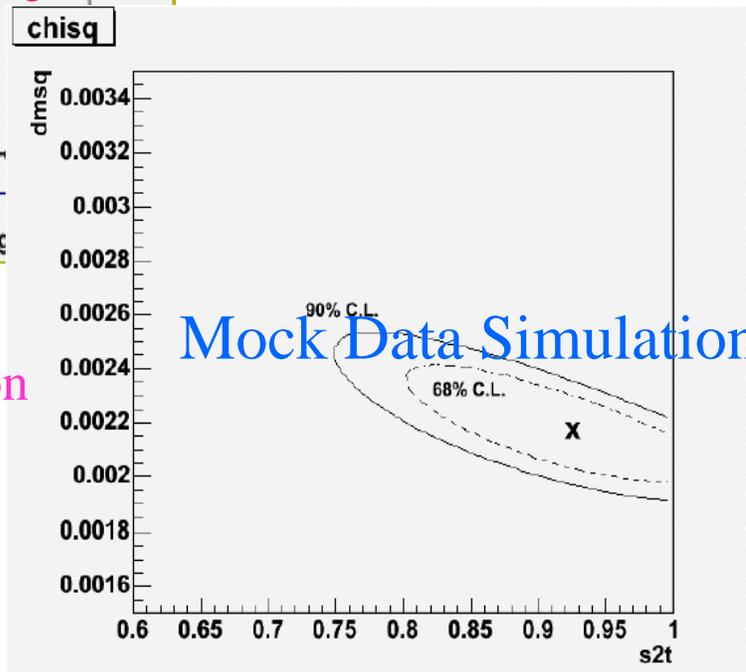
Typical $\Delta\chi^2$ plots for BMPT and cross-section parameters

Results of Mock Data Challenge (simulated 7.4×10^{20} protons



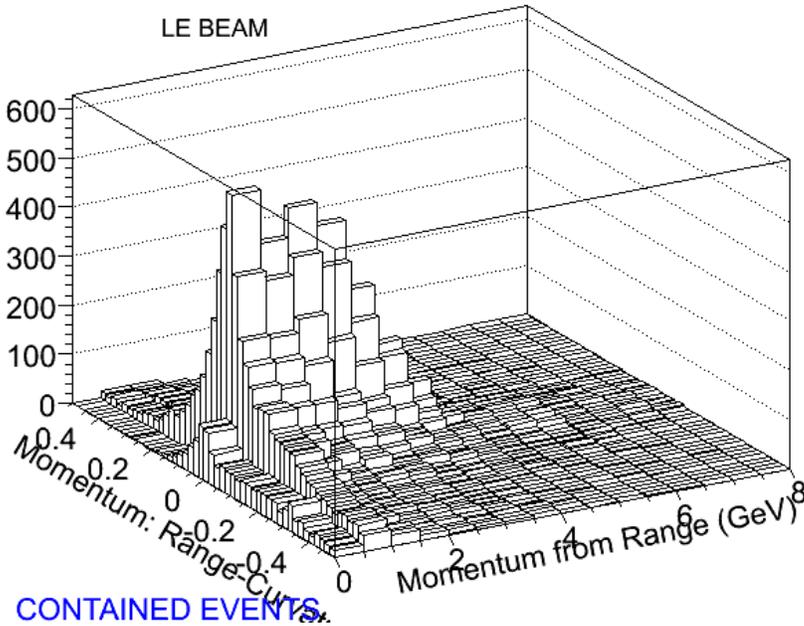
Fit to reconstructed far energy spectrum

Best fit results for oscillation parameters



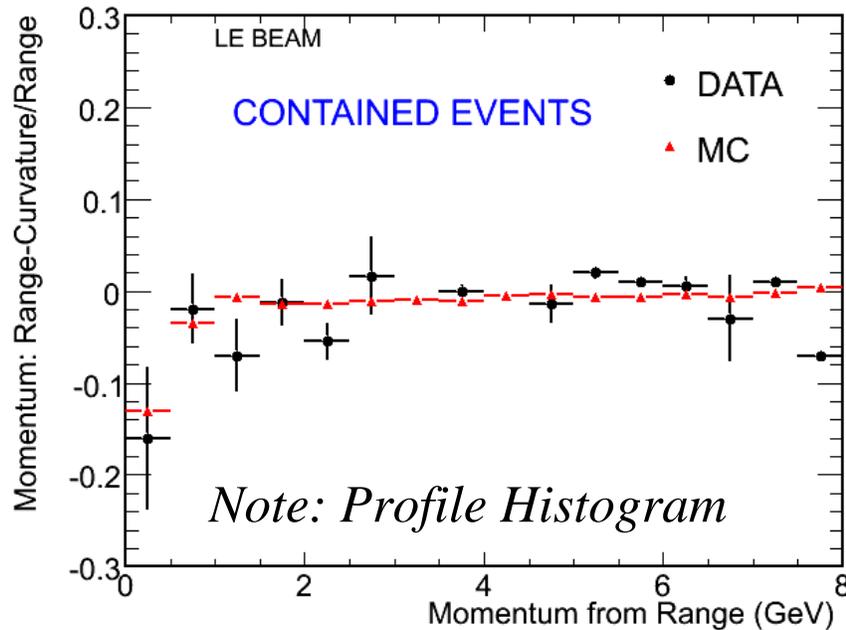
	Challenge Value	Fitted Value
Δm^2	2.1×10^{-3}	2.2×10^{-3}
$\text{Sin}^2(2\theta_{23})$	0.88	0.93

Range and Curvature Momentum Comparison



Near Detector

Far Detector



Comparison in momentum regime where events are contained.

Builds confidence in magnetic field map and calibration

Example -MINOS CC Disappearance Analysis

Blind analysis - only <50% of data in open sample for comparisons.

Remainder modified by “blinding function”.

Steps in analysis

Select neutrino events

Classify as CC events

Likelihood-based procedure using pulse height, event length.

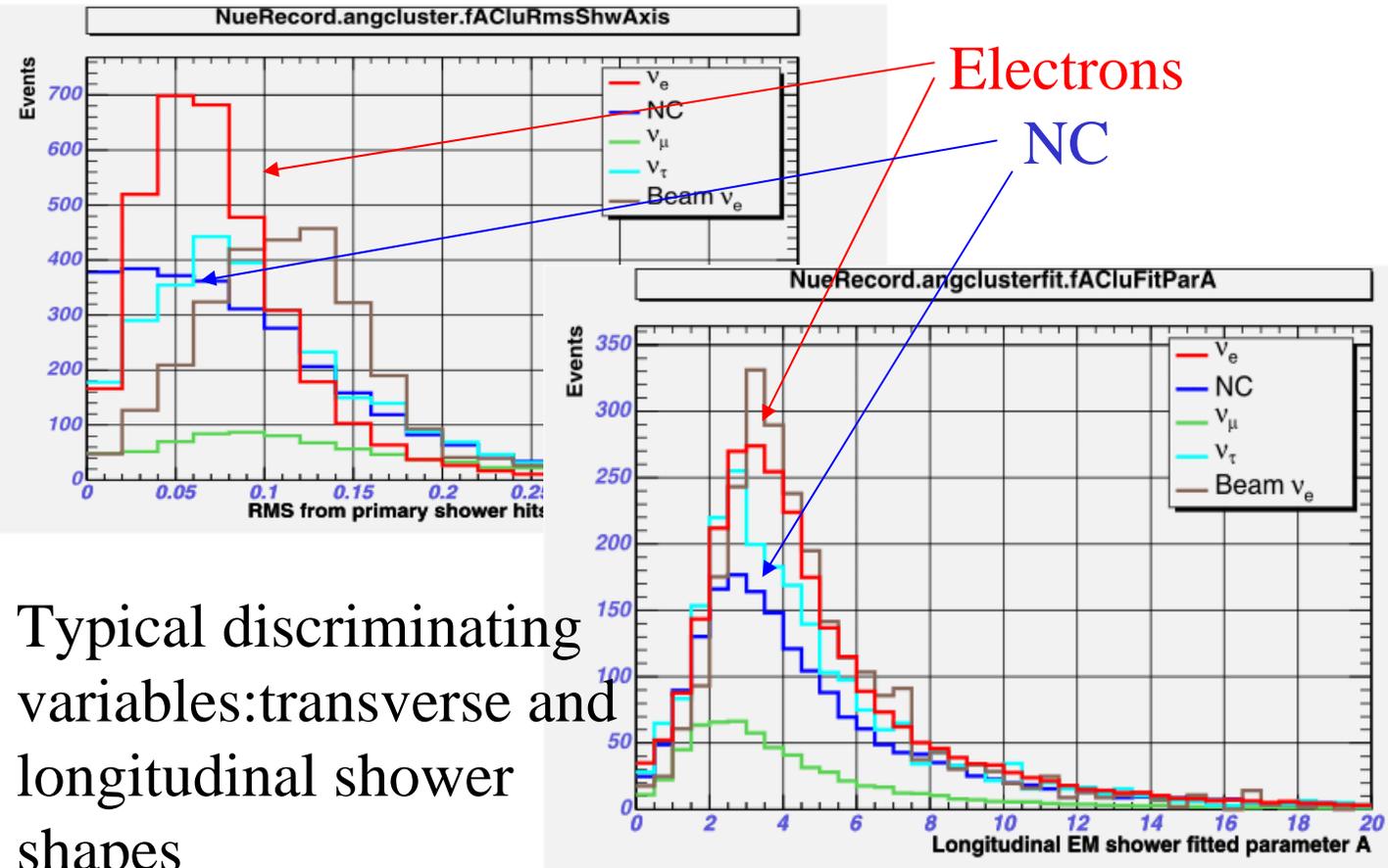
Check with neural net based procedure - good agreement. Also controlled scanning checks.

