

Two Lectures on Heavy Ion Physics

CTEQ

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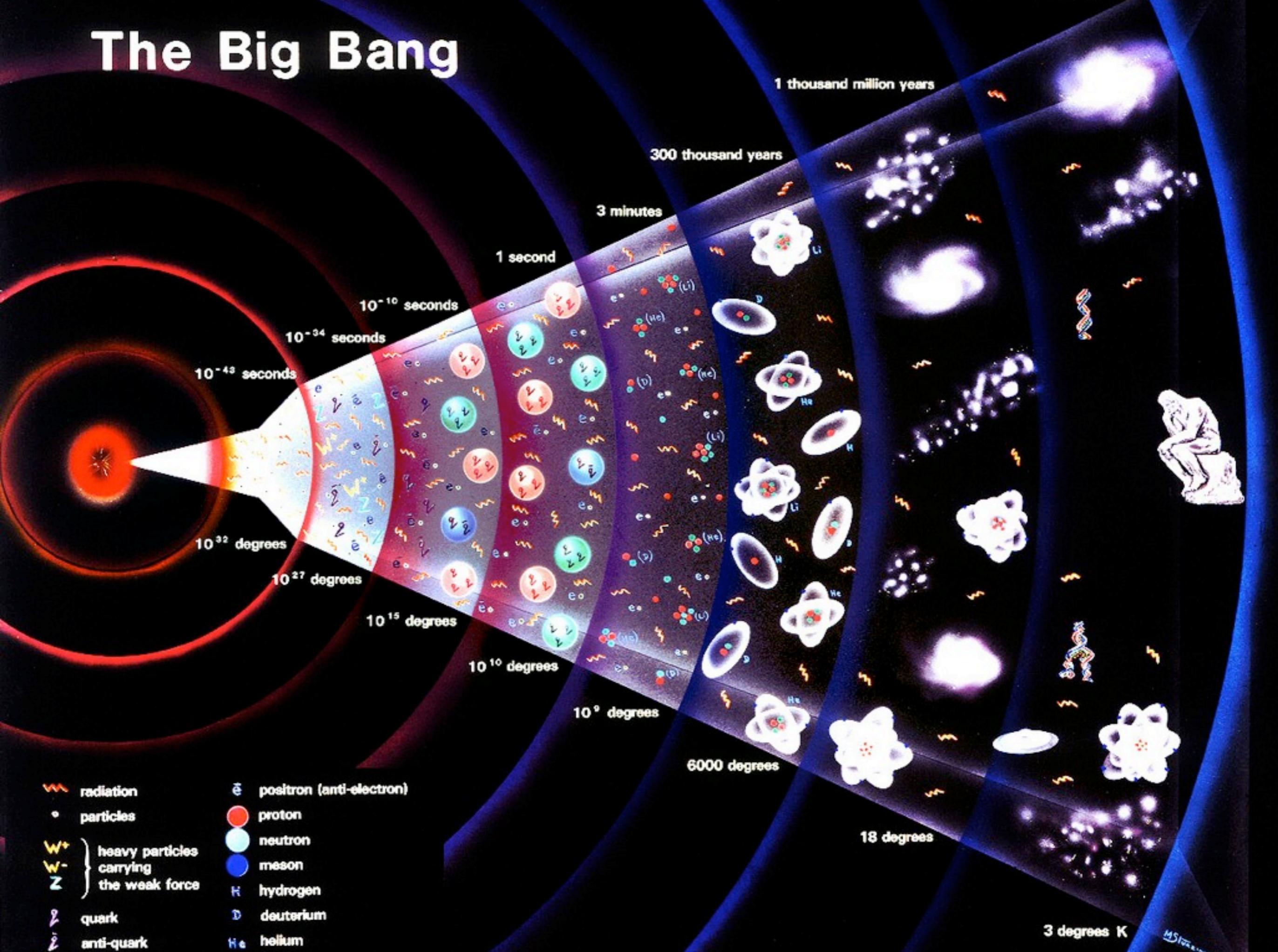


CTEQ Summer School

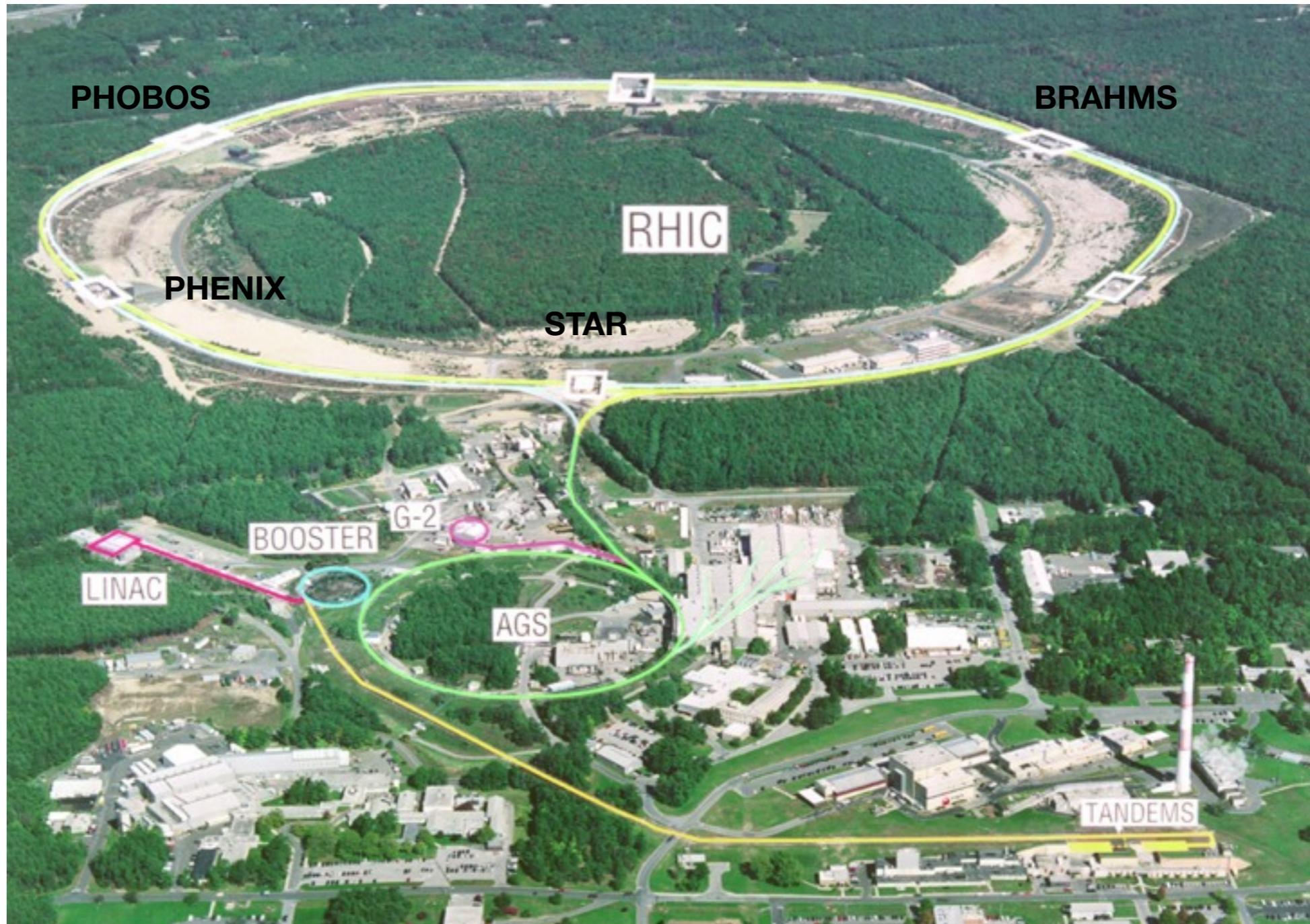
July 2, 2009

Madison, WI, USA

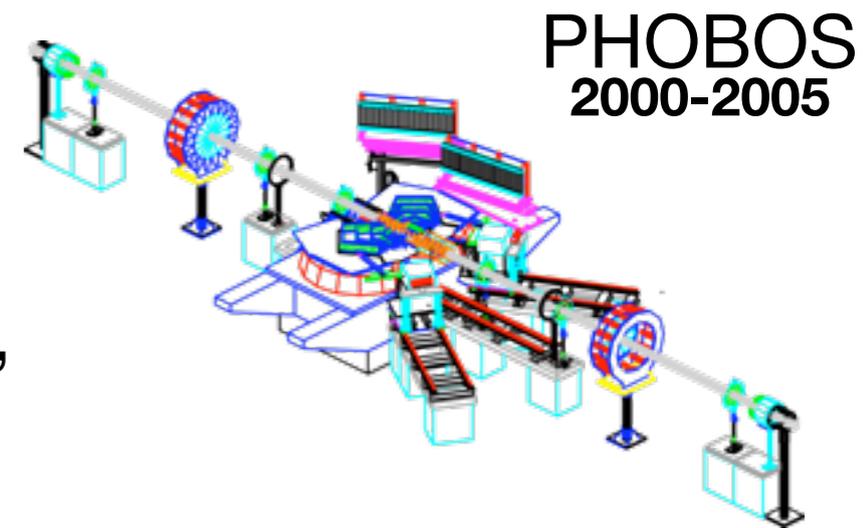
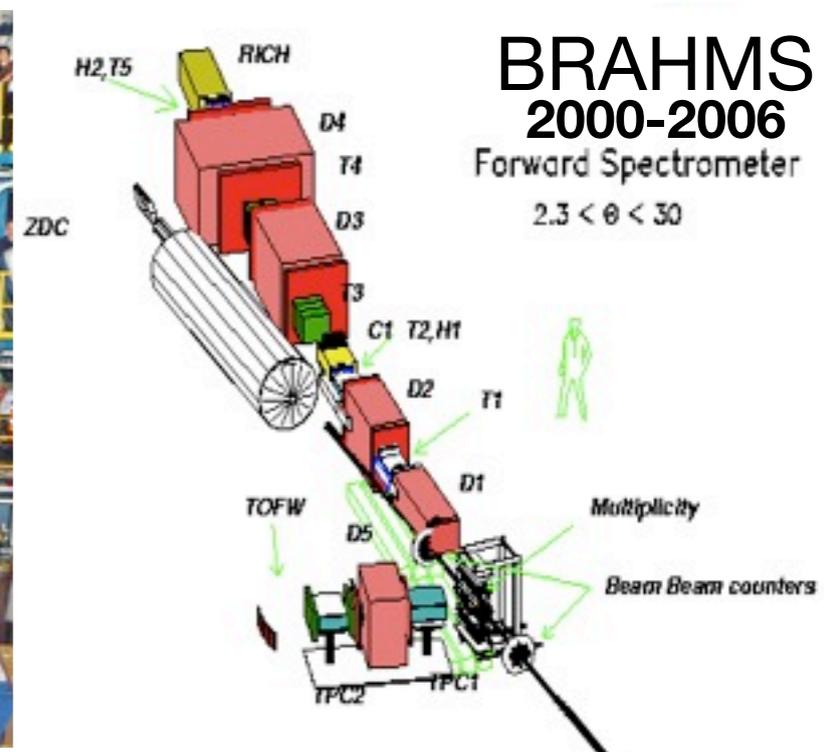
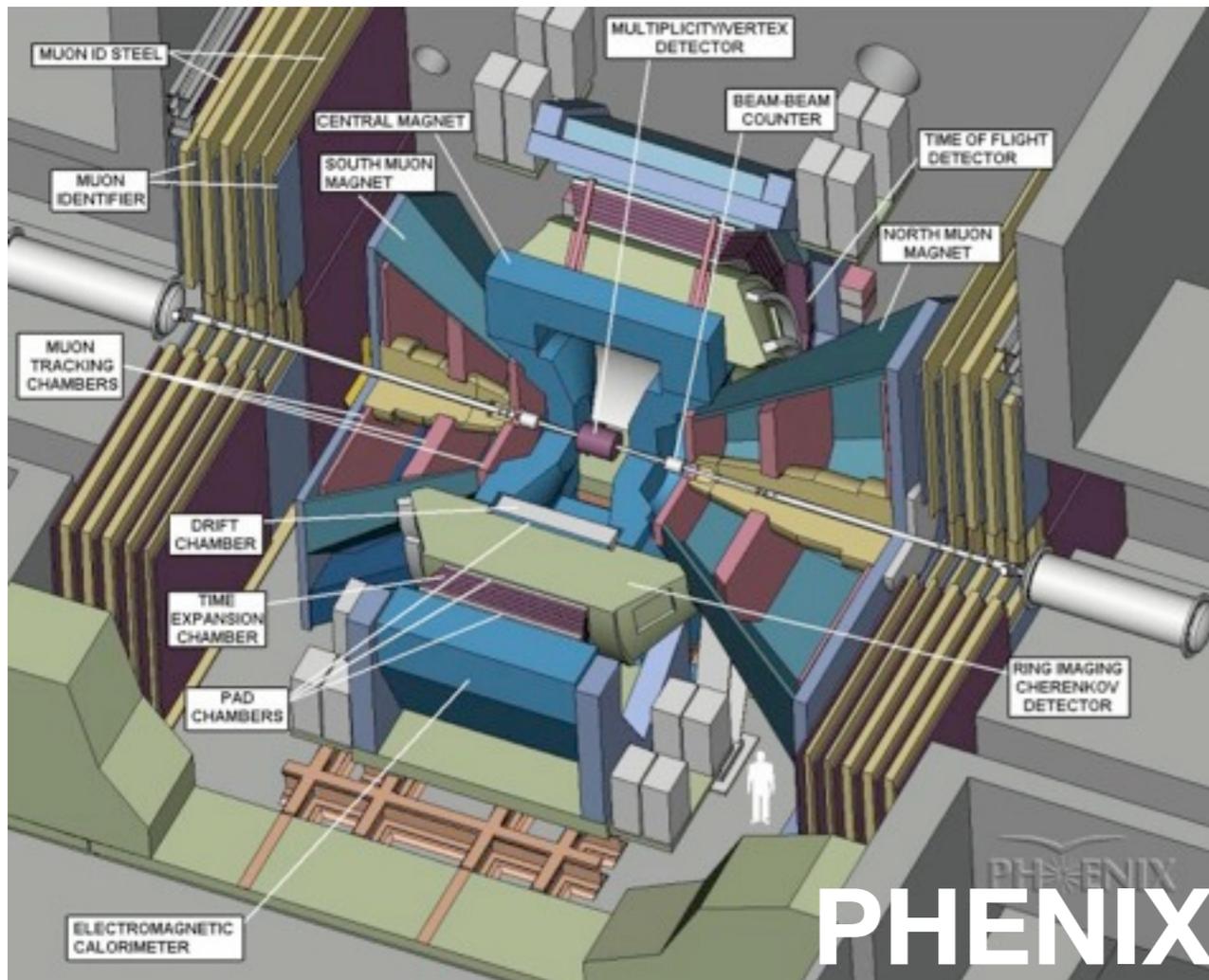
The Big Bang



Relativistic Heavy Ion Collider (RHIC)



RHIC Experiments to Scale



RHIC program has large and small detectors,
sufficient overlap to make cross-checks.

Billions of events since 2000:

p+p, d+Au, Cu+Cu, Au+Au, 19.6-200 GeV

Heavy Ion Collisions: Soft Physics

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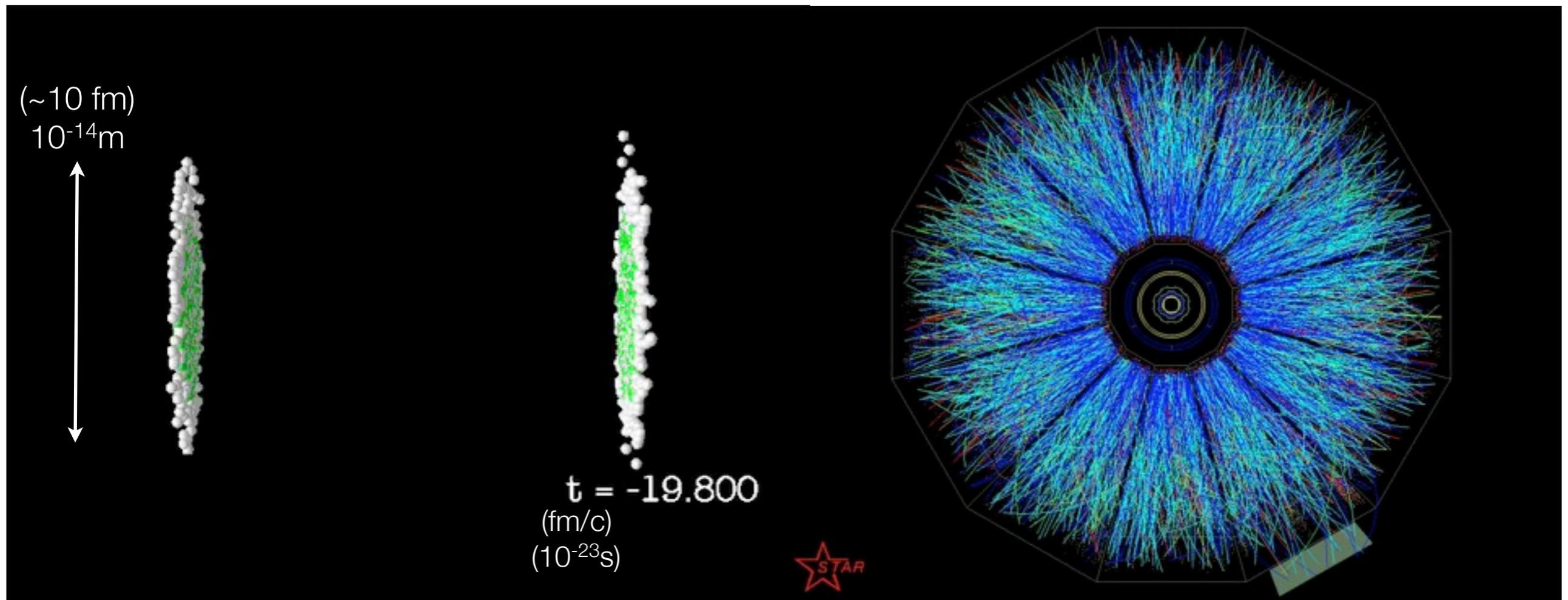
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RHIC Physics in a Nutshell

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Collisions of Ions

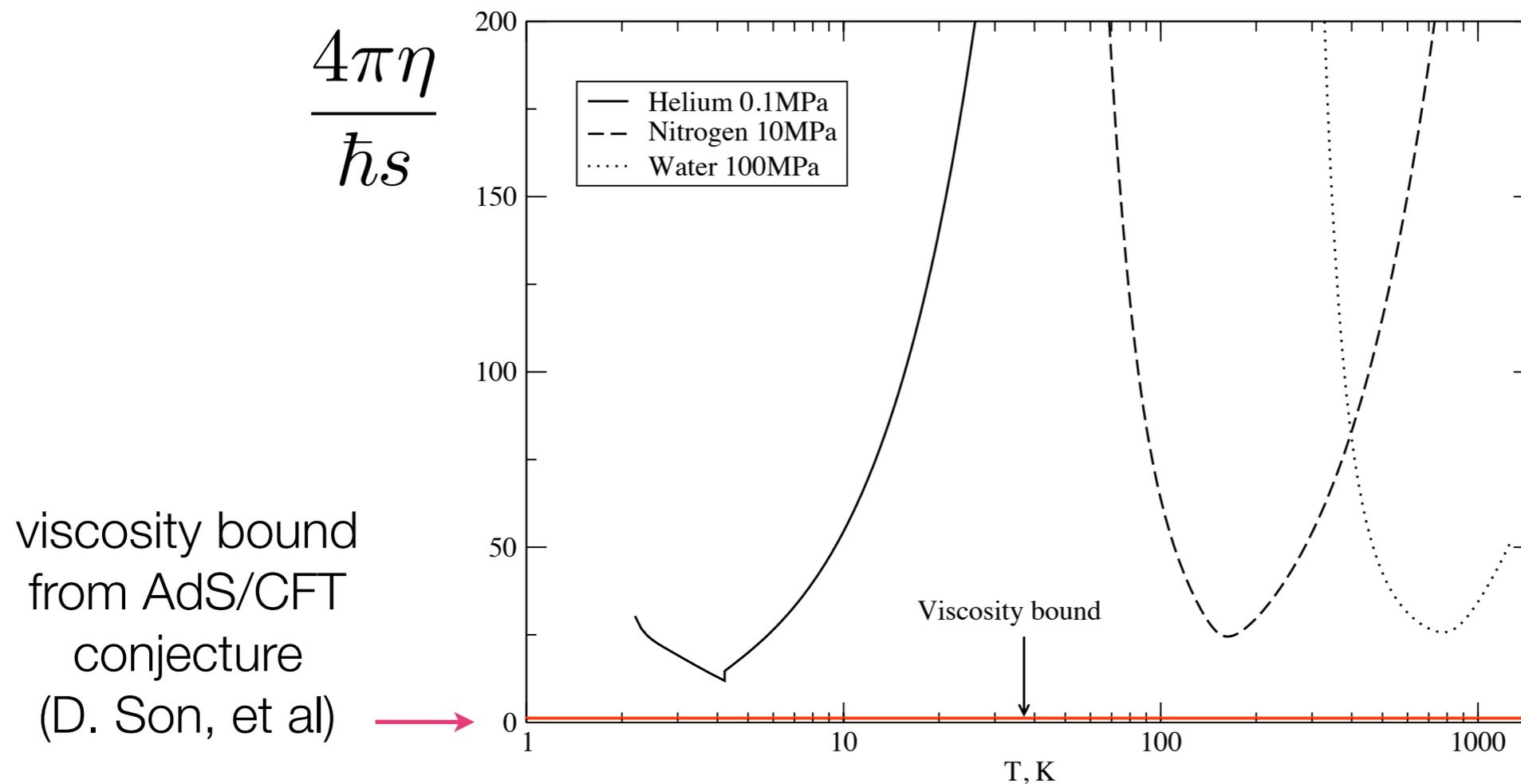
1000's of Particles



RHIC physics takes place in **space-time**
Need to “rewind” dynamical evolution to study
QCD at high temperature and density

Bottom line: The most-perfect fluid?

We say RHIC collisions behave like a “**perfect liquid**”
(or at least a “near perfect fluid”)



$$\frac{\eta}{s} \geq \frac{\hbar}{4\pi k_B}$$

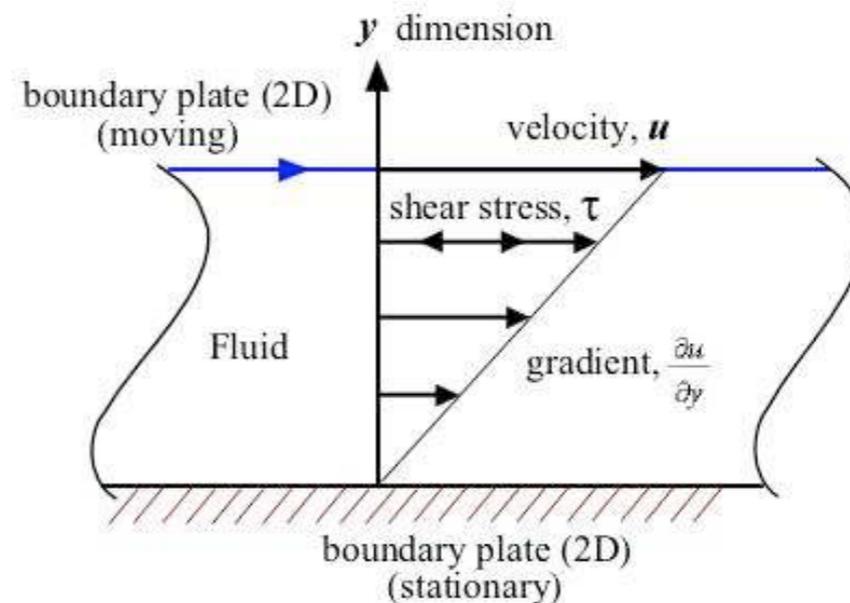
local thermalization: **strong coupling** → low viscosity

These conclusions originally emerged from “soft” physics

What is viscosity?

- From wikipedia (via Dam Son), shear viscosity is the flow of momentum between adjacent fluid layers

- Deviation from local equilibrium
- Frictional (i.e. dissipative) force



- Simple way to think of it is as presence of a fundamental length scale in the dynamical evolution

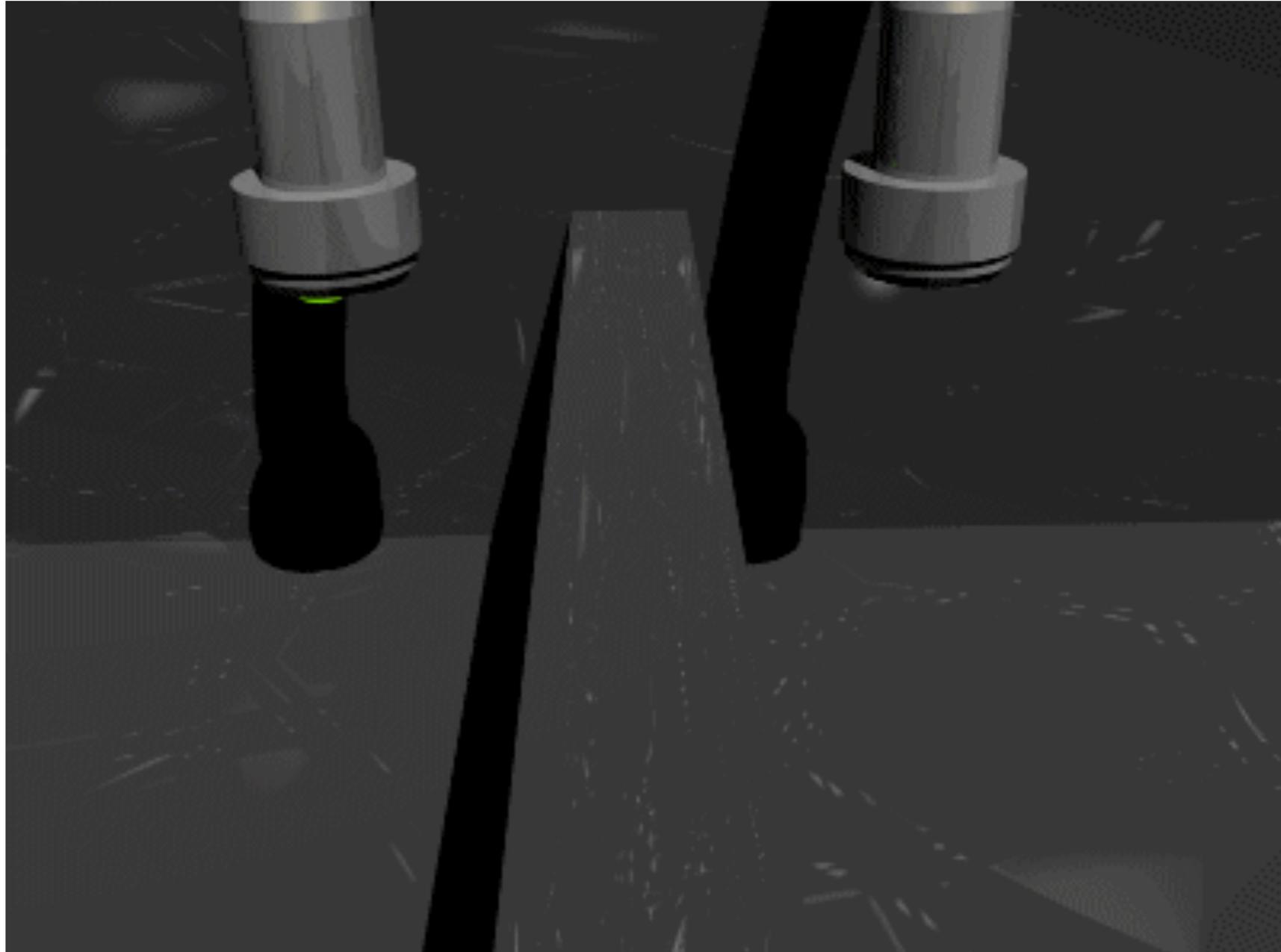
- Scale over which energy/momentum can “escape” from a fluid cell

- Ideal gas has infinite shear viscosity! $\eta = \rho v l$ $l = \frac{1}{n\sigma}$ $\eta \sim \frac{1}{\sigma}$

Viscosity Illustrated

en.wikipedia.org/Viscosity

low
resistance
to flow



propagates
ripples

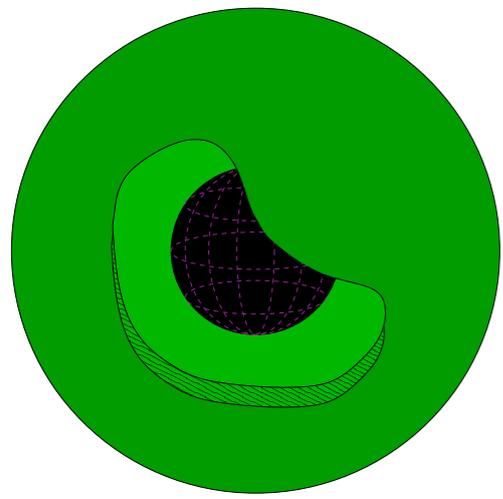
high viscosity (e.g. honey)

low viscosity (e.g. water)

The least-perfect fluid

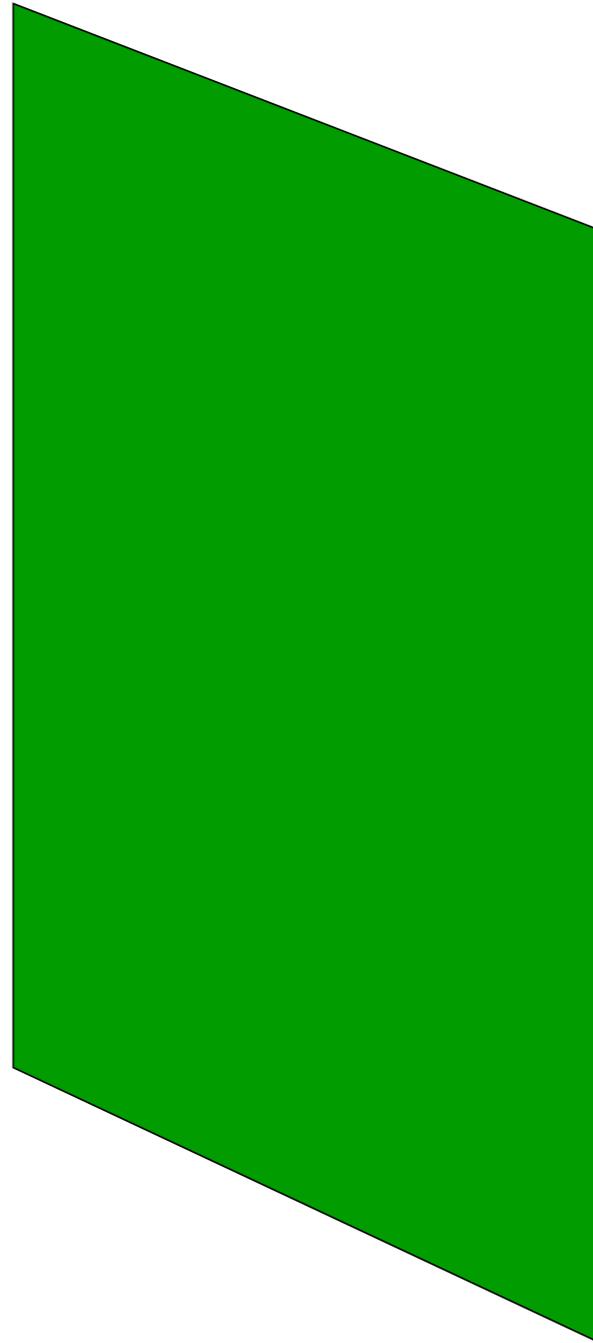


The Pitch Drop Experiment: <http://www.physics.uq.edu.au/pitchdrop/pitchdrop.shtml>

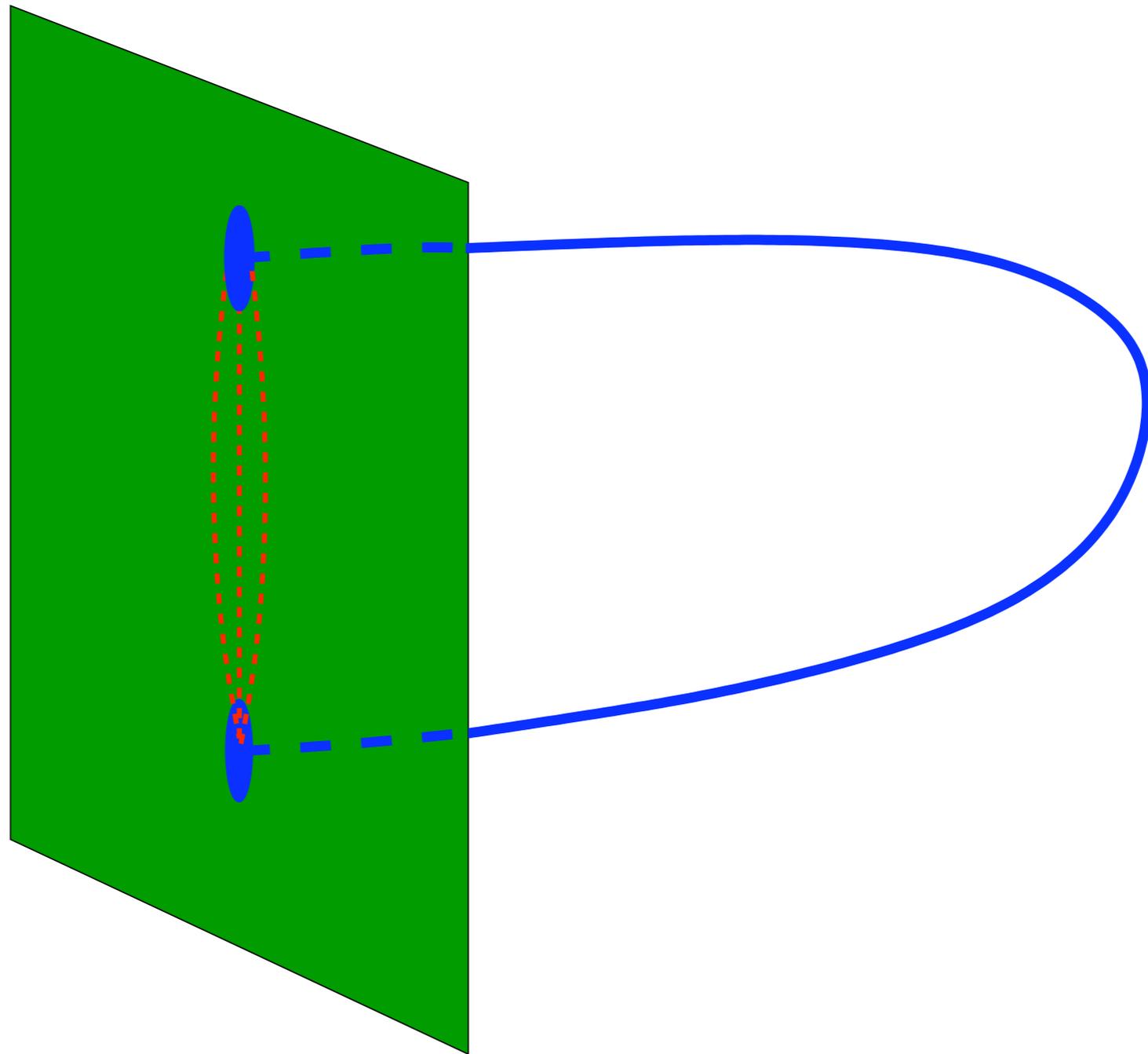


AdS/CFT calculations put
a black hole in the 5th
dimension (thermalization)

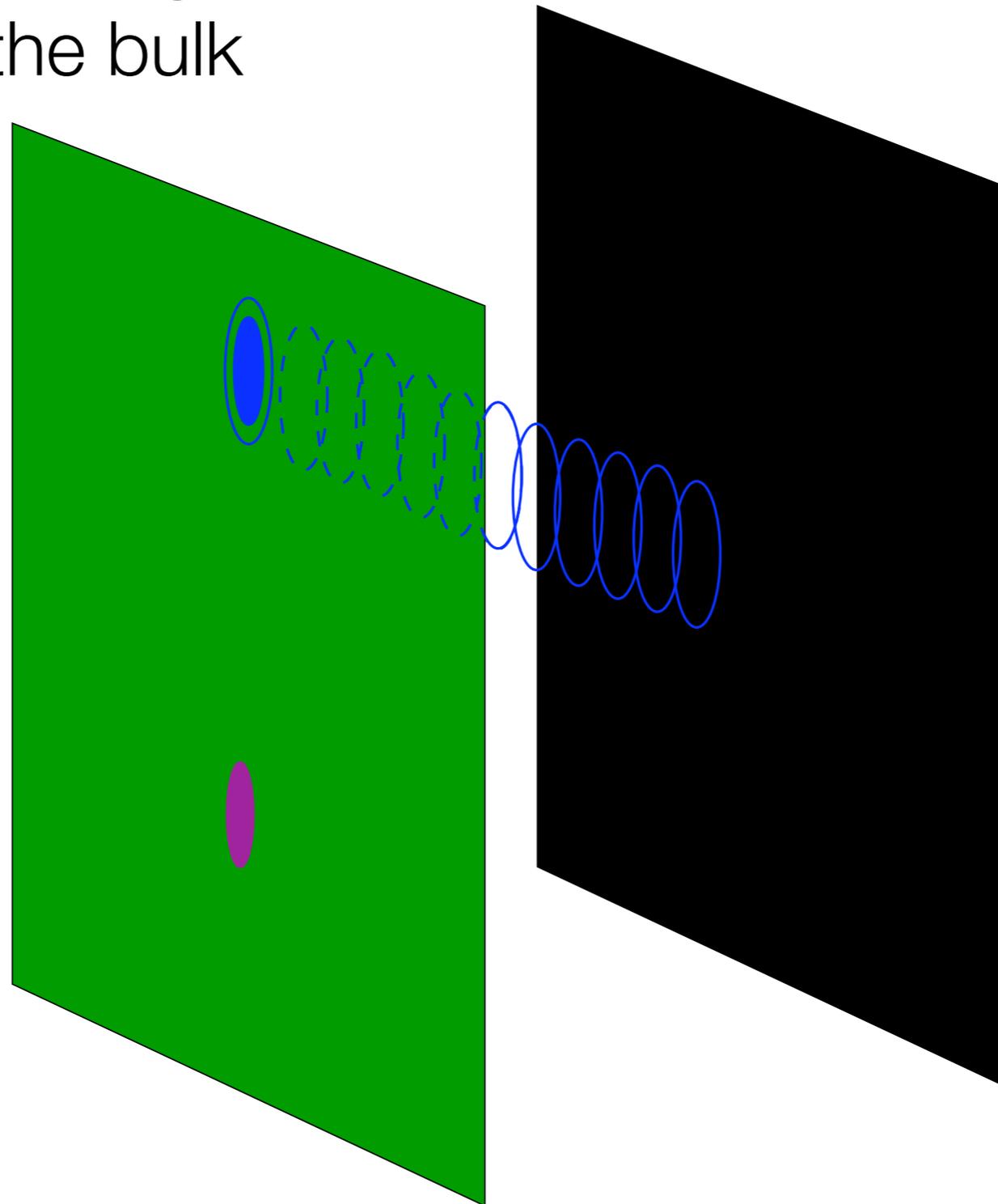
“unwrap”
the spheres
for clarity...



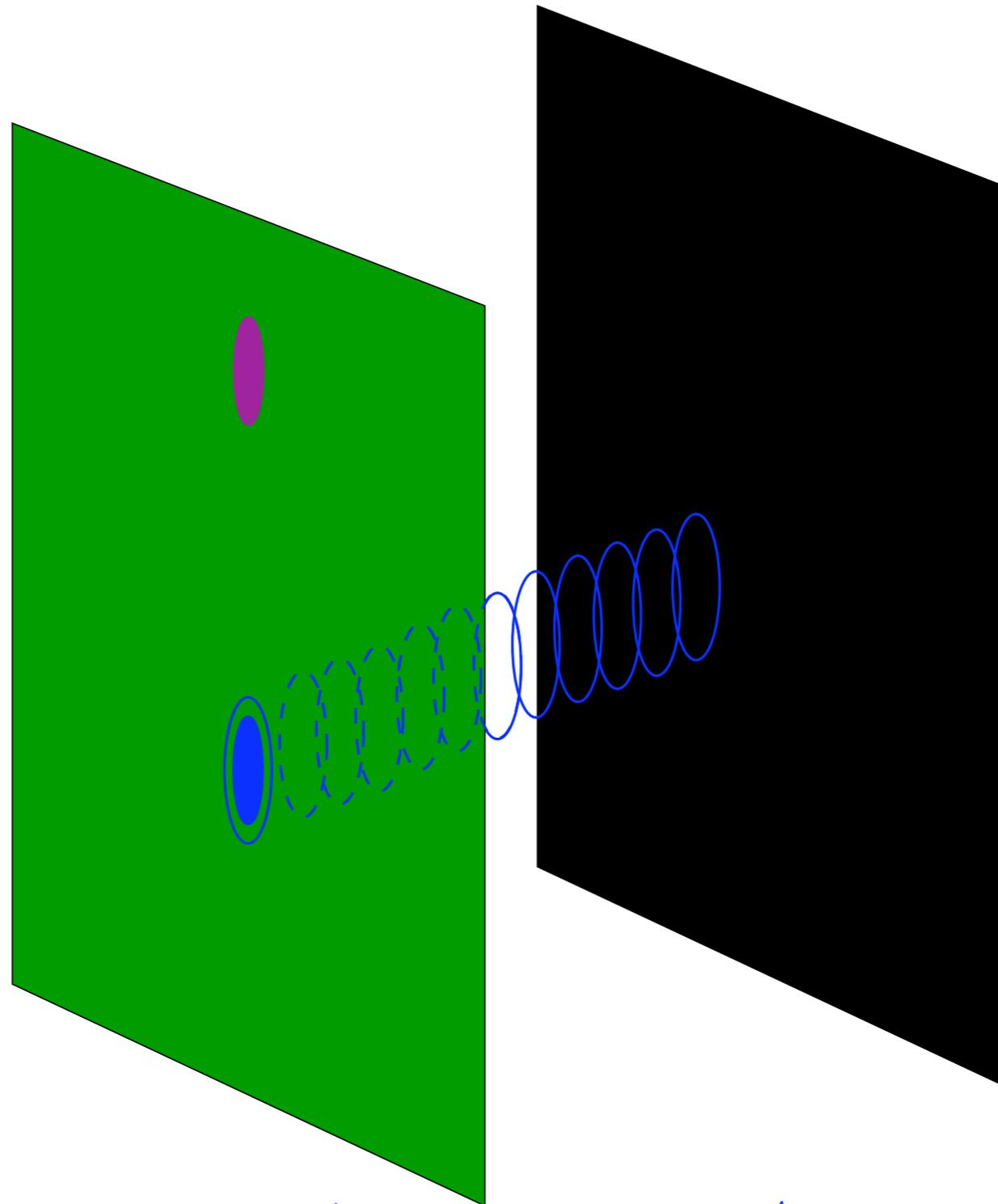
“QCD” in AdS/CFT is a string trailing into the bulk



A quark fires a graviton
into the bulk



And scatters off
of the black hole



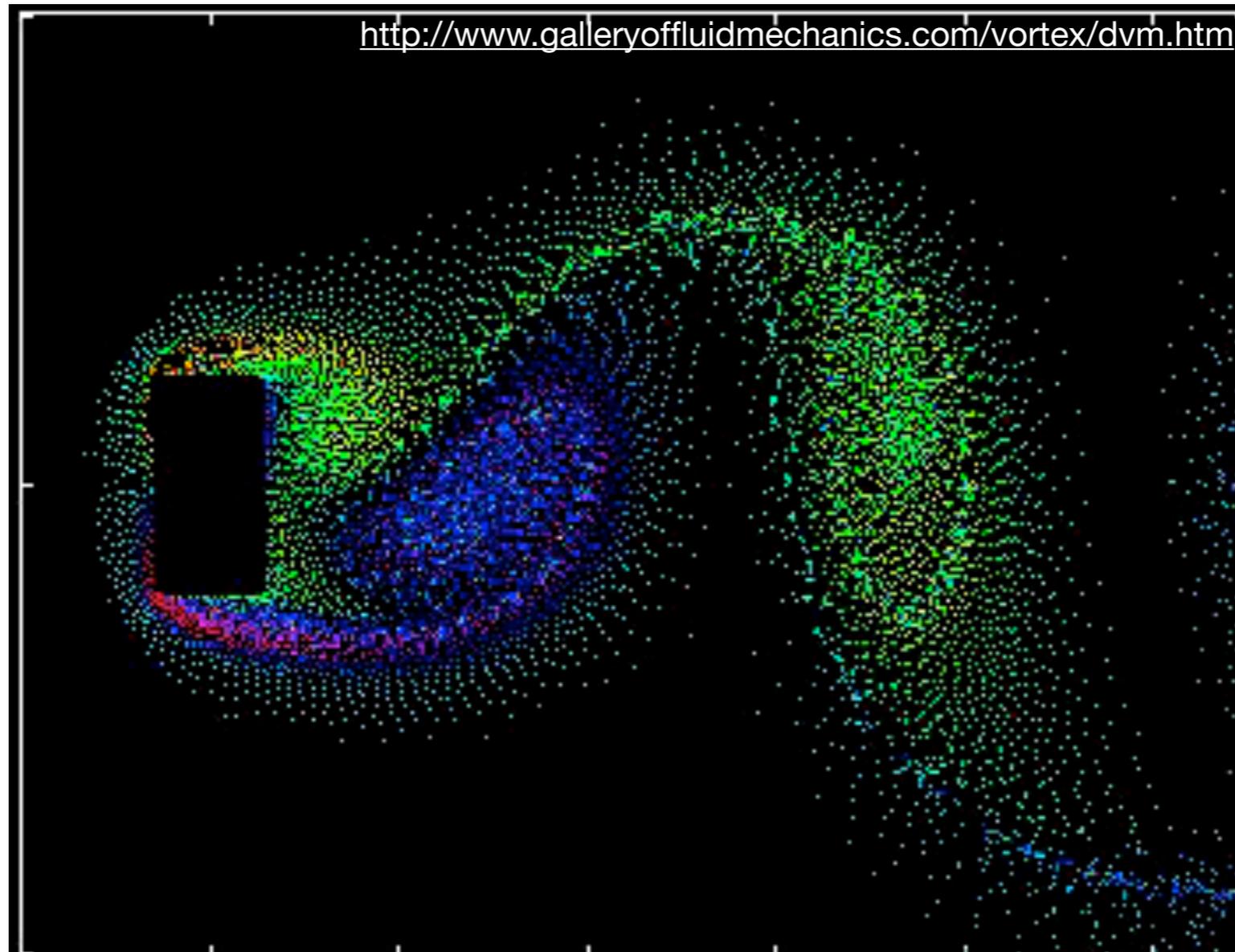
Black hole
absorption
is dissipative

$$\eta = \frac{A}{16\pi G}$$

$$S = \frac{A}{4G}$$

$$\frac{\eta}{S} = \frac{1}{4\pi}$$

Viscous Hydrodynamics



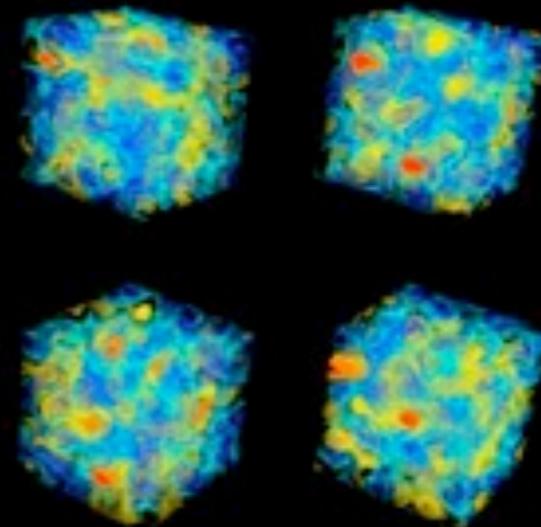
Viscosity introduces new dimensions to hydrodynamic phenomena

Dynamical Regimes of Hot QCD

Strongly
Coupled

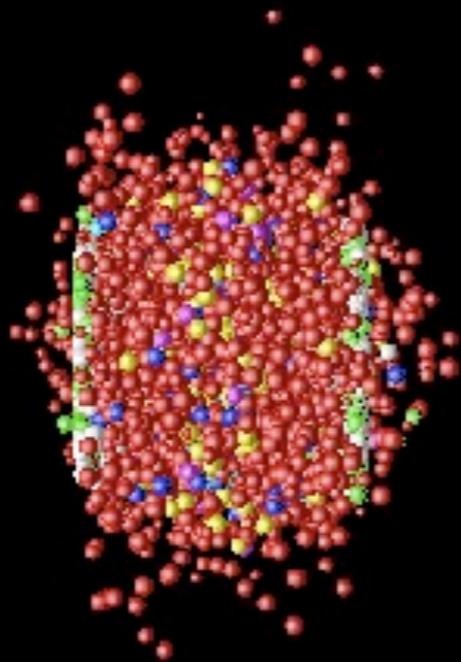


Near-Perfect Fluid?