Extrapolate (using near) to far CC spectra to extract oscillation parameters.

Notes: NuMI/MINOS MC used to extrapolate far/near

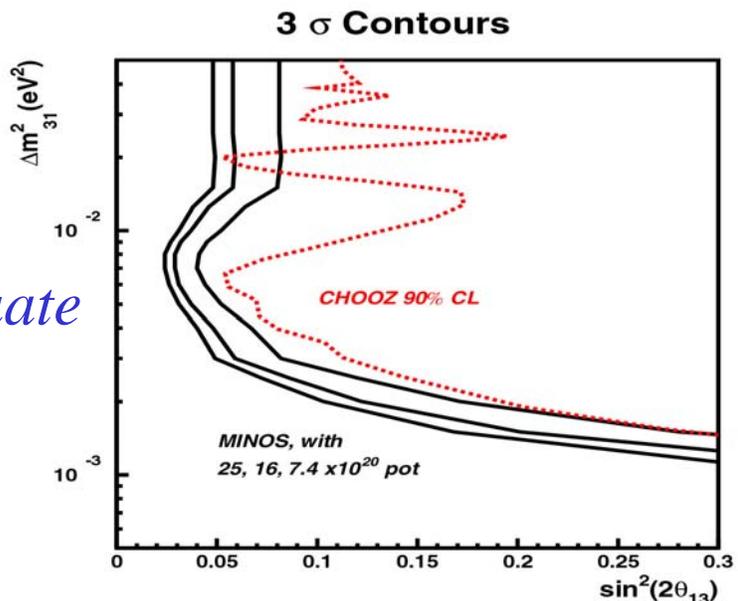
Fit χ^2 will include systematic errors.

MINOS electrons and their sensitivity



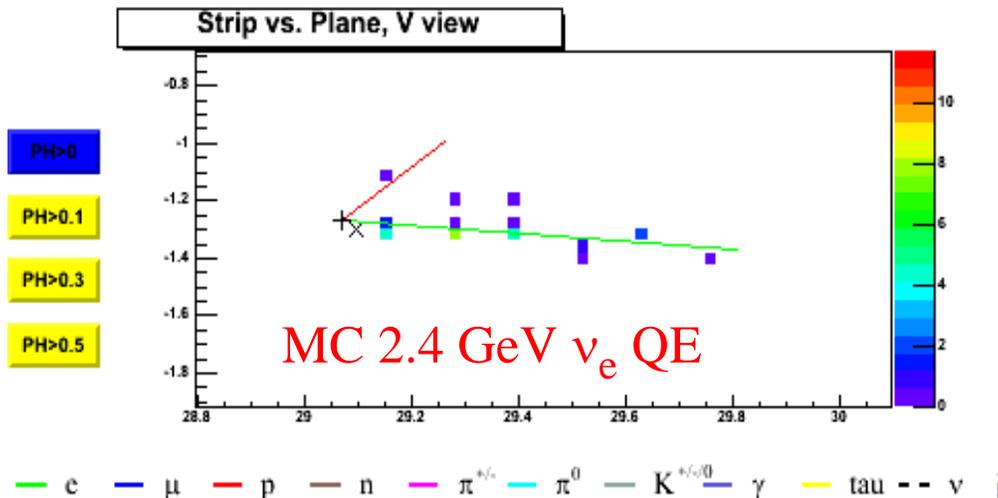
Typical discriminating variables: transverse and longitudinal shower shapes

Can improve CHOOZ limit by ~ 2 with adequate protons



Electrons and $\sin^2 2\theta_{13}$ in MINOS

As an iron calorimeter, MINOS is not optimized for low-energy e^- detection (1-3 GeV)



With appropriate attention, progress can be made on limits for $\sin^2 2\theta_{13}$

Principal background ($\sim 2/3$) is from mis-identified π^0 NC events:

$$\nu_{\mu} + N \rightarrow \nu_{\mu} + N' + \pi^0, \quad \pi^0 \rightarrow \gamma\gamma$$

Understand Off-axis Mathematically

Model beam with $\theta^* = \pi/2$

$$P_T = P^* \sin \theta^*$$

$$P_L = \gamma P^* (1 + \cos \theta^*) \cong E_\nu$$

We fix $\theta_{lab} \equiv P_T / P_L$

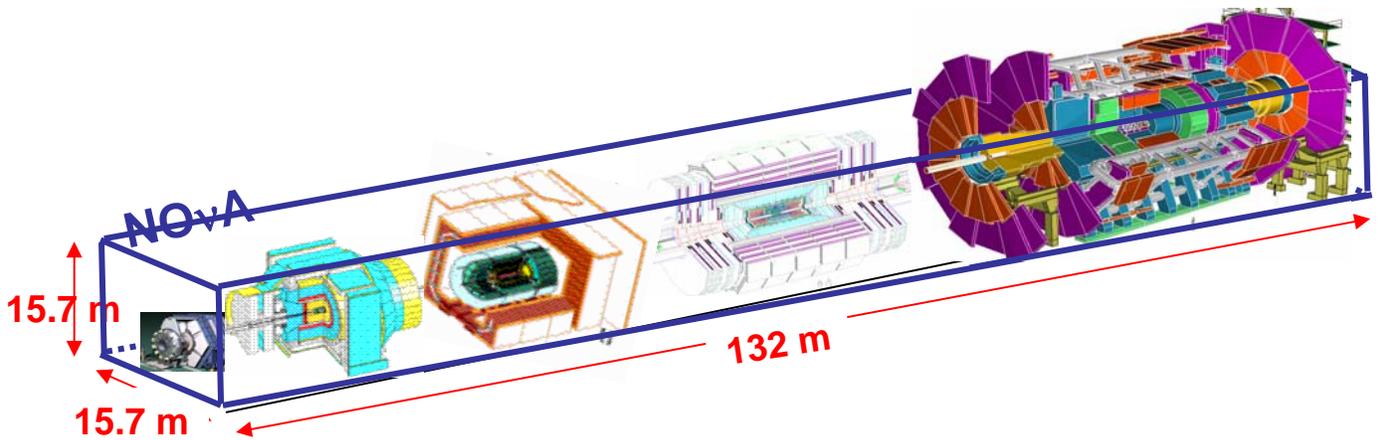
At $\pi/2$ $\Delta P_T \cong 0$ as vary θ^*

$$\Delta P_L = \frac{\Delta P_T}{\theta_{lab}} \cong 0$$

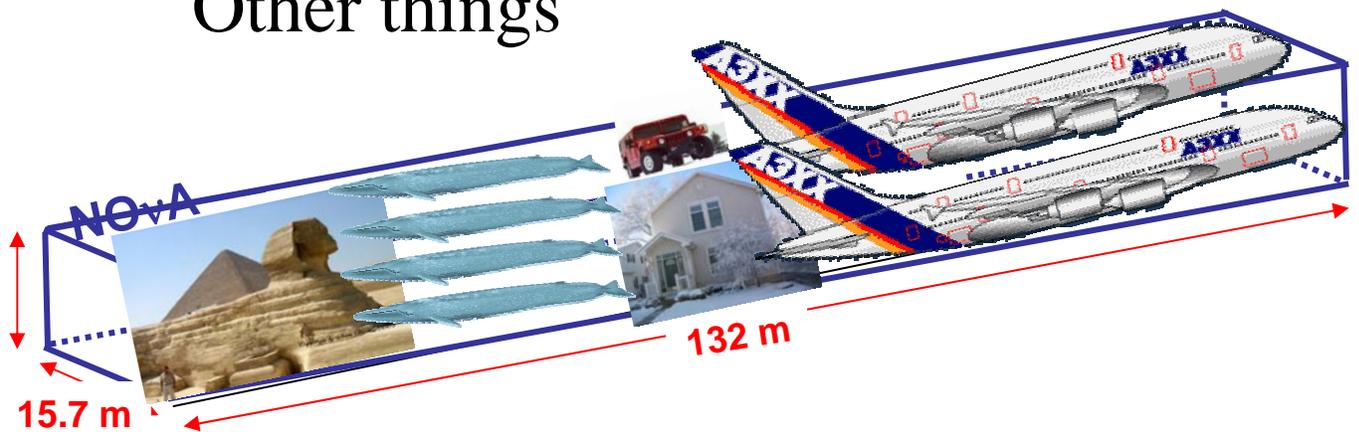
Constraint on angle means γ compensates θ^* to fix P_L (at fixed θ_{lab})

Two Views of the very Large

BaBar, CDF, D0, CMS, & ATLAS



Other things



Next Order and CP Violation

- Do not neglect Δ_{12}
- Add back CP phase δ
- Terms group as

$$\left| e^{-i(\delta + \Delta_{32})} A_{atm} + A_{solar} \right|^2$$

(Before we looked only at A_{atm}^2)

Giving

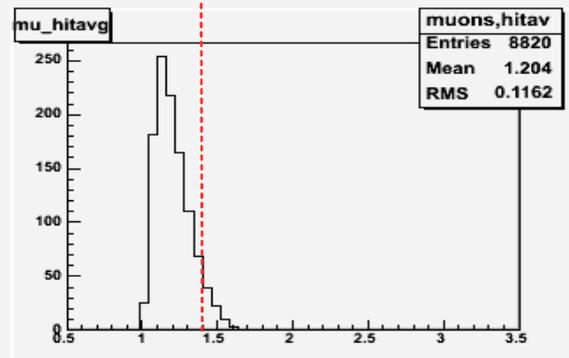
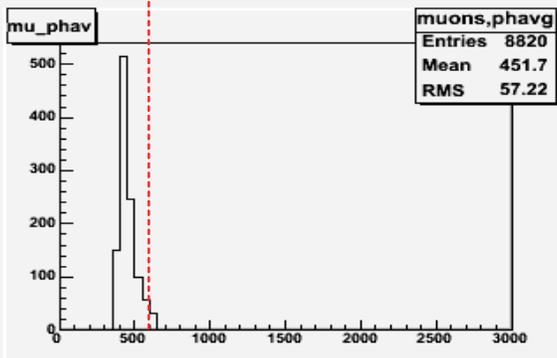
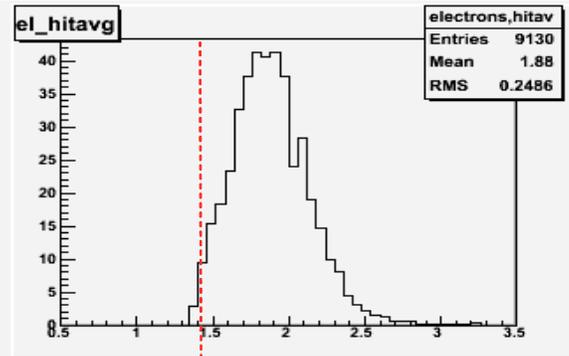
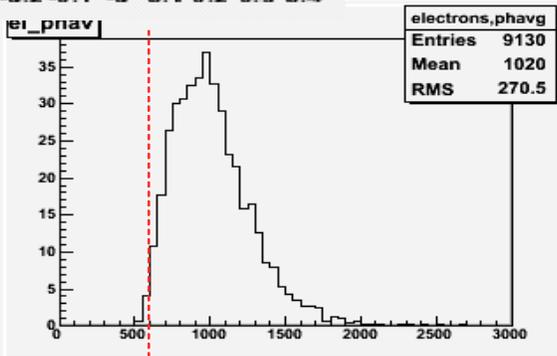
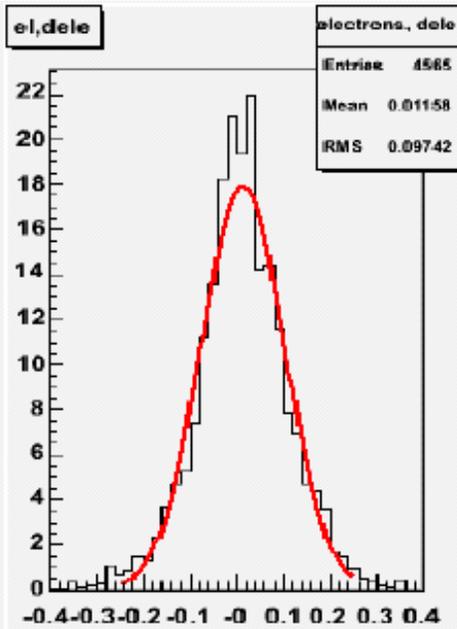
$$A_{atm}^2 + A_{solar}^2 + P_{interference}$$