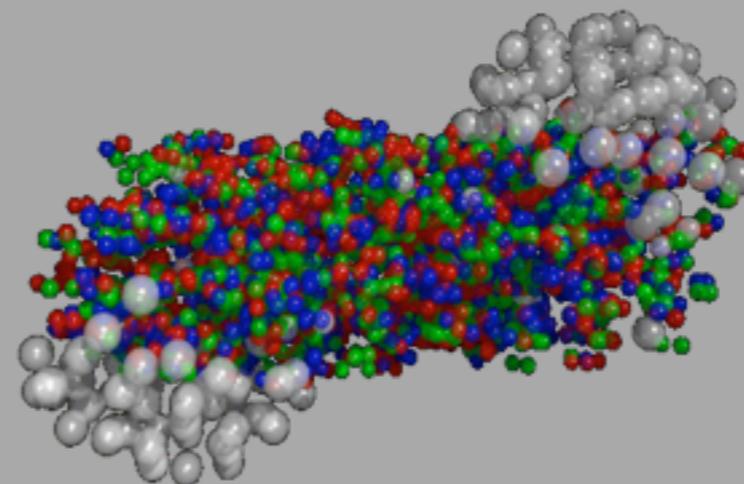
Lattice QCD?

Weakly
Coupled



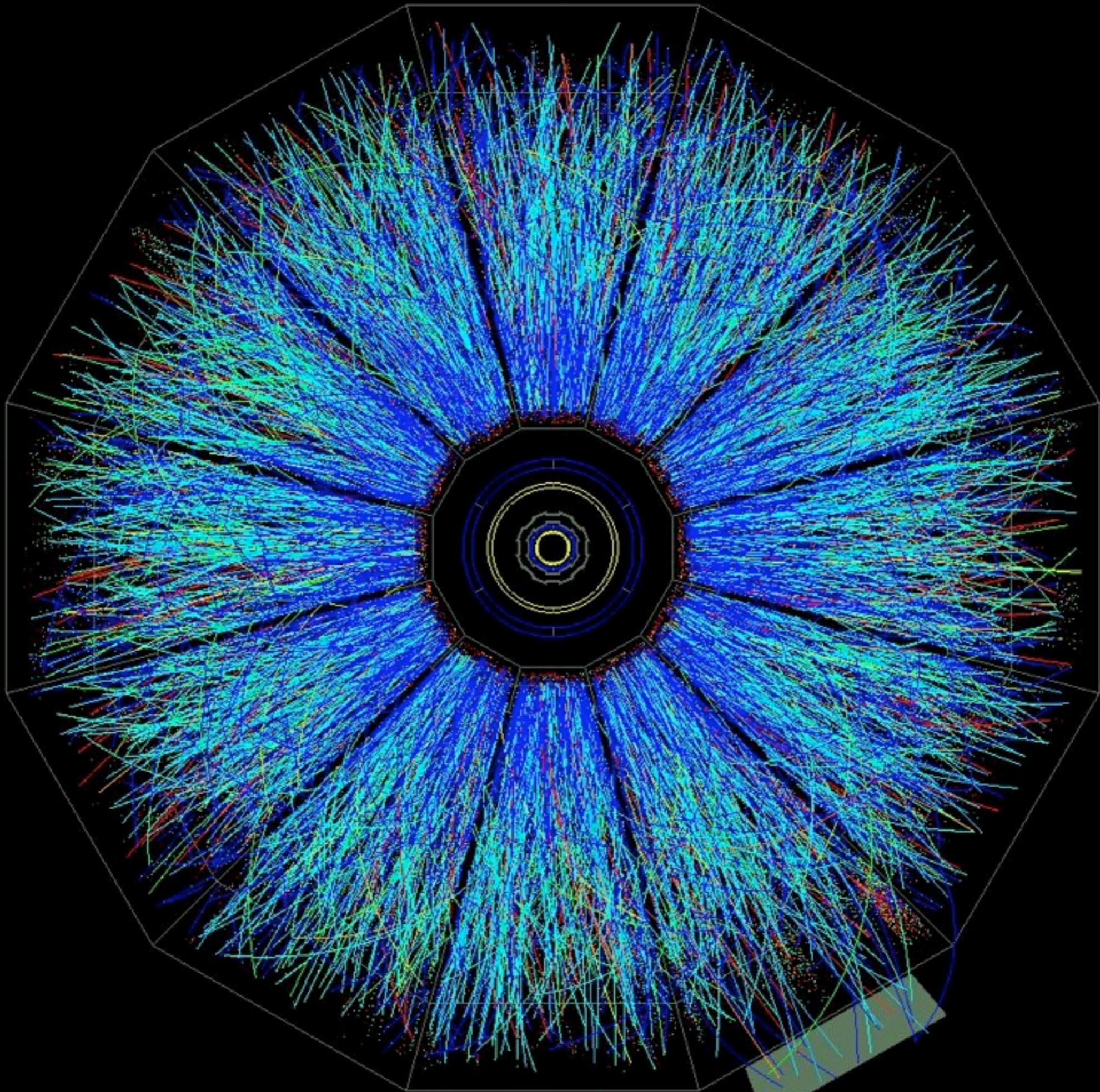
pQCD Cascade

Short Times

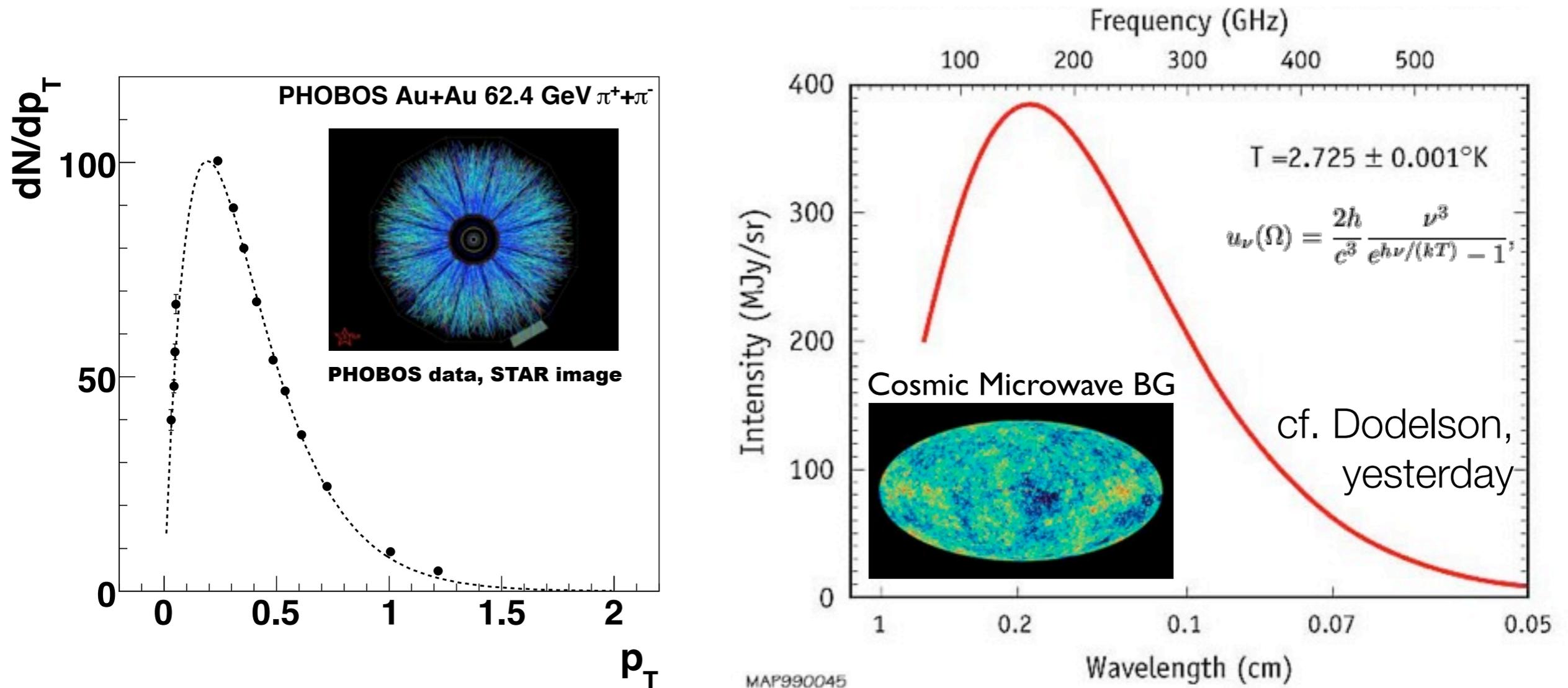


Hadron Cascade

Long Times



Strong Blackbody Radiation



It is often overlooked that the spectra of particles emerging from a heavy ion collision is nearly blackbody but with hadrons instead of photons: **thermal system**

Hagedorn Temperature



1

Rolf Hagedorn predicted
bound state spectrum
rises indefinitely →
Singularity at
limiting temperature
 $T_H \sim 170 \text{ MeV}$

$$\rho(m) \sim m^a e^{m/T_0}$$

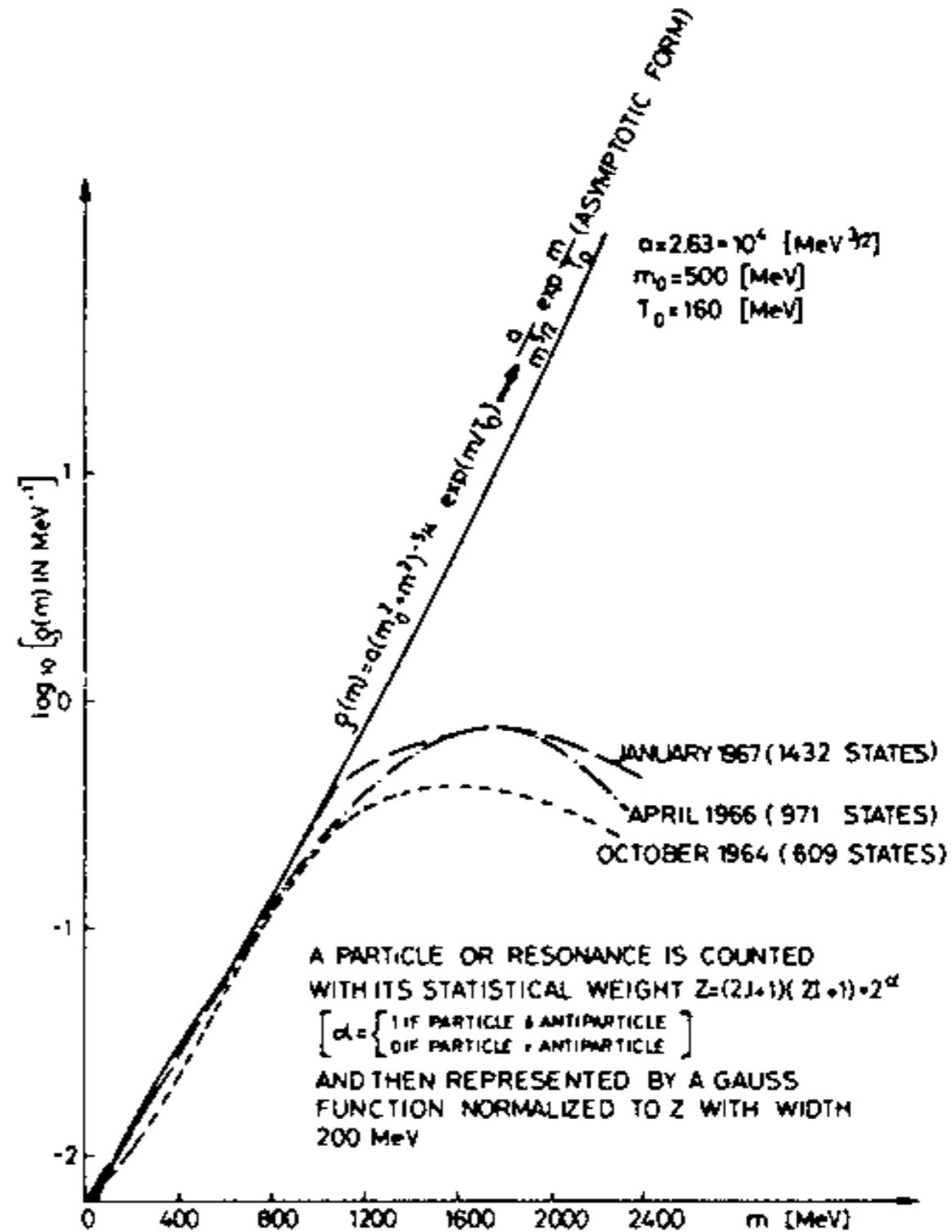
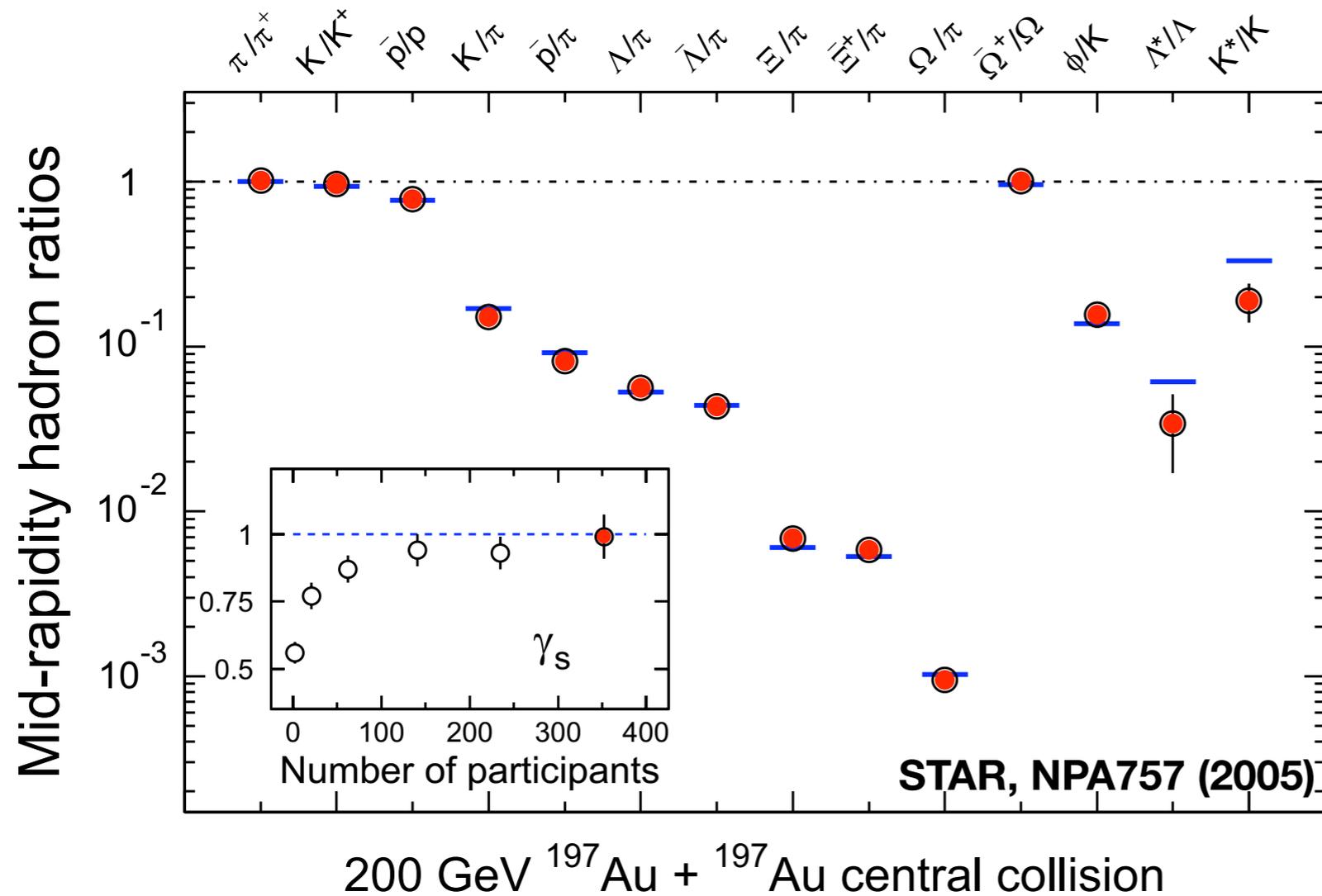


Fig. 3.1: The predicted and the experimental mass spectrum as it evolved from 1964 to 1967.



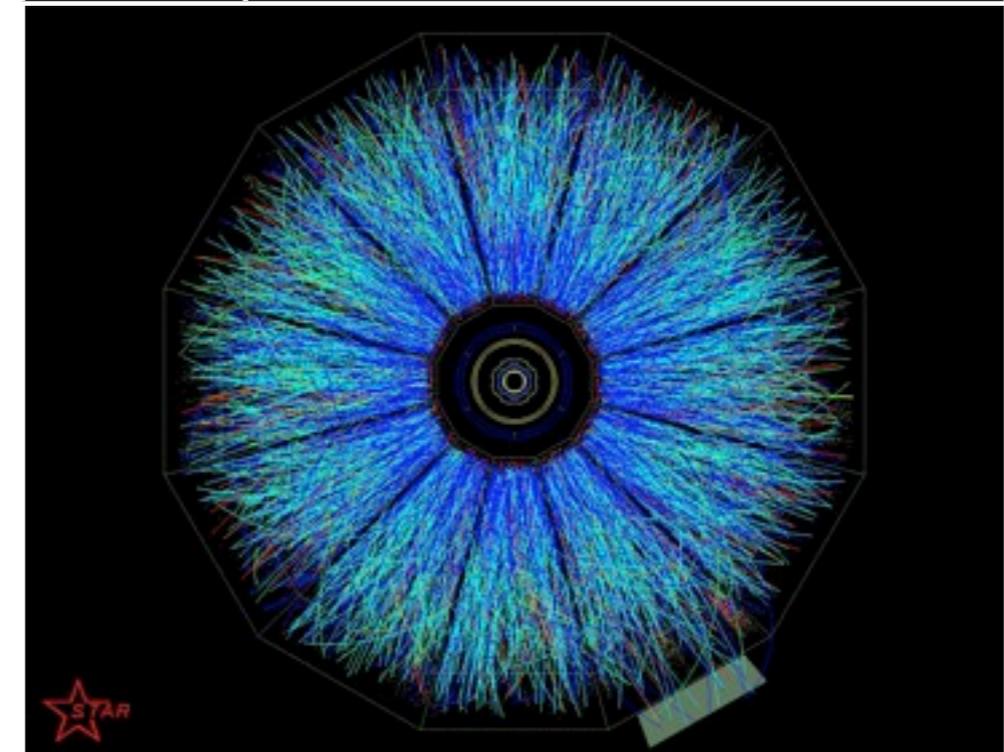
Strong Blackbody Radiation

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$$N_i \propto V \int \frac{d^3p}{(2\pi)^3} \frac{1}{e^{(\sqrt{p^2+m^2}-\mu_B)/T} \pm 1}$$

T	Chemical freezeout temperature
μ_B	Baryochemical potential (more matter than antimatter)



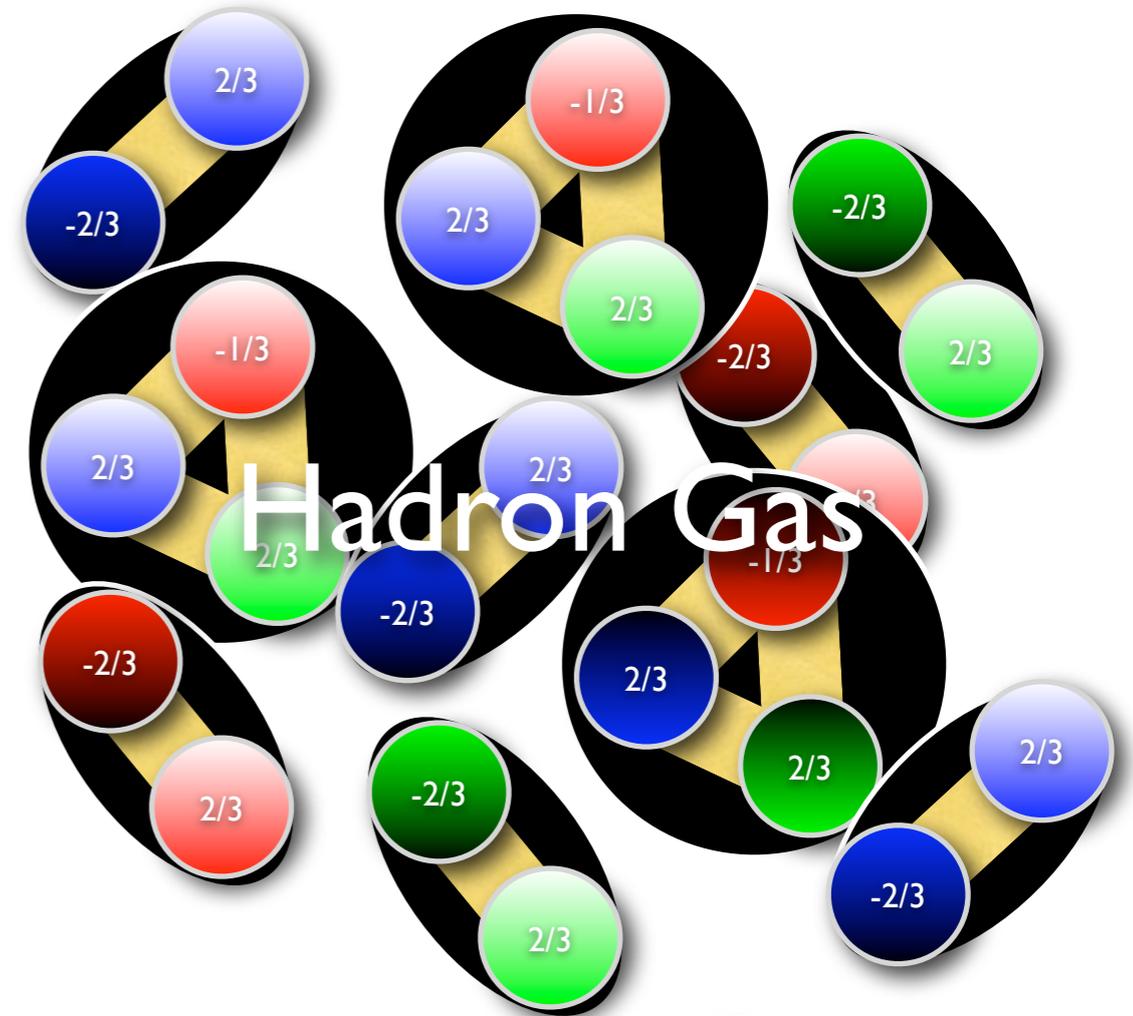
All hadron species apparently emitted from a thermal source $T_{ch} = 163 \pm 4$, $\mu_B = 24 \pm 4$

System decouples @ $T_{ch} \sim T_H$

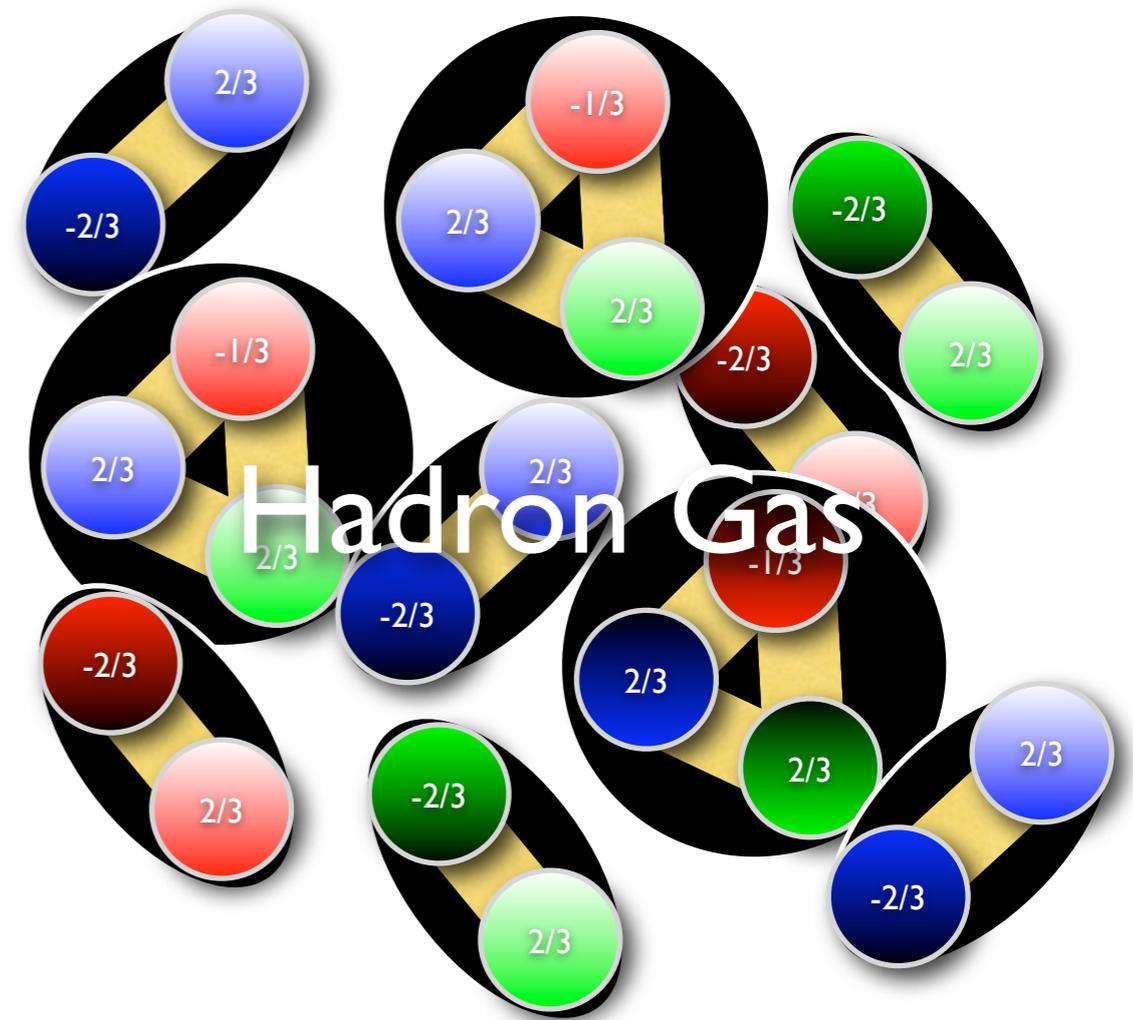
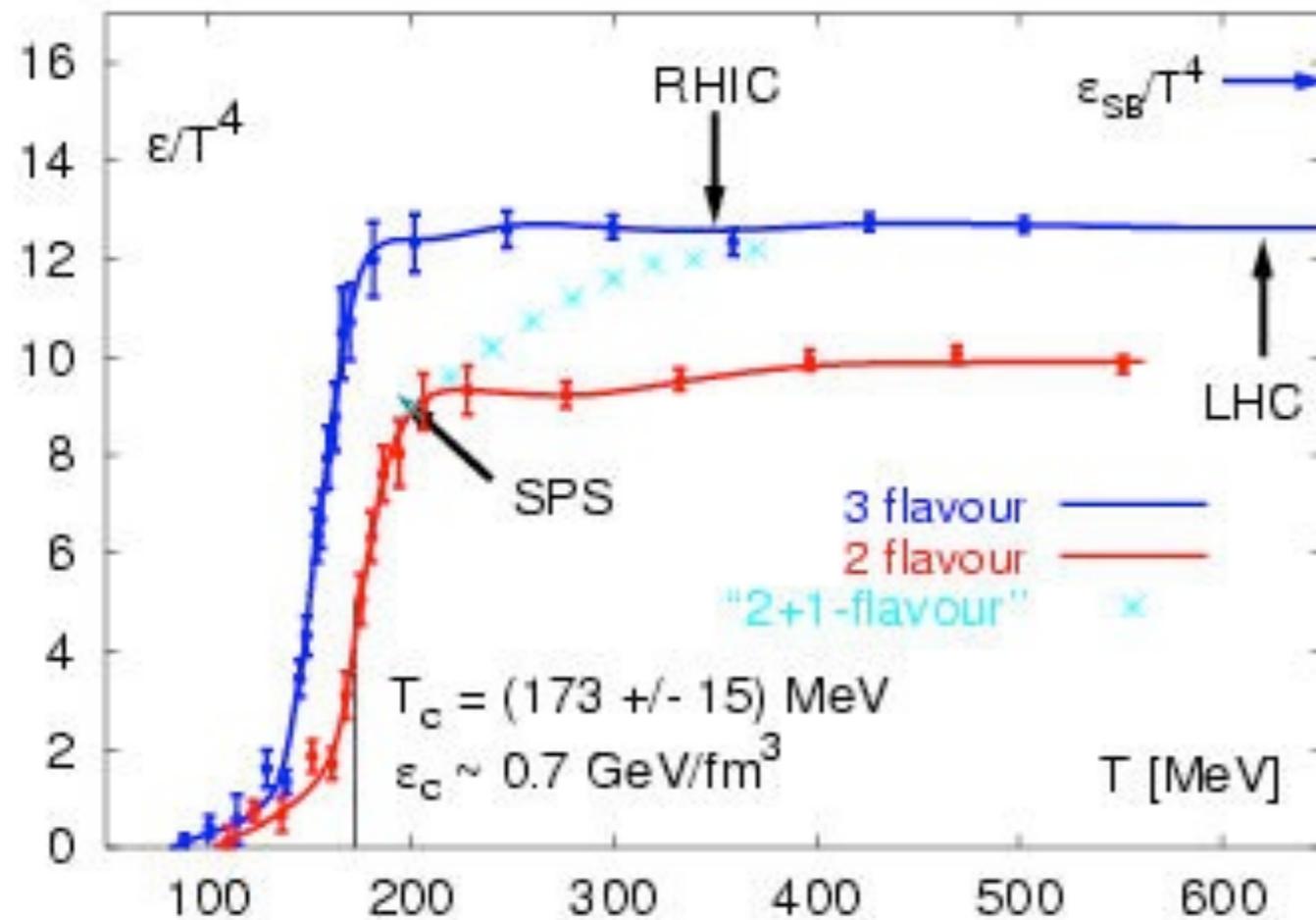
The Hagedorn Limit

Heating up a Hagedorn gas excites higher-mass resonances
→ T_H is the “limit” temperature

Nominally, no way to ever reach
higher temperatures in a
strongly-interacting system!



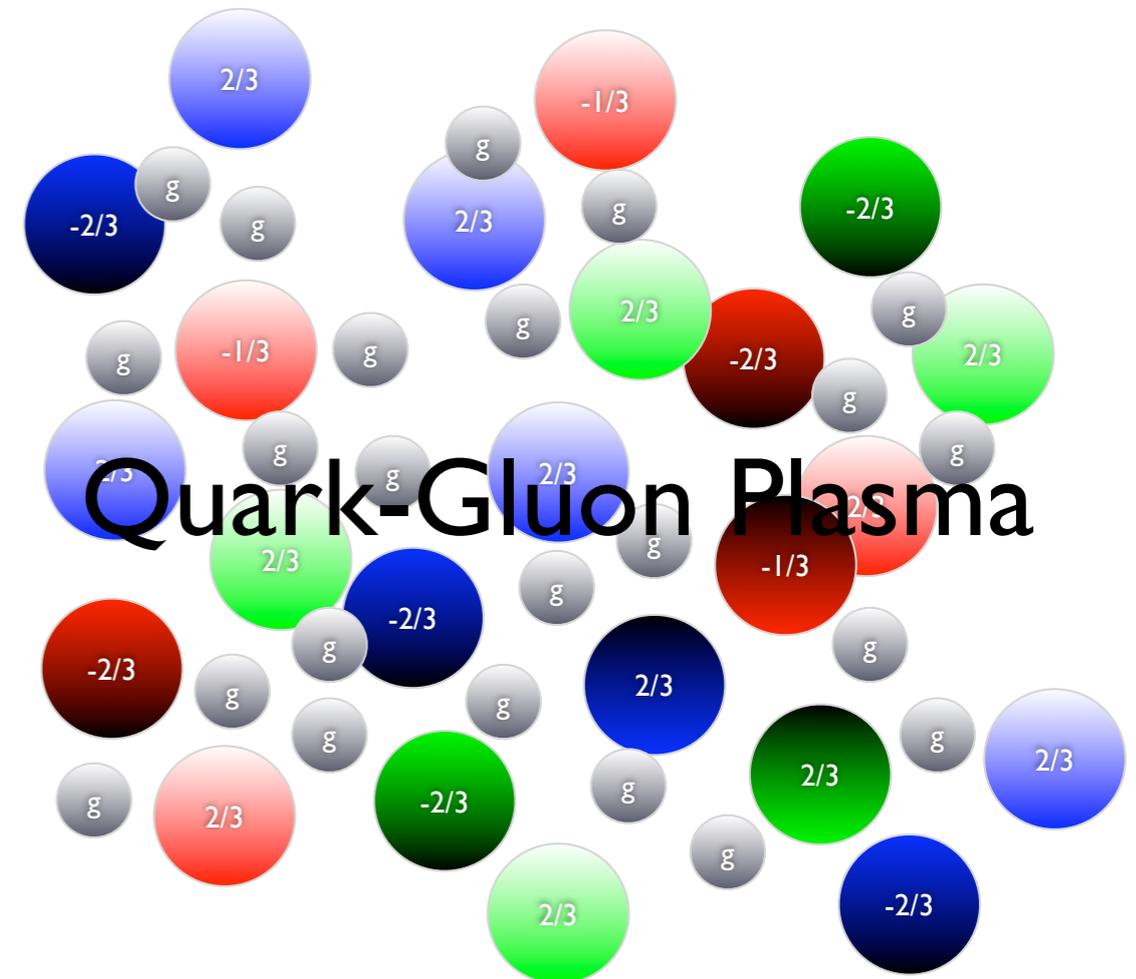
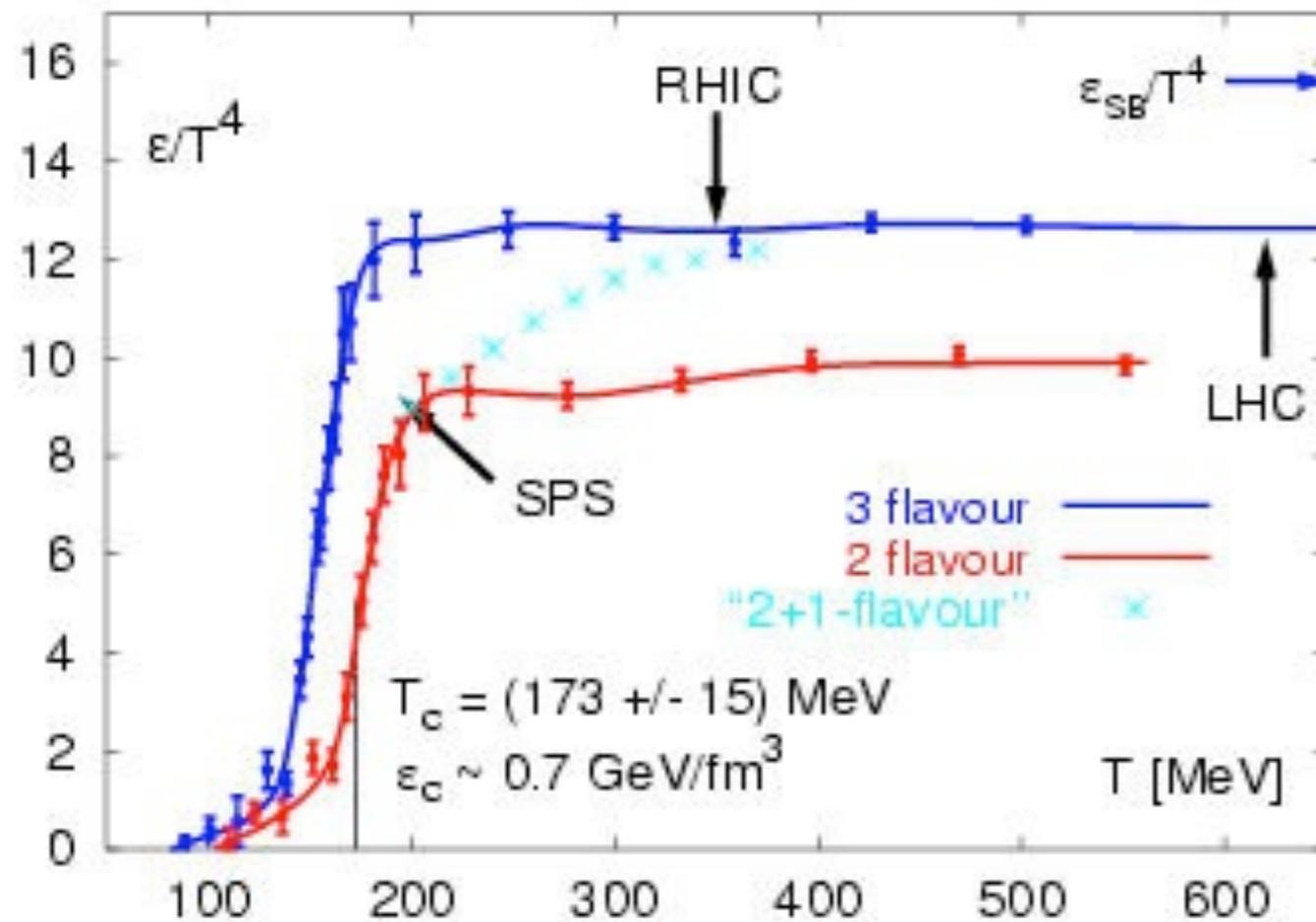
Quark Gluon Plasma



Lattice predicts a phase transition at $T_c \sim 170 \text{ MeV} \sim T_H$
($\epsilon_c \sim 700 \text{ MeV/fm}^3 > \epsilon_N \sim 500 \text{ MeV/fm}^3$)
from hadronic degrees of freedom to quark/gluon

Quark Gluon Plasma

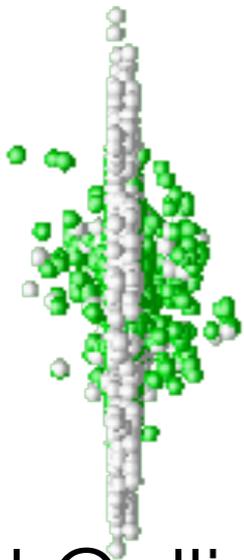
#DOF



Lattice predicts a phase transition at $T_c \sim 170 \text{ MeV} \sim T_H$
($\epsilon_c \sim 700 \text{ MeV}/\text{fm}^3 > \epsilon_N \sim 500 \text{ MeV}/\text{fm}^3$)
from hadronic degrees of freedom to quark/gluon

Probing Heavy Ion Dynamics with Soft Physics

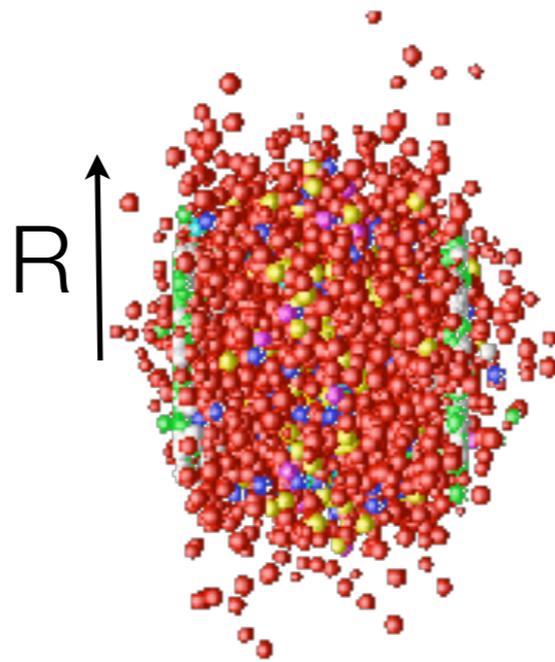
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Initial Collisions
(Nuclear geometry,
Baryon stopping)

**Initial Boundary
Conditions
(Space-time)**

Multiplicities

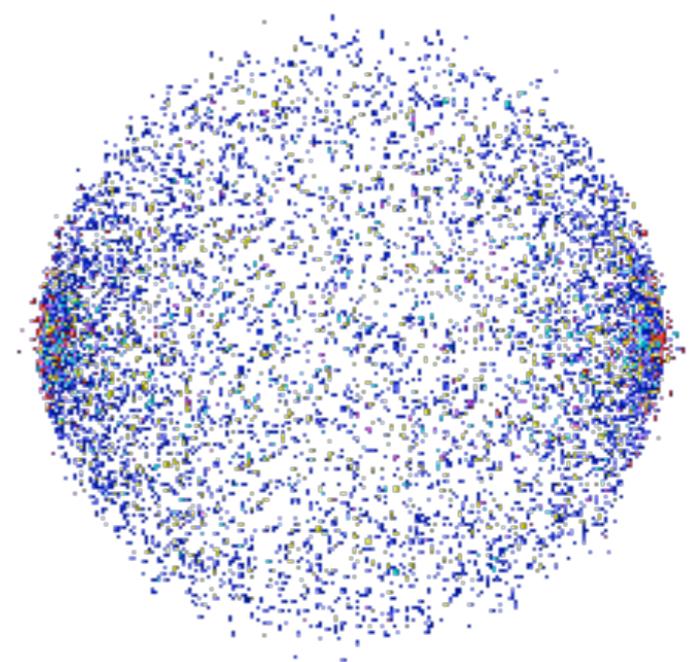


Dynamical evolution
 $\tau \sim O(R)$

$$\partial_{\mu} T^{\mu\nu} = 0$$

$p(\epsilon)$ (e.g. lattice)

Spectra, Flow



Freezeout
(Hadronization)