NOvA performance and event classification variables

Active detector gives good resolution

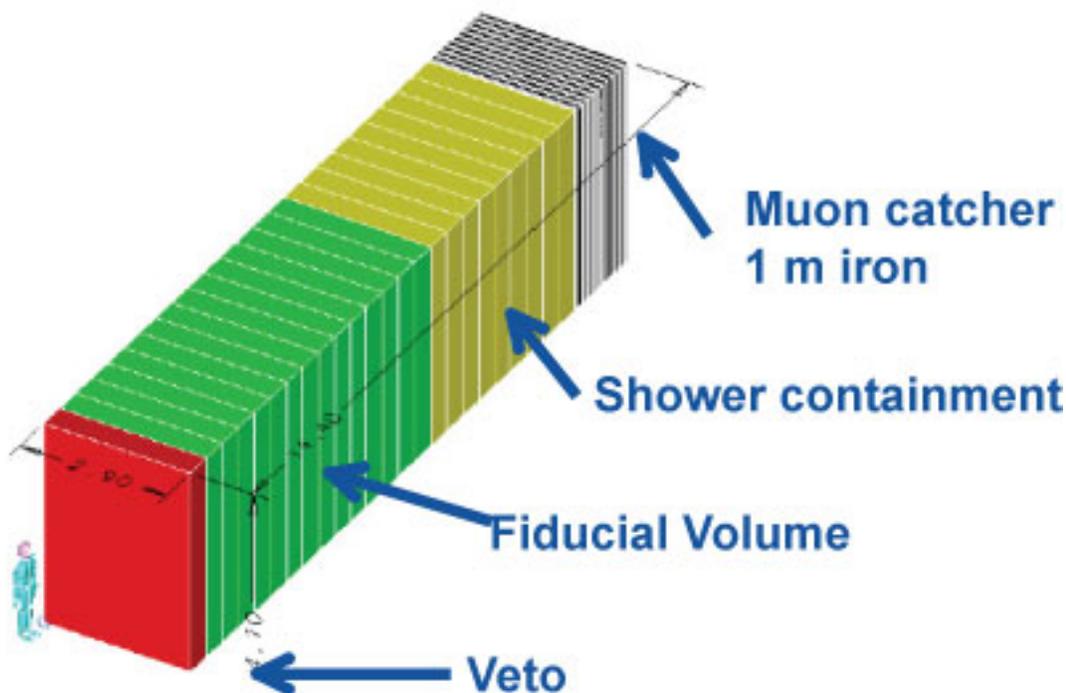
$$\frac{\Delta E}{E} = \frac{9\%}{\sqrt{E}}$$

Overall $\nu_{\mu} \rightarrow \nu_e$ efficiency
 $\sim 24\%$

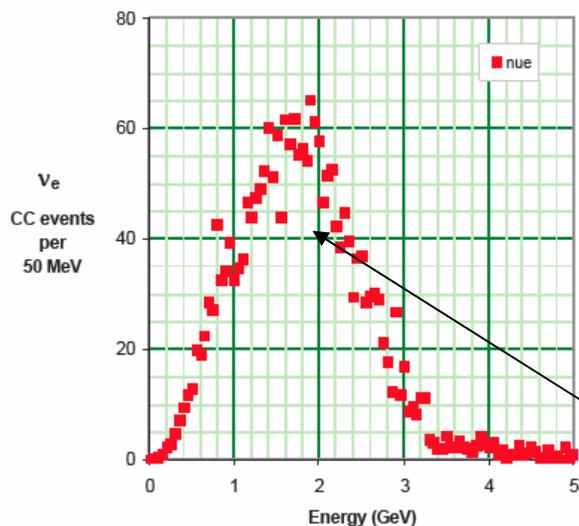


e- μ discrimination using average pulse height/plane and $\langle \text{hits} \rangle / \text{plane}$

NOvA also has a near detector as a crucial component



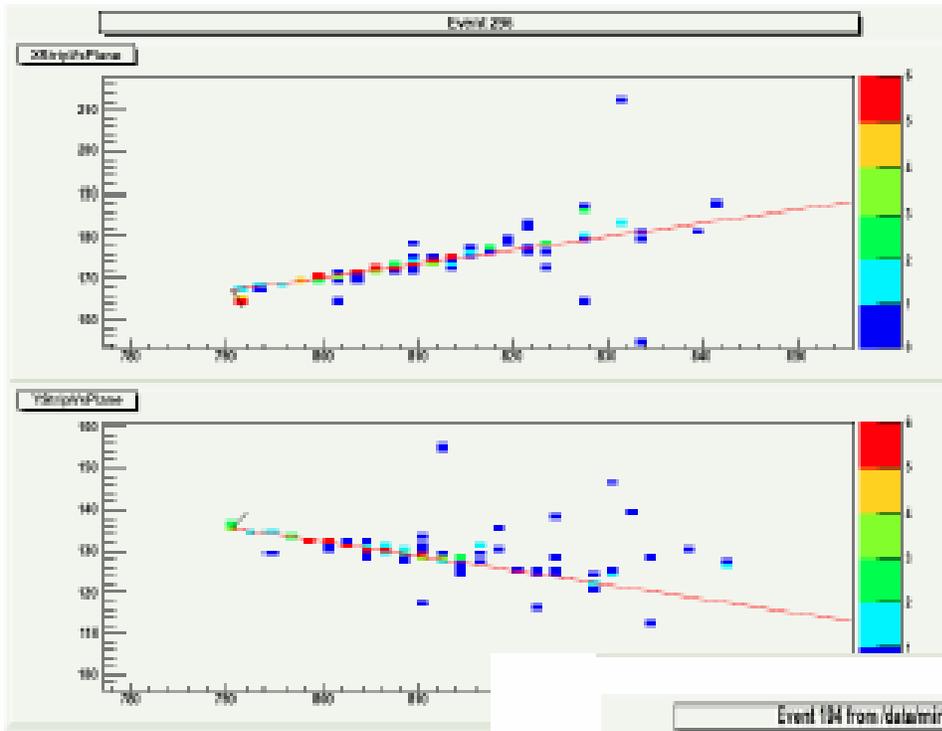
Can move to different locations in underground tunnel to measure off-axis effects and calibrate far spectra



Can also operate early on surface to understand response

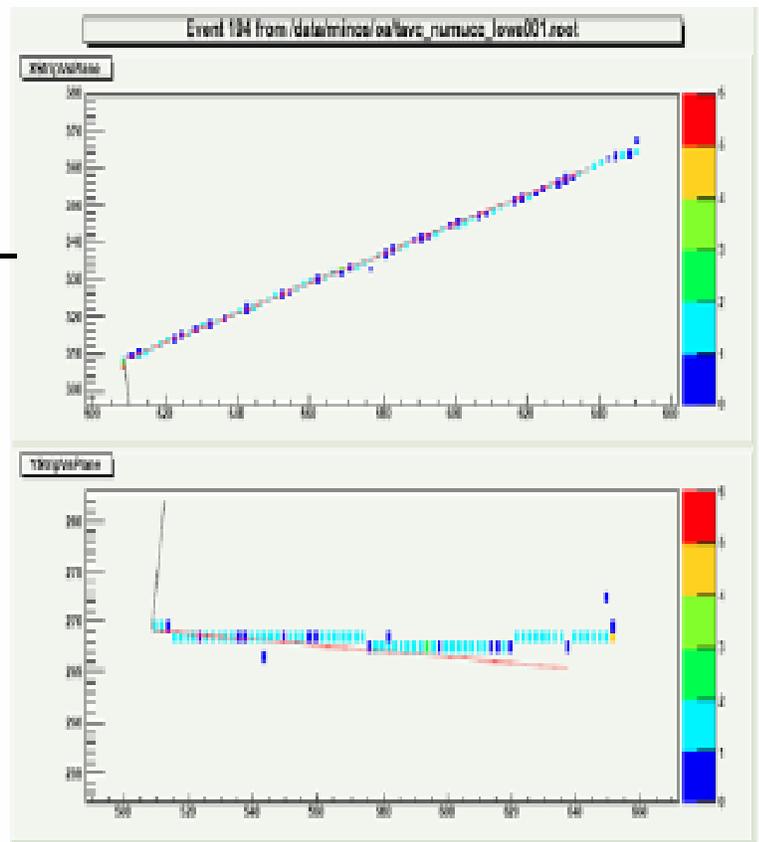
Sample ν_e spectrum from “off-axis neutrino test beam”

Representative NOvA MC Events



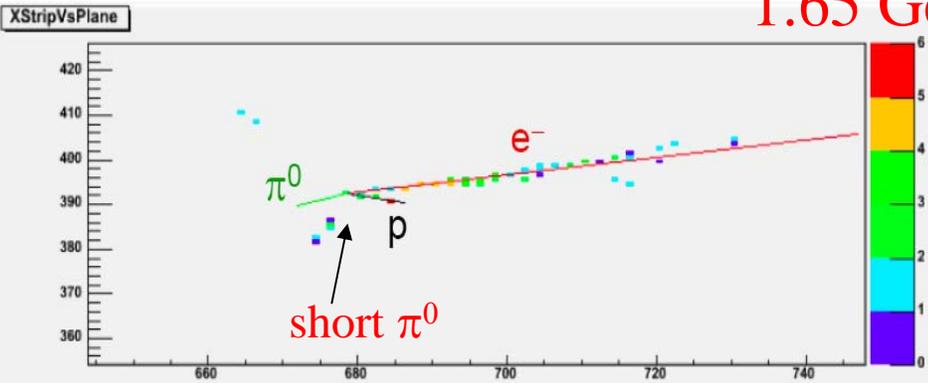
2 GeV ν_e QE
- 60 planes

2.2 GeV ν_μ QE -
140 planes

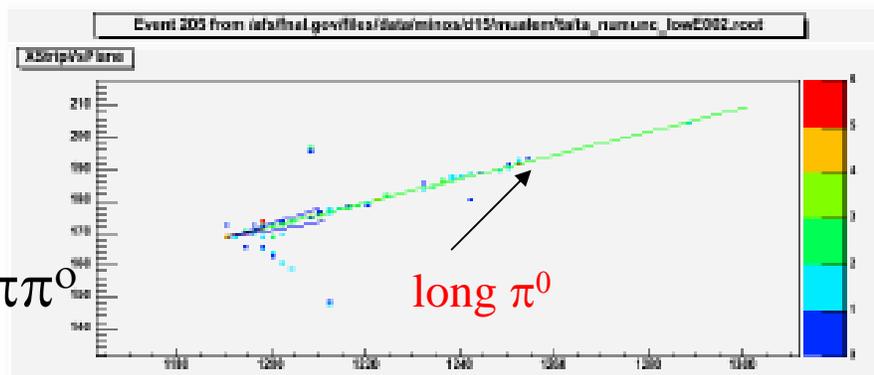
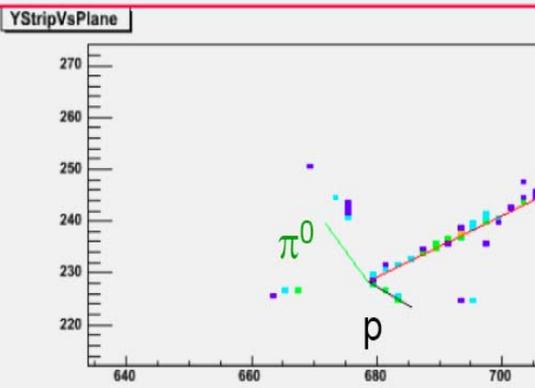


NOvA MC events showing successful and unsuccessful π^0 rejection.

1.65 GeV $\nu_e N \rightarrow e p \pi^0$

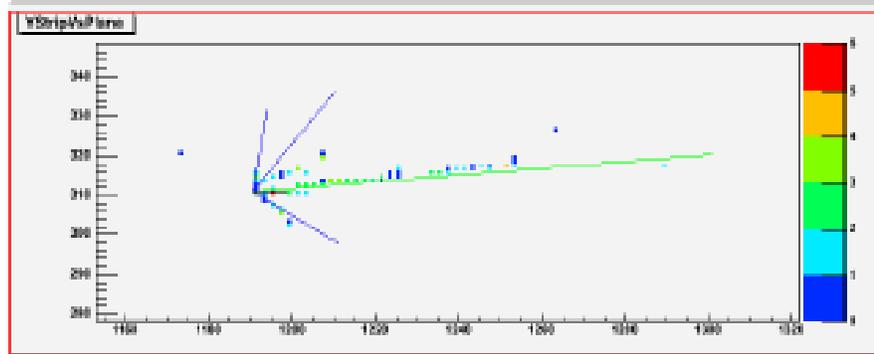


55 planes
Success!!



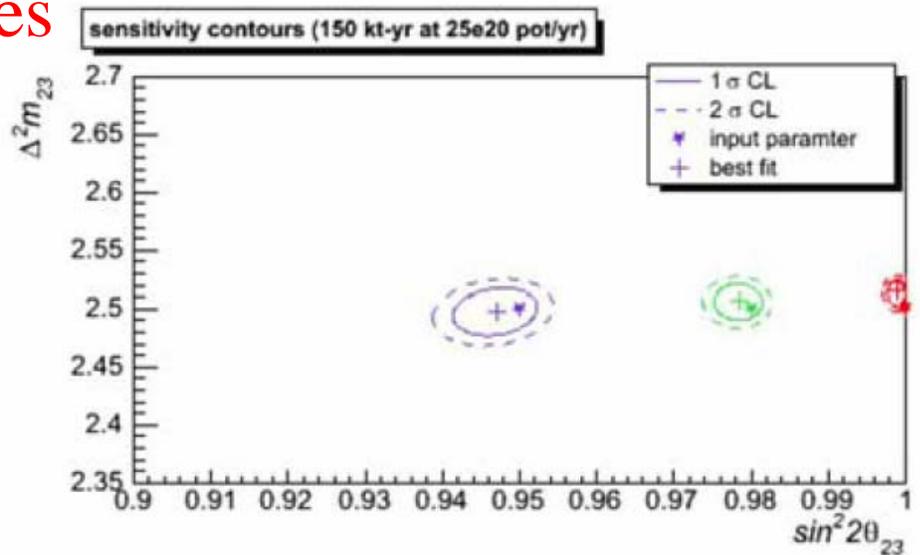
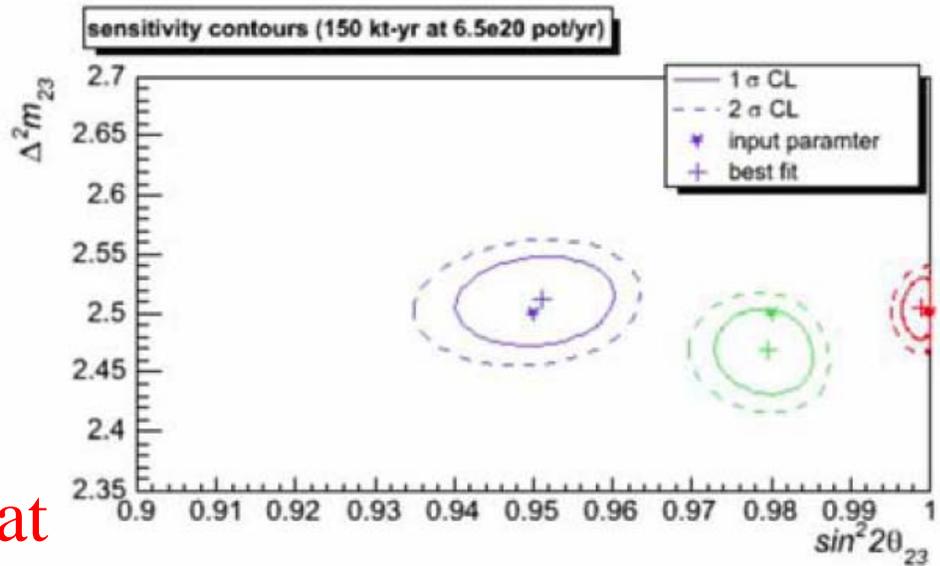
4.95 GeV $\nu N \rightarrow \nu \pi \pi \pi^0$

70 planes
Failure!!