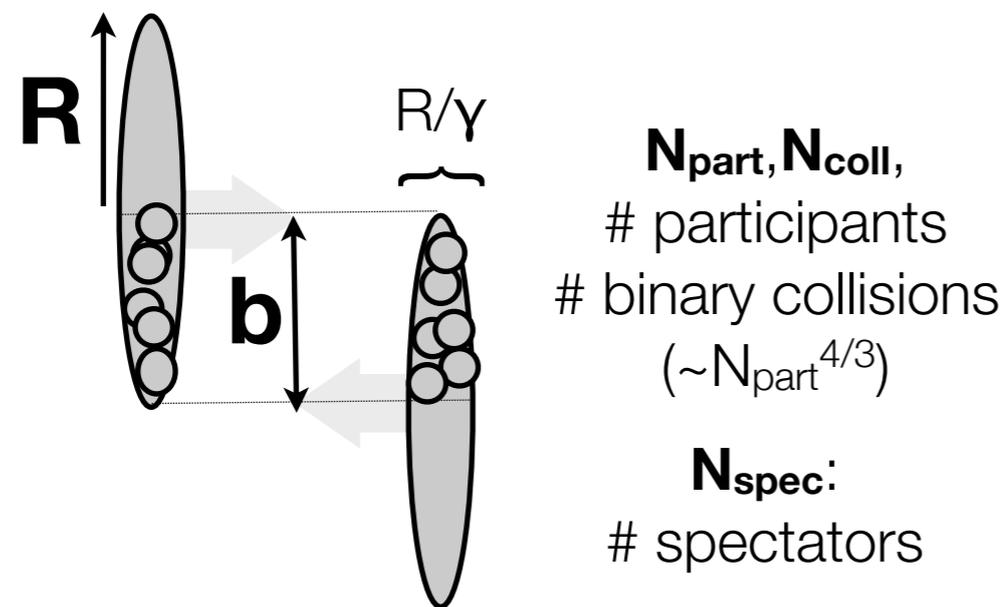
**Final Boundary
Conditions (Local)**

$$T = T_H \text{ or } T_{th}$$

Yields, Correlations

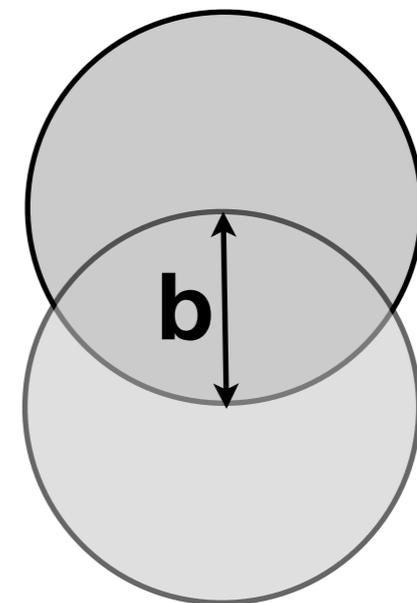
Nuclear Geometry

Knowing collision geometry is essential for all heavy ion results



$$\epsilon = \frac{\sigma_y^2 - \sigma_x^2}{\sigma_y^2 + \sigma_x^2}$$

“eccentricity”



Transverse and longitudinal scales are quite different:
spatial, temporal, momentum (via $\Delta p = h/\Delta R$)

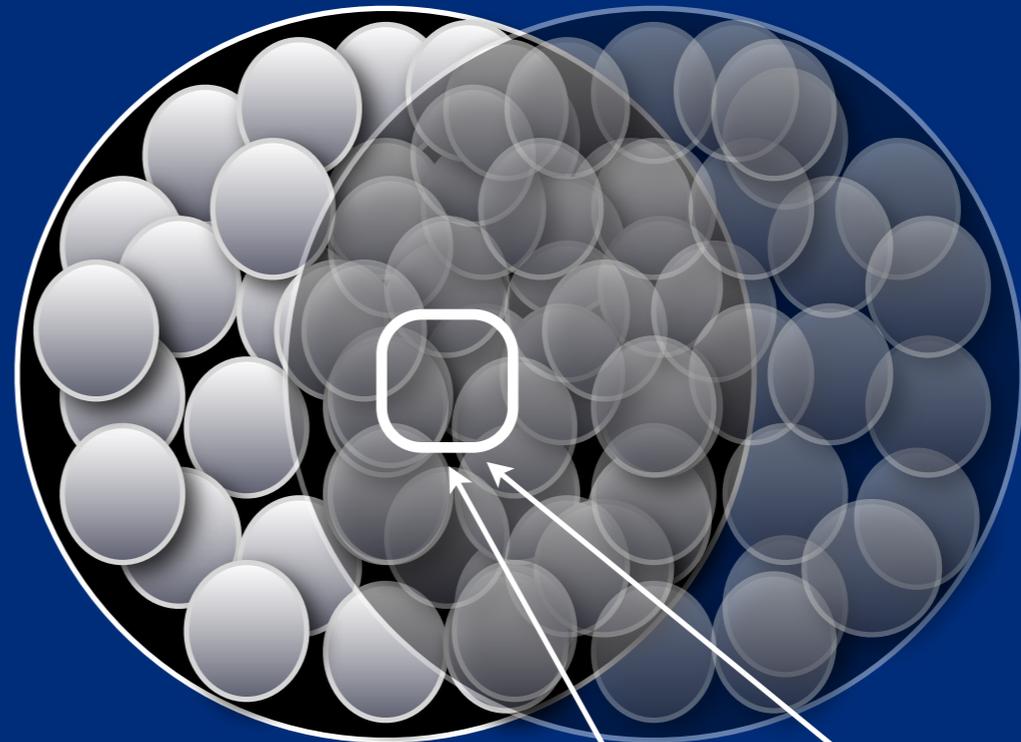
The Initial State

“Glauber model” is used to establish initial state (energy or entropy density) at one impact parameter, and evolve to other impact parameters

$$n_{part}(\vec{x}, b)$$

$$n_{coll}(\vec{x}, b)$$

more on this in the next lecture



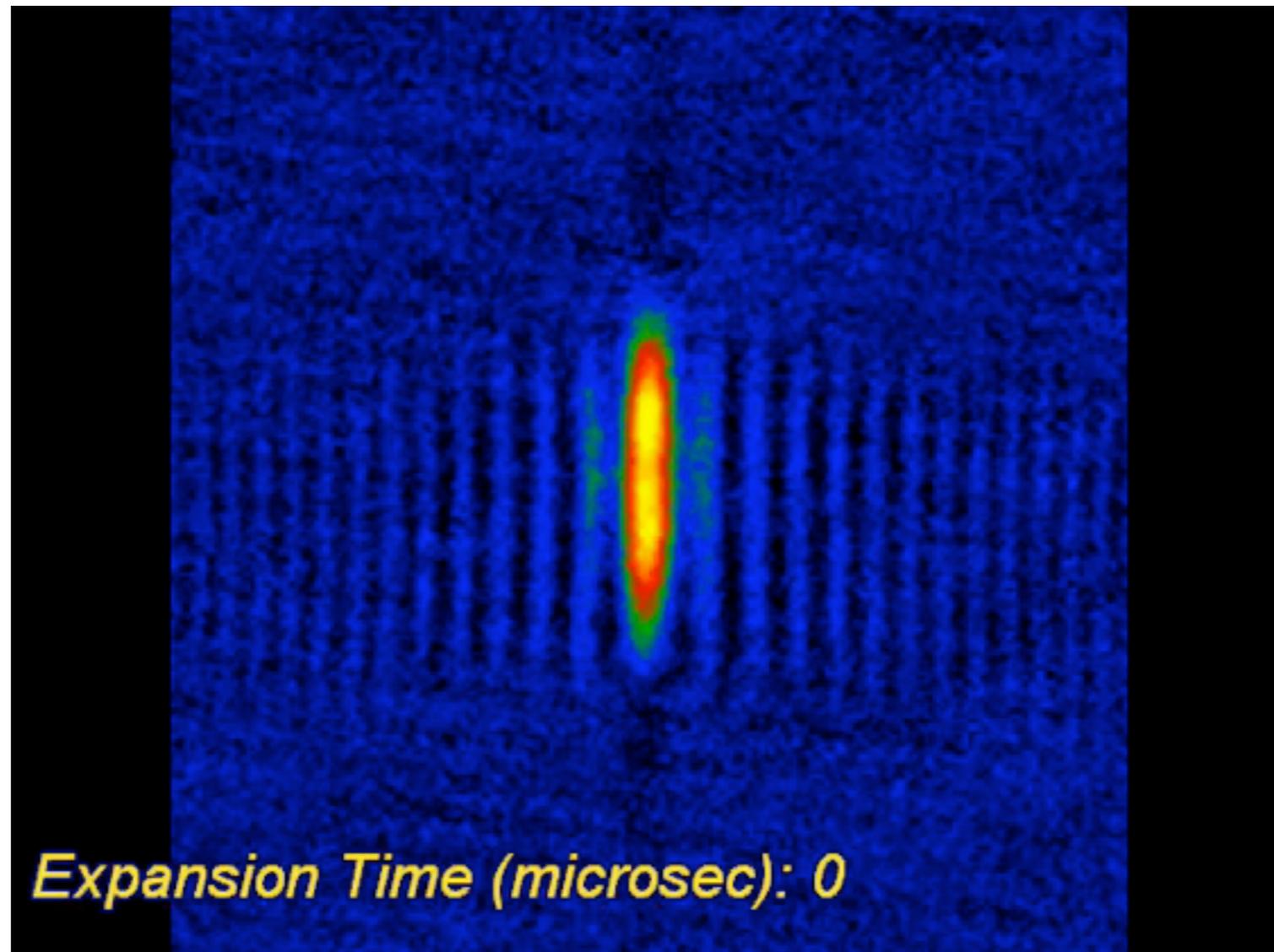
$$\frac{d^3 s}{d\eta d^2 \vec{x}} = n_{pp} \left\{ (1-x) \frac{n_{part}(\vec{x})}{2} + x n_{coll}(\vec{x}) \right\}$$

Participant density Collision density

Hydrodynamic Expansion

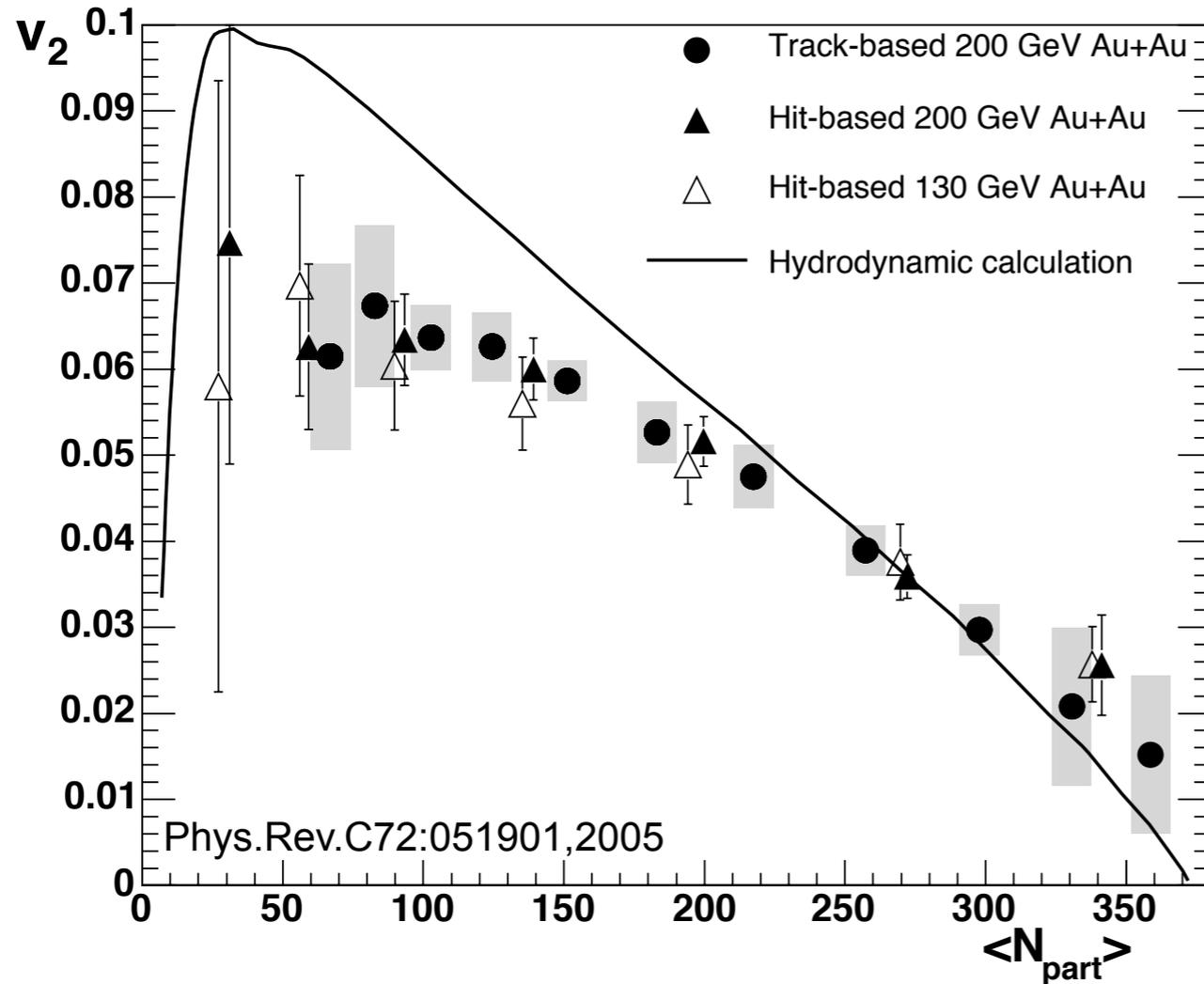
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laser flash photography of trapped Li ions (J. Thomas, Duke)

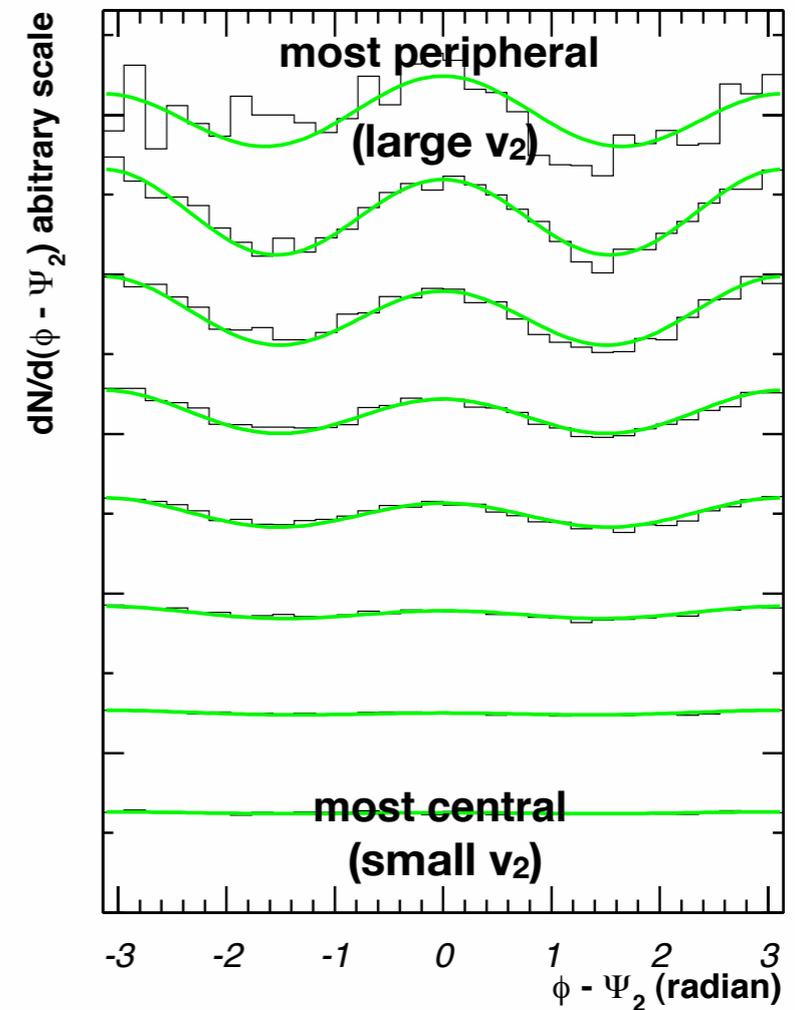


Hydrodynamic evolution is very general:
initial spatial asymmetries \rightarrow momentum asymmetries

Elliptic Flow (v_2)



PHOBOS Au+Au 130 GeV



$$\frac{1}{N} \frac{dN}{d\phi} = 1 + 2v_1 \cos(\phi - \Phi_R) + 2v_2 \cos(2[\phi - \Phi_R]) + \dots \quad \text{compared with Hydro (Huovinen)}$$

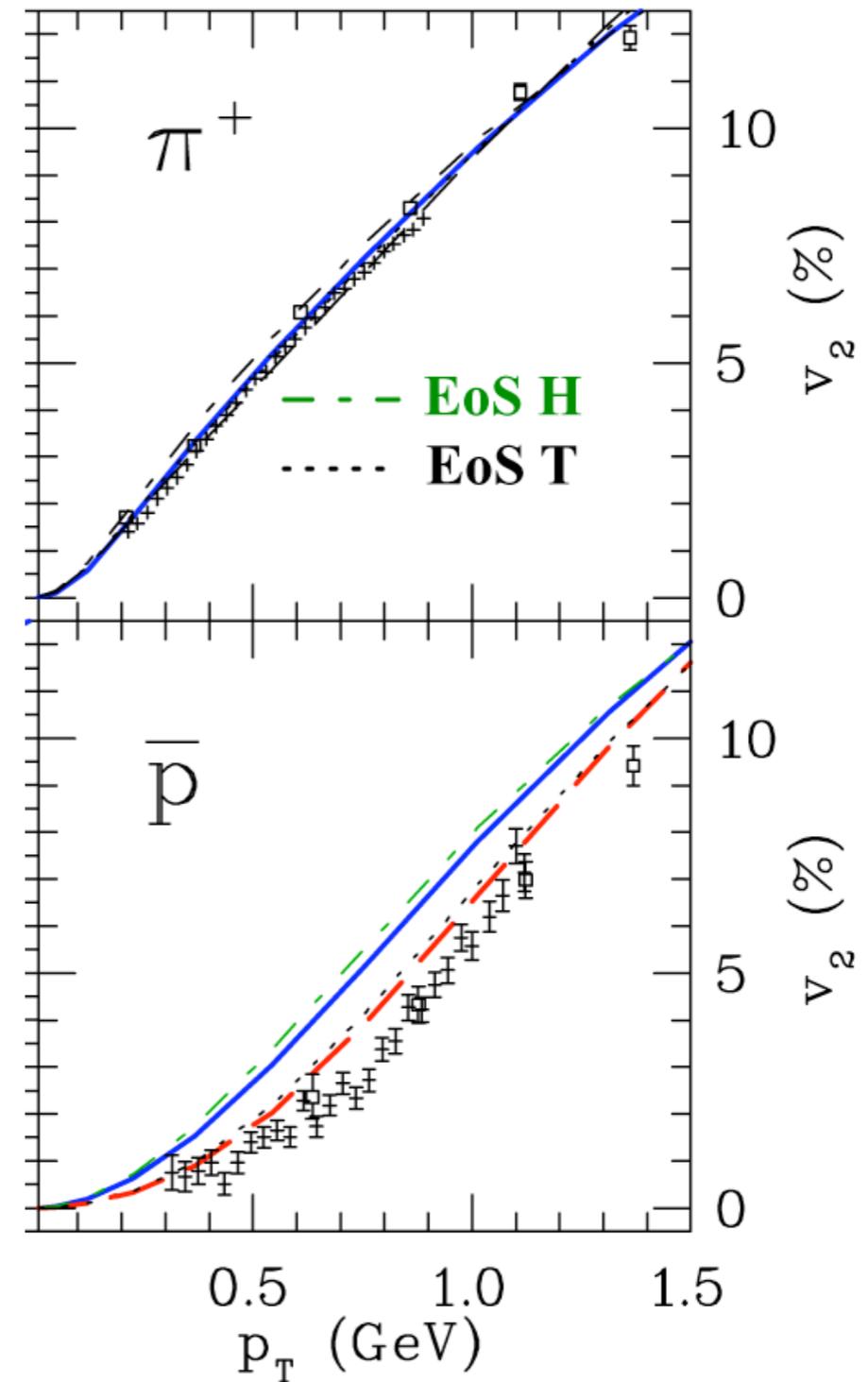
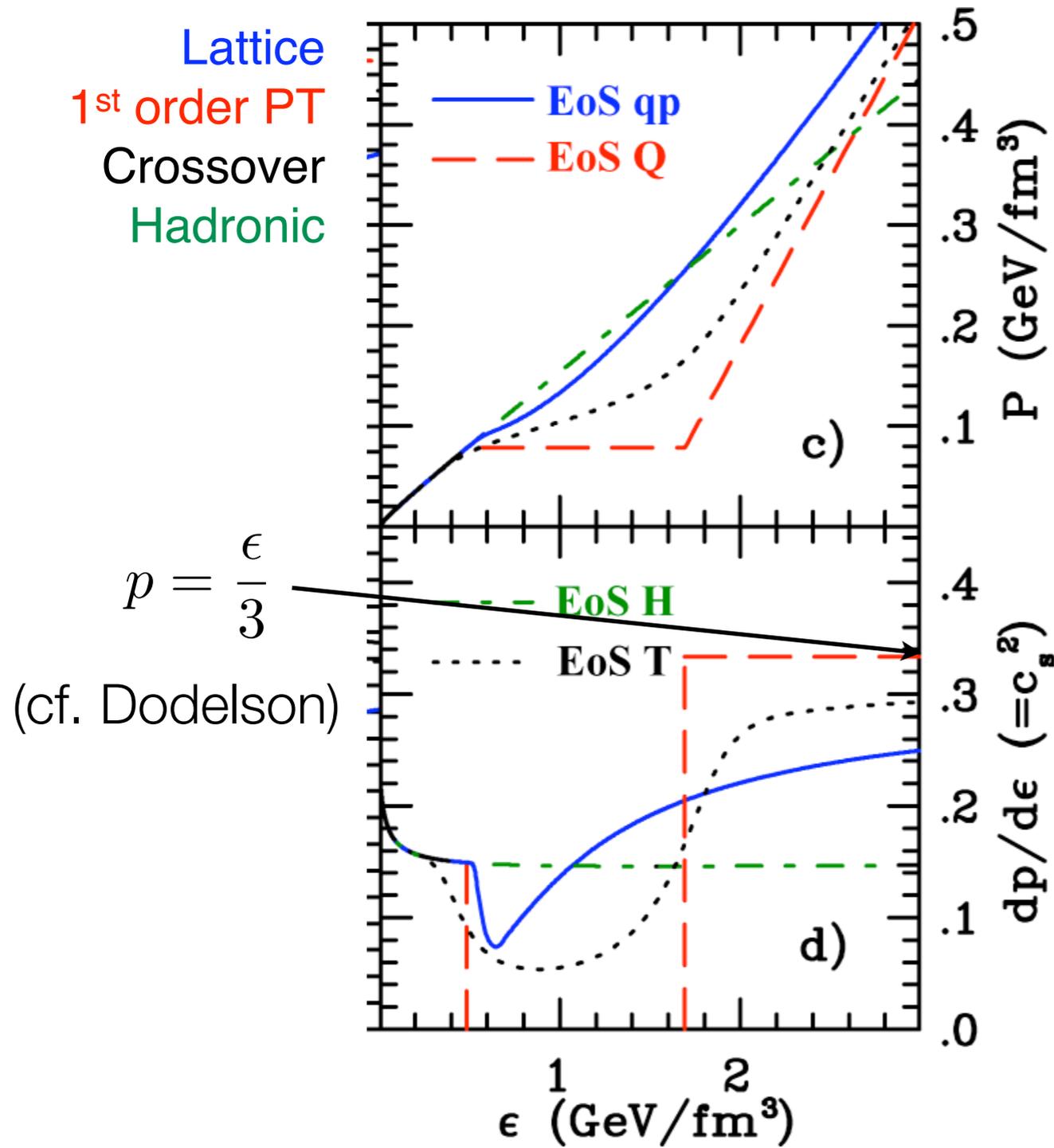
Hydro (boost invariant)

$\epsilon_0 \sim 30 \text{ GeV}/\text{fm}^3$
 $\tau_0 < 0.6 \text{ fm}^3$

$\epsilon \sim 0.5 \text{ GeV}/\text{fm}^3$
 $\tau \sim 1 \text{ fm}^3$

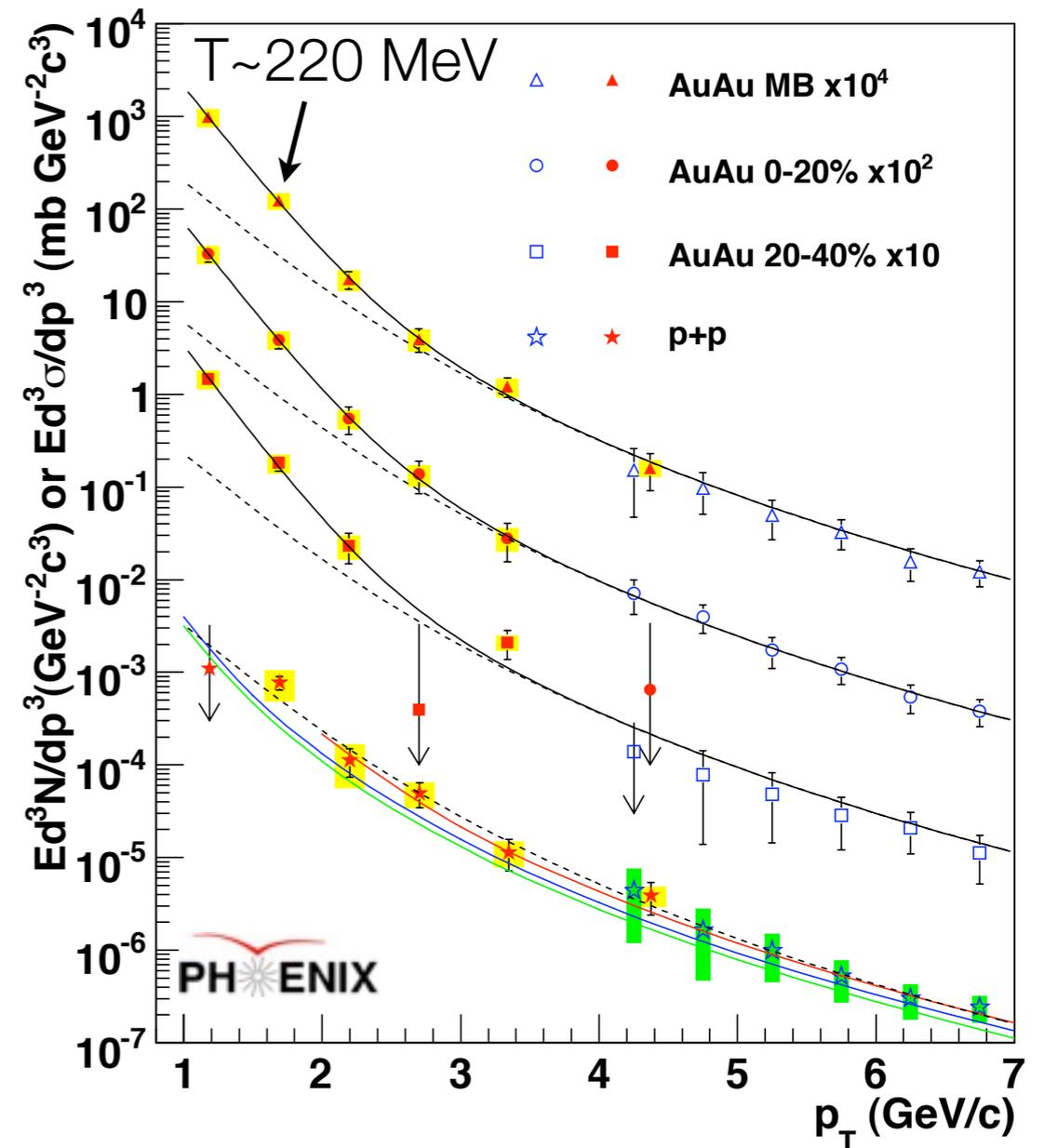
Hadronic scales

Equation of State ($v_2(p_T)$)



Implications of Hydro

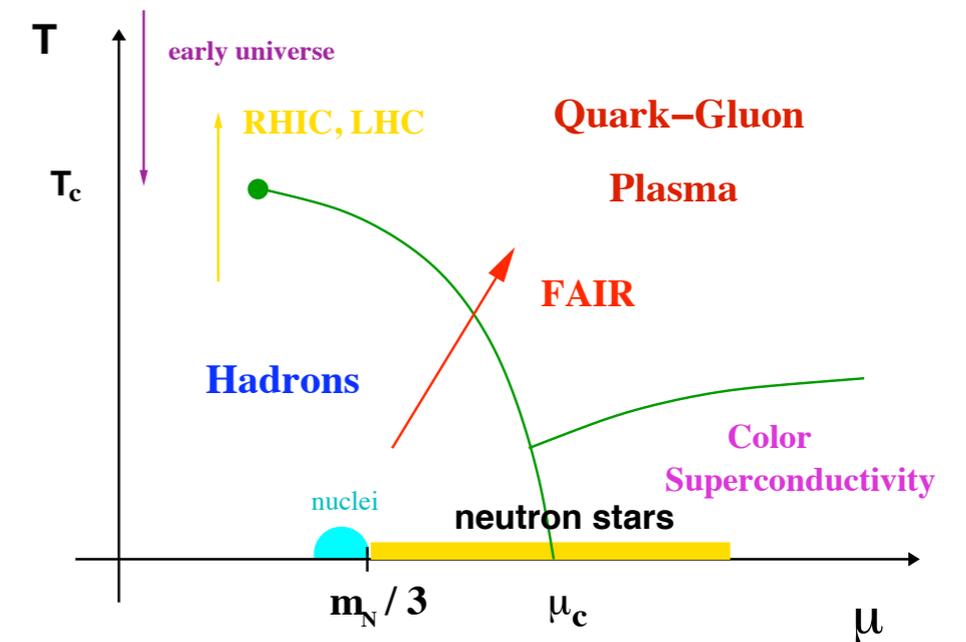
- **Hydro is not “just another” model**
- **The flowing medium is quite special**
 - Local thermalization (strong coupling)
 - “Lattice”-like equation of state (1st order PT)
- **Initial density so large ($30 \times \epsilon_p$) that can't justify a hadronic cascade**
- **Temperature exceeds T_H**
 - High temperature seen directly by direct thermal photons from PHENIX
- **This suggests that the medium may well be the “quark gluon plasma”**
 - What points to quark and gluon DOFs?



Studying Matter in Laboratory

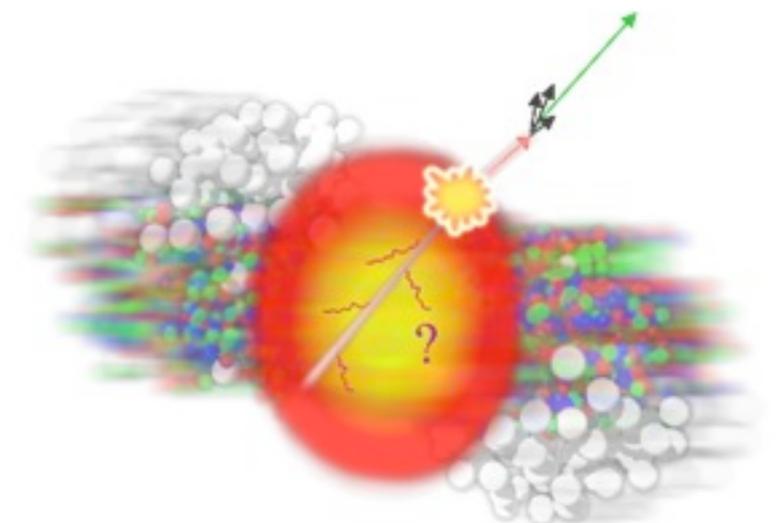
- **Changing initial conditions -- “soft” physics**

Ideal	Temperature	Matter	Density
Practical	Energy	Size	Shape



- **Probing it microscopically -- “hard” probes (next lecture)**

Ideal	Microscopy	Tomography	
Practical	Jets	Photons	Quarkonia



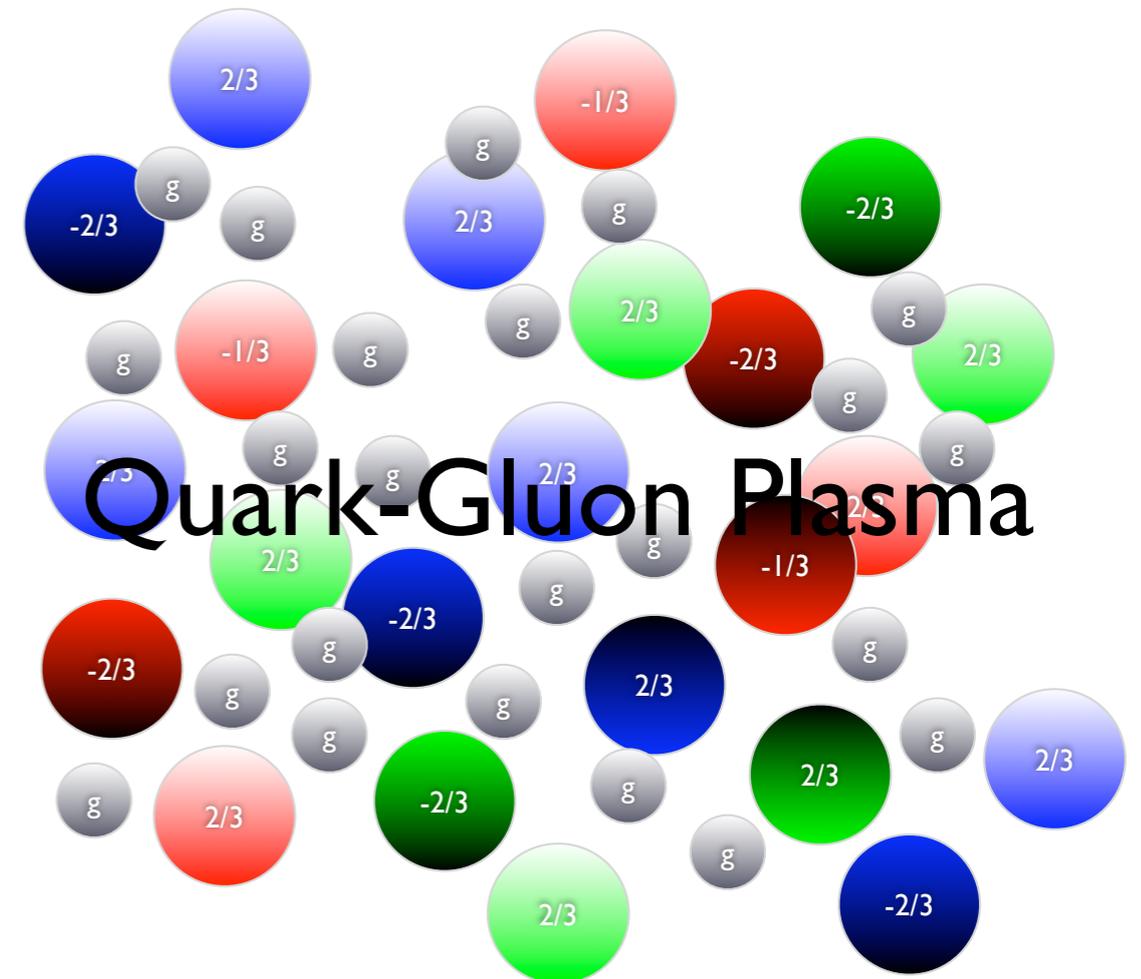
The Big Question: Thermalization at RHIC

Does the system truly thermalize? Everywhere?

What are the conditions (energy, density, size) for thermalization?

How fast does the system thermalize?

What thermalizes?



Degrees of Freedom: What Thermalizes?

Parton distributions
Nuclear Geometry
Nuclear shadowing

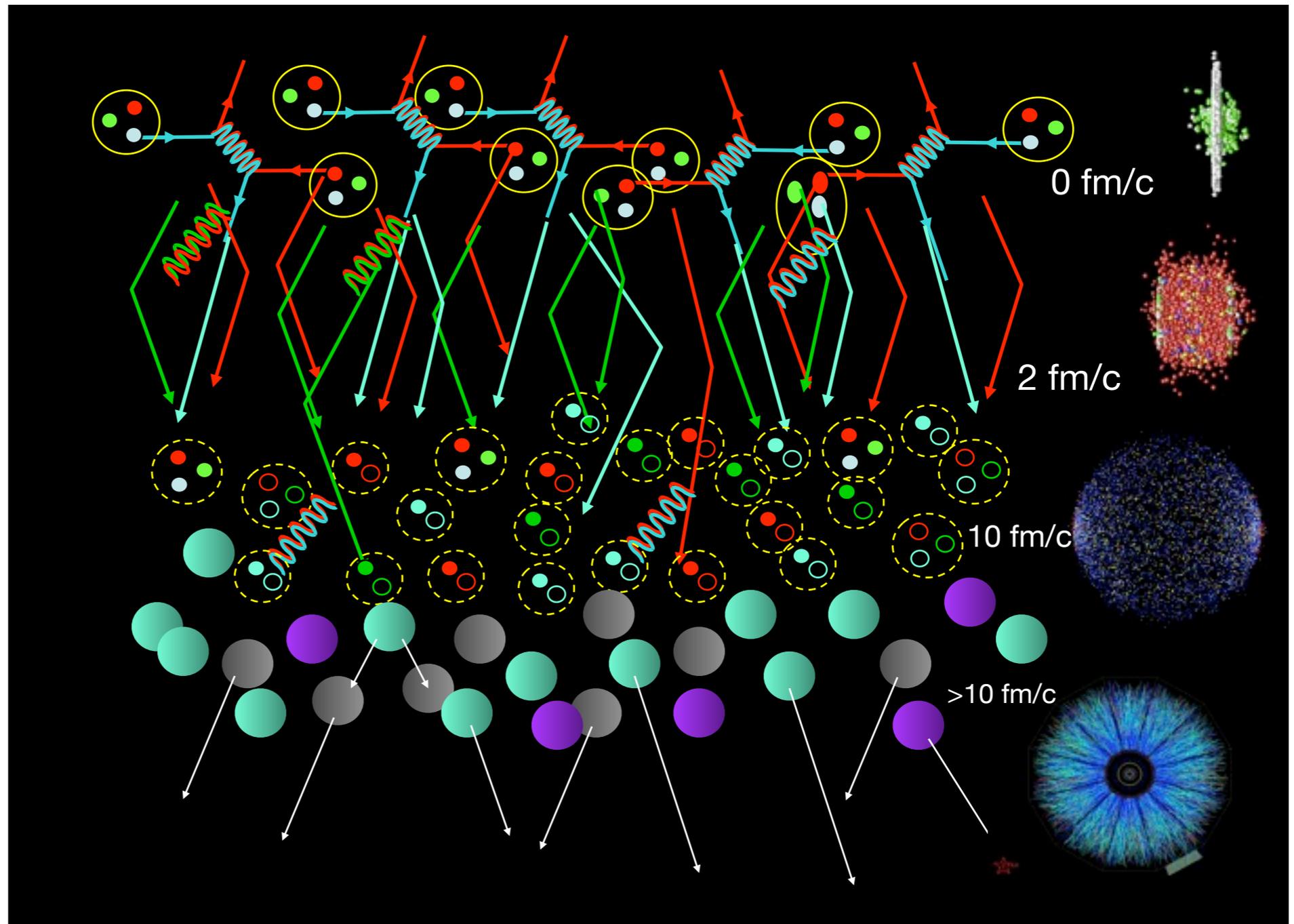
Parton production
& reinteraction

Chemical Freezeout &
Quark Recombination

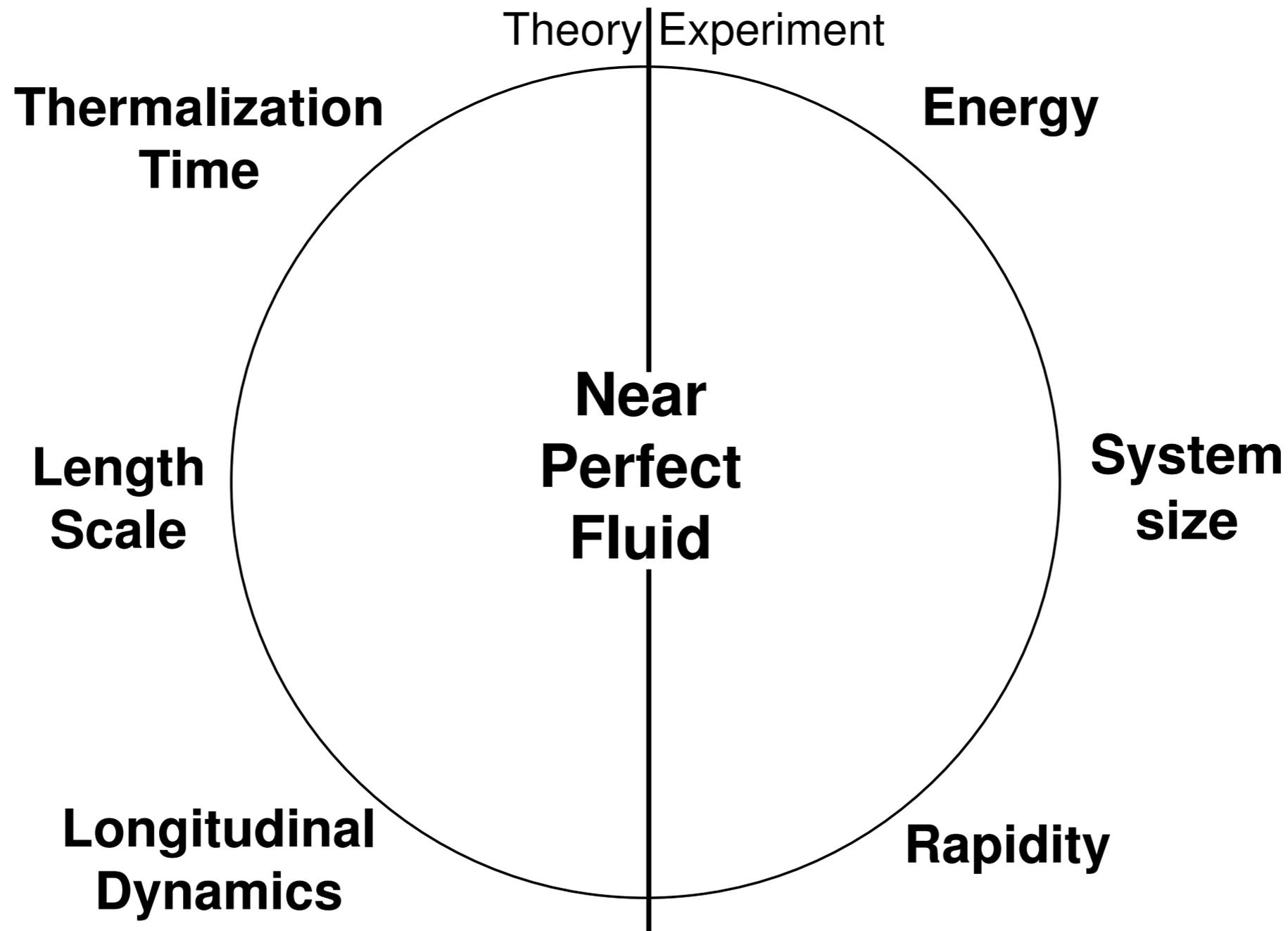
Jet Fragmentation
Functions

Hadron Rescattering

Thermal Freezeout &
Hadron decays



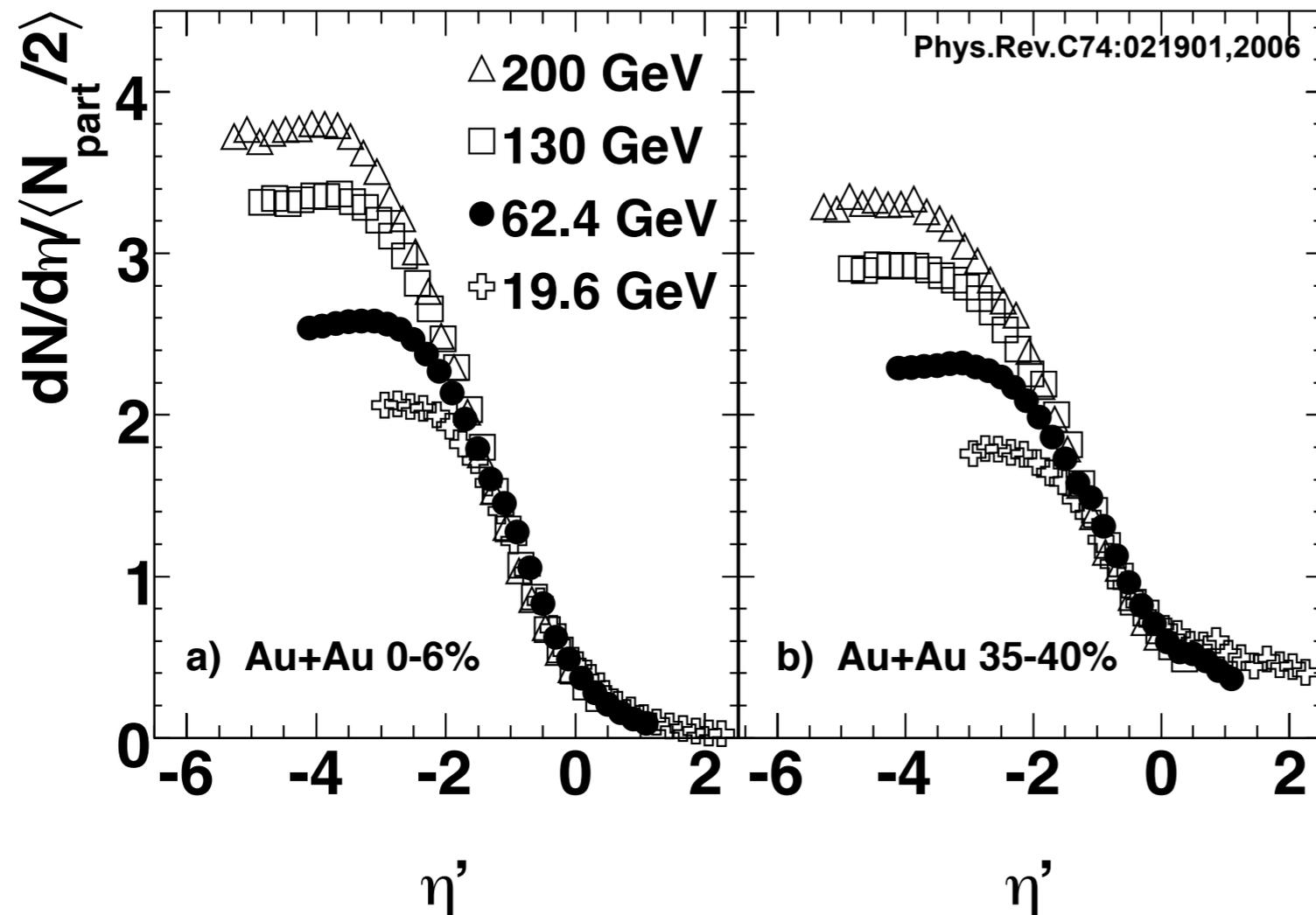
The Edge of Liquidity



Can we make what we see at RHIC “turn off”?

Extended Longitudinal Scaling (Limiting Fragmentation)

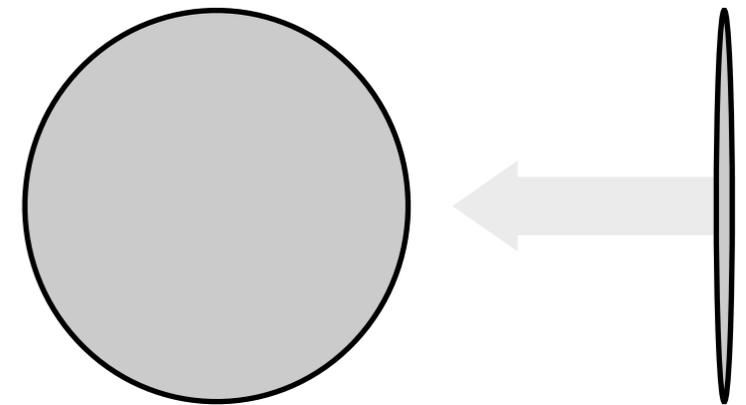
- From rest frame of one projectile: yields invariant at fixed geometry (i.e. same $b/2R$ or $N_{\text{part}}/2A$)



$$y = \tanh^{-1}(\beta)$$

$$\sim -\log(\tan(\theta/2)) \equiv \eta$$

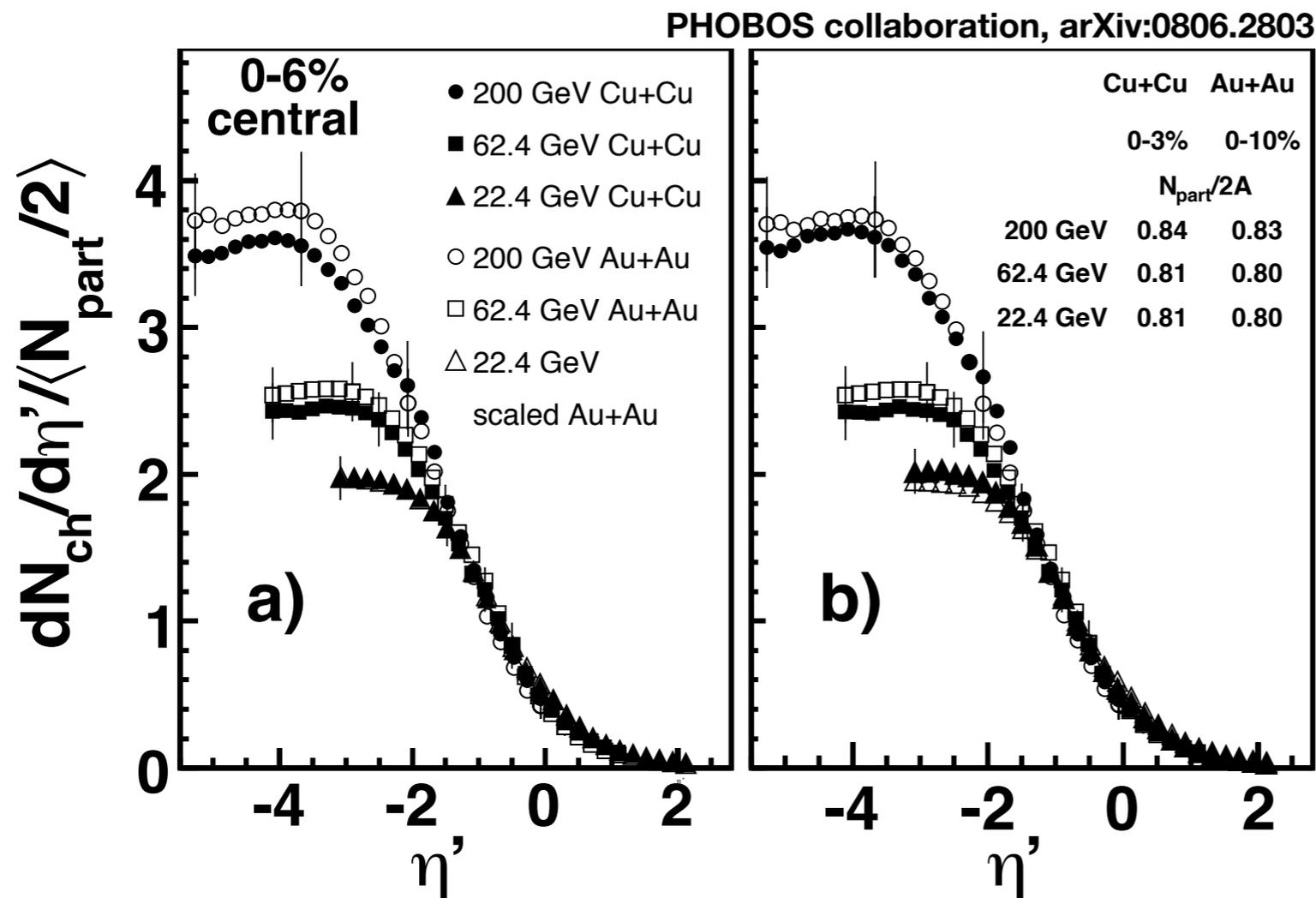
$$\eta' = \eta - y_b \sim \log(x_F)$$



Midrapidity physics at lower energy = forward physics at higher energy
Energy and rapidity dependence is smooth

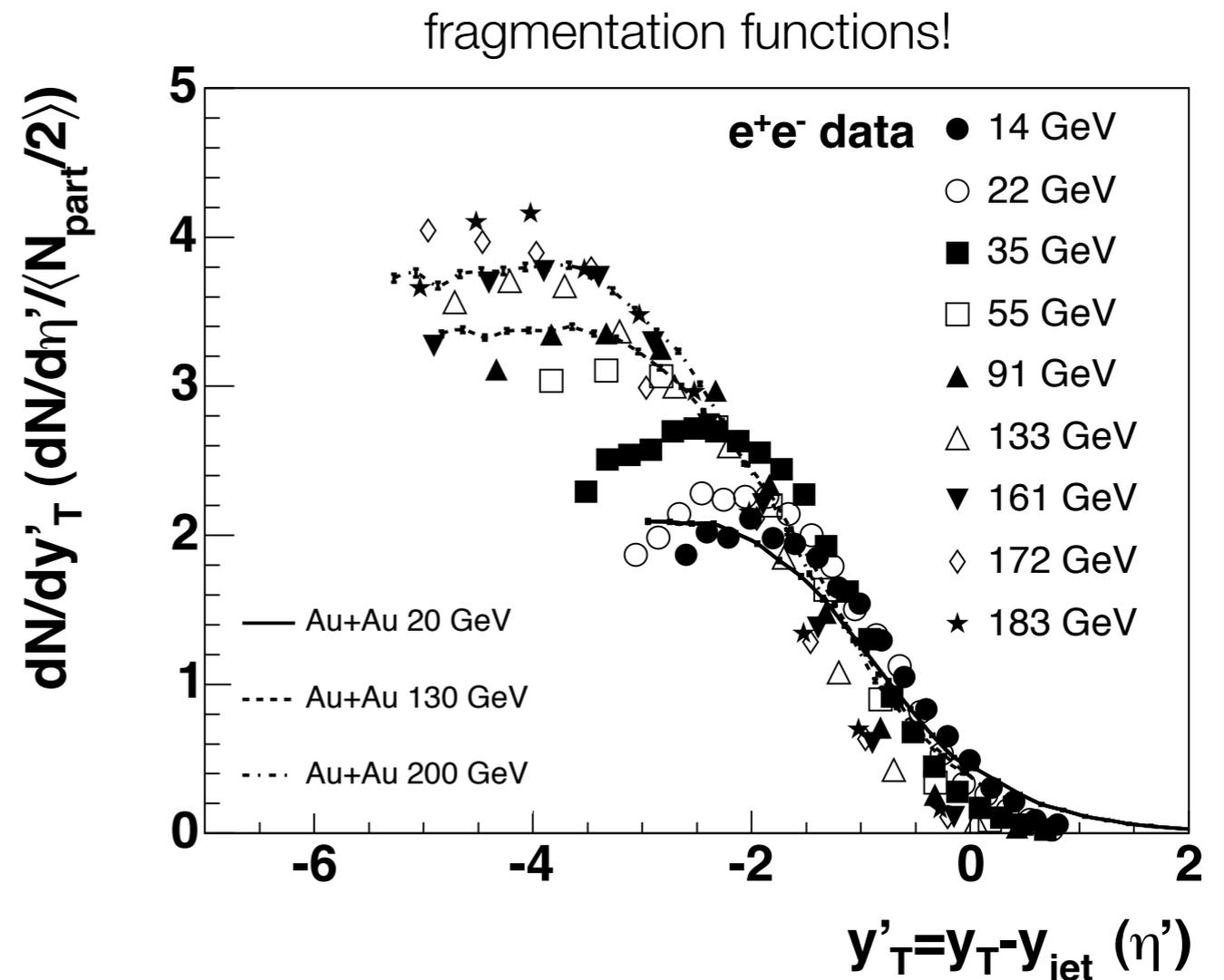
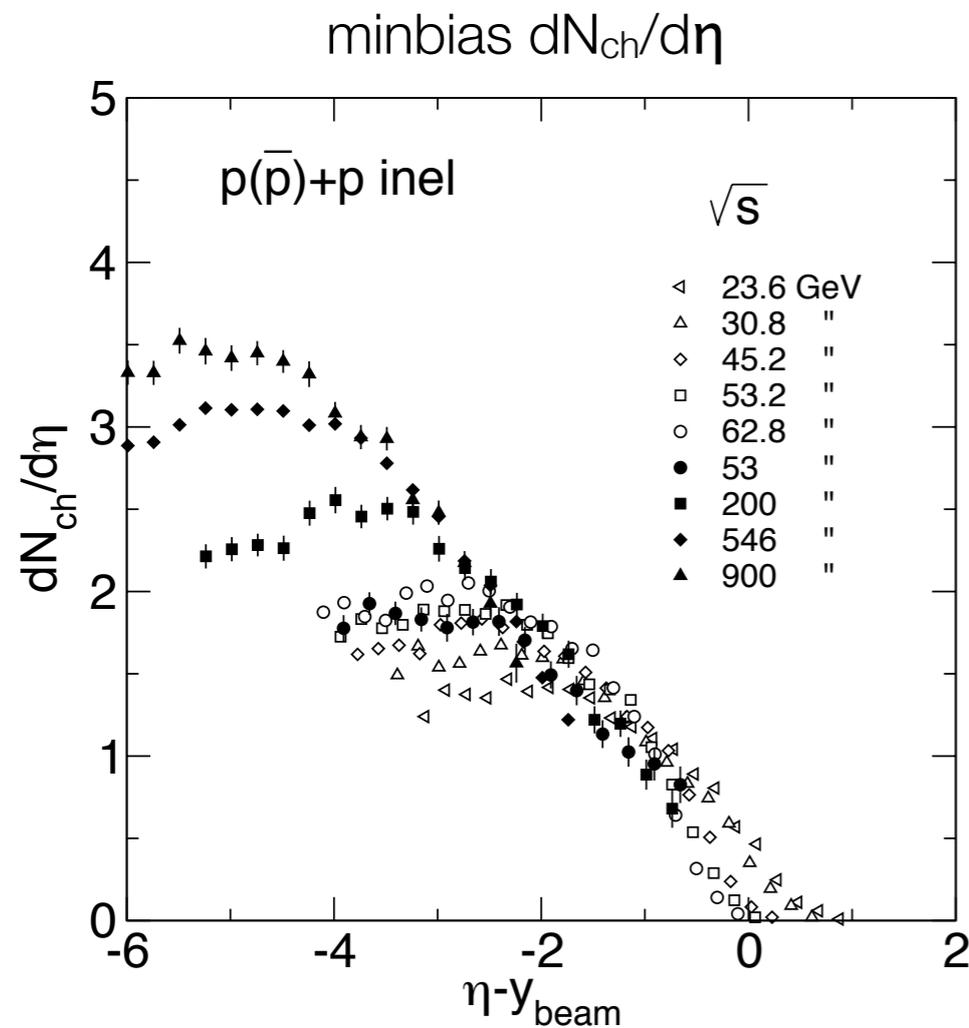
Extended Longitudinal Scaling (Limiting Fragmentation) PHOBOS Q

- No change in shape moving from Au+Au → Cu+Cu



Change the nuclear size by x3: Au+Au → Cu+Cu
 No change in shape for same centrality ($b/2R$)

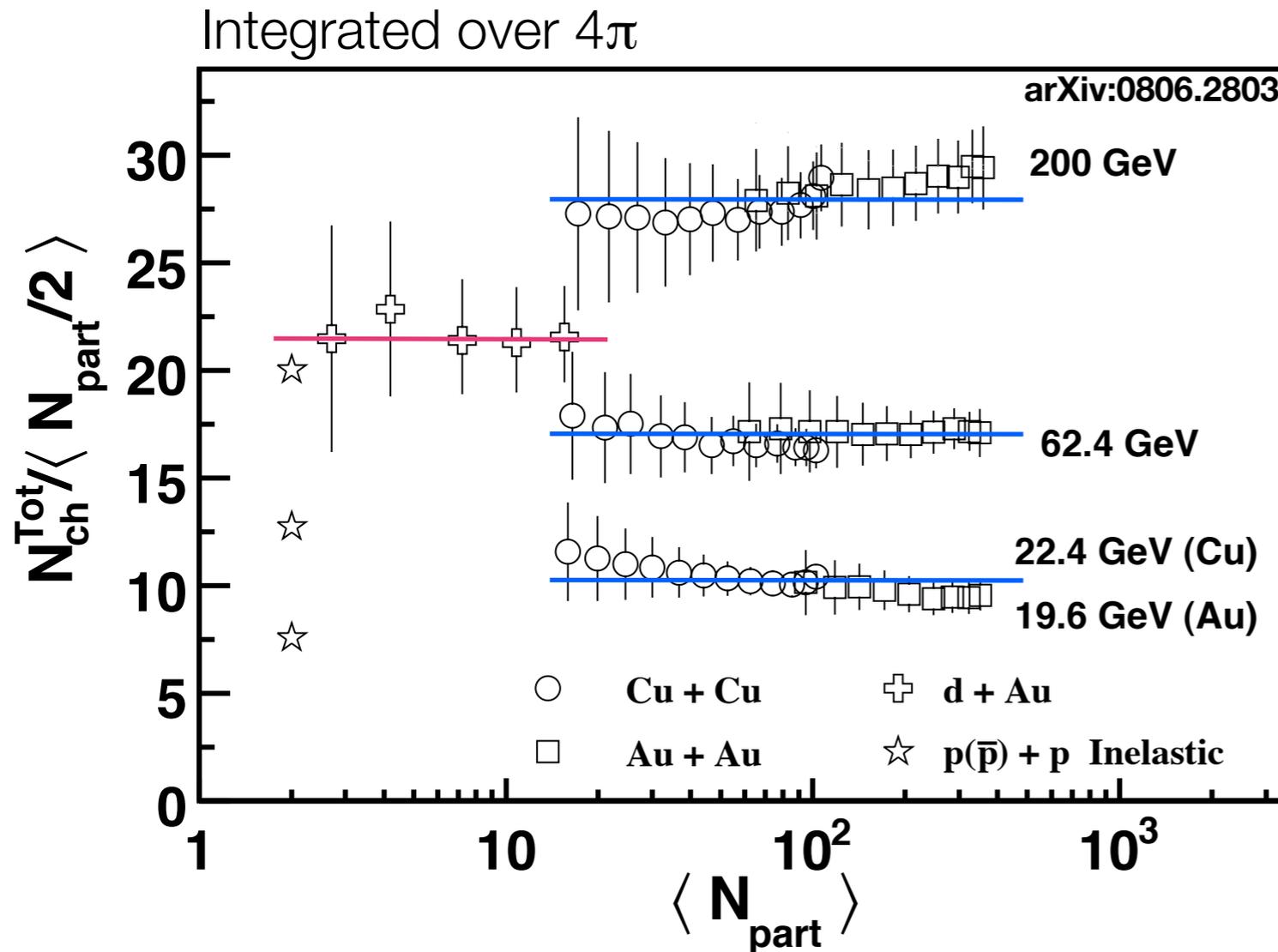
Longitudinal Scaling in Elementary Systems ($p+p$, e^+e^-)



Even “small” systems ($p+p$ and e^+e^-) show same feature,
and magnitude for e^+e^-

Ubiquity \rightarrow 1) **trivial** or 2) **deep**

When does the system get too small?



Cu+Cu & low energy data
fit into systematics
established by RHIC Au+Au:
when is system “too small”?

Total multiplicity (4π)
shows “wounded nucleon”
scaling to $N_{part}=20$

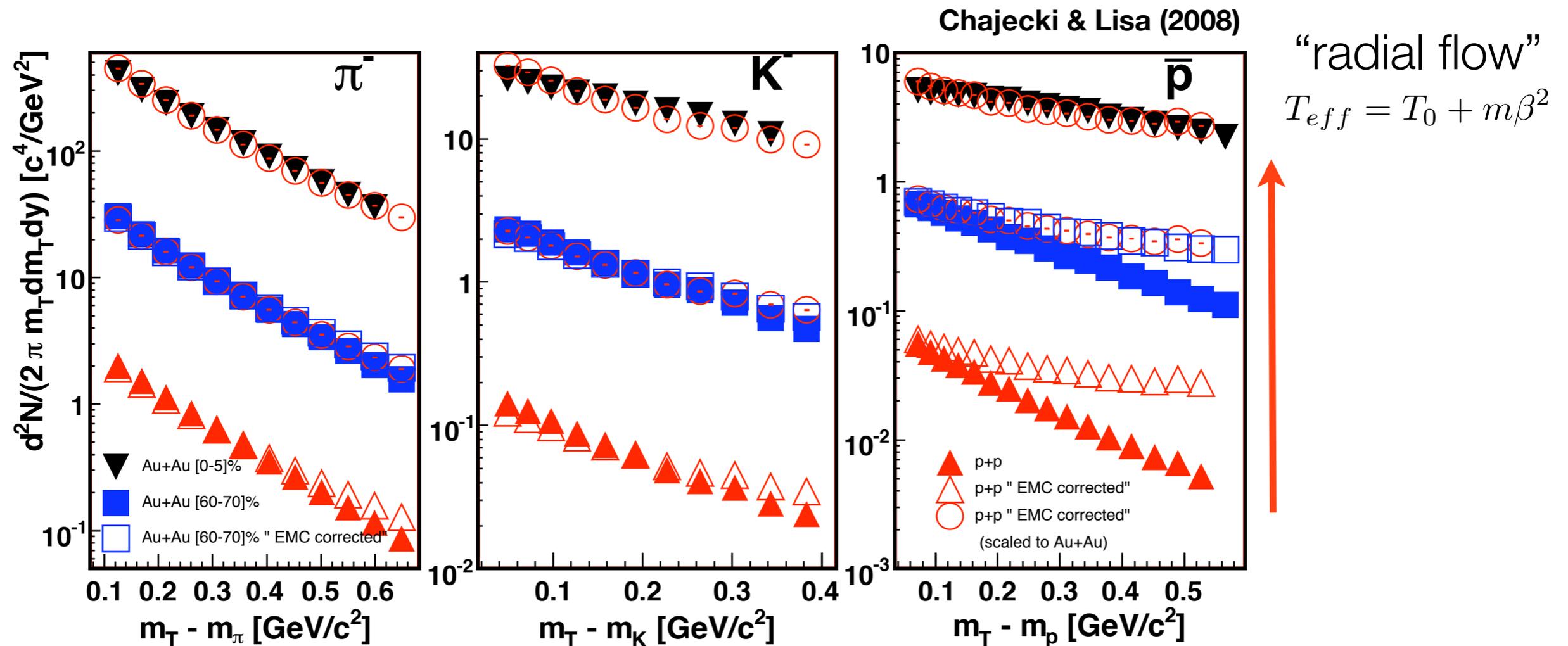
Suggests no change
in overall degrees of freedom
with **system size**

p+p/d+Au \rightarrow A+A explained by
“leading-particle” effects

$$S \propto sV \propto N_{tot} \propto N_{part}$$

(Fermi-Landau model)

Small ~ Large?

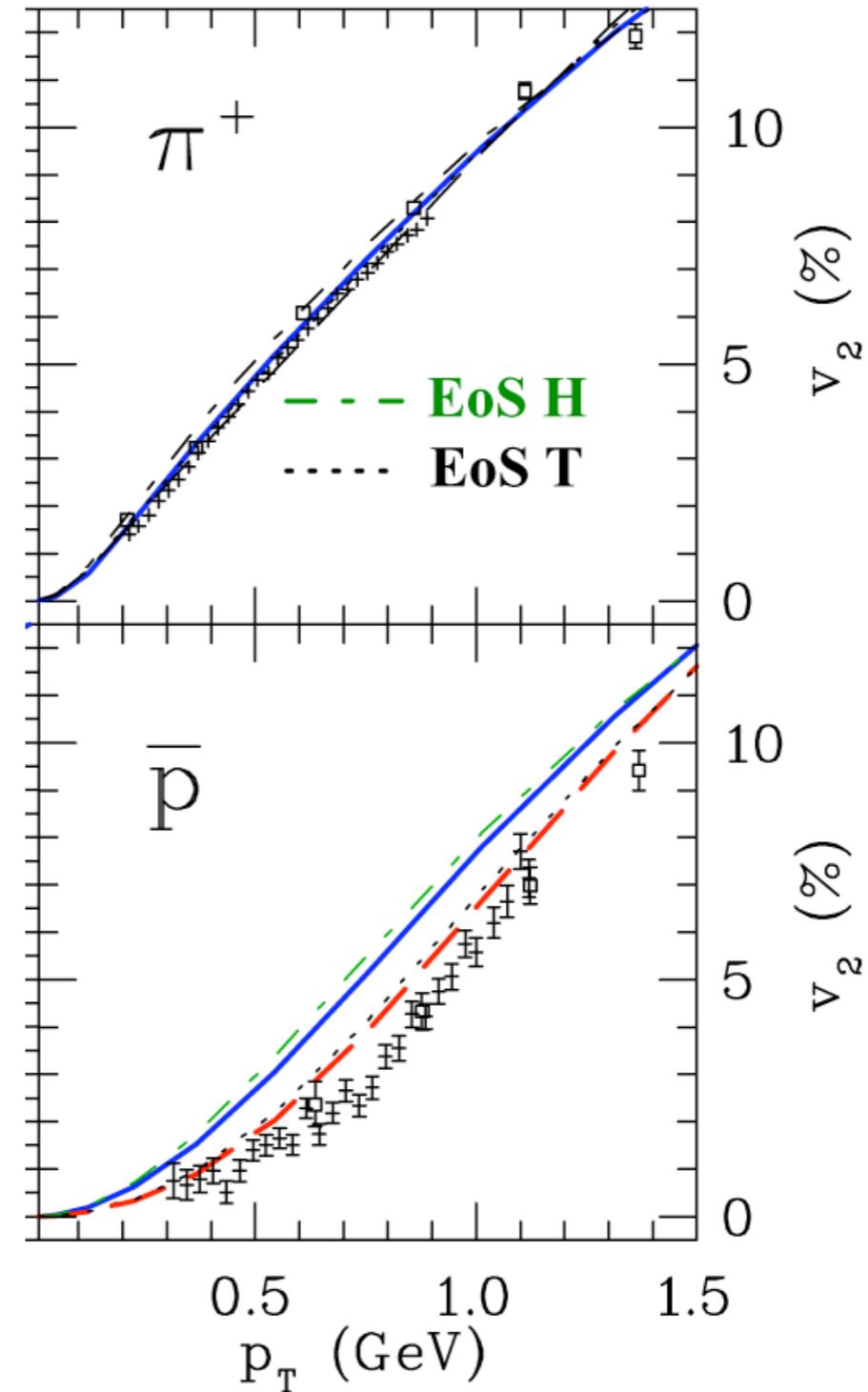
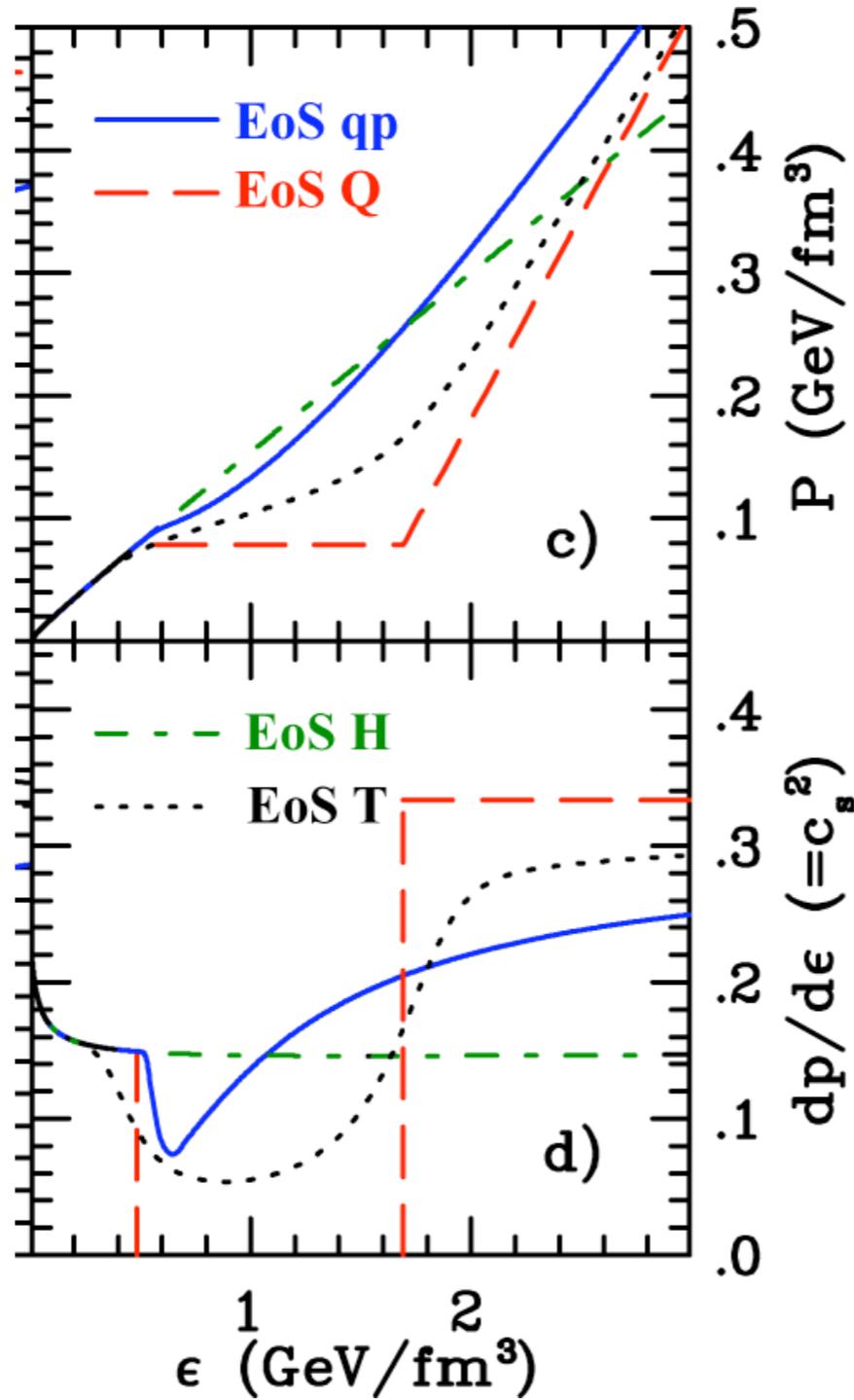


p+p spectrum, “undistorted” by energy-momentum conservation, compares well to A+A: radial expansion in p+p?

Does p+p behave like a thermalized, flowing system, like A+A? (if so, what does this suggest about “underlying event”?)

Recap: EOS Degrees of Freedom

Lattice
 1st order PT
 Crossover
 Hadronic



Degrees of Freedom: What Flows?

Parton distributions
Nuclear Geometry
Nuclear shadowing

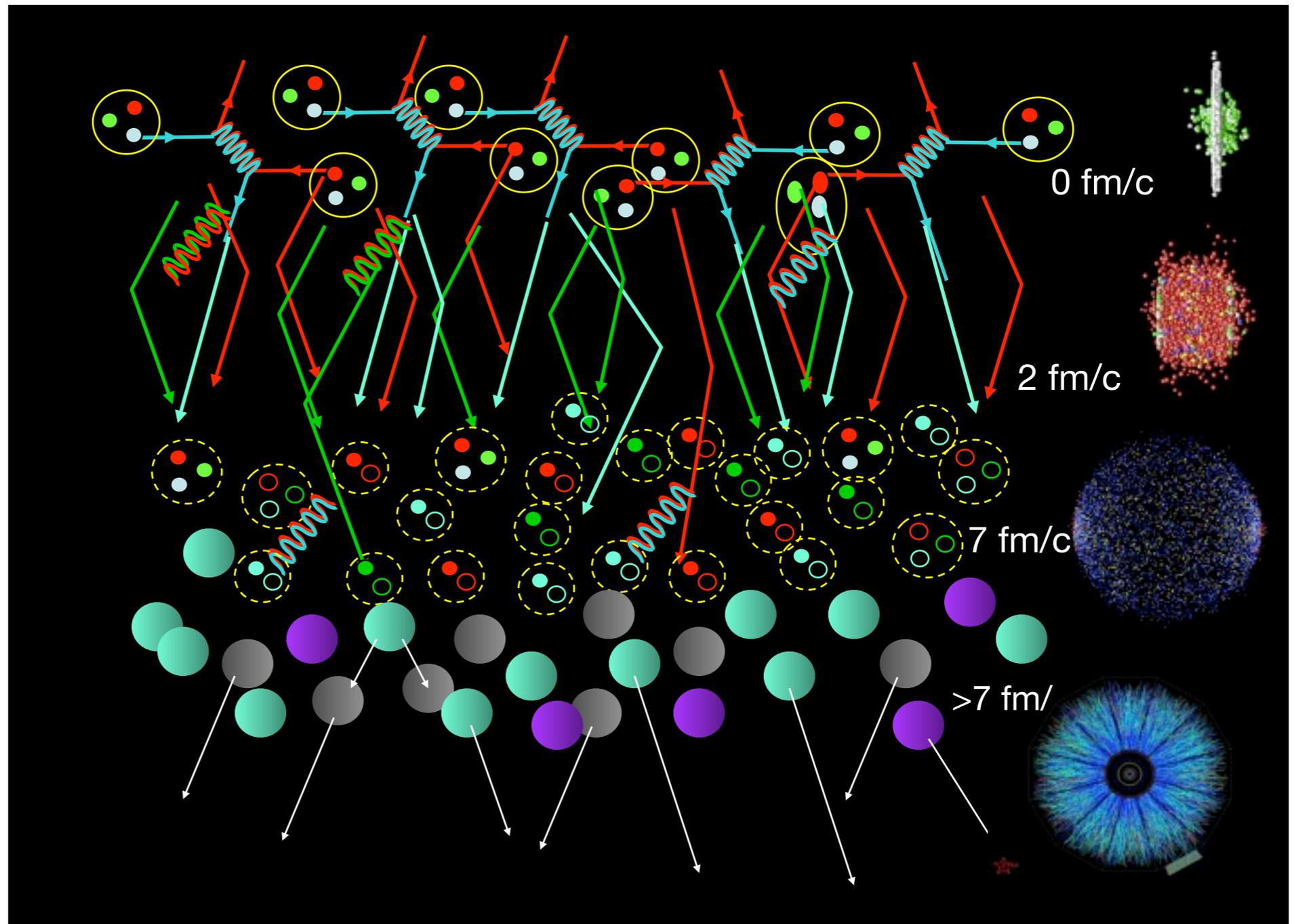
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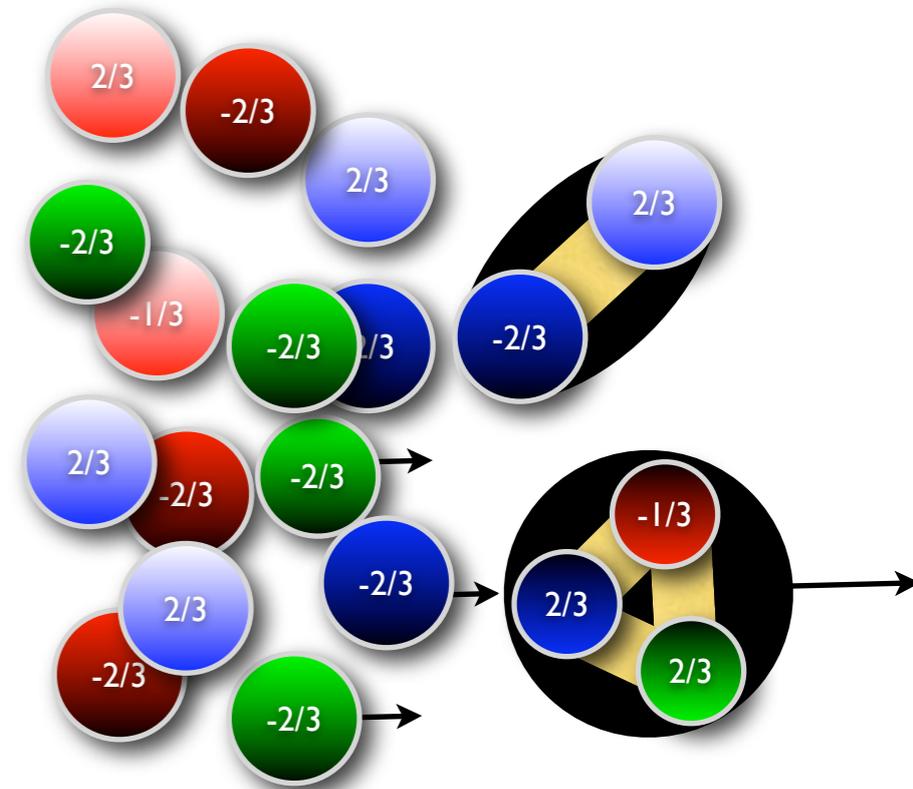
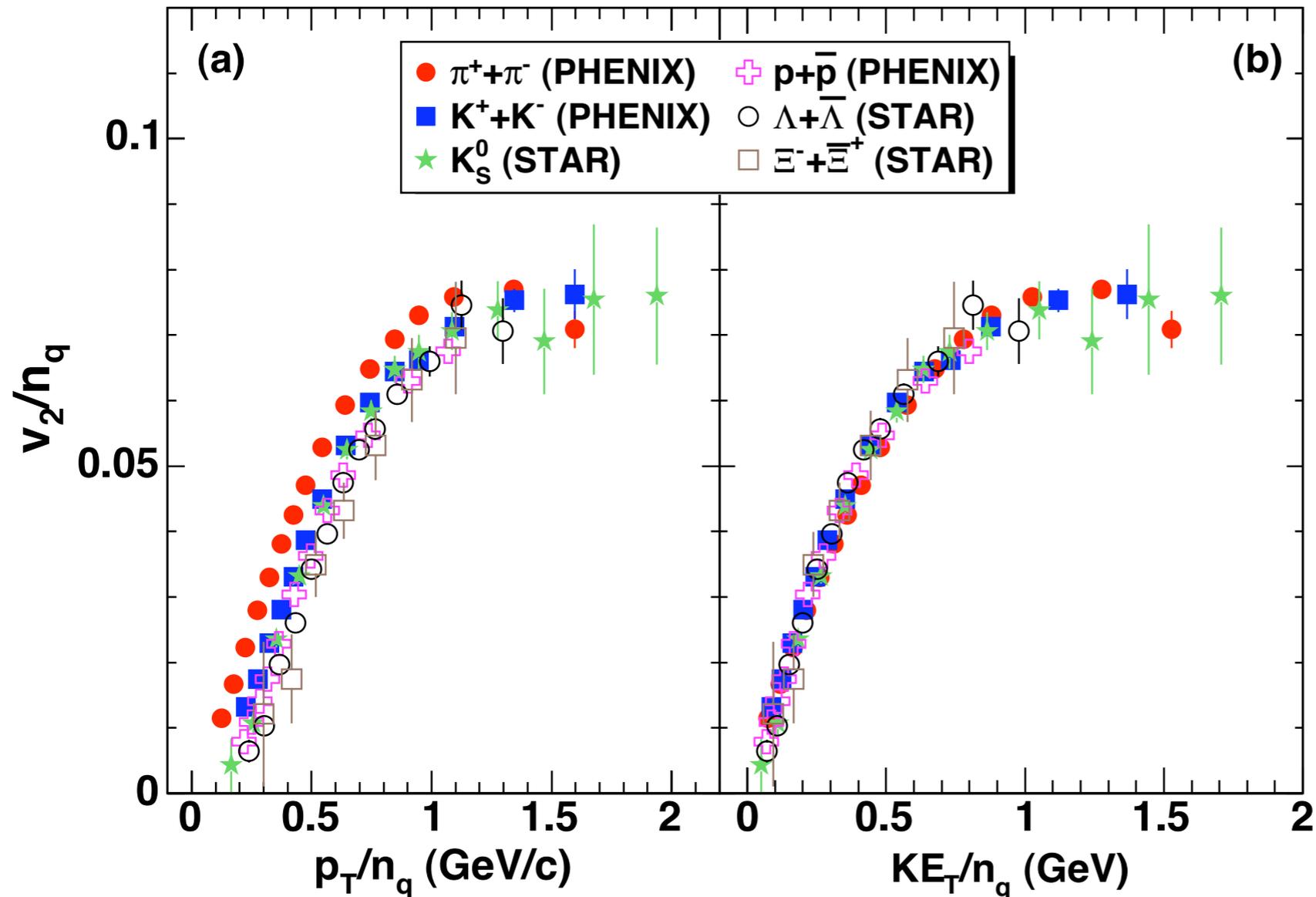
Hadron Rescattering

Thermal Freezeout &
Hadron decays



Constituent Quark Scaling

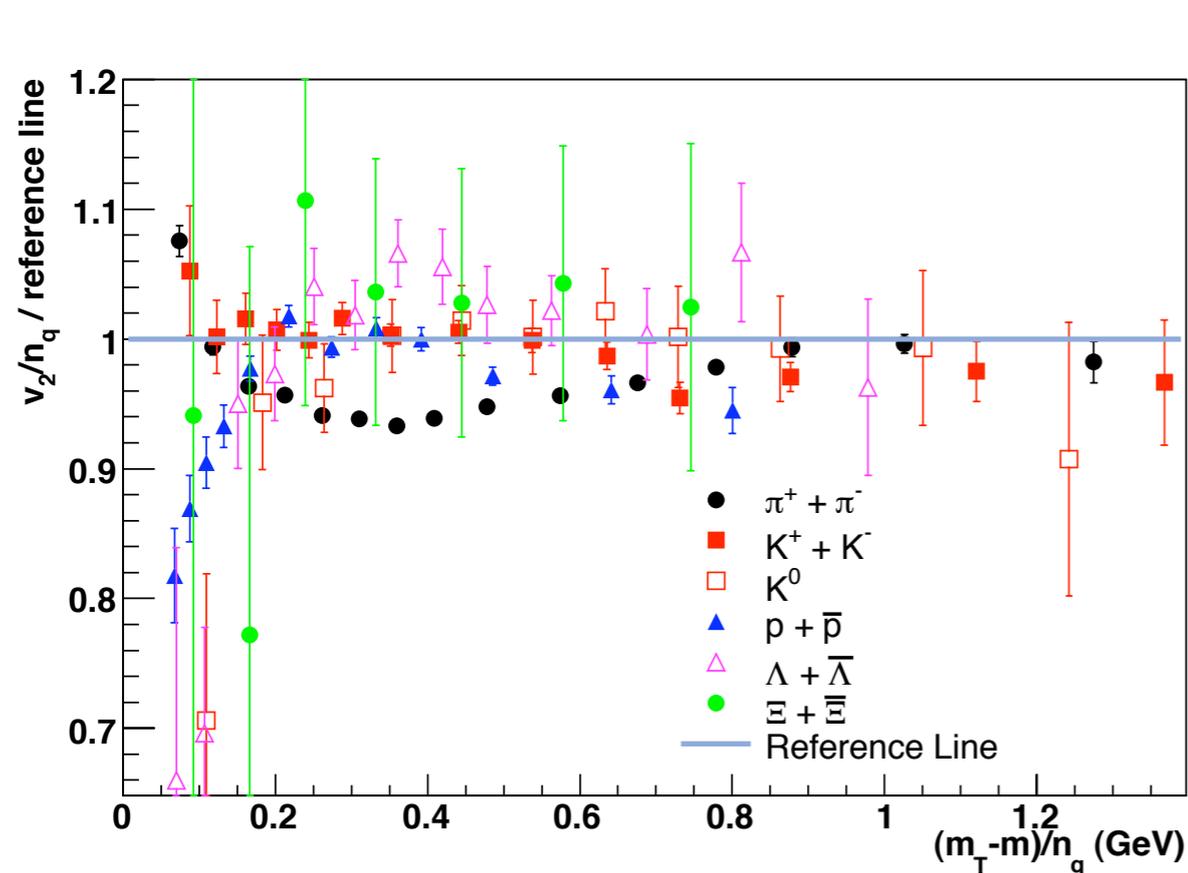
Phys.Rev.Lett.98:162301,2007



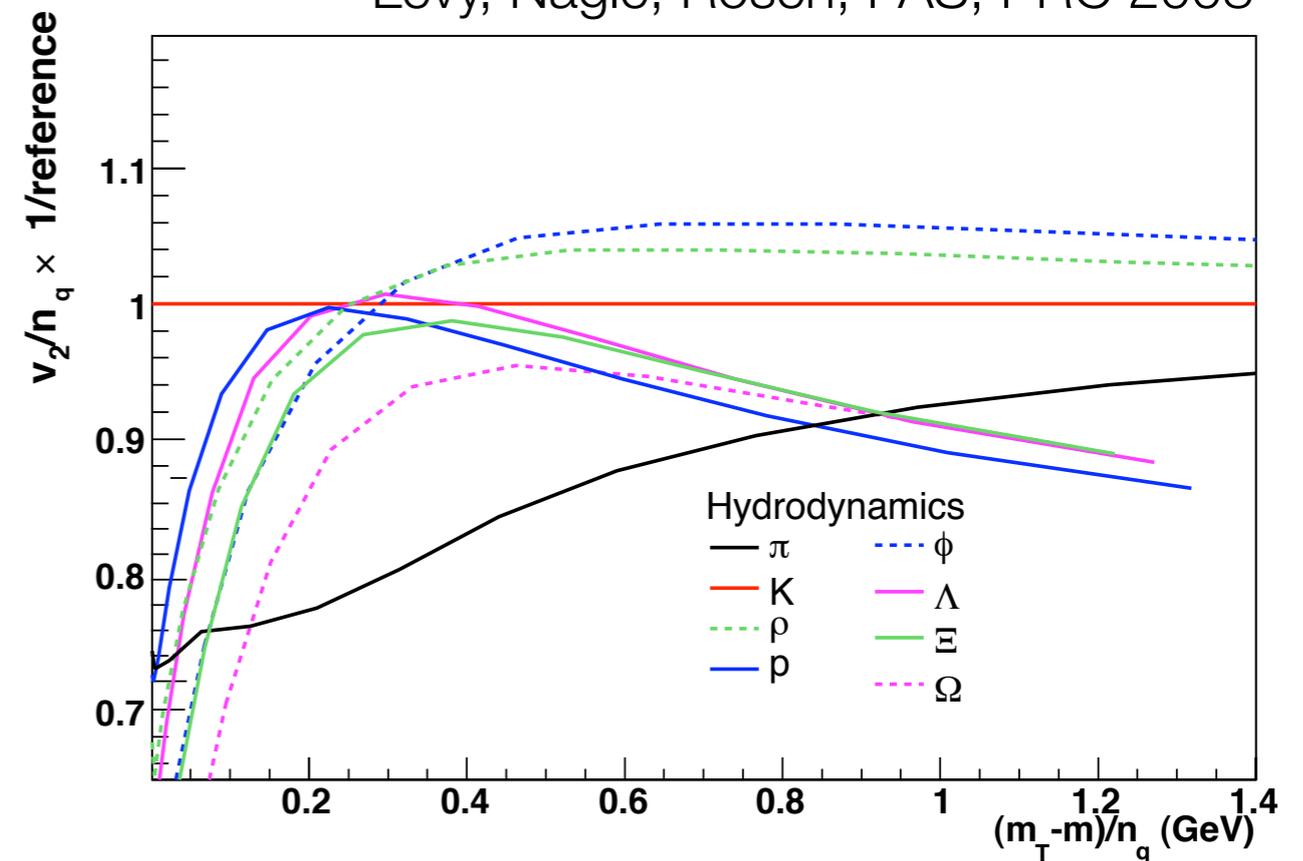
All particles flow as if frozen out from a flowing soup of constituent quarks

$$KE_T = m_T - m$$

Quasi-Particles in the sQGP



Levy, Nagle, Rosen, PAS, PRC 2008

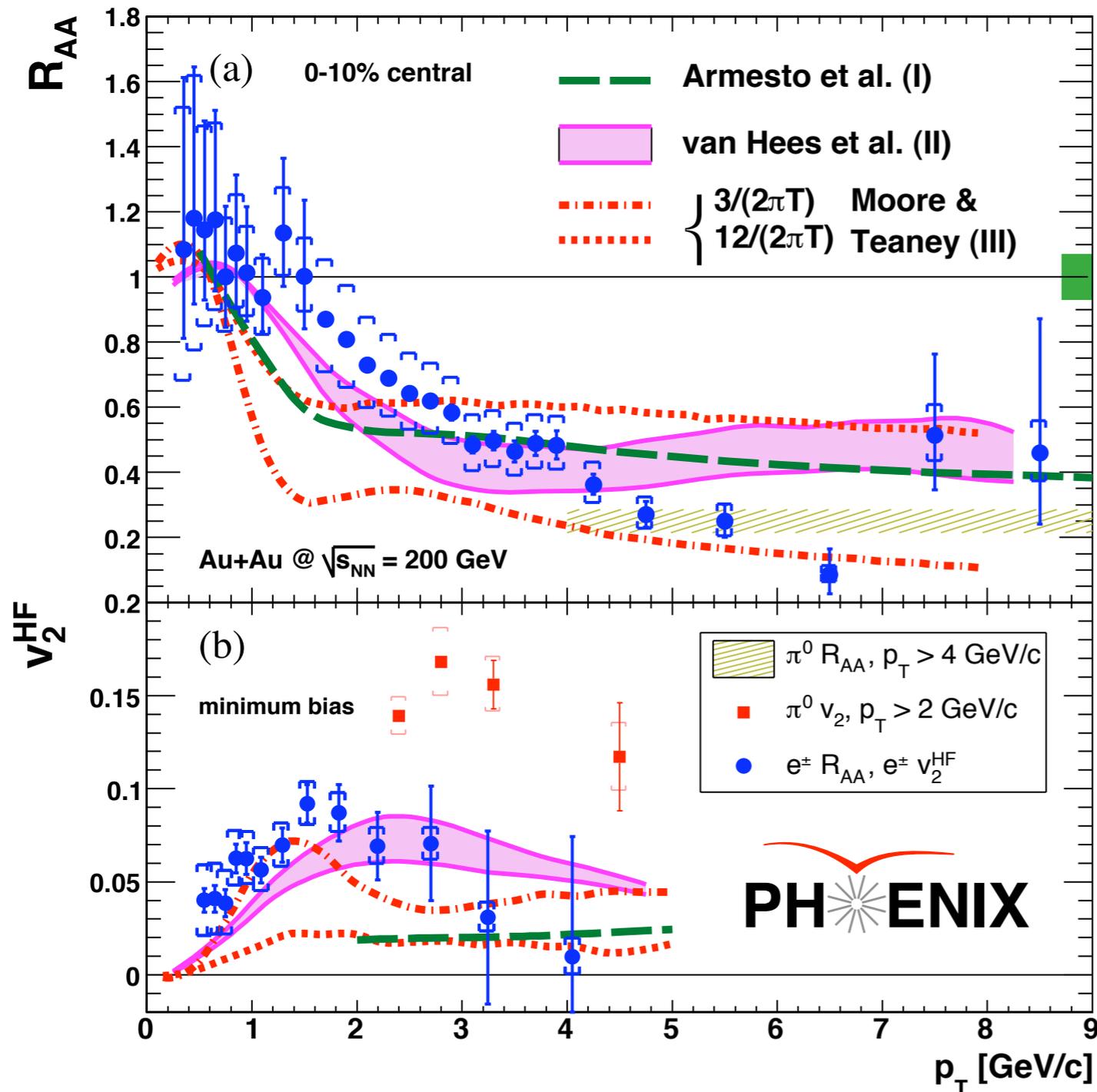


Data agree with n_q scaling substantially better than ideal hydro

“perfect fluid” \neq good quasiparticles:
can we harmonize two scenarios?