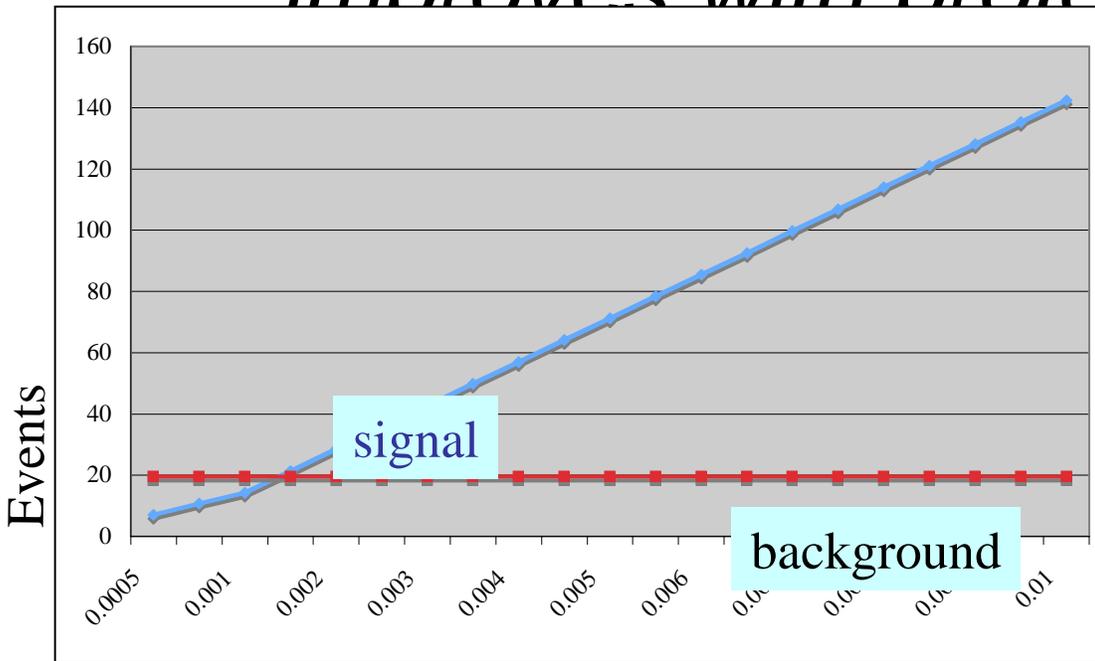
Sufficient running of NOvA
can make do other precision
measurements.

5-year run at
two intensities



Excellent results for $\sin^2 \theta_{23}$ with expected intensities.

Measurement is limited by statistics - experiment improves with protons



0.01

0.05

0.1

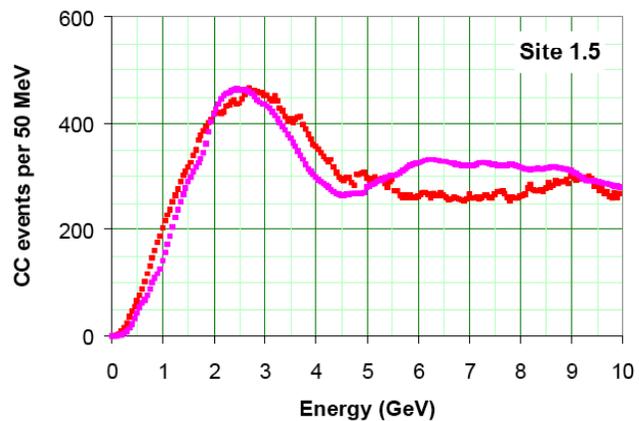
$$\sin^2 2\theta_1$$

For $\sin^2 2\theta_{13} = 0.1$:

ν : S=142.1, B=19.5

$\bar{\nu}$: S= 71.8, B=12.1

5 yrs at $6.5E20$ pot/yr,
efficiencies included



Background ν_e from near
detector

Graphic - R. Bernstein

The Fine Art of Small Masses

- **Direct searches**
 - Look for deviations at very end of β -decay spectrum.
 - Time-of-flight spread in Supernova !
 - Cosmological “constraints”
 - **Don't exceed allowed density**
 - Accelerator-based studies of π , τ decays.
- Current state is limits only.
- **Resort to indirect techniques**
 - Oscillations and mixing!

Decays

Astronomy

ν_e	$< 3 - 7 \text{ eV}$	$< 20 \text{ eV}$ (SN 1987A)
	$< 1 \text{ eV}$ (“double- β ”, model-dependent)	
ν_μ	$< 0.170 \text{ MeV}$	$\sum m_\nu < 8 \text{ eV}$ (cosmology)
ν_τ	$< 24 \text{ MeV}$	

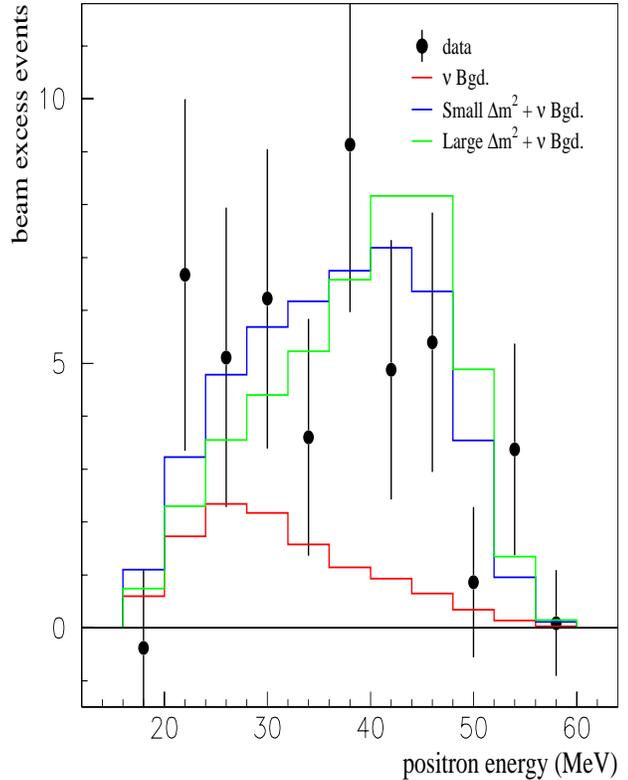
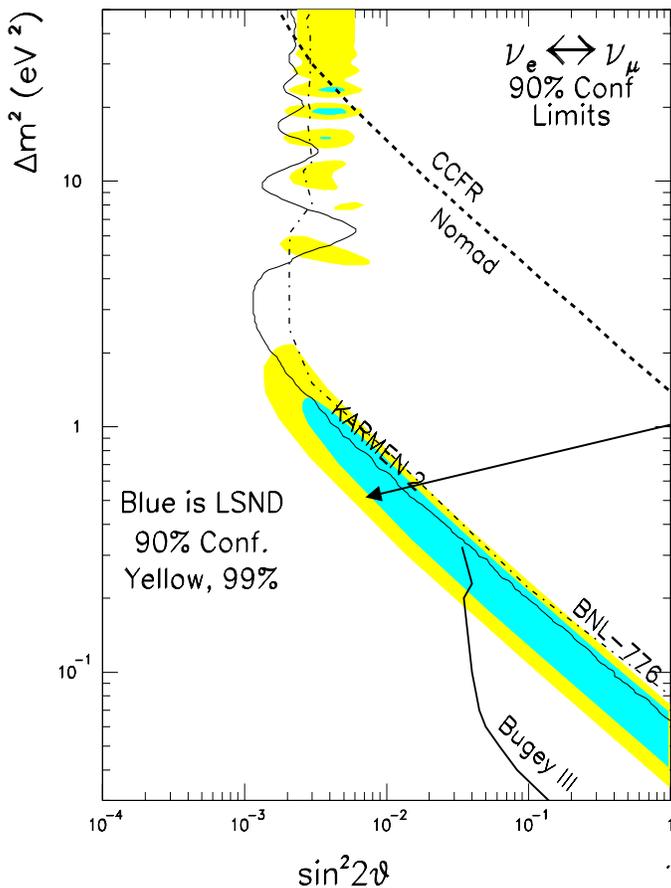
Hints of Oscillations # 3

LSND - Stopped Pion Decay



Use stopped π, μ decays
as source of anti - ν_μ

Detect positron
and neutron from CC.

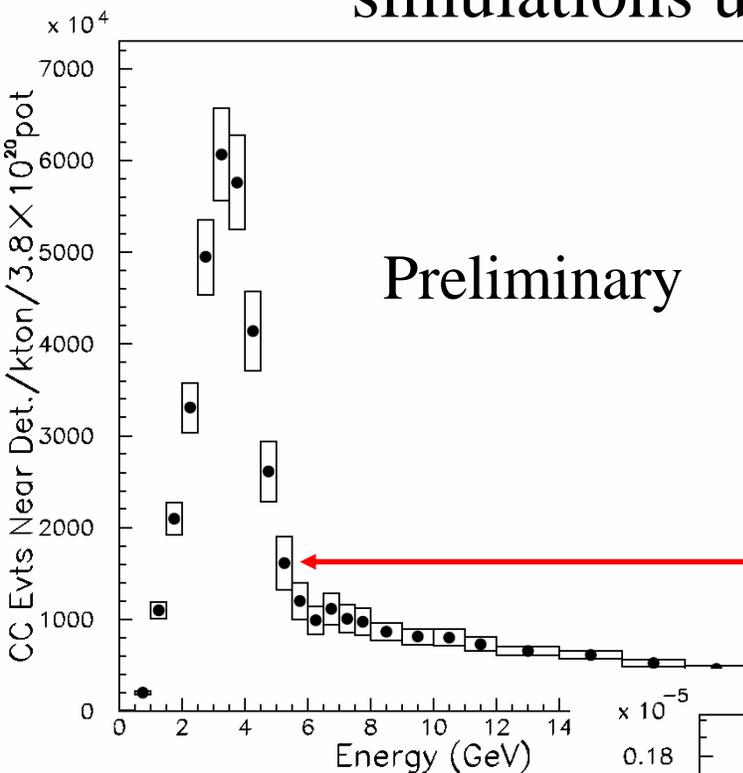


Recent analysis shows compatible with other experiments - but close.

Fermilab MiniBOONE
experiment is checking

Two detectors to reduce uncertainty

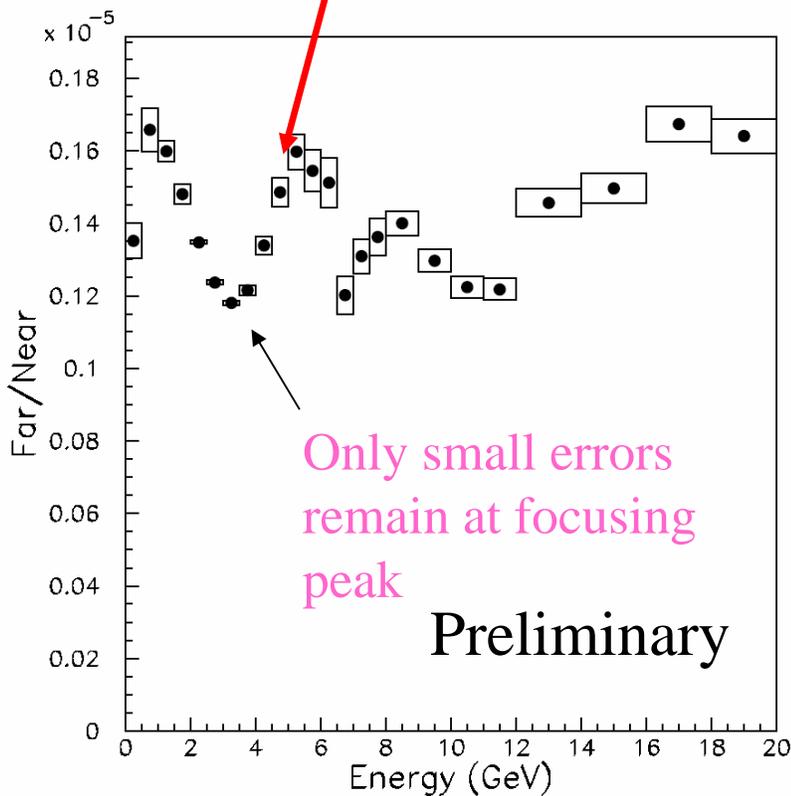
Examples from MINOS
simulations using FLUKA



Hadron production
uncertainty dominates
this plot

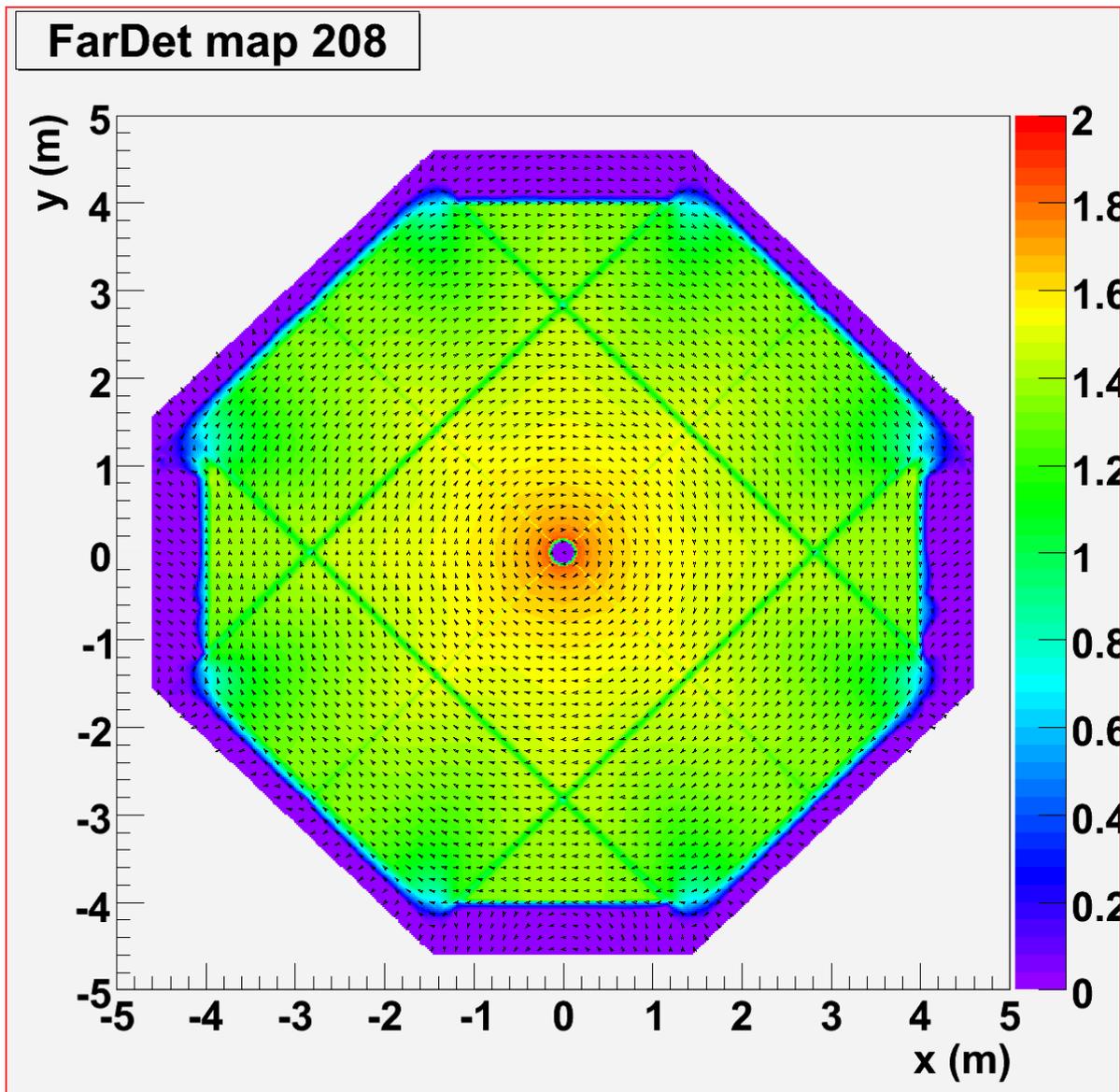
At 5 GeV, uncertainty of
 $\sim 15\% \rightarrow \sim 4\%$

In far/near ratio,
hadron uncertainty
largely cancels



Tracking must go on in magnetized iron

Determine field map by finite element analysis, steel properties.