(or give up on constituent quarks, or the perfect fluid...)

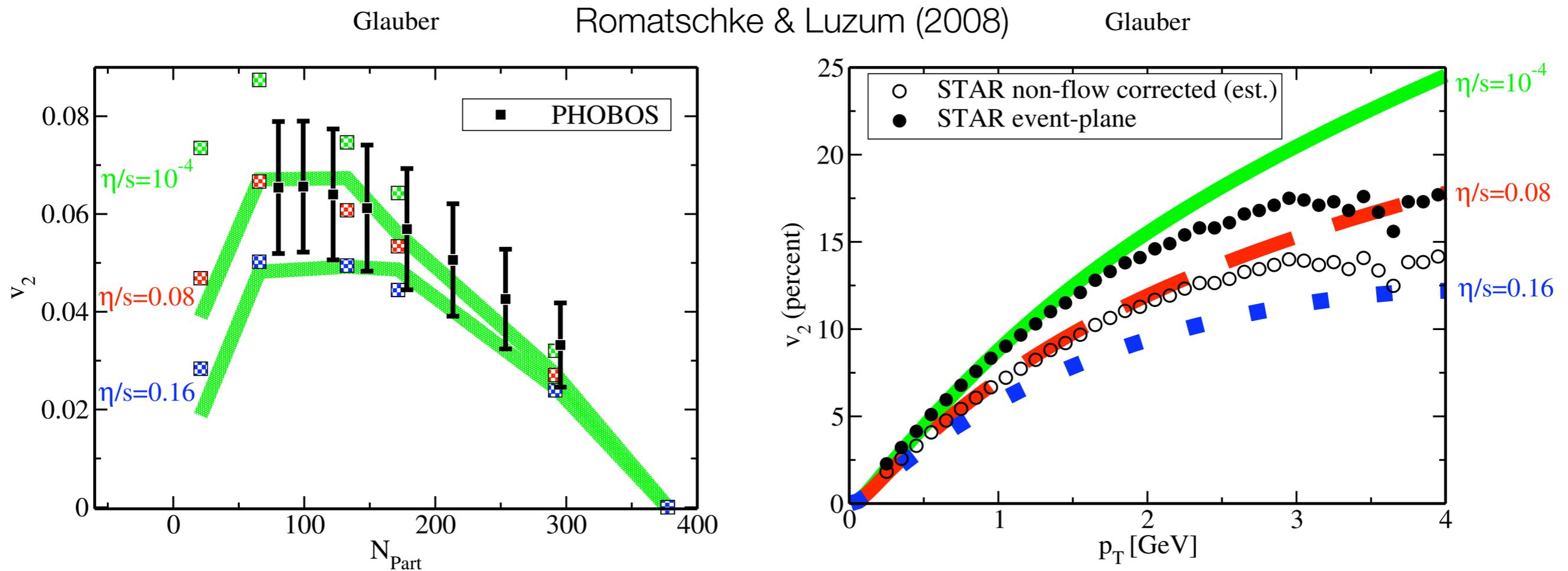
Direct Measurements of Viscosity: Heavy Flavor CTEQ



Charm suppression and flow (via non-photonic electrons) are correlated and reflect diffusion of heavy quark in medium (i.e. viscosity)

Models suggest non-zero viscosity, at or near AdS/CFT bound: challenge to quasi-particle interpretations of RHIC medium?

Viscous Hydrodynamics



Recent advances in implementing viscous hydrodynamics

With Glauber initial conditions, RHIC data saturating viscosity bound
(off by a factor of 2 in “Color Glass Condensate” approach)

Lessons from Soft Physics @ RHIC

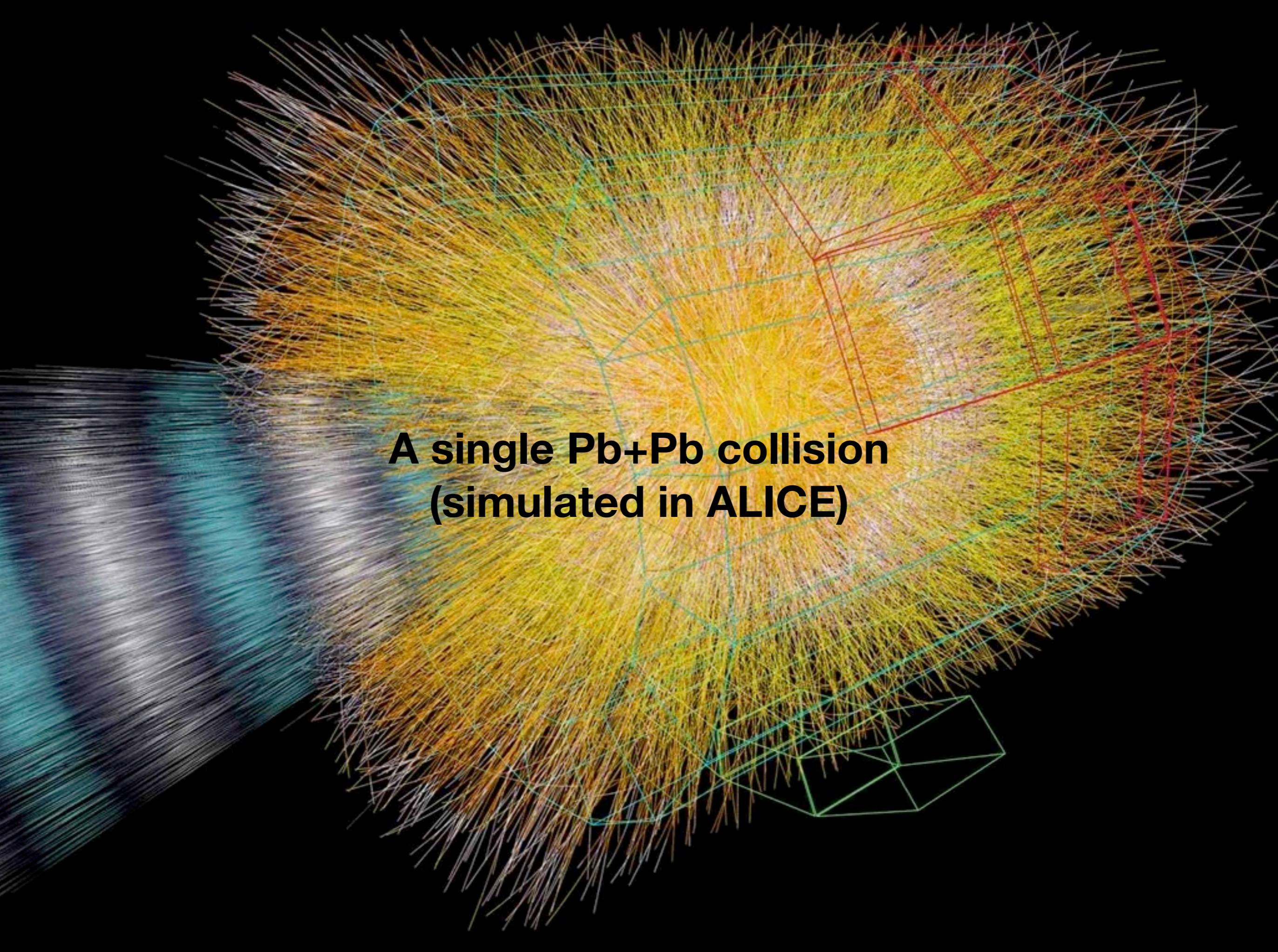
- **System is manifestly thermalized in final state**
 - Appears to freezeout at or near Hagedorn temperature
- **Hydro implies local thermalization in the initial state (and possibly presence of phase transition)**
 - System much hotter than T_H
- **No deviations in soft observables vs. rapidity, energy, size**
 - Extended longitudinal scaling ubiquitous in all systems
- **Viscosity appears to saturate bound from string theory**
 - How to harmonize with observation of constituent quark scaling, suggesting “quark” DOFs just before freezeout?

The Next Chapter: The LHC



Parameter	Units	Nominal	Early Beam
Energy per nucleon	TeV/n	2.76	2.76
Initial Luminosity L_0	$\text{cm}^{-2} \text{s}^{-1}$	$1 \cdot 10^{27}$	$5 \cdot 10^{25}$
No. bunches/bunch harmonic		592/891	62/66
Bunch spacing	ns	99.8	1350
β^*	m	0.5 (same as p)	1.0
Number of Pb ions/bunch		$7 \cdot 10^7$	$7 \cdot 10^7$
Transv. norm. RMS emittance	μm	1.5	1.5
Longitudinal emittance	eV s/charge	2.5	2.5
Luminosity half-life (1,2,3 expts.)	H	8, 4.5, 3	14, 7.5, 5.5

RHIC: 200 GeV/N Au+Au
 LHC: 5500 GeV/N Pb+Pb
 (14 TeV p+p)

The image depicts a simulated heavy-ion collision. On the left, a beam of particles, represented by a dense, multi-colored spray of lines (blue, cyan, and white), enters from the left. This beam impacts a target, resulting in a large, dense, and roughly spherical cloud of particles on the right. The particles in this cloud are primarily yellow and orange, with some white and grey lines interspersed. A grid of thin, colored lines (green, red, and blue) is overlaid on the particle cloud, forming a series of rectangular and hexagonal shapes that represent different regions or sectors within the collision volume. The background is solid black, which makes the particle tracks stand out prominently.

**A single Pb+Pb collision
(simulated in ALICE)**

Does the system still thermalize?

(or will running coupling increase viscous effects?)

What are the conditions (energy, density, size) for thermalization?

(Will p+p collisions clearly show it?)

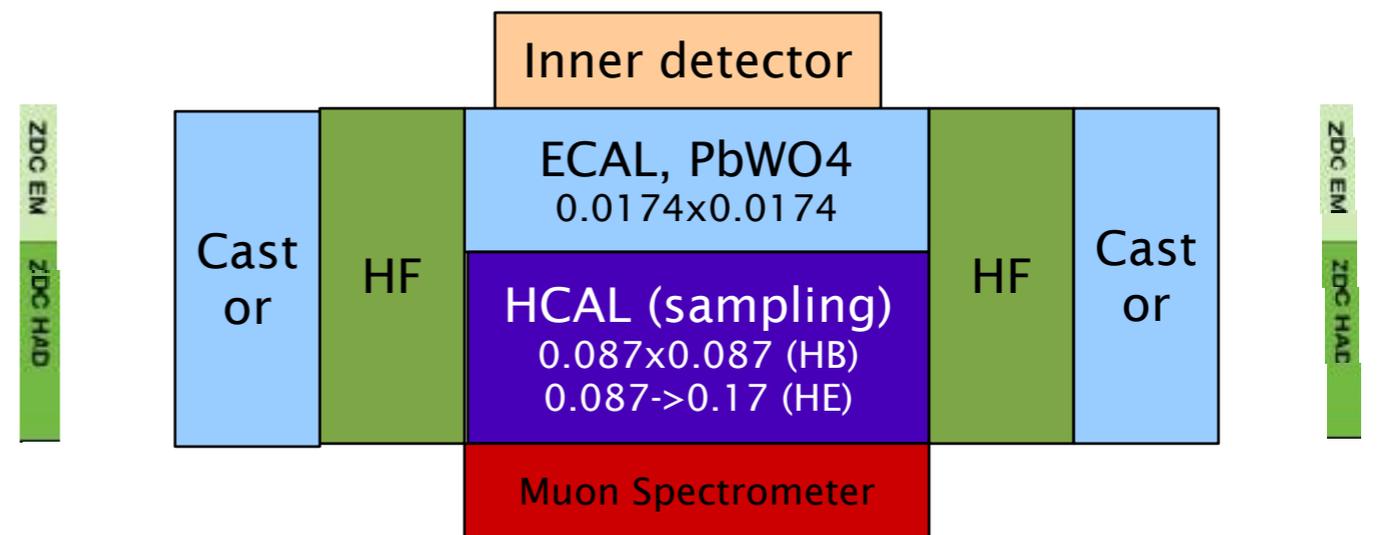
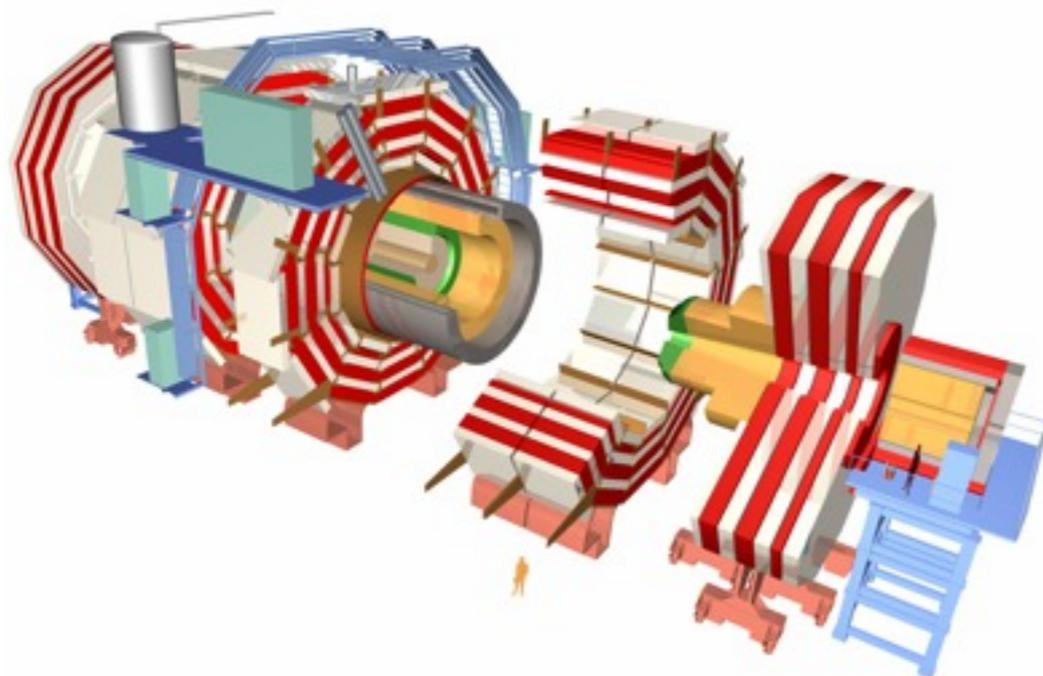
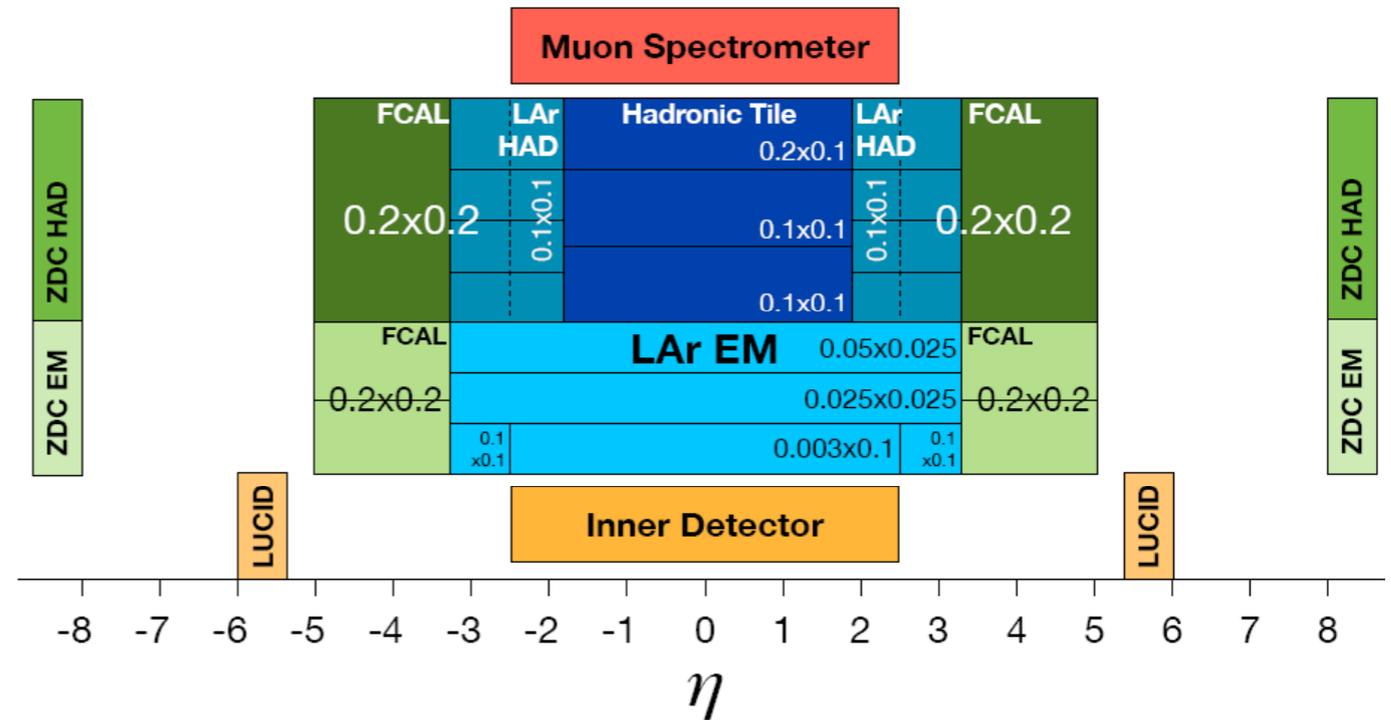
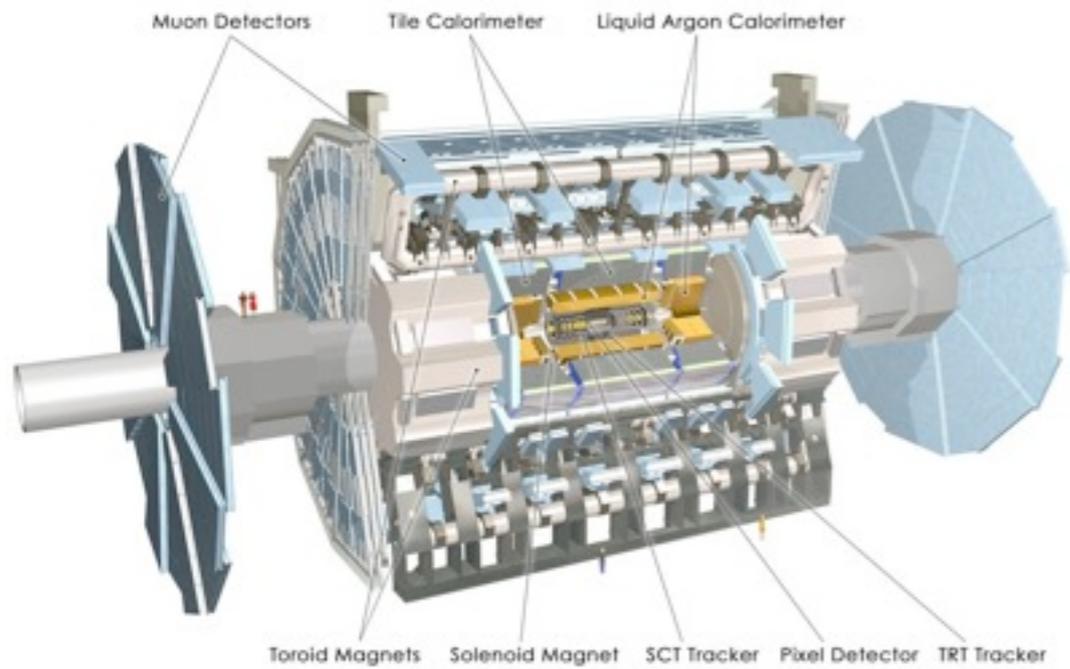
How fast does the system thermalize?

(initial longitudinal scale is 30x smaller!)

What thermalizes?

(Will degrees of freedom change?)

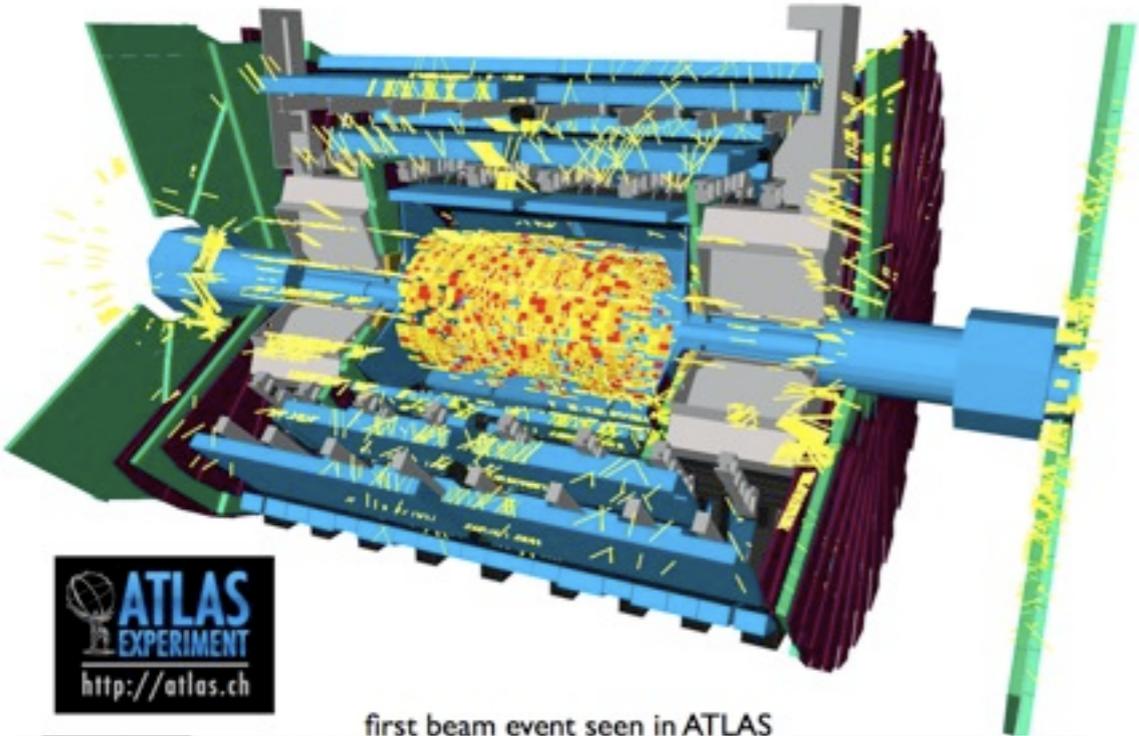
ATLAS & CMS @ LHC



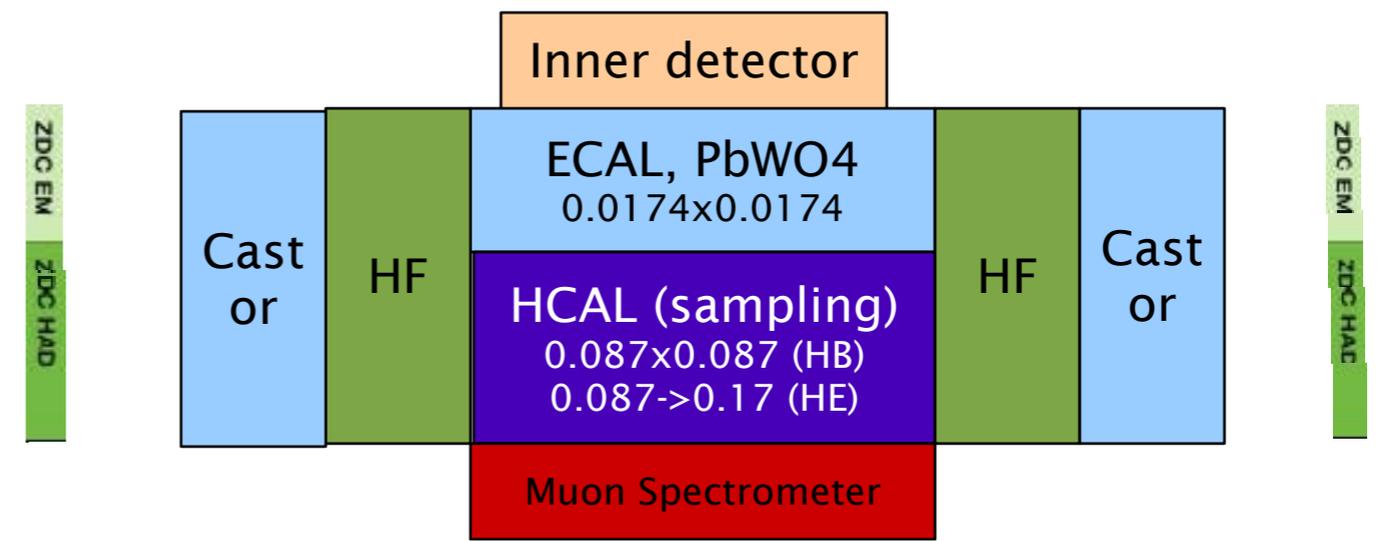
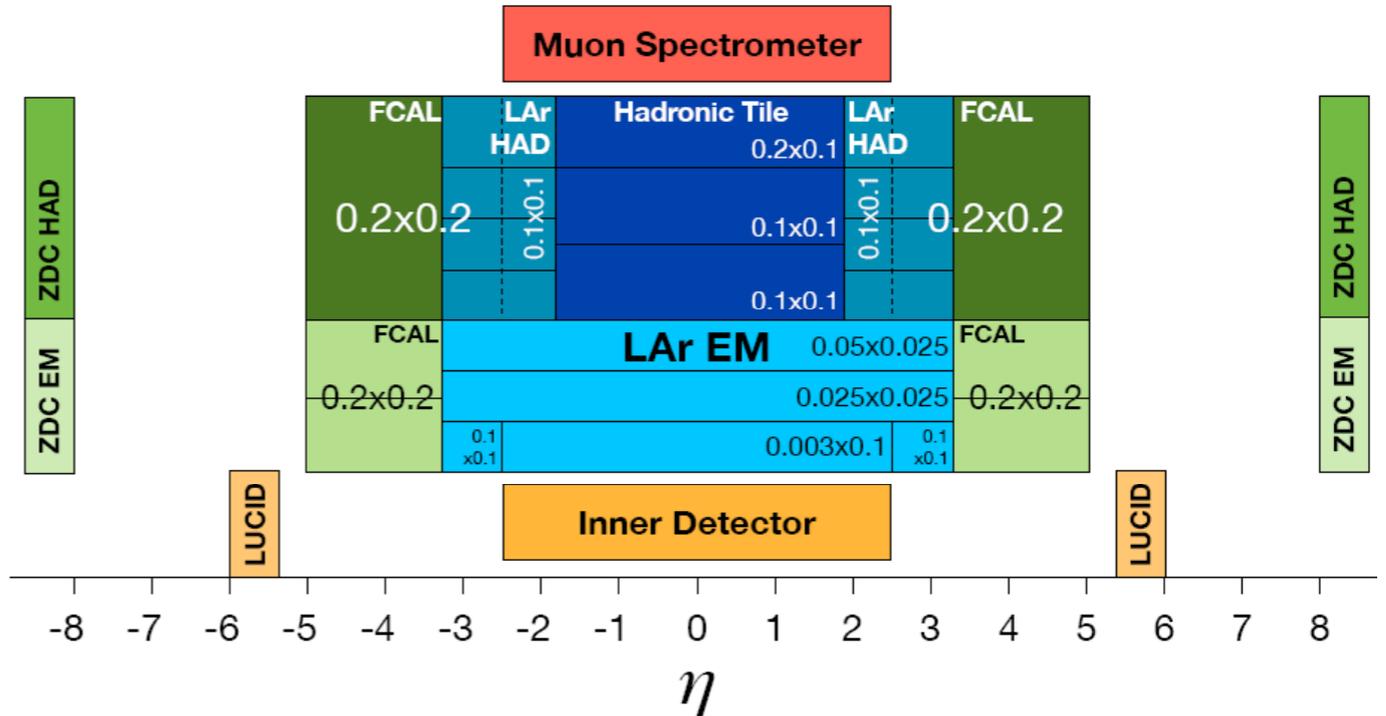
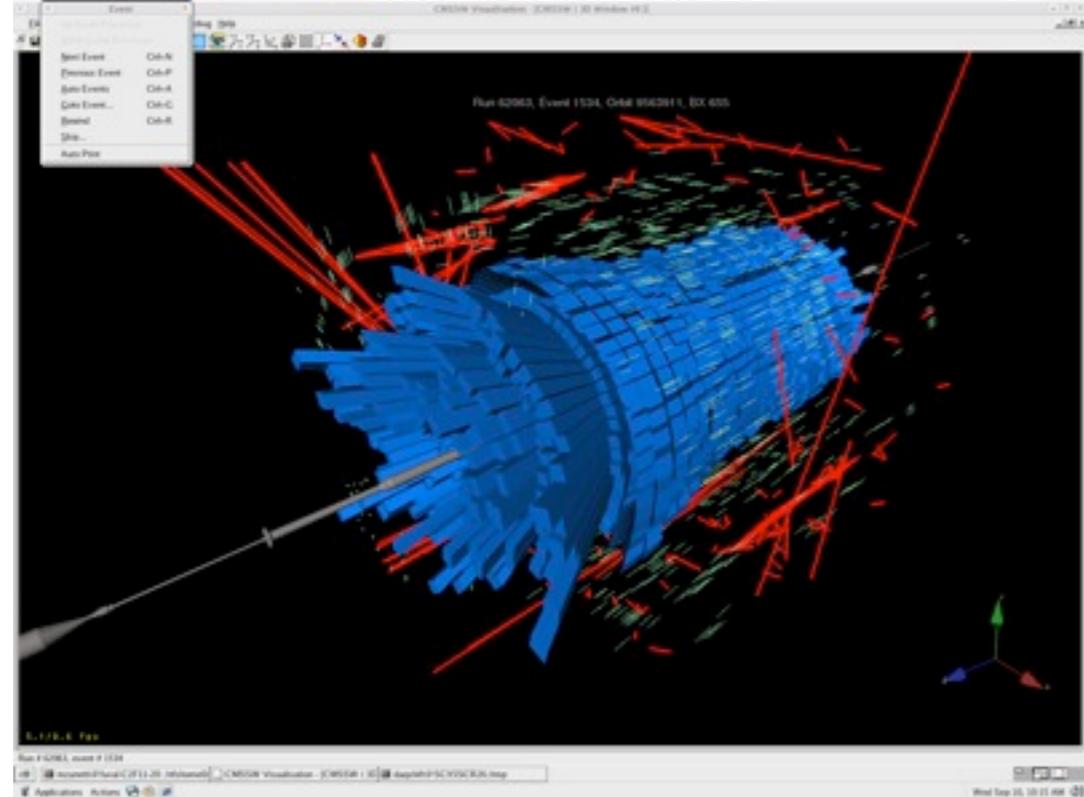
Enormous acceptance for soft physics

ATLAS & CMS @ LHC

9/10/08!



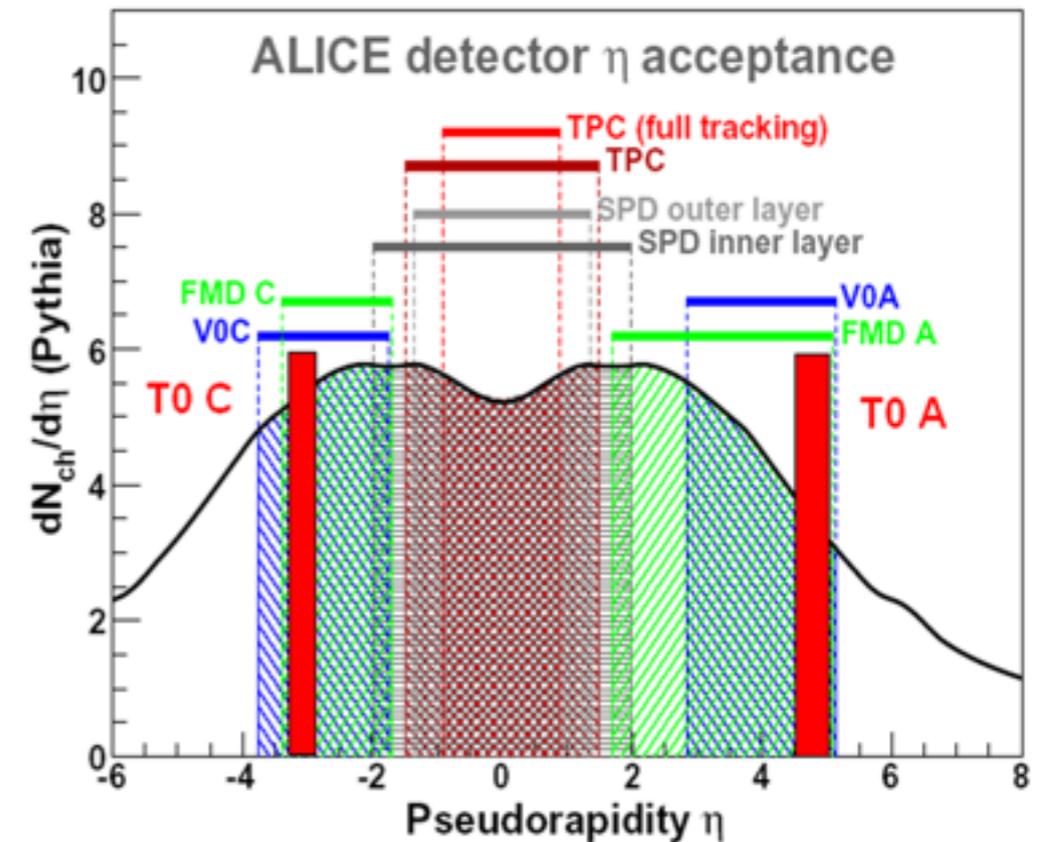
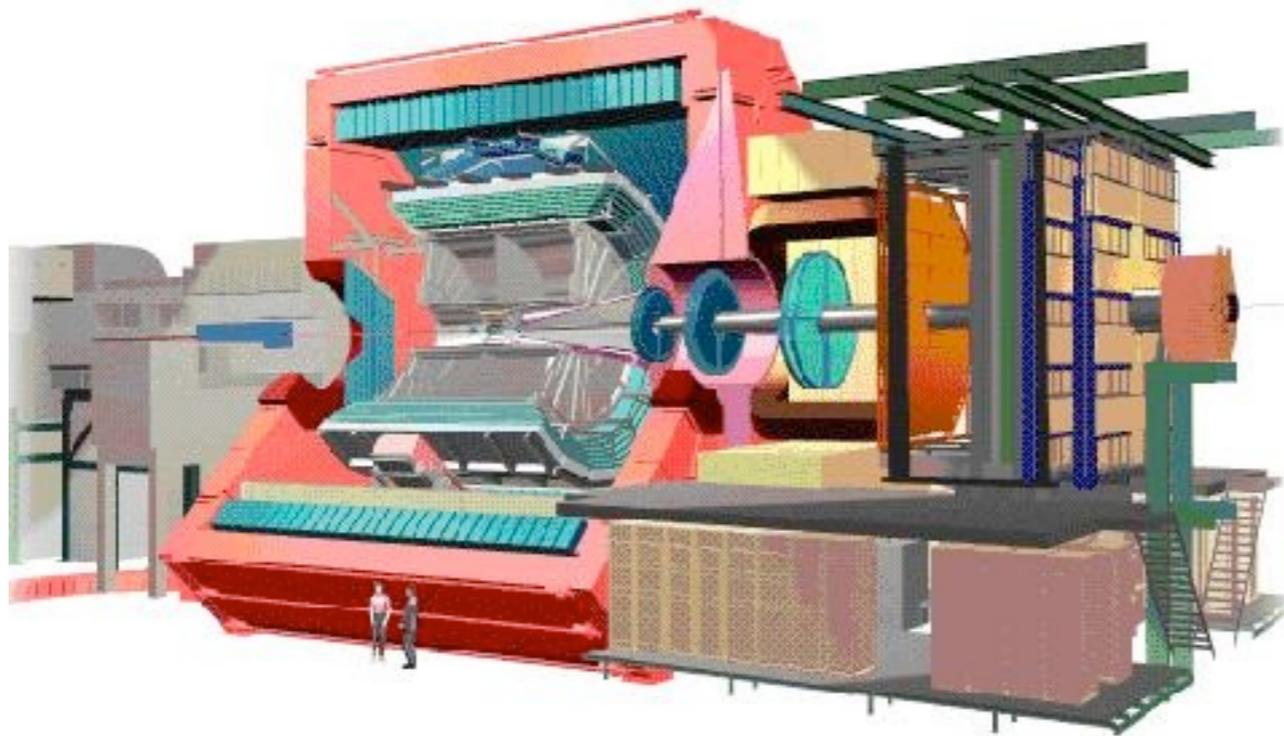
first beam event seen in ATLAS