MINOS Far Detector B-field (Tesla)



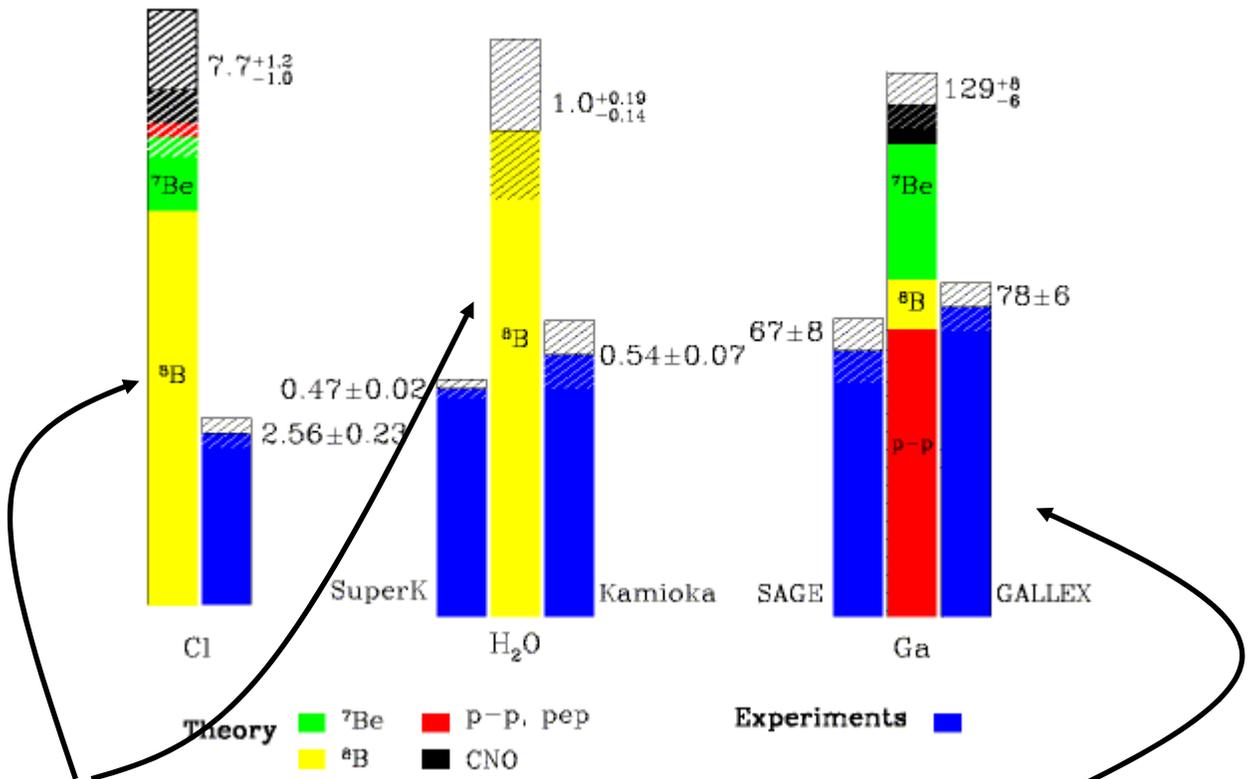
Solar Neutrinos - Observations

Difficult experiments using 3 different technologies have detected the solar neutrino radiation

Sensitive to different mixes of channels

All have reported significant deficits

Total Rates: Standard Model vs. Experiment
Bahcall-Pinsonneault 98



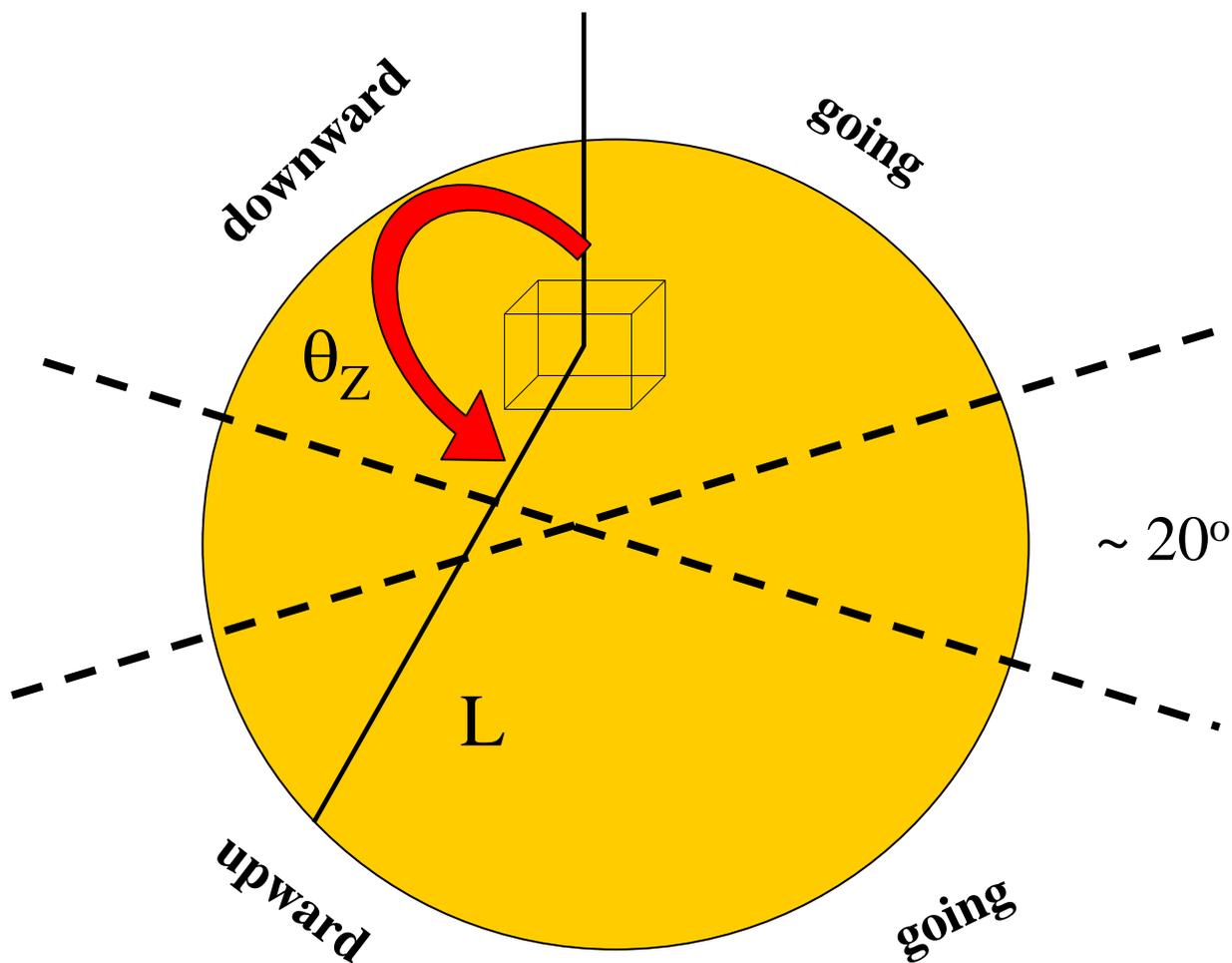
Correcting B increases Be

- not seen.

pp alone, no room for Be



Zenith Angle as a Clue to Oscillation



ν path lengths vary from 20 - 13,000 km.

Wide energy ν energy spectrum with
 ~ 1 GeV "typical"