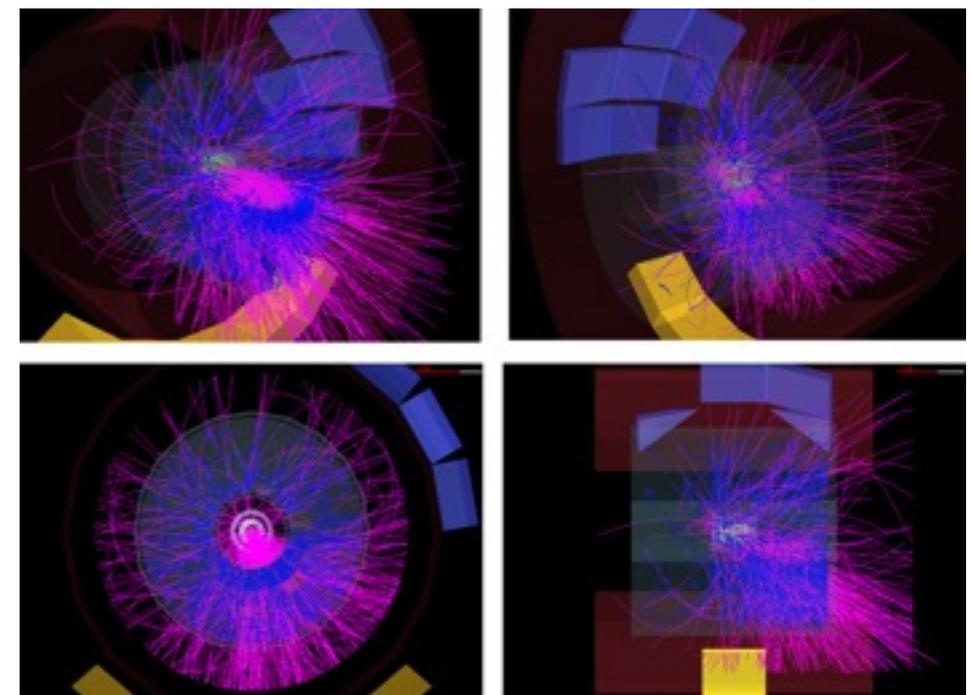
Enormous acceptance for soft physics

ALICE @ LHC

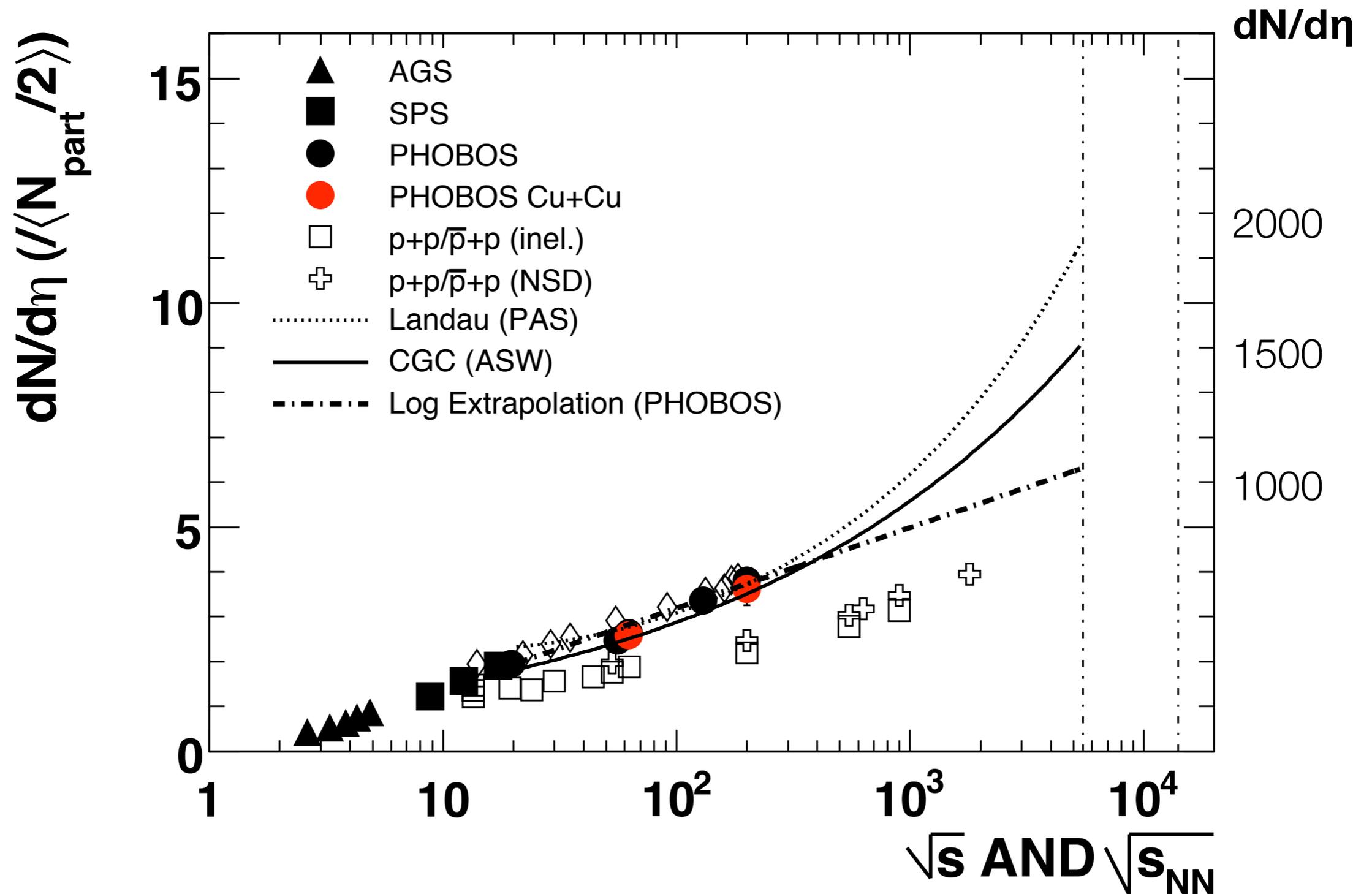
CTEQ



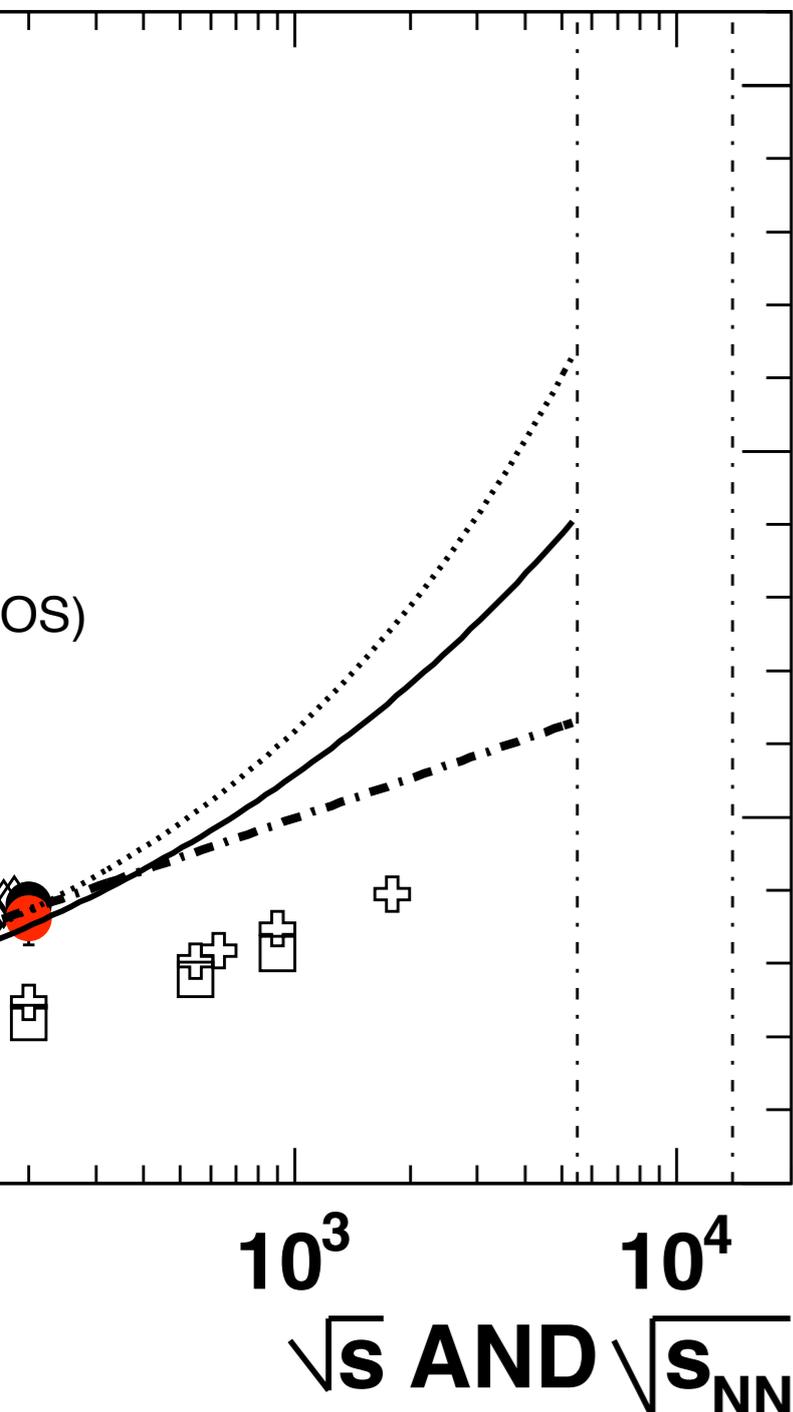
Optimized for tracking & PID
in high multiplicity
heavy ion events (TPC), with new
high p_T triggering abilities (EMCAL)



Predictions for the LHC: Multiplicity



Extrapolating to the LHC



Landau hydrodynamics: application of 1950's Fermi-Landau model & 1970's Carruther's approach

$$\frac{dN}{dy} = K s^{1/4} \frac{\exp(-y^2/2L)}{\sqrt{2\pi L}} \rightarrow \frac{s^{1/4}}{\sqrt{\ln(s)}} \quad L = \ln \left[\frac{s}{4m_p^2} \right]$$

Color Glass Condensate:

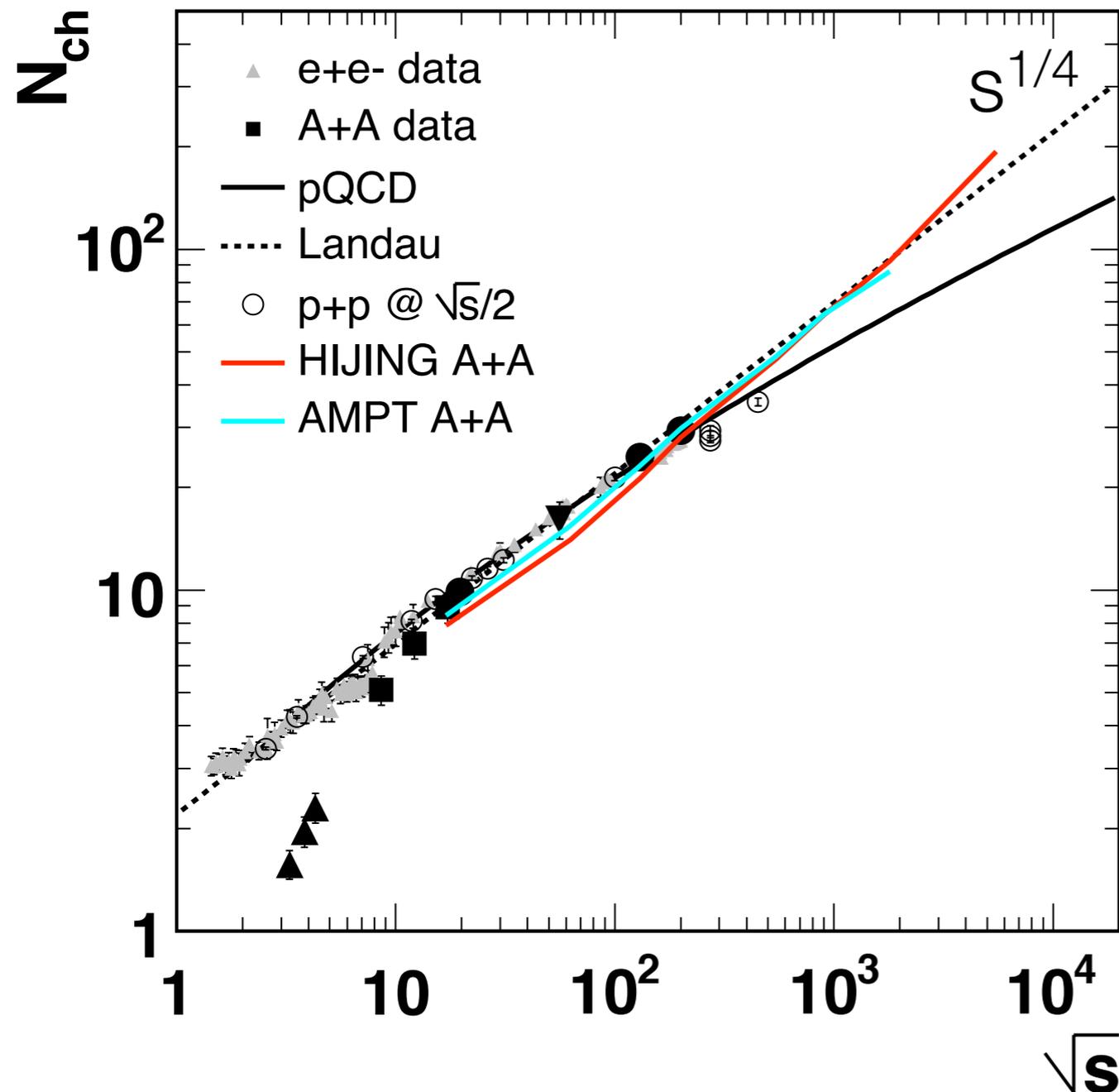
QCD at *very* low x . involves parametrization of gluon PDF and simple hadronization model

$$\frac{dN}{d\eta} = K s^\delta$$

Logarithmic rise: assumption that current trends apply at larger energies

Total Multiplicity @ LHC using data & models

CTEQ

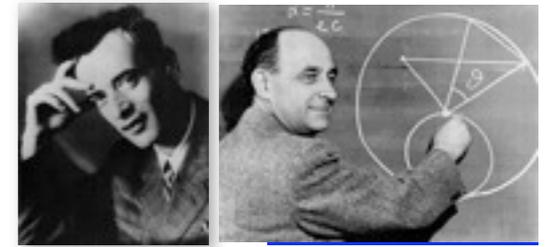


On a log-log plot,
all data is approximately
power-law

Where data exist, various
models can approximate it:
only the LHC provides test
(unfortunately total yield
is tough to measure!)

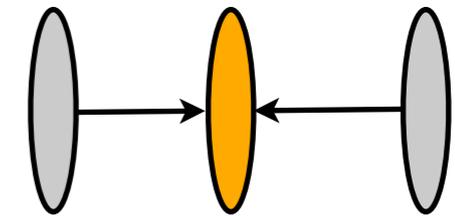
My favorite is Fermi-Landau
(power law rise)

Fermi-Landau Model for Total Multiplicity



- **Assume complete thermalization in Lorentz-contracted volume**

- Total energy scales as beam energy: $E = \sqrt{s}$
- Total volume scales as $V = V_0/2m\sqrt{s}$
- Energy density is *quadratic* in beam energy: $\epsilon = E/V = (2m/V_0)s$



- **Blackbody radiation: $p = \epsilon/d$ ($d = \#$ spatial dimensions, e.g. 3)**

- Related to tracelessness of stress-energy tensor
- When $d=3$, entropy density $\sigma \sim \epsilon^{3/4}$ (or $d/(d+1)$)

$$T_{\mu\nu} = \begin{pmatrix} \epsilon & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix}$$

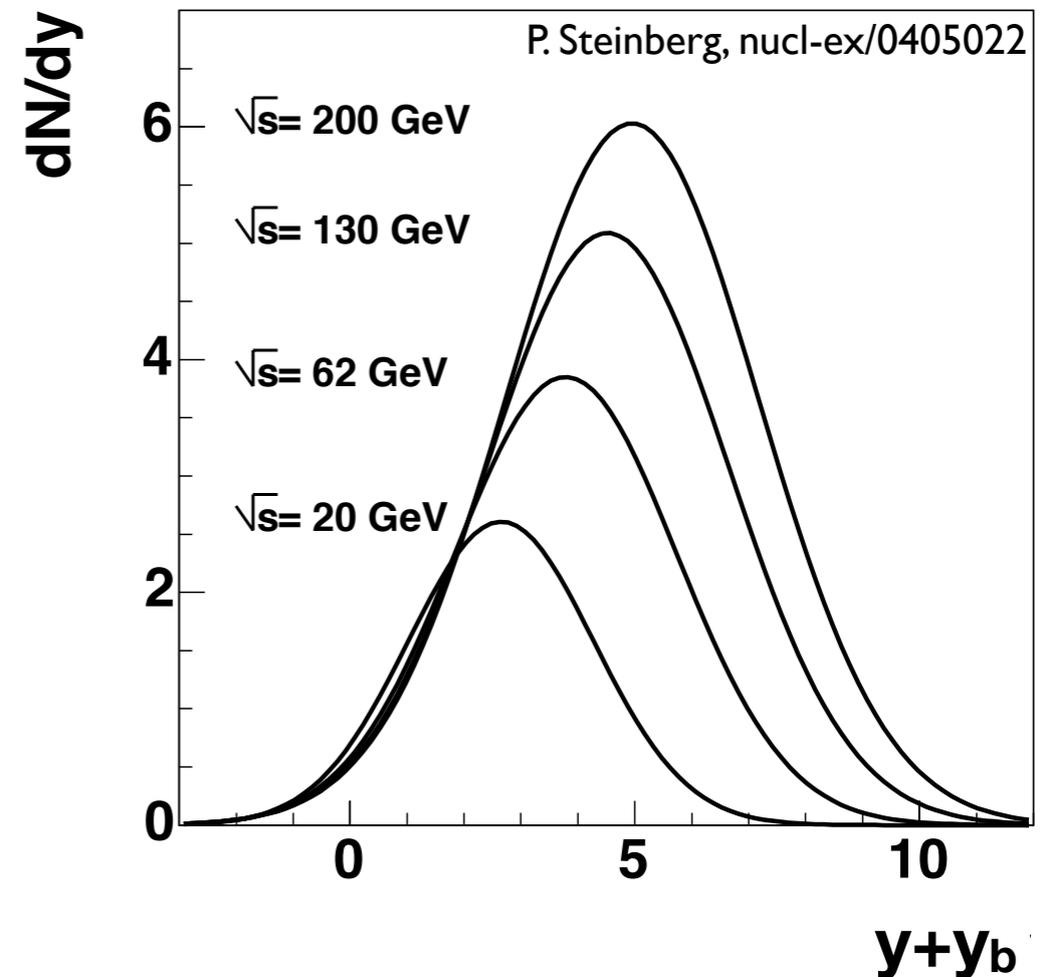
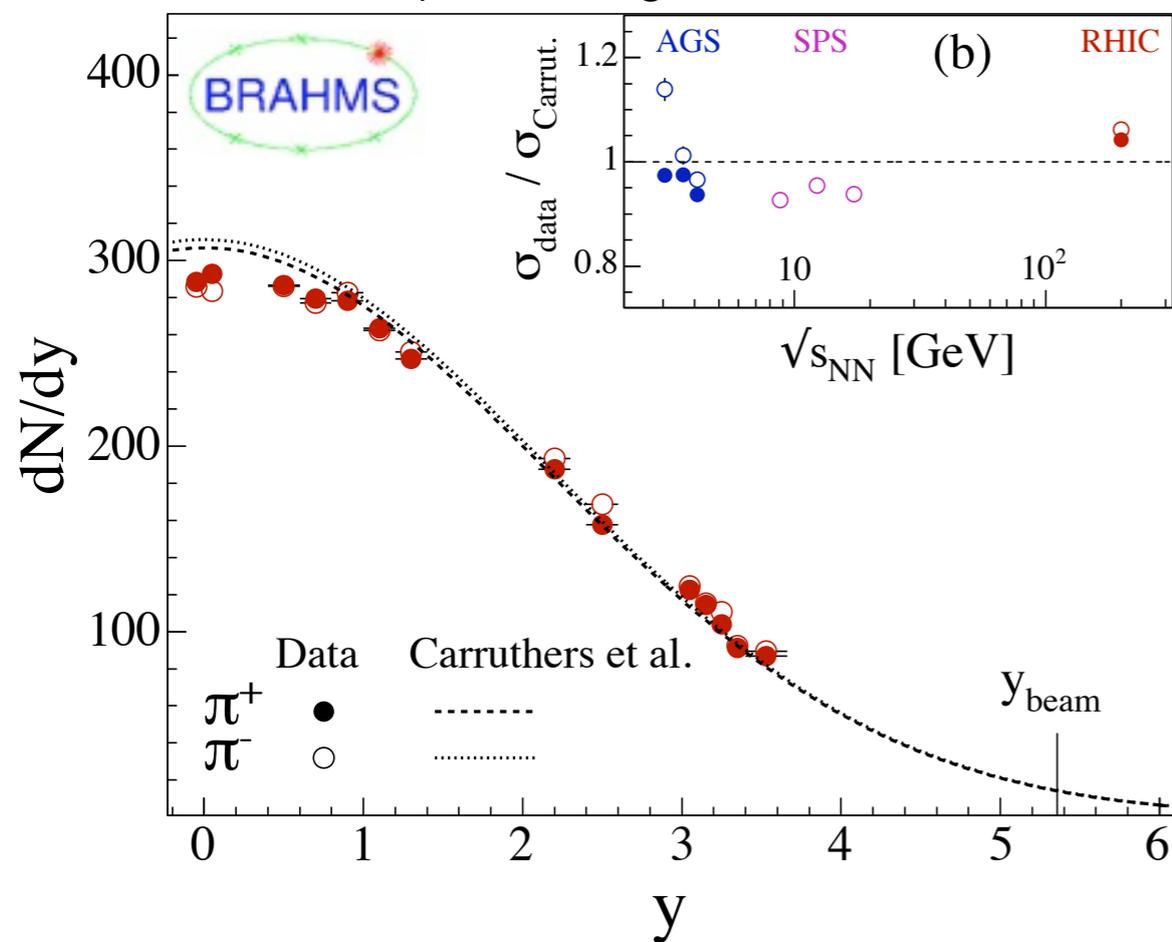
- **Total entropy $S = \sigma V \sim s^{3/4}/s^{1/2} \sim s^{1/4} = \sqrt{E}$
(or $s^{0.5(d-1)/(d+1)}$)**

- Total multiplicity assumed linear with N_{tot}
- N_{tot} assumed linear with N_{ch}

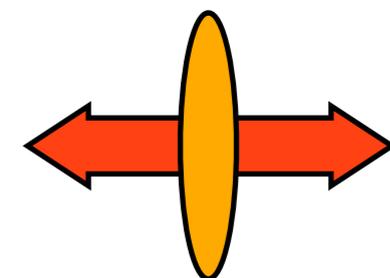
- **Putting it all together: $N_{\text{ch}} \sim s^{1/4}$**

Gaussian Rapidity Distributions

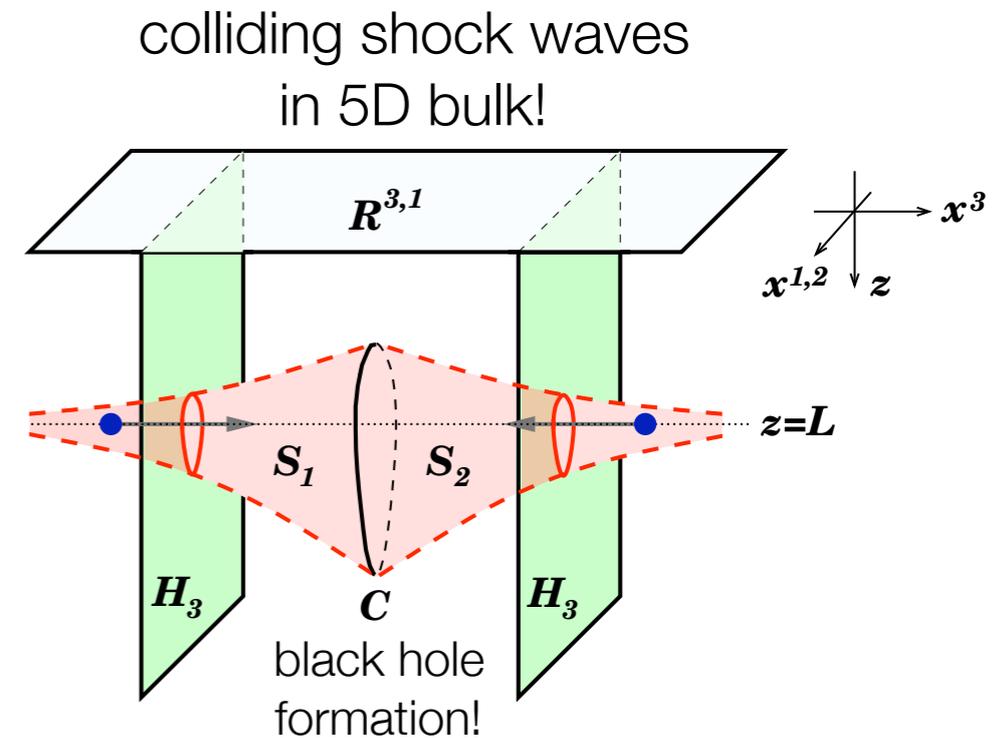
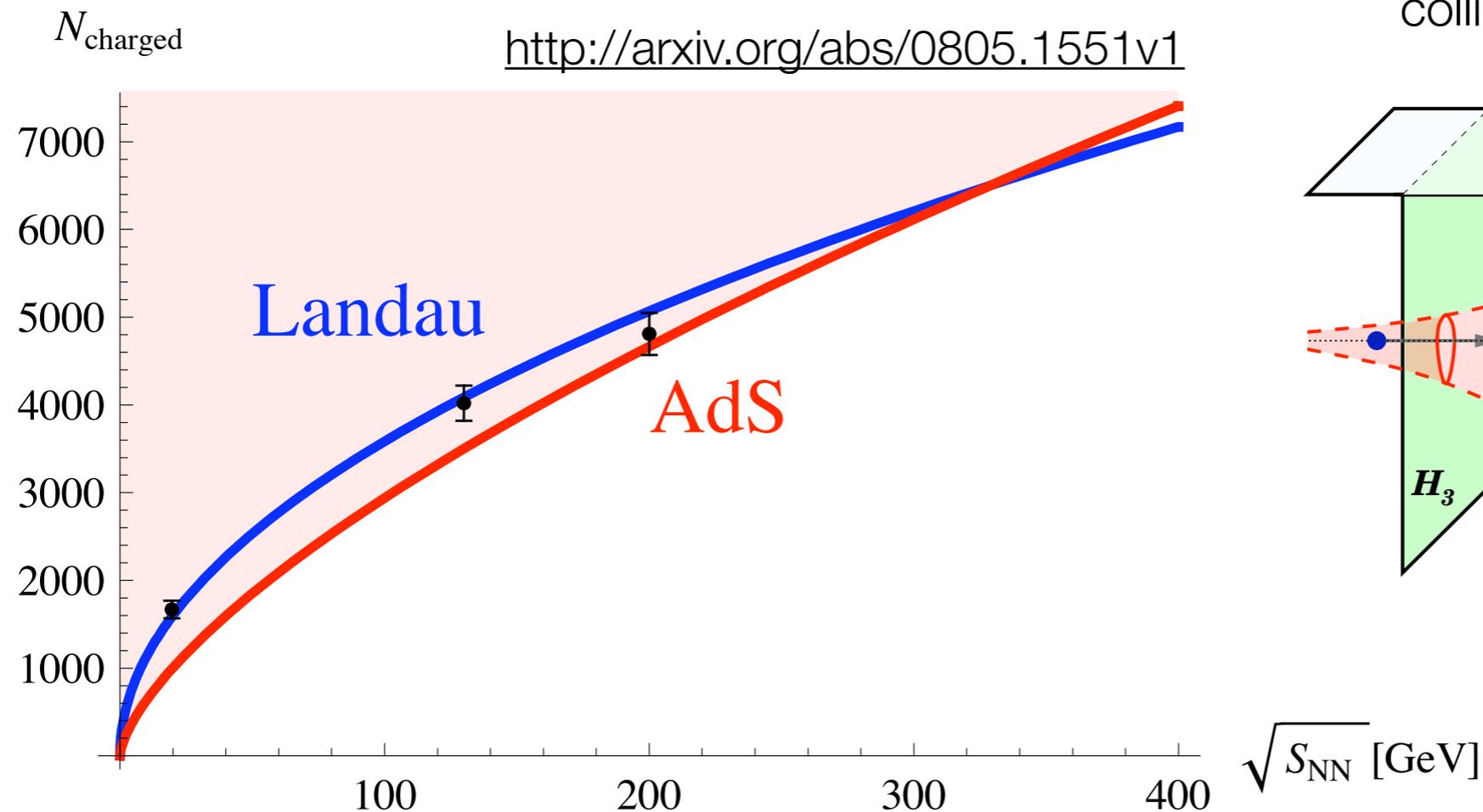
<http://arxiv.org/abs/nucl-ex/0403050v2>



Hydrodynamics maps Lorentz contracted pancake into a Gaussian in rapidity with variance $\sigma^2 = \ln(s/m_P)$. Agrees with data on dN/dy and naturally produces longitudinal scaling seen in data (all of it...)



Total Multiplicity from AdS/CFT



AdS/CFT now in the game: using colliding shock waves, and calculating area of trapped surface, Gubser et al predict $N_{\text{ch}} \sim s^{1/3}$

Substantially faster than measured rise of multiplicity,
1/3 implies $d=5$: true by construction?! Anyway, we'll test!

Soft Observables: From RHIC to LHC

- **In the soft sector, RHIC has observed:**
 - **Elliptic flow** suggesting early thermalization (near perfect fluid)
 - **Extended longitudinal scaling** (trivial or deep?)
 - Intriguing **connections to elementary systems**
 - **Statistical (& constituent quark) freezeout**
- **Empirical and theoretical trends have emerged from the extensive data set, requiring testing at the LHC**
 - “**Day 1**” measurements in 2009 (p+p) and 2010 (A+A)
 - Lots will be learned very quickly -- stay tuned!
- **The collider era of heavy ion physics is upon us!**
 - We need all the help we can get to understand what we see

Help Wanted (Lattice? AdS/CFT?)

- **Even if prediction of η/s is “true”: still lots to do**
- **Strongly-coupled sector of QCD is a real challenge**
 - Lattice is a powerful technique, but it requires analysis and interpretation
 - Data is simpler than expected, and thus more complicated to explain
- **Many basic properties of baryon still not understood: how do soft interactions proceed?**
 - Is there a fundamental (& simple) way to understand hadron structure?
 - pQCD works for jets, etc., but is inapplicable for small momentum transfers
 - What are colliding shock waves on gauge theory side, as interaction occurs?
- **How & when does thermalization occur?**
 - Again, describing system as a black hole doesn't explain the formation of the black hole
 - Attempts have been made (e.g. Shuryak, Kovchegov, Nastase), but have not yet led to novel predictions etc.

Heavy Ion Collisions: Hard Probes

CTEQ

Peter Steinberg

Brookhaven National Laboratory
Upton, NY



CTEQ Summer School

July 2, 2009

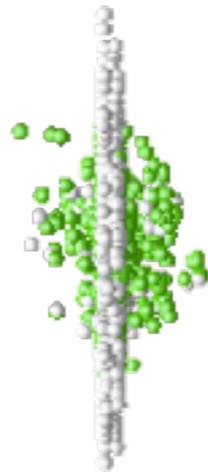
Madison, WI, USA

The promise of hard probes

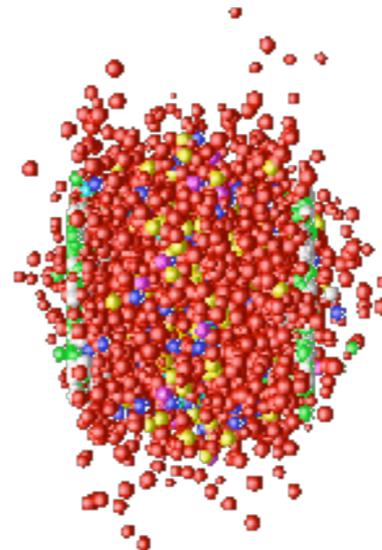
CTEQ



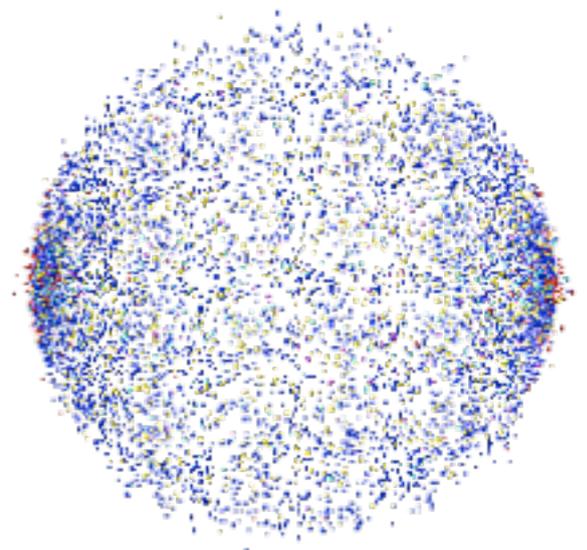
Colliding
nuclei



Hard
scattering

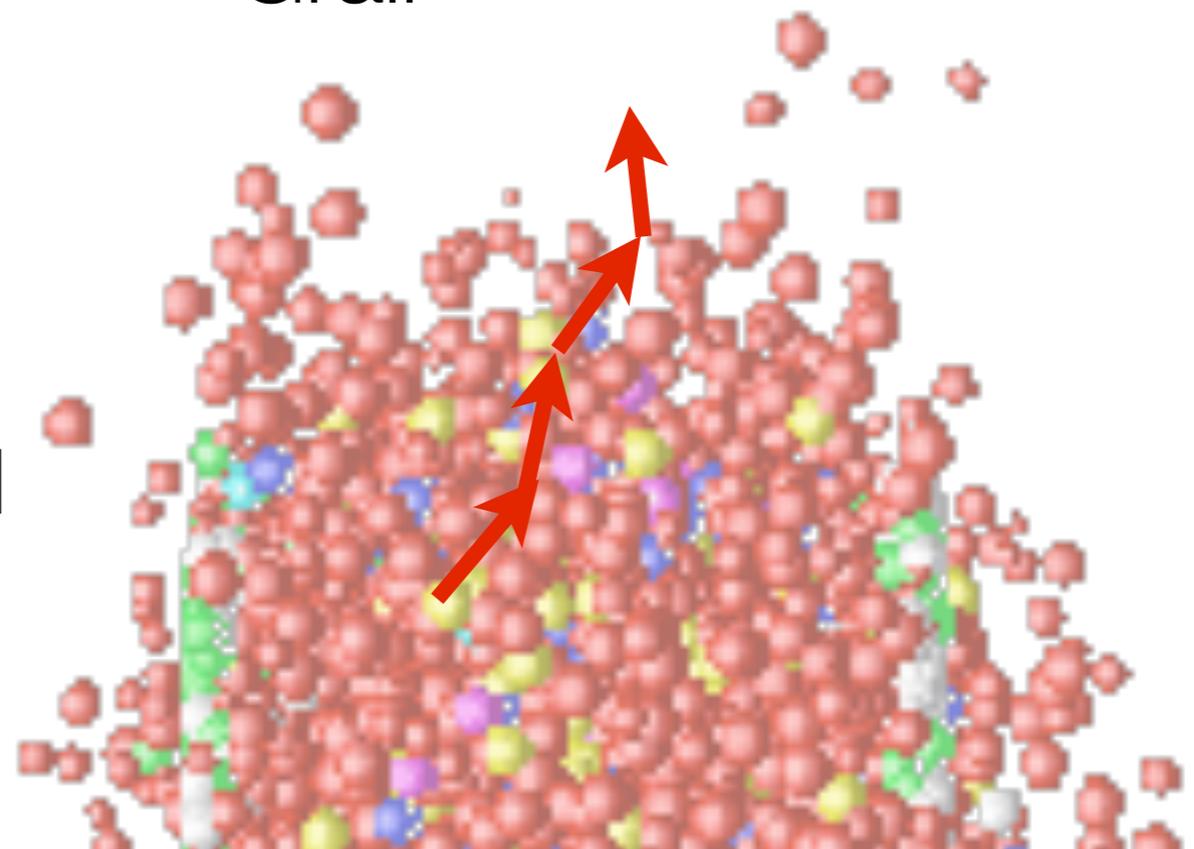


Holy
Grail

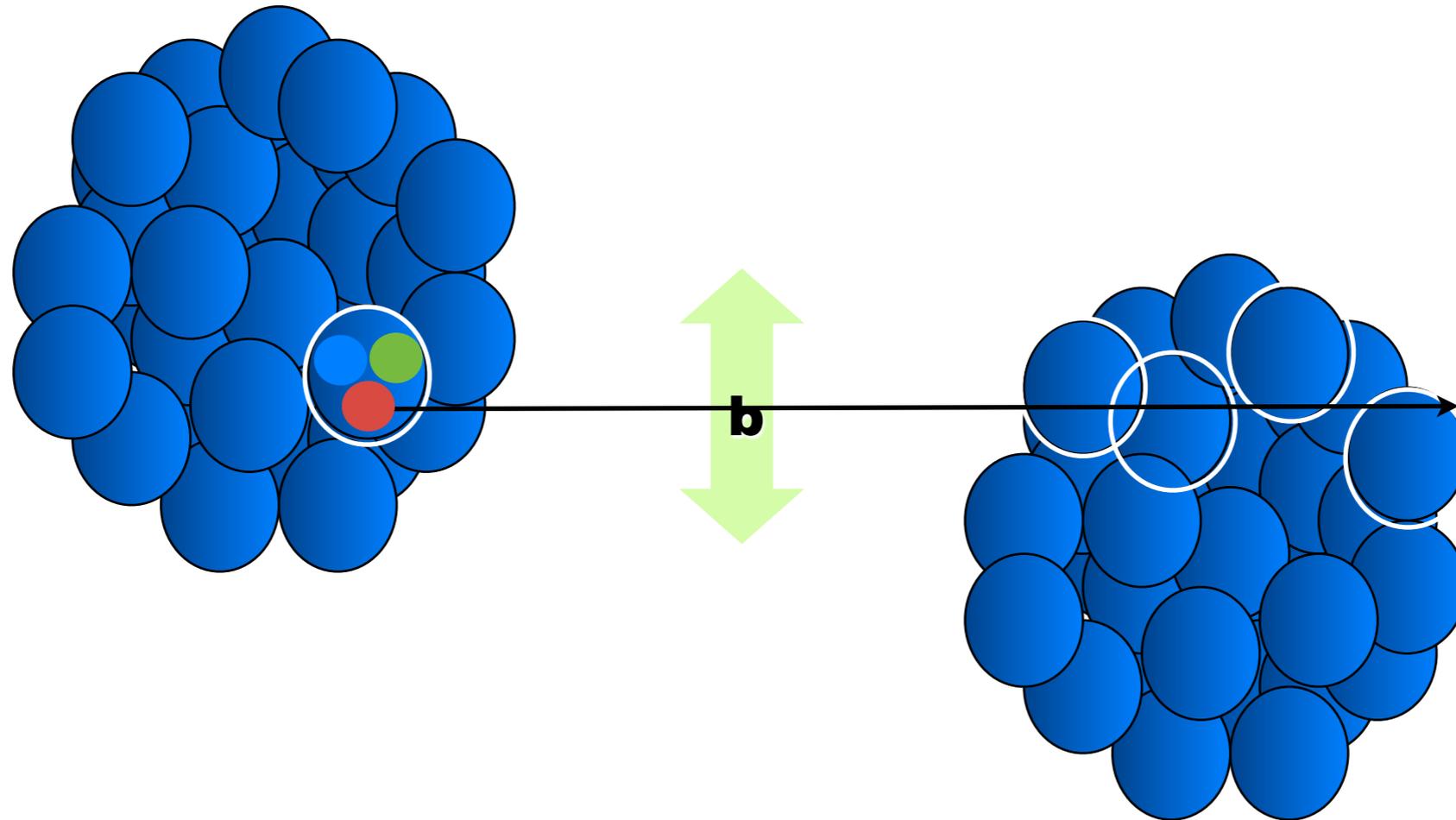


Final hadrons

To understand the microscopic properties of the sQGP, need a well-defined “probe” that can scatter off of its constituents

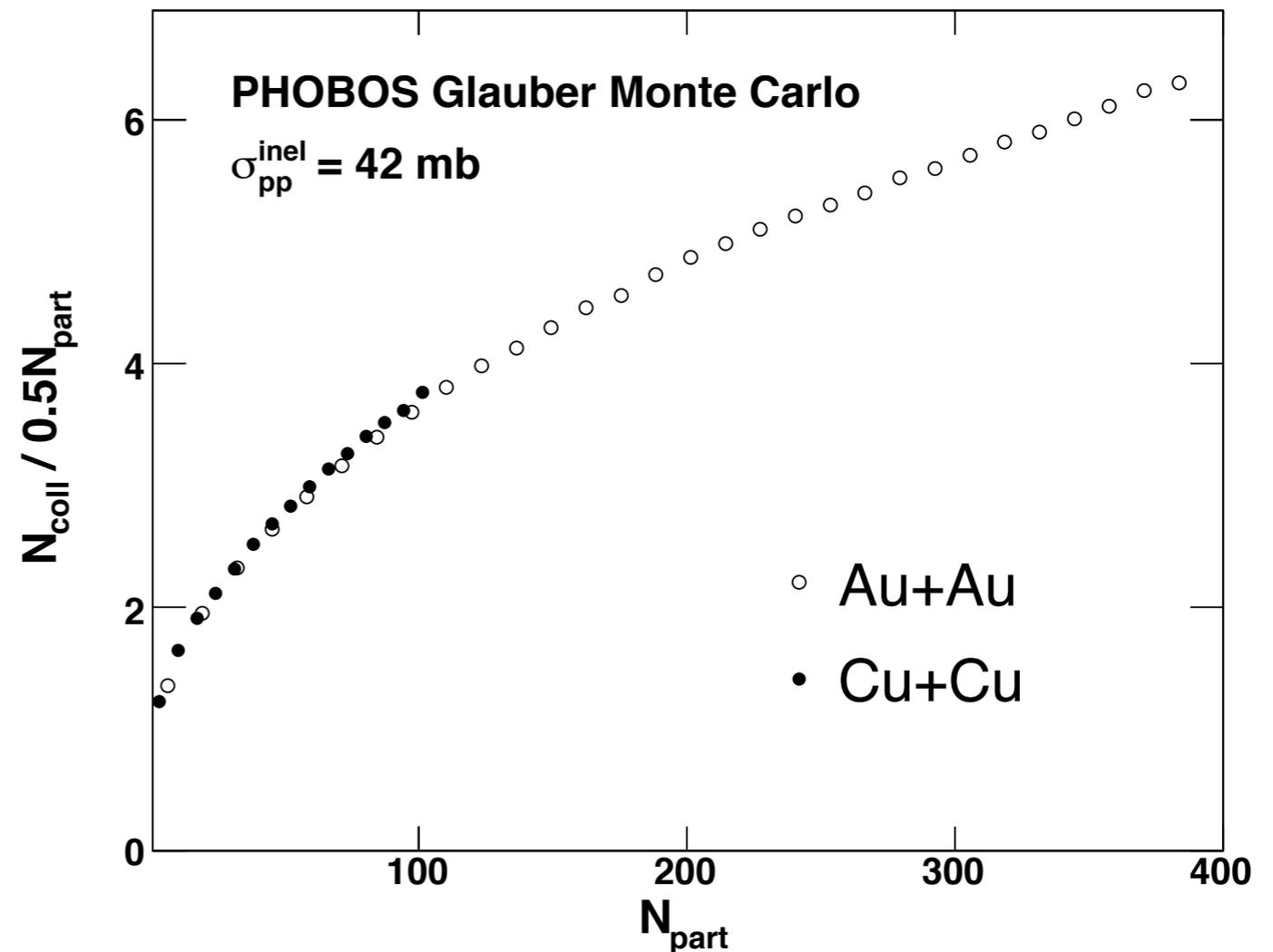
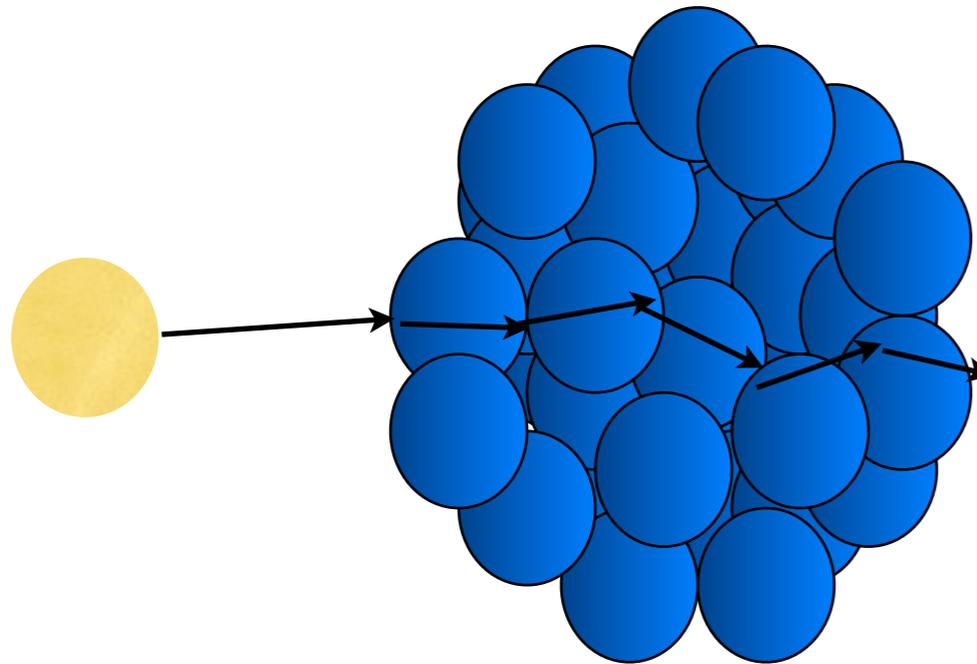


Participants vs. Collisions



In principle, rates of hard processes should be sensitive to the “microscopic” structure, N_{coll} , iff pQCD factorization holds true in $A+A$

Participants vs. Collisions



The number of binary collisions scales as $\sim N_{\text{part}}^{4/3}$.

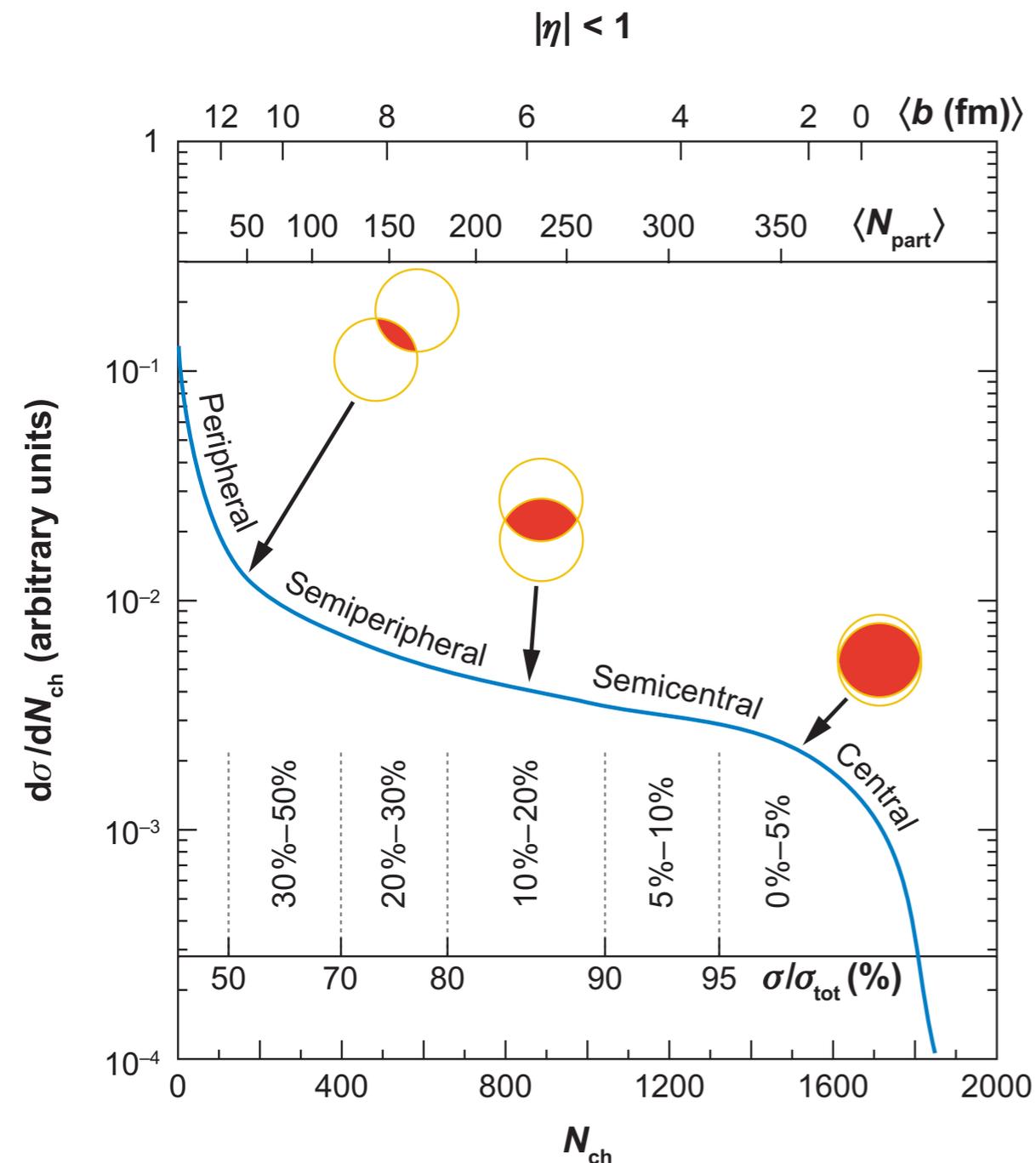
Nuclear “thickness” (v) scales like $N_{\text{coll}} / (N_{\text{part}} / 2)$

Extrapolating hard processes to A+A

	A+A luminosity	HI total cross section	# of binary subcollisions (Glauber!)	yield in proton-proton	
what you measure in Pb+Pb	$N_{A+A}^X = \mathcal{L} \sigma_{A+A} \langle N_{coll} \rangle N_{p+p}$				
	$= \mathcal{L} \sigma_{A+A} \langle N_{coll} \rangle \frac{\sigma_{p+p}^X}{\sigma_{p+p}^{tot}}$				ratio of subprocess to total
	$= \mathcal{L} \frac{\sigma_{A+A}}{\sigma_{p+p}^{tot}} \langle N_{coll} \rangle \sigma_{p+p}^X$				results from p+p reference
	$= \mathcal{L} (nb^{-1}) A^2 \sigma_{p+p}^X (nb)$				pocket formula!

10^4 increase in “cross section” (but 10^7 decrease in \mathcal{L})

Estimating Centrality



Any quantity which varies monotonically with impact parameter can be used to get a data-driven estimate of

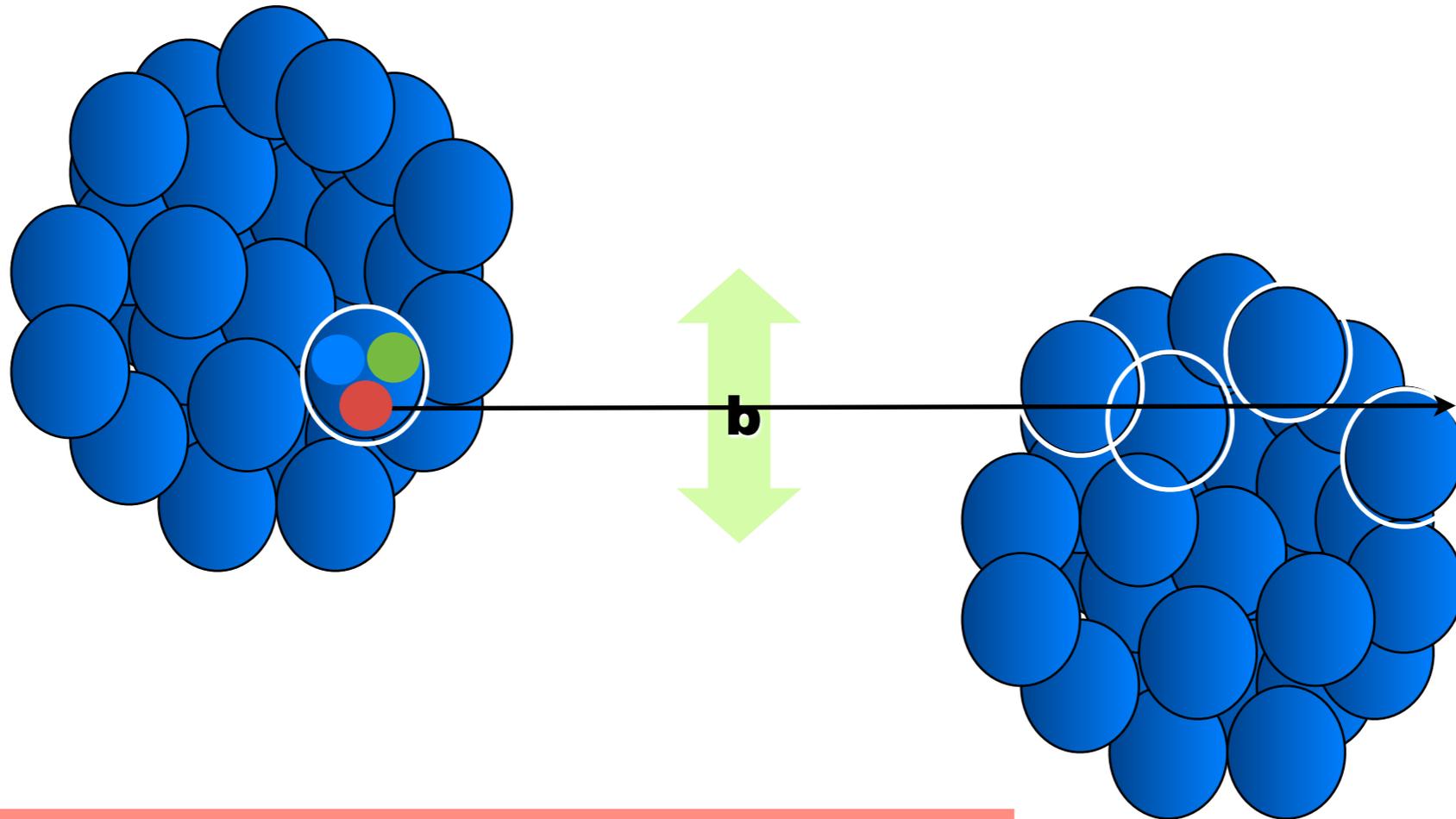
- Number of participants (N_{part})
- Number of binary collisions (N_{coll})
- Impact parameter (b)

by calculating a “percentile bin” in the variable, and the similar percentile in the geometric quantity.

For this, we use the “Glauber model”

Nuclear Modification Factor

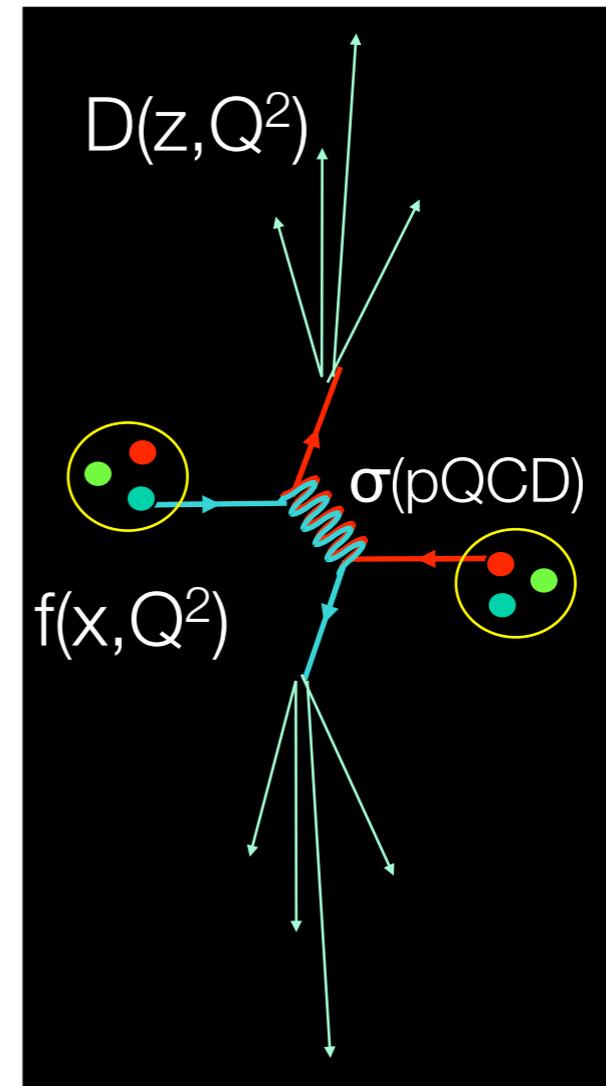
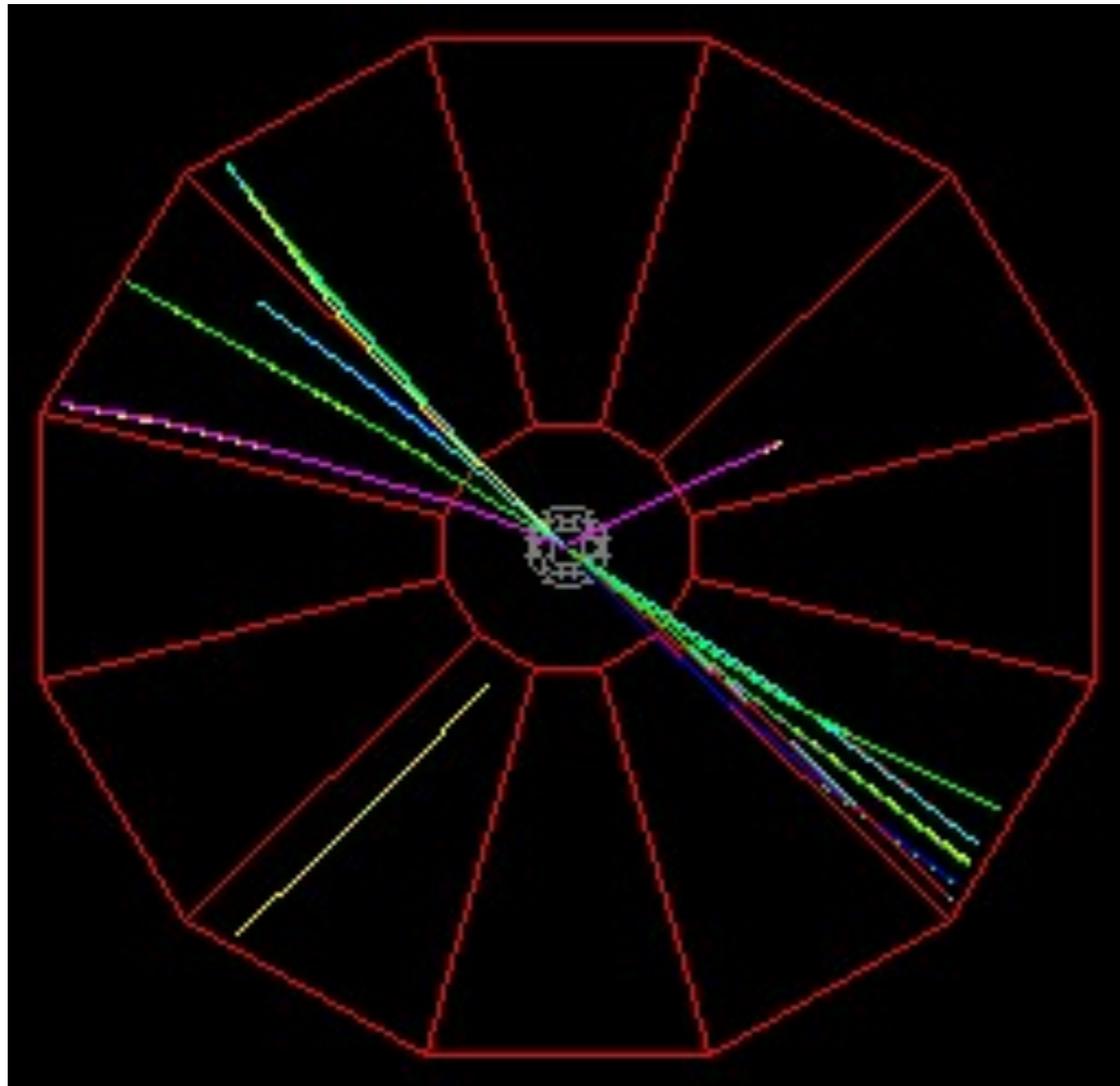
CTEQ



$$R_{AB} = \frac{1}{N_{AB}^{coll}} \frac{dN_{AB}}{dp_T}$$

Yield per collision
relative to p+p
< 1 implies
“nuclear effect”
(or no factorization)

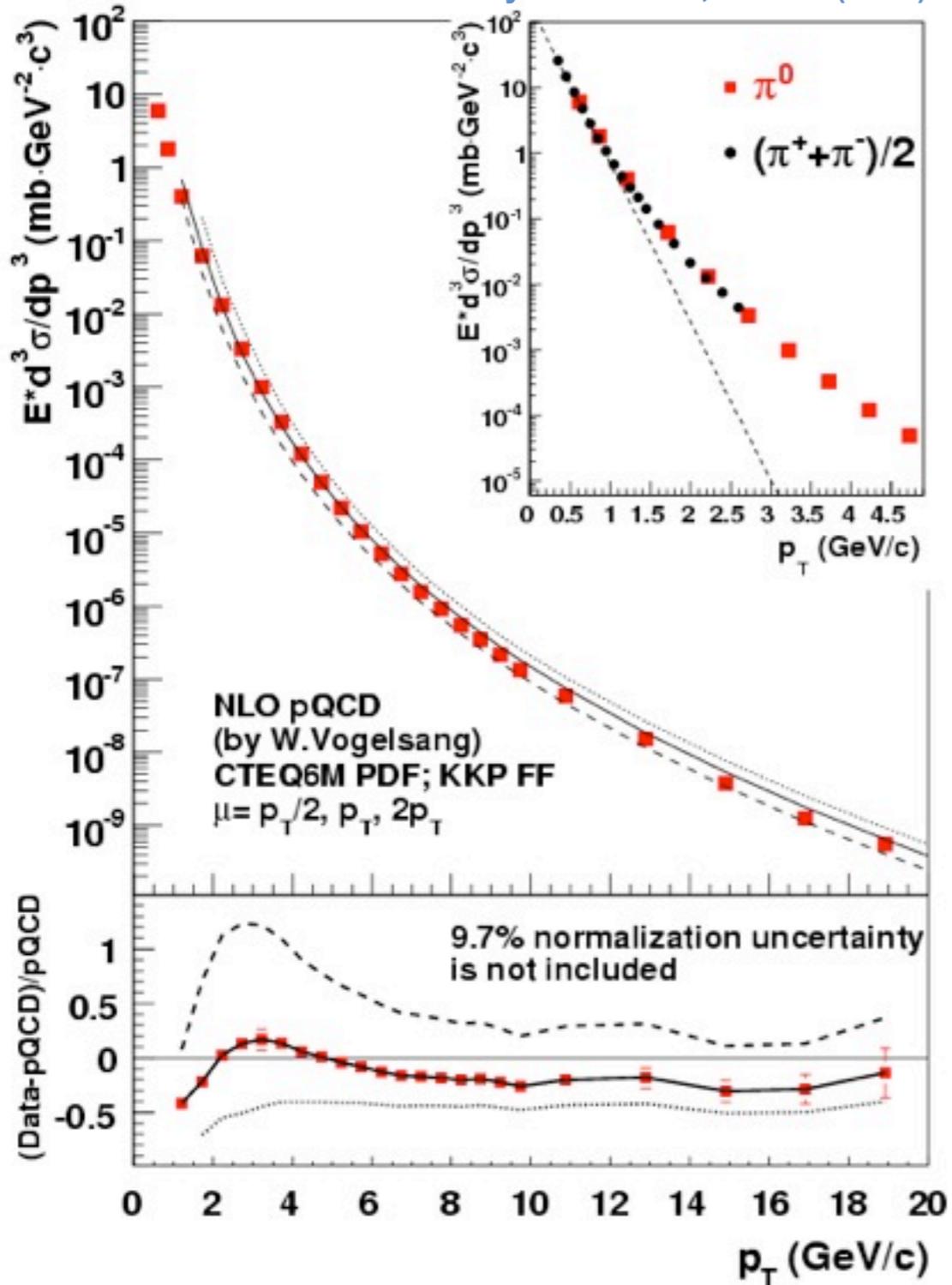
Jets in p+p



Jets are a “hard” phenomenon, allow pQCD description

pQCD @ RHIC

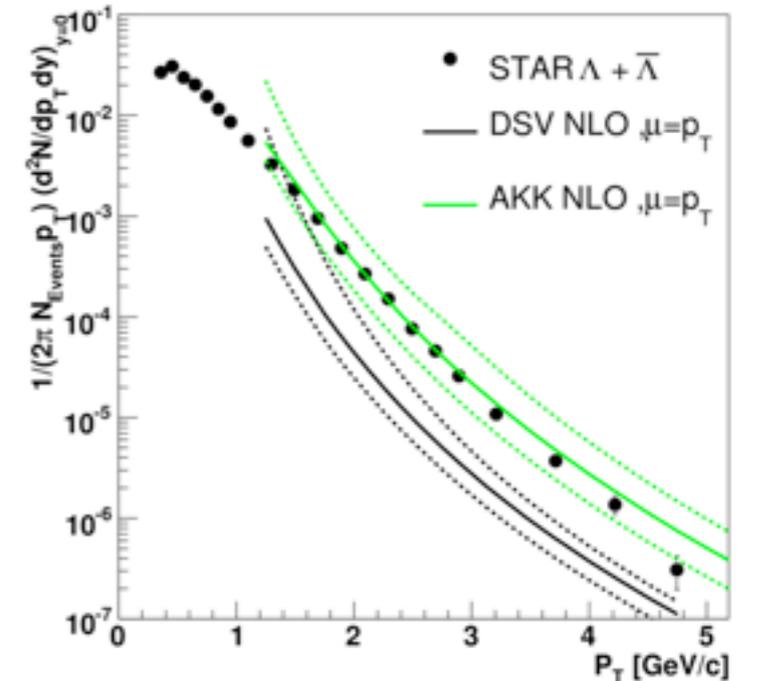
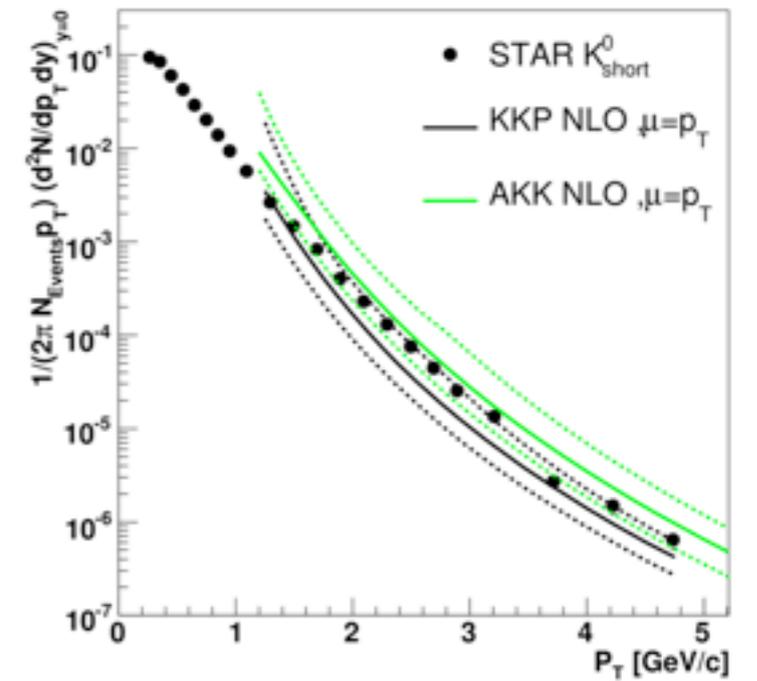
Phys. Rev. D 76, 092002 (2007)



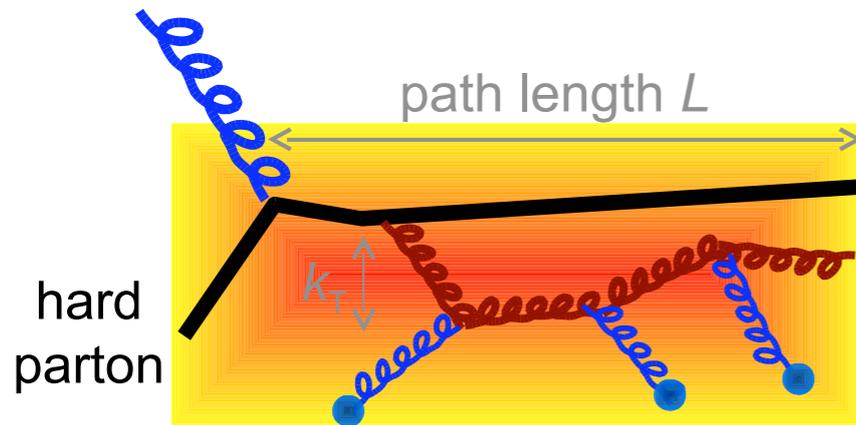
← Great for pions

Less so for K & Λ →

Phys. Rev. C 75, 064901 (2007)

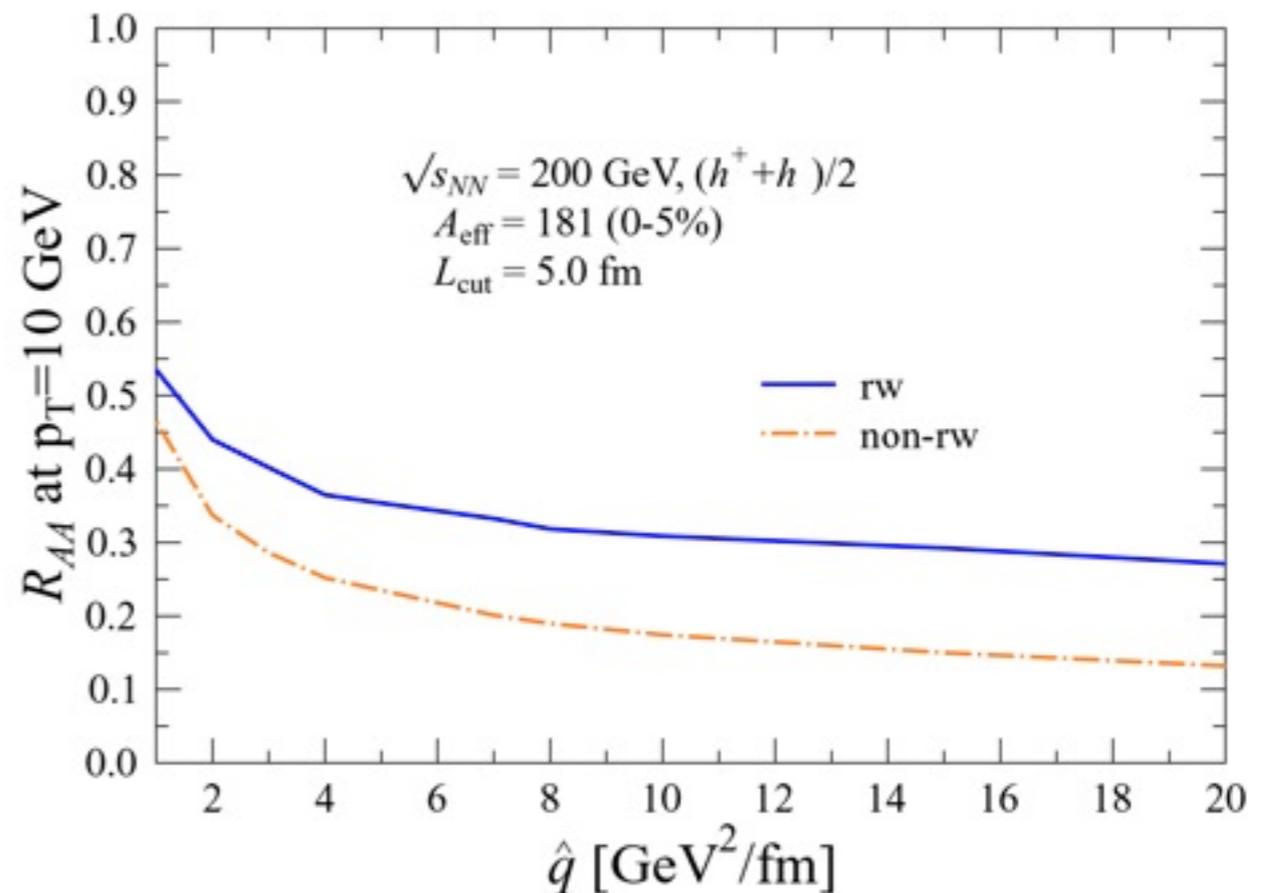
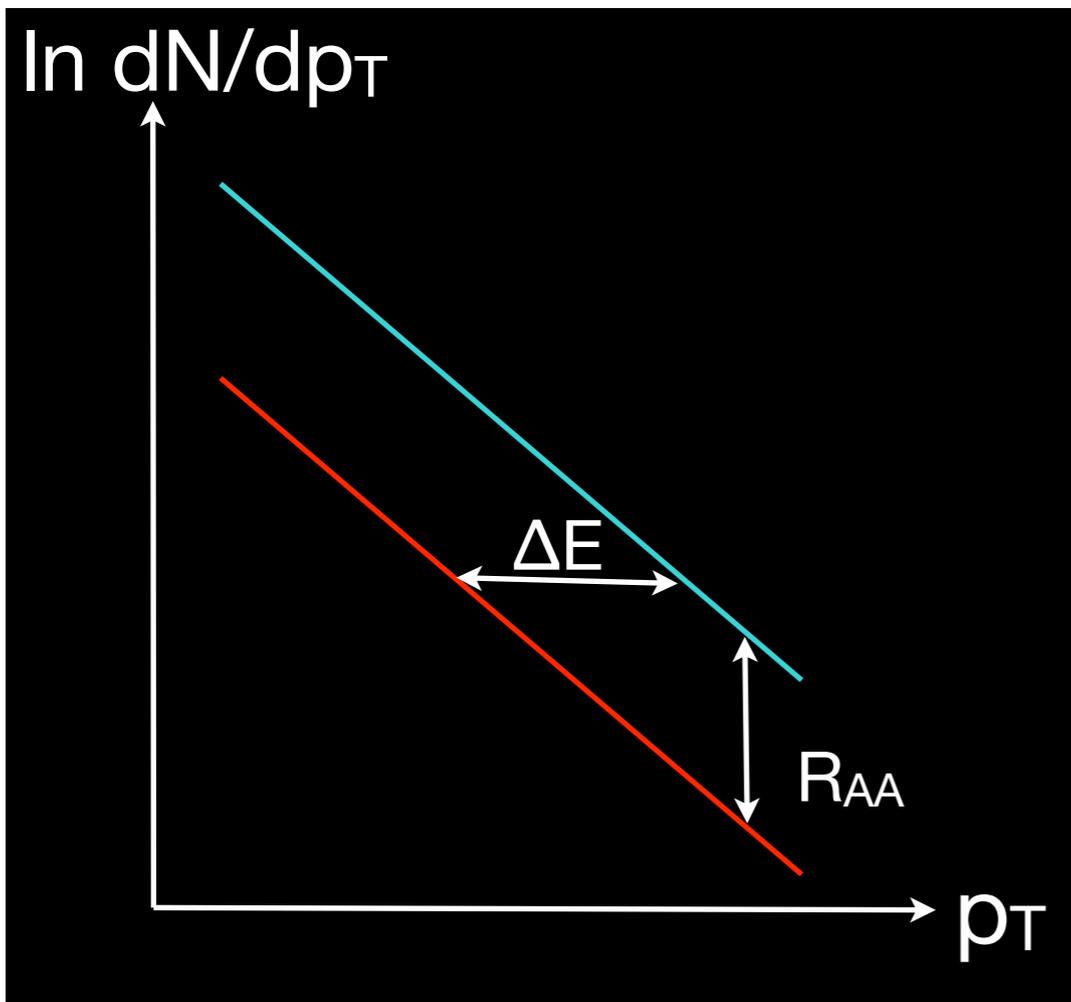


“Jet Quenching”



$$\langle \Delta E \rangle \approx \int_0^{\omega_c} d\omega \omega \frac{dI}{d\omega} \propto \alpha_s C_R \omega_c \propto \alpha_s C_R \hat{q} L^2$$

$\hat{q} = \frac{\langle k_T^2 \rangle}{\lambda}$
 Transport coefficient (GeV²/fm)
 sensitive to gluon density
 (not deconfinement!)



Modified Fragmentation

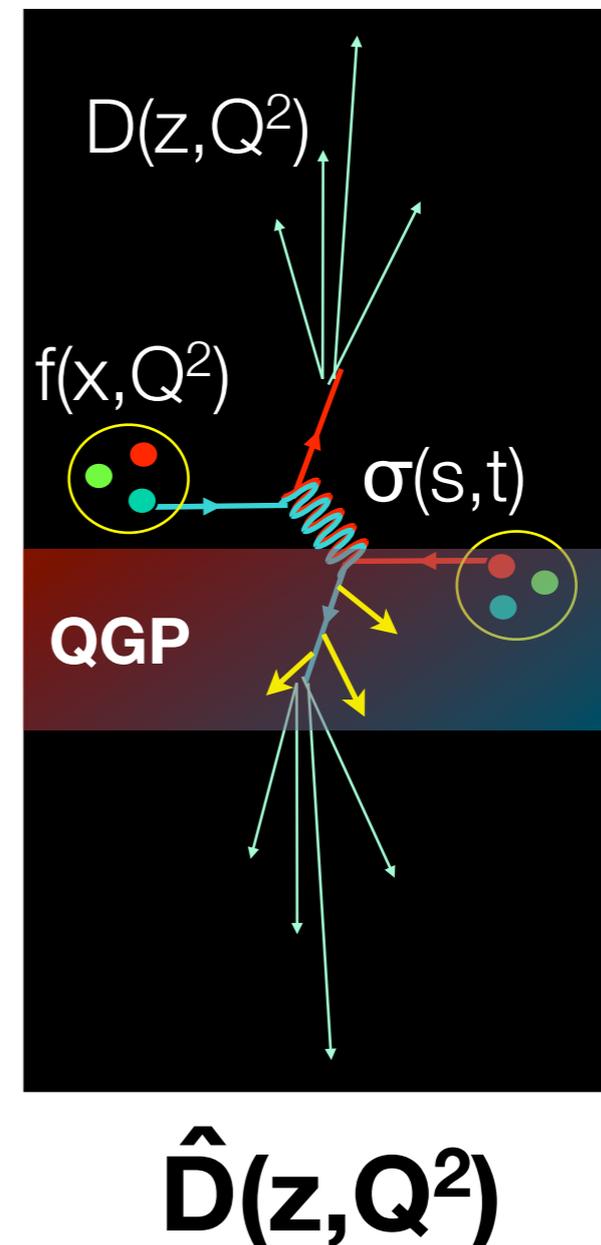
- **The typical scenario**

- Hard parton produced in medium according to pQCD cross sections
- Energy loss: generation of additional radiation, elastic scattering
- Vacuum fragmentation of attenuated parton

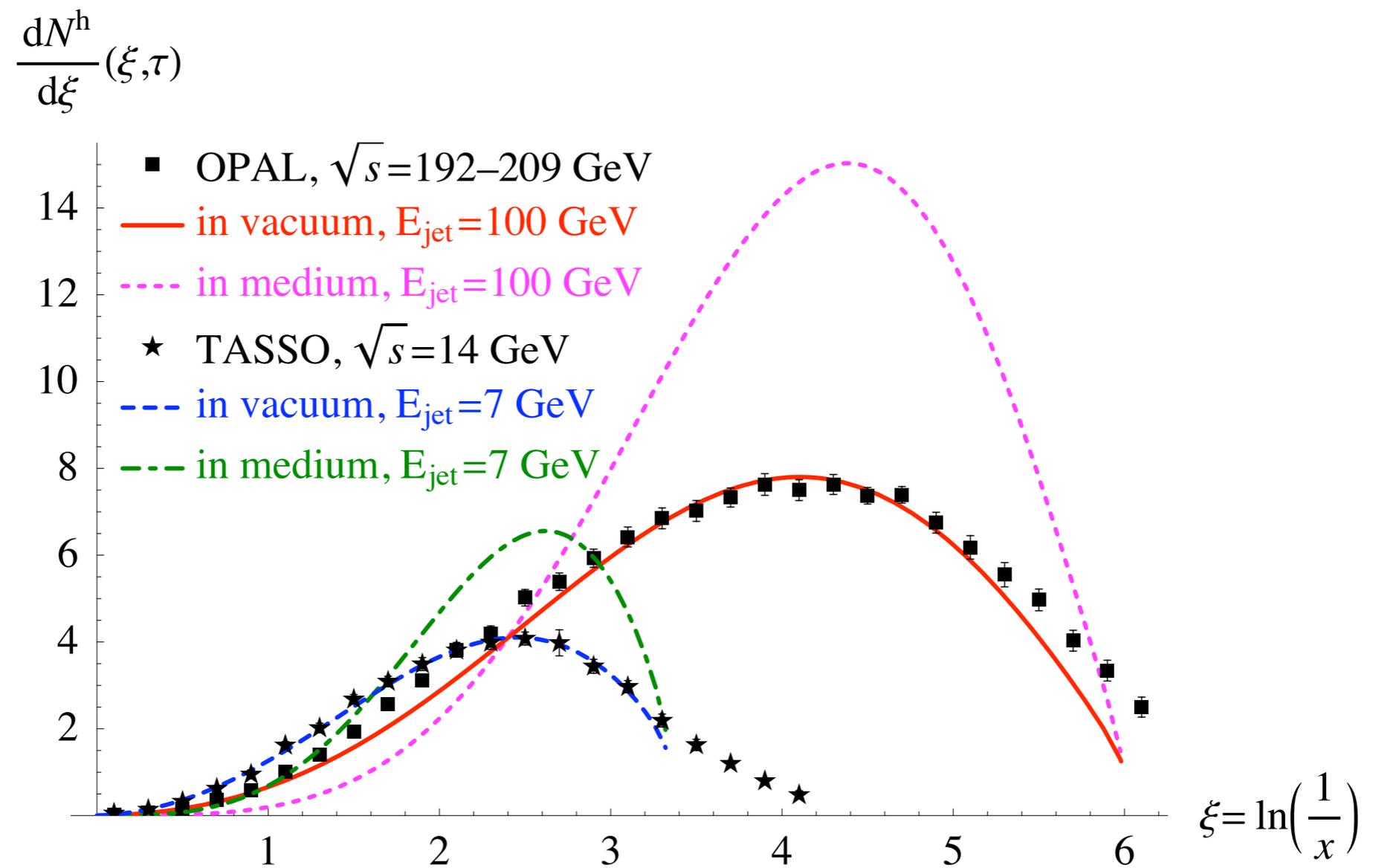
- **“Modified fragmentation”**

- **Variety of implementations**

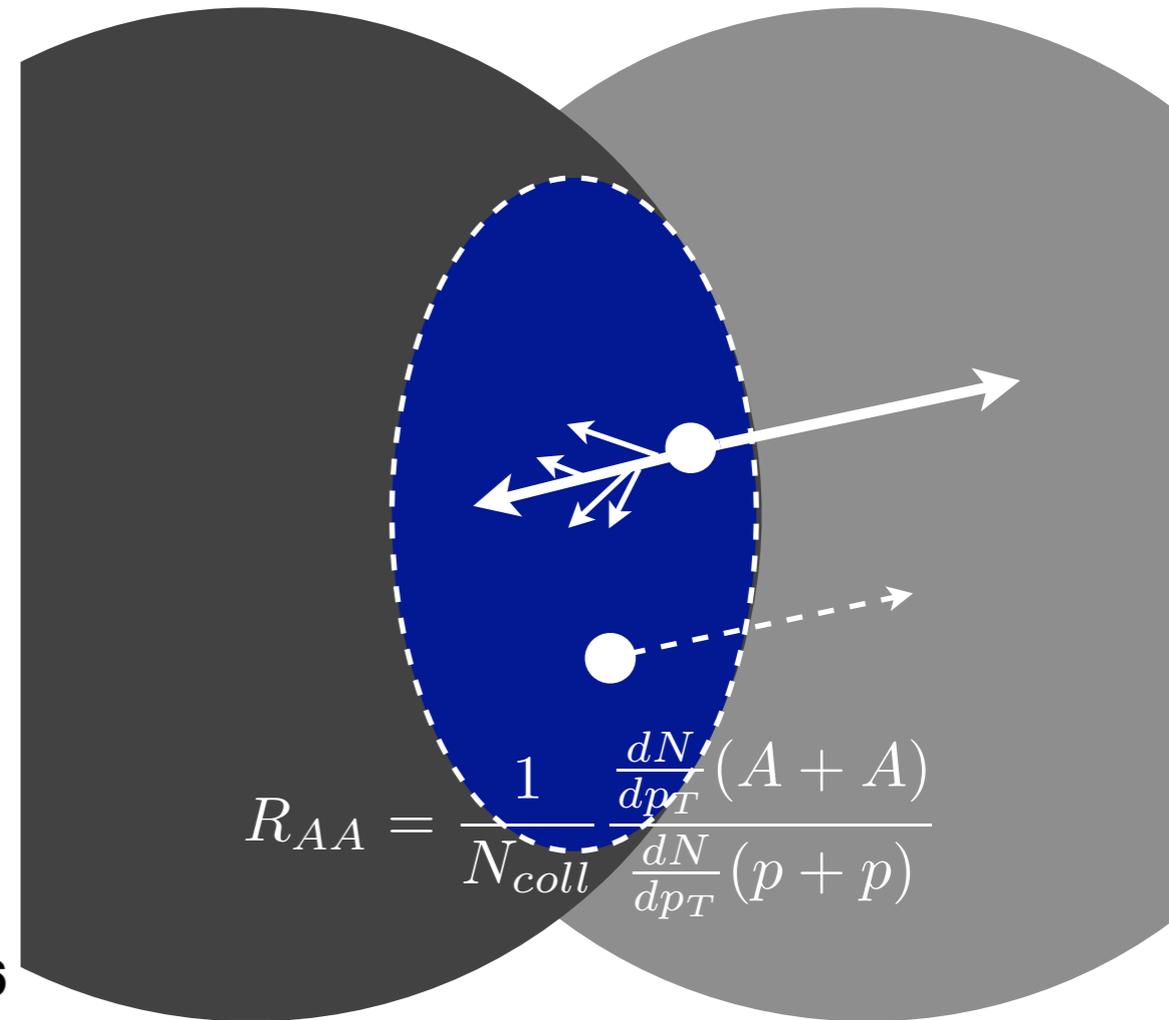
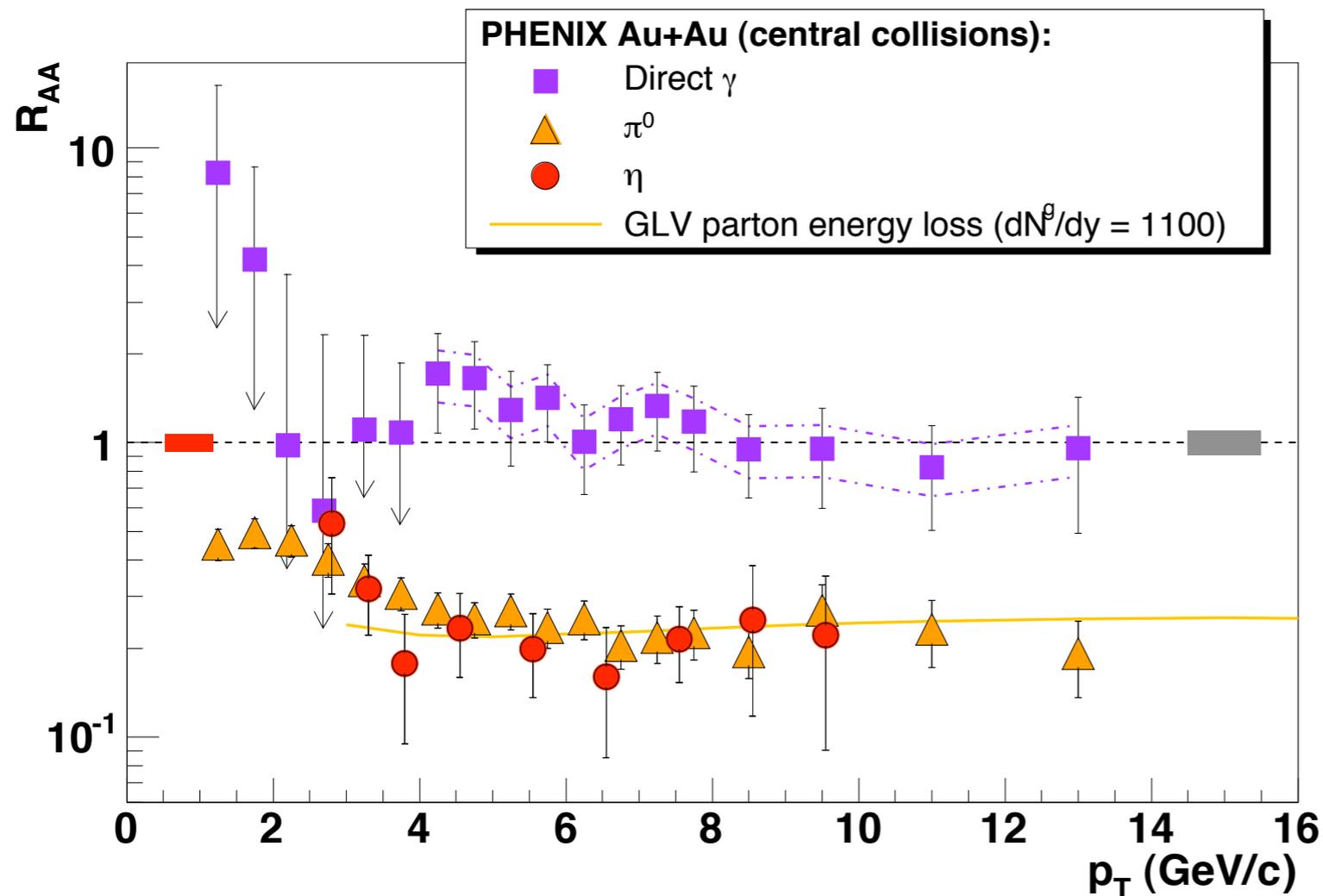
- Many soft gluons (ASW/BDMPS)
- Few hard gluon (DGLV)
- Higher twist (HT)
- Thermal field theory (AMY)
- Even AdS/CFT (LRW)



Modified Fragmentation

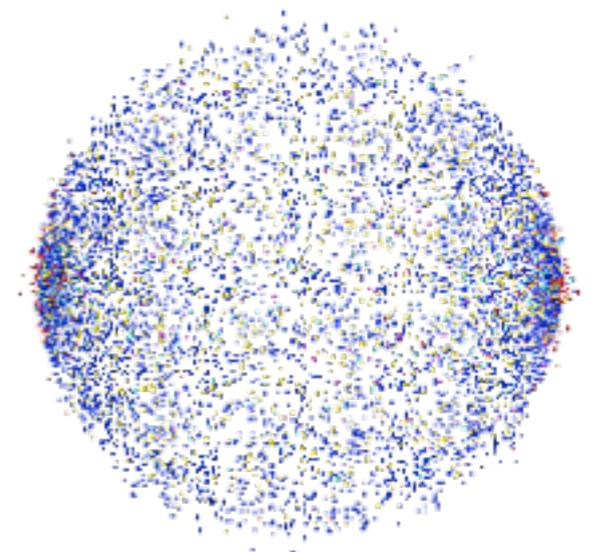
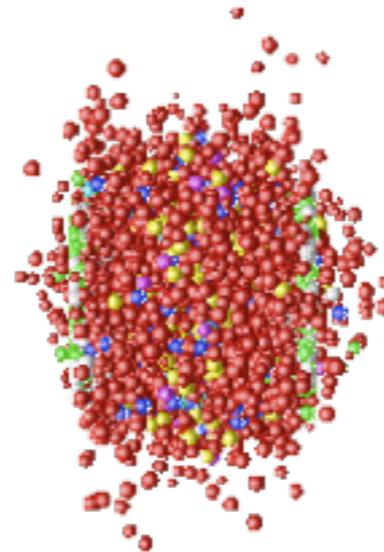
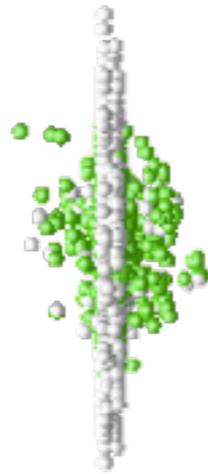
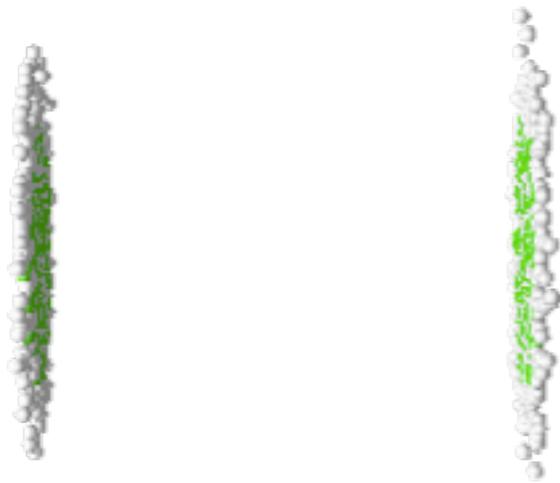


High p_T Suppression



High p_T particles are strongly suppressed relative to $p+p$ spectrum \times binary collisions (N_{coll}). Photons not. pQCD energy loss calculations describe light hadrons. Photons appear to be unaffected by medium.

When does suppression occur?



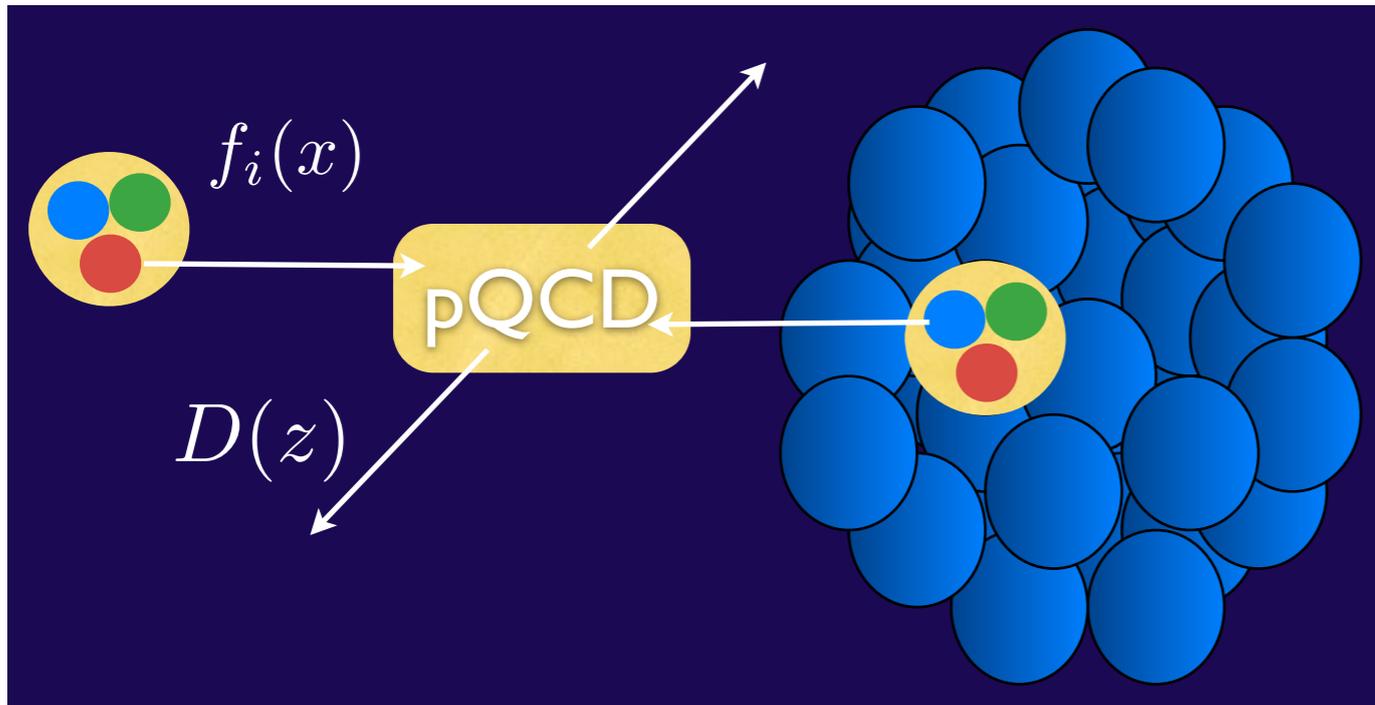
Some sort of strong “shadowing” phenomenon in the initial nuclear parton distributions...

...so there are simply fewer hard scatterings in the reaction?

Or does something occur in the strongly interacting stage??

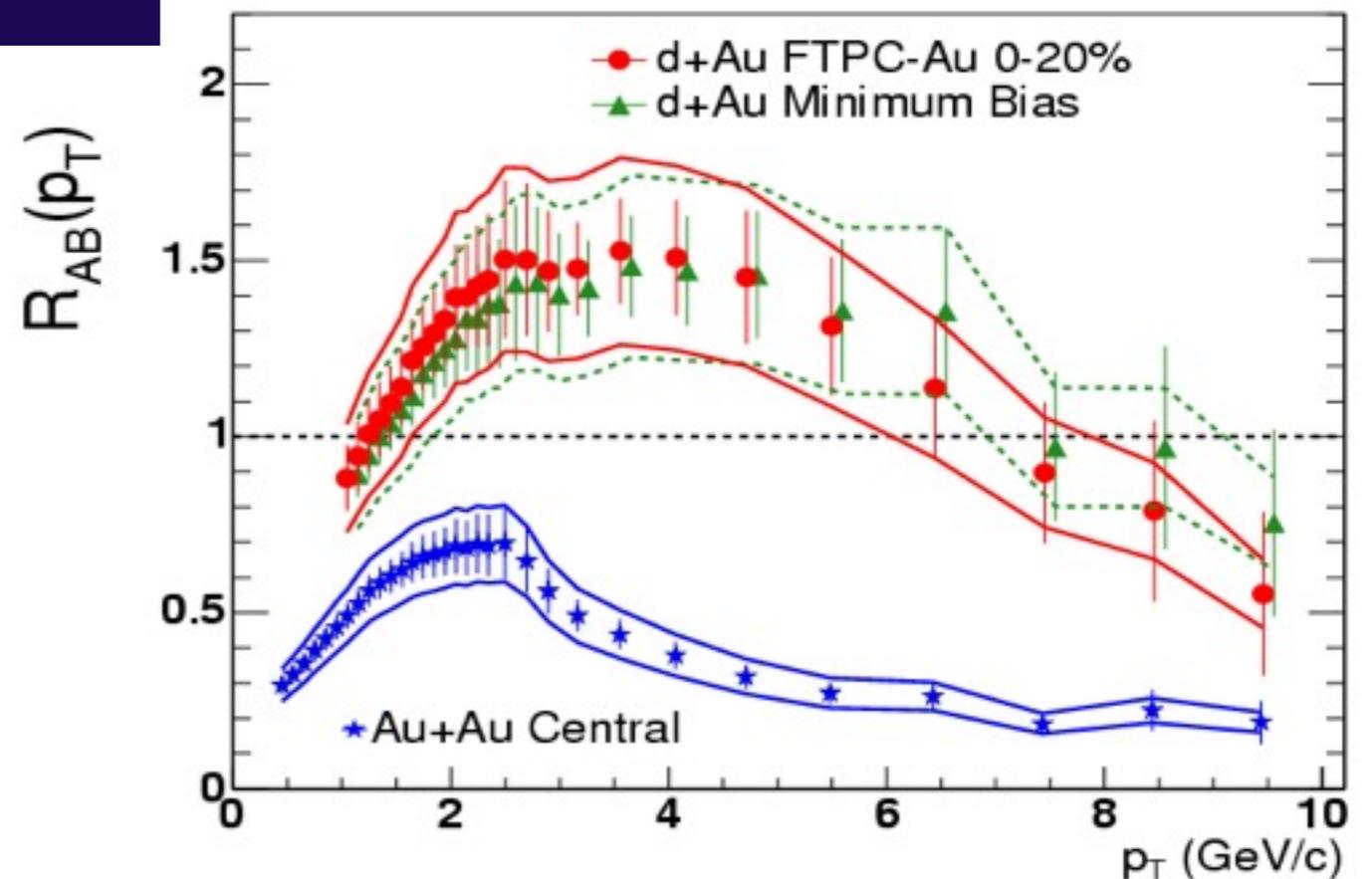
Collisions of the outgoing particles with the background of soft hadrons?

Initial vs. Final State

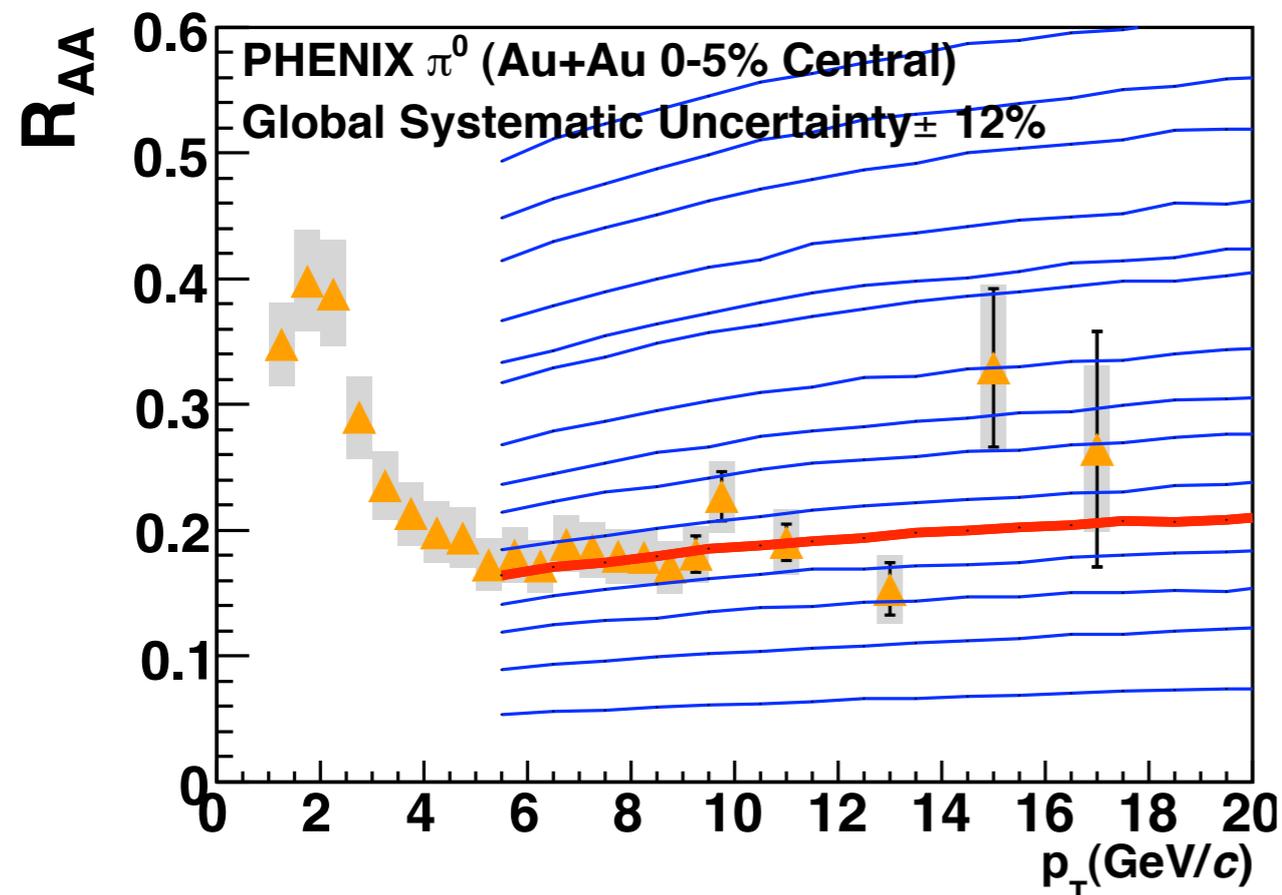


p+A collisions provide some access to the shadowing from the nuclear wave-function (reduction in initial flux)

d+Au not suppressed (except at high p_T ...) while Au+Au shows large effect at all p_T : evidence for final state



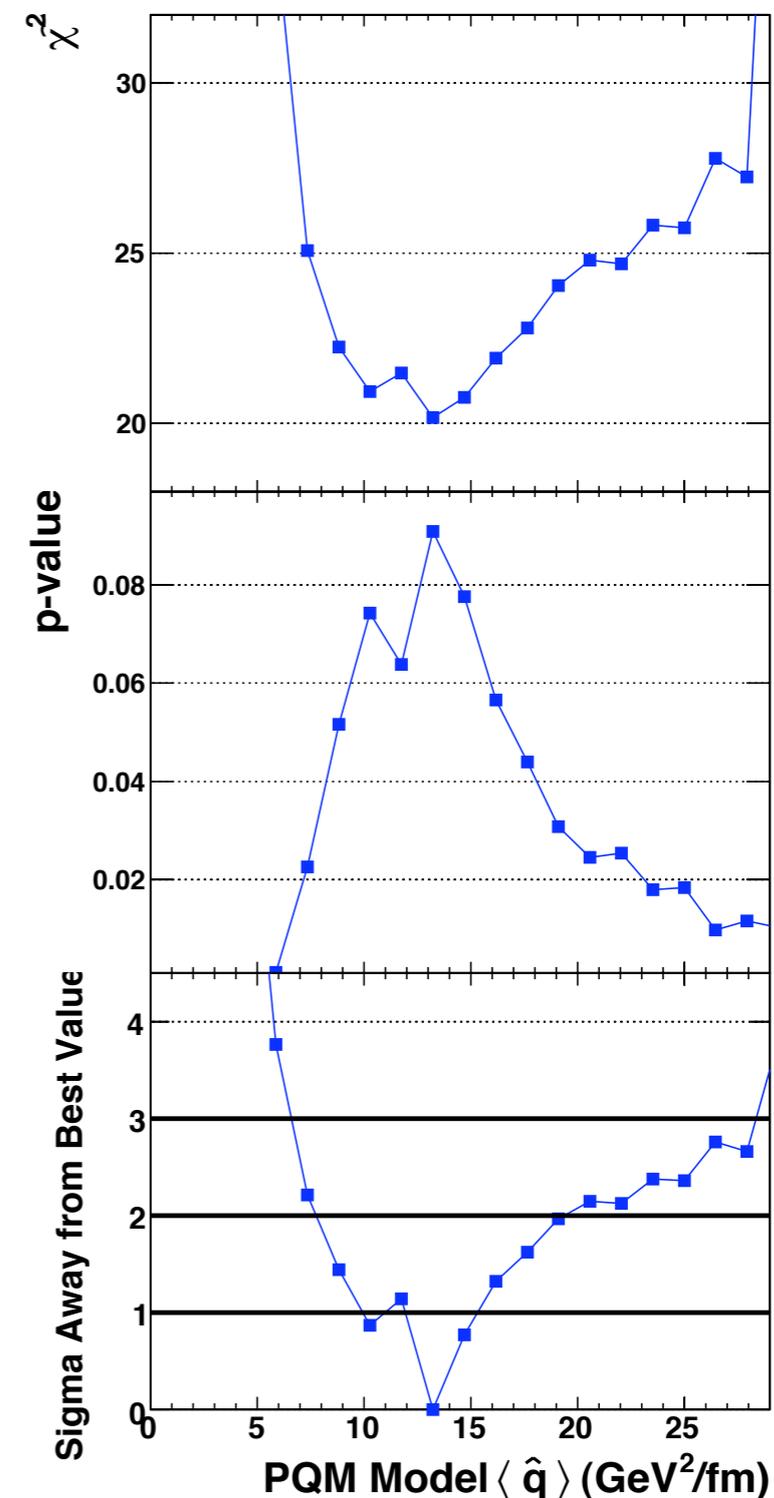
Estimating Stopping Power (PHENIX)



$$\Delta E \propto \hat{q} \propto \langle p_T^2 \rangle / \lambda$$

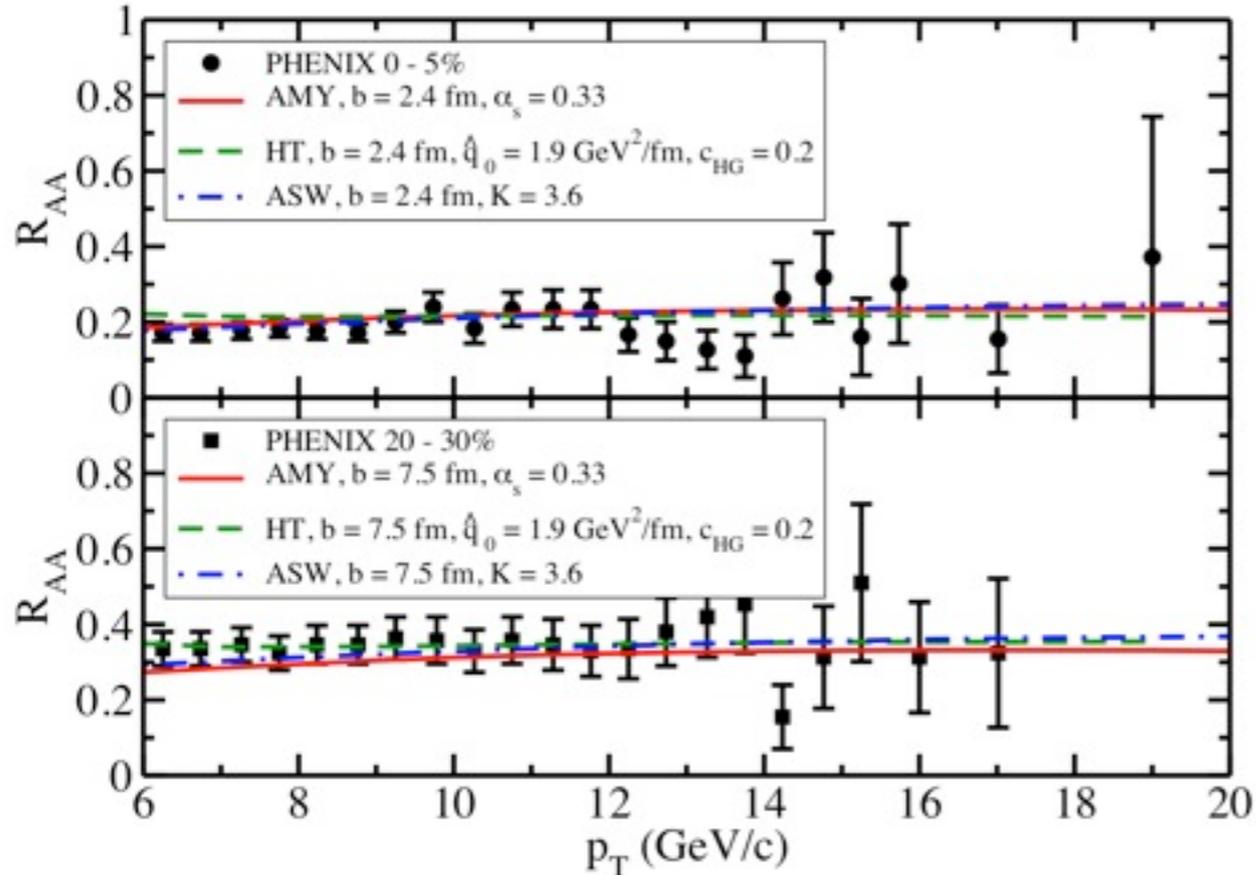
PHENIX χ^2 fits to PQM
indicate $6 < \hat{q} < 24 \text{ GeV}^2/\text{fm}$.

First “real” quantitative constraints!



Comparison of approaches - singles

Bass, et al, arXiv:0808.0908v3

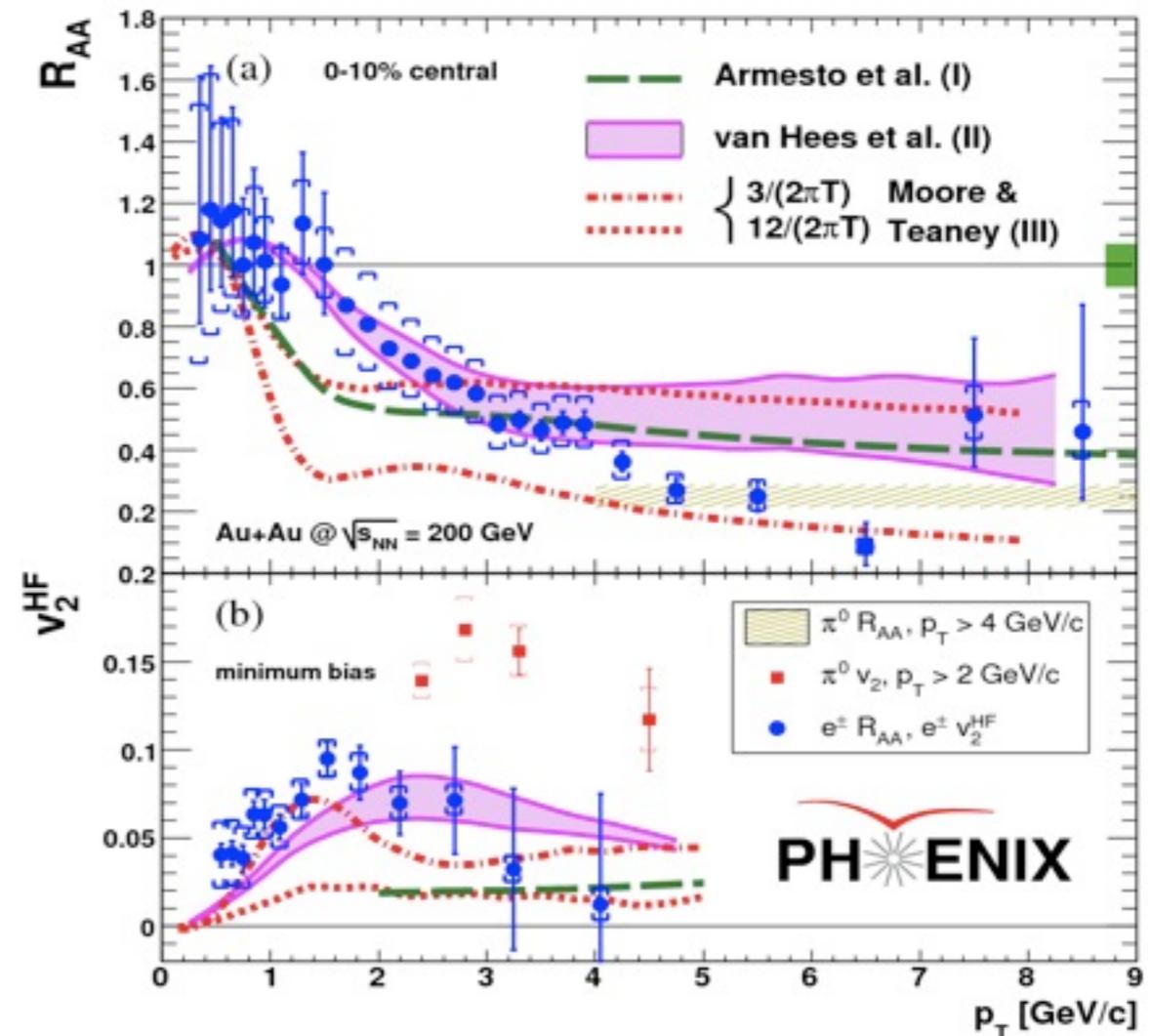
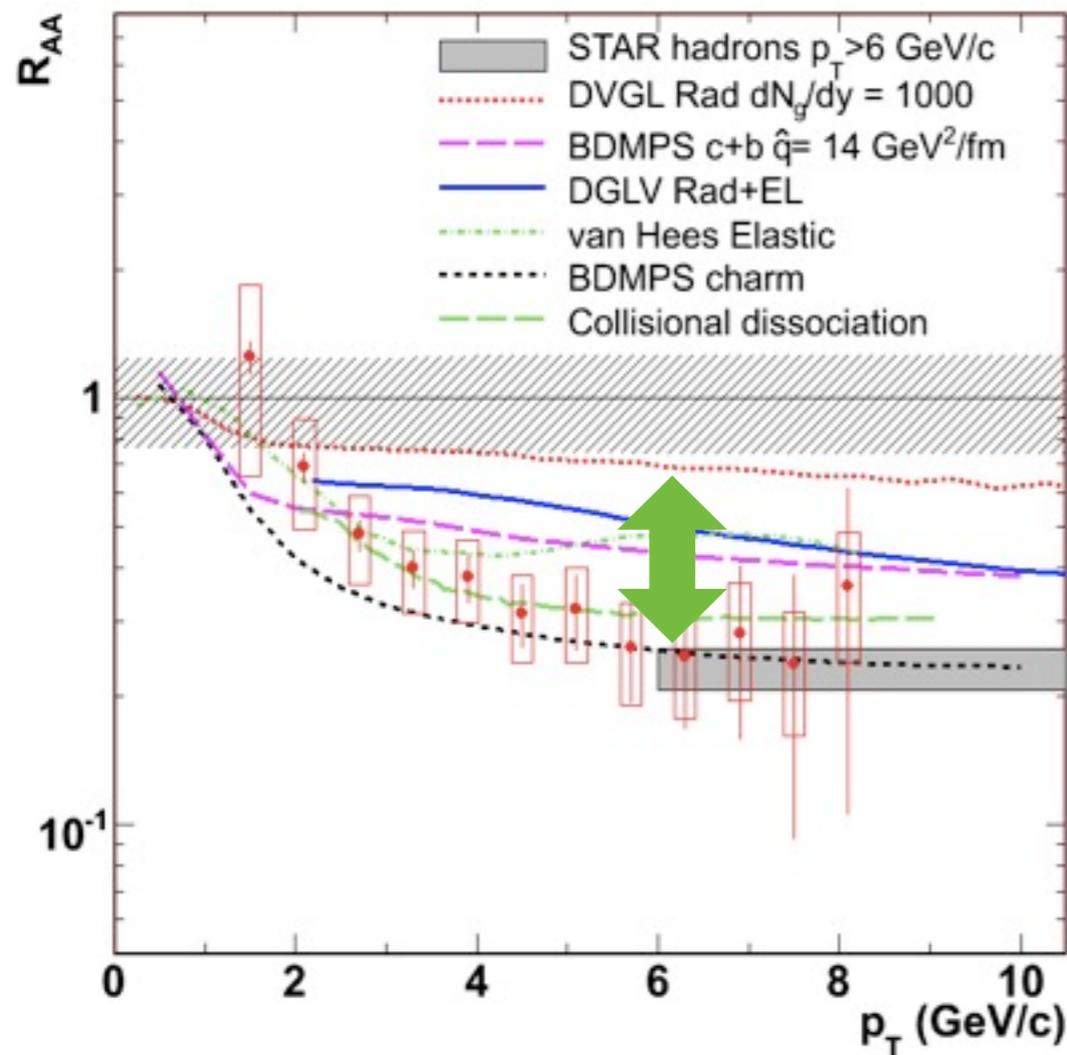


$\hat{q}(\vec{r}, \tau)$ scales as	ASW	HT	AMY
	\hat{q}_0	\hat{q}_0	\hat{q}_0
$T(\vec{r}, \tau)$	10 GeV ² /fm	2.3 GeV ² /fm	4.1 GeV ² /fm
$\epsilon^{3/4}(\vec{r}, \tau)$	18.5 GeV ² /fm	4.5 GeV ² /fm	
$s(\vec{r}, \tau)$		4.3 GeV ² /fm	

Very different values
for transport coefficient...

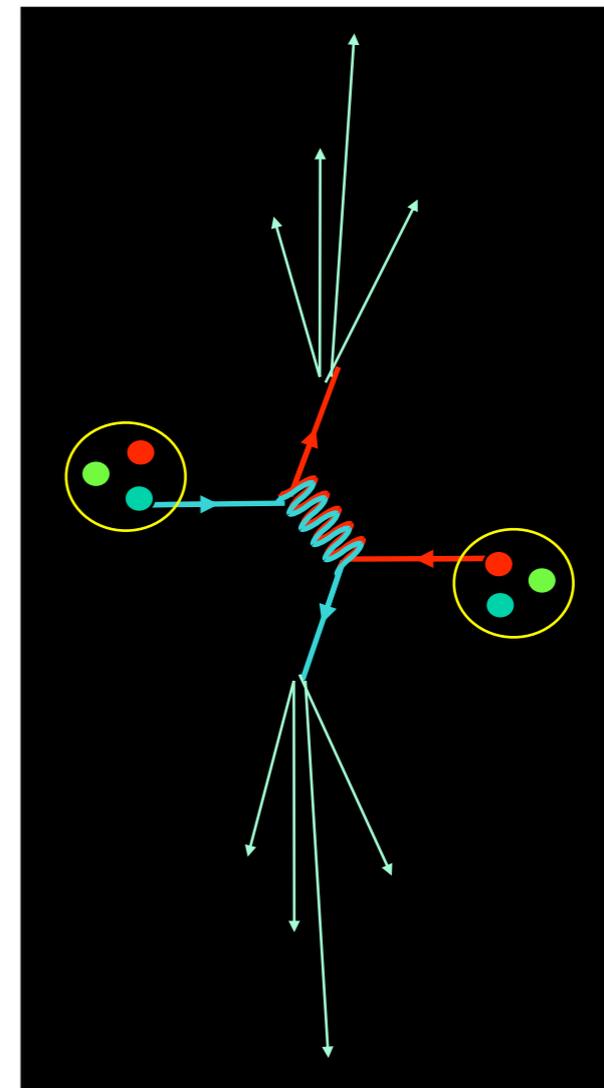
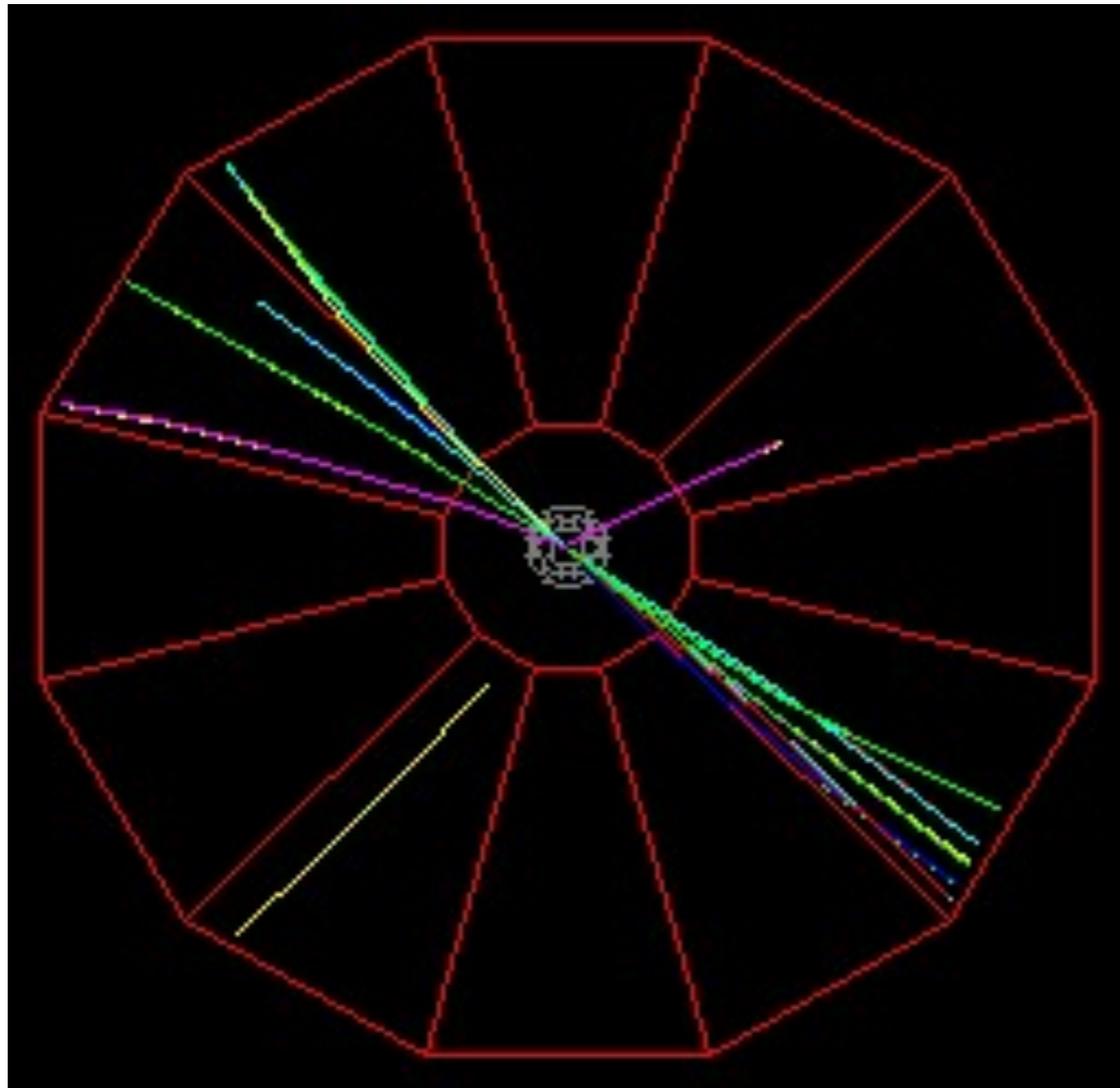
Variety of jet quenching implementations all give the “same” result for high p_T suppression - dominated by geometry

Heavy Flavor: a fly in the ointment



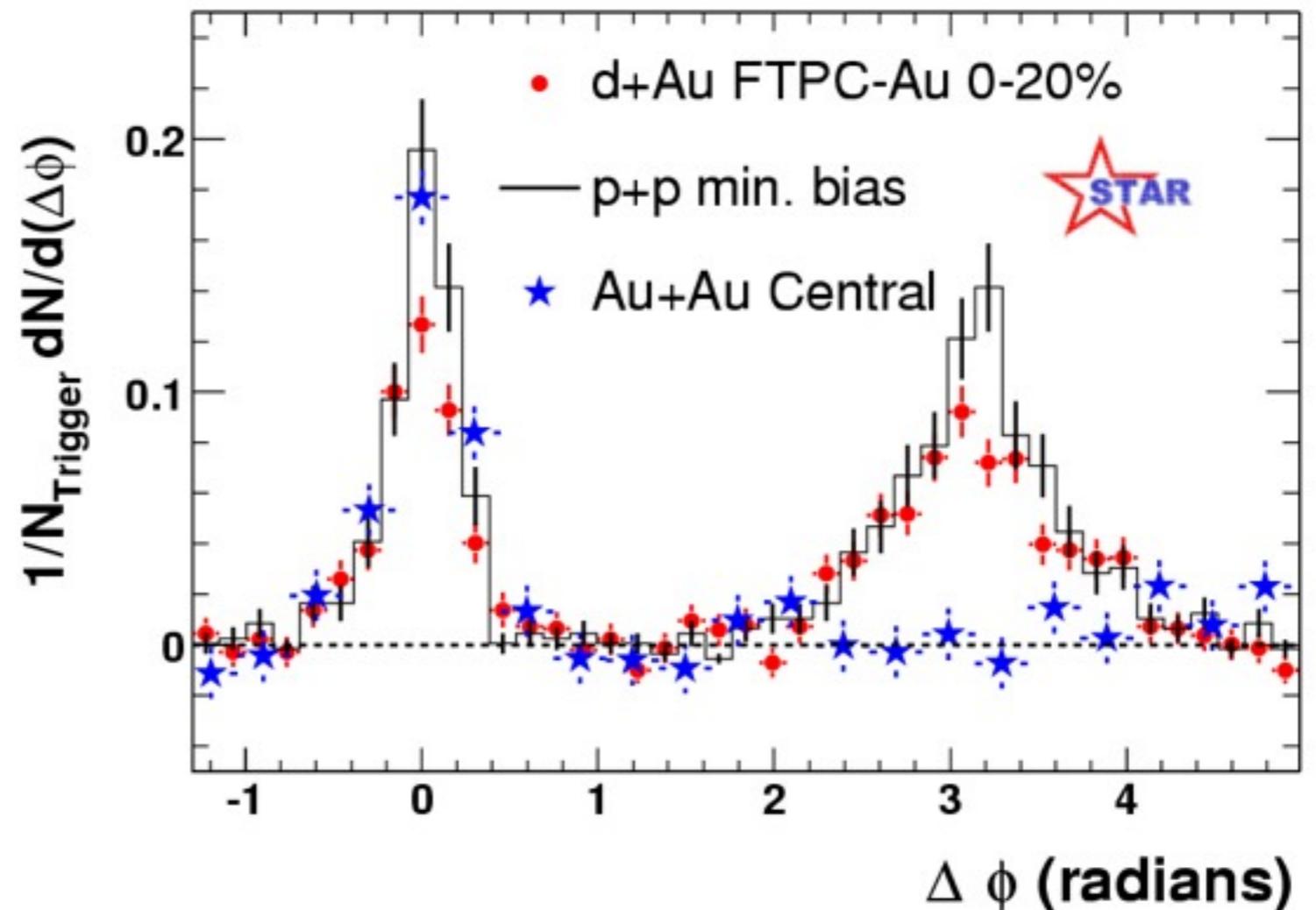
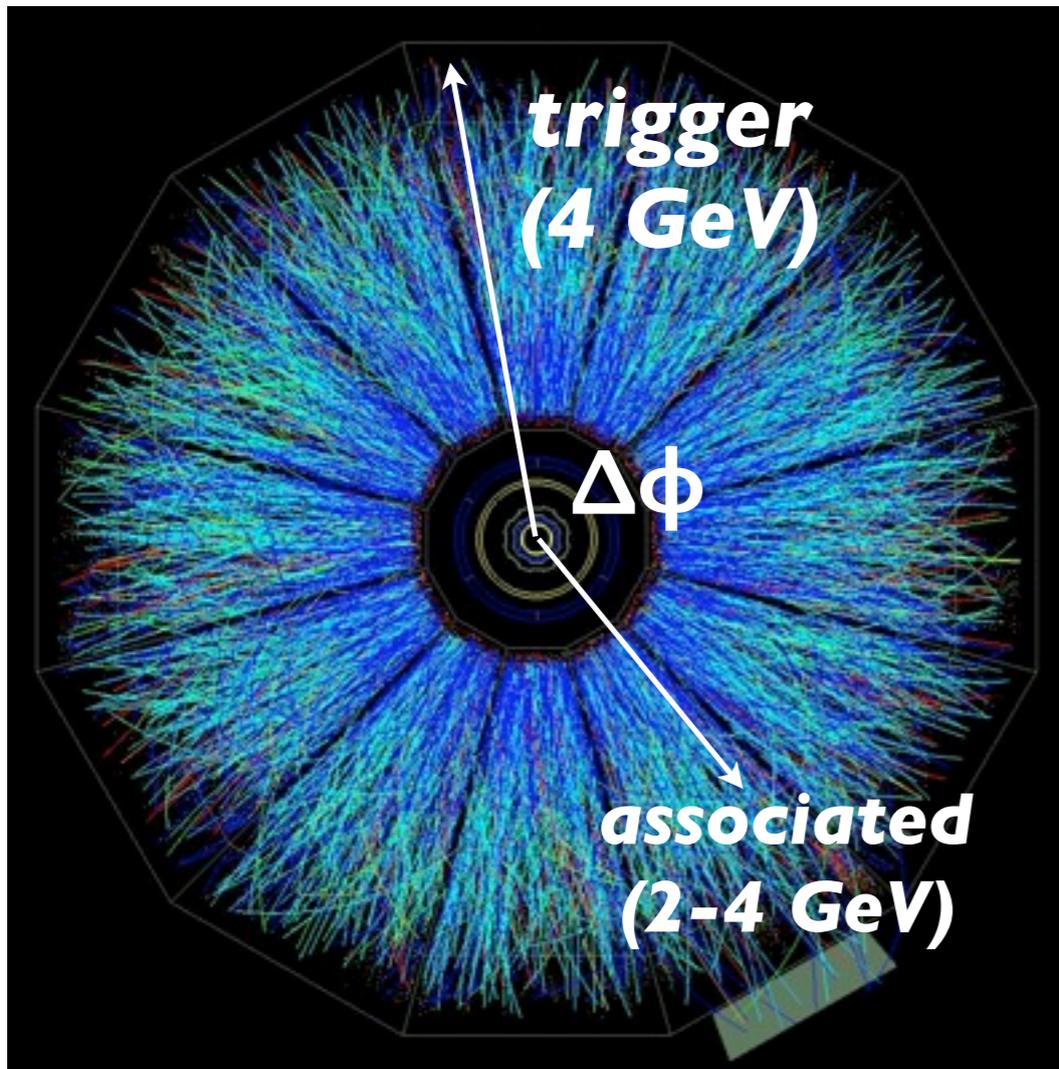
Heavy flavor (c,b) measured via semi-leptonic mode.
 Suppression at high p_T is very similar to other hadrons:
 surprising since mass should reduce radiation (“dead-cone”)

Correlations (jet-medium interactions)

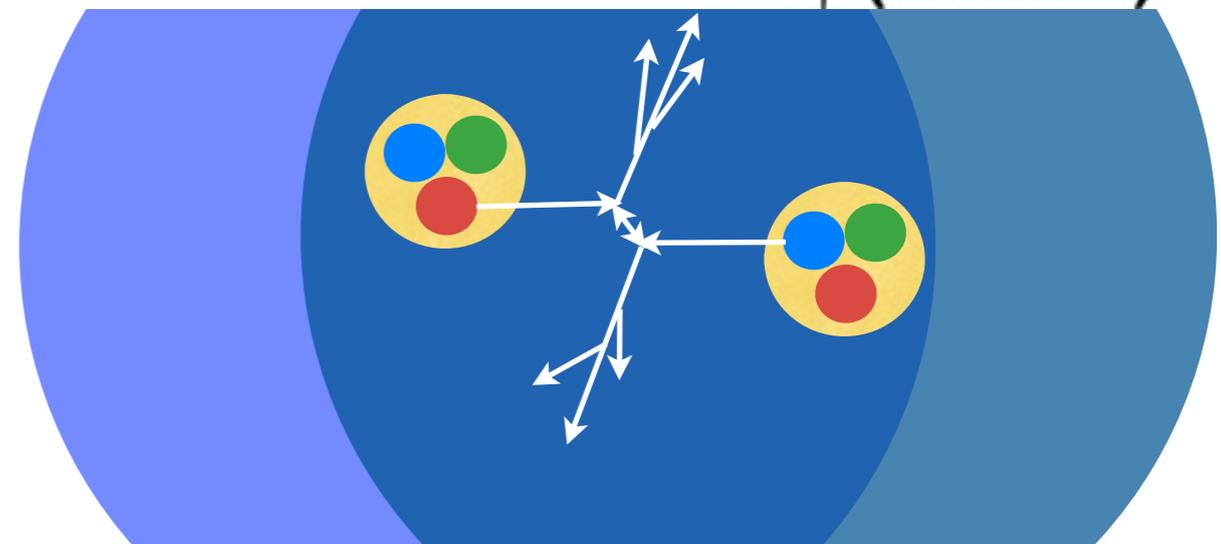


Jets are multi-particle phenomena:
2+ high p_T particles (quark and/or photon)
quarks fragment into multiple hadrons

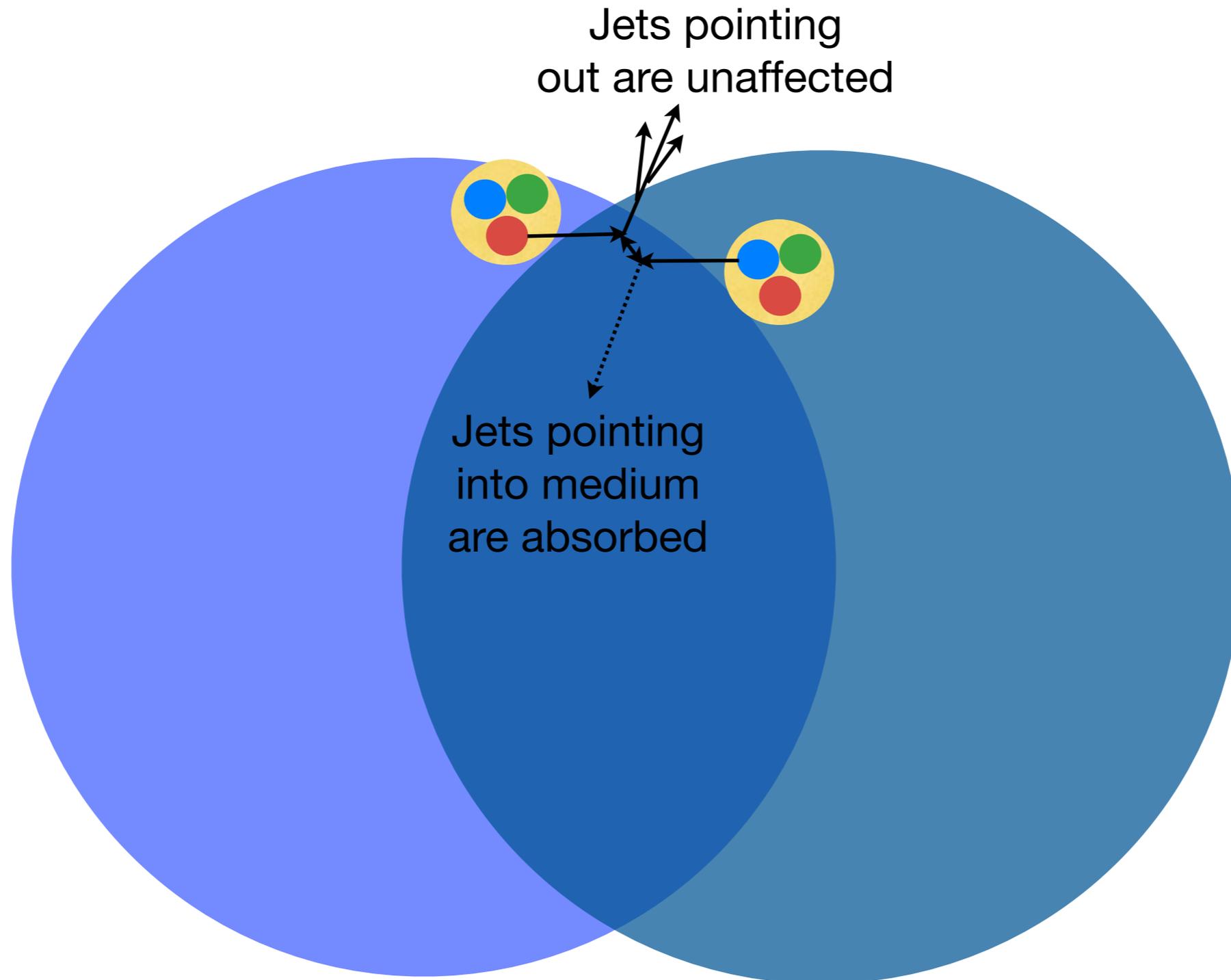
“Back-to-Back” Disappearance (b2bd)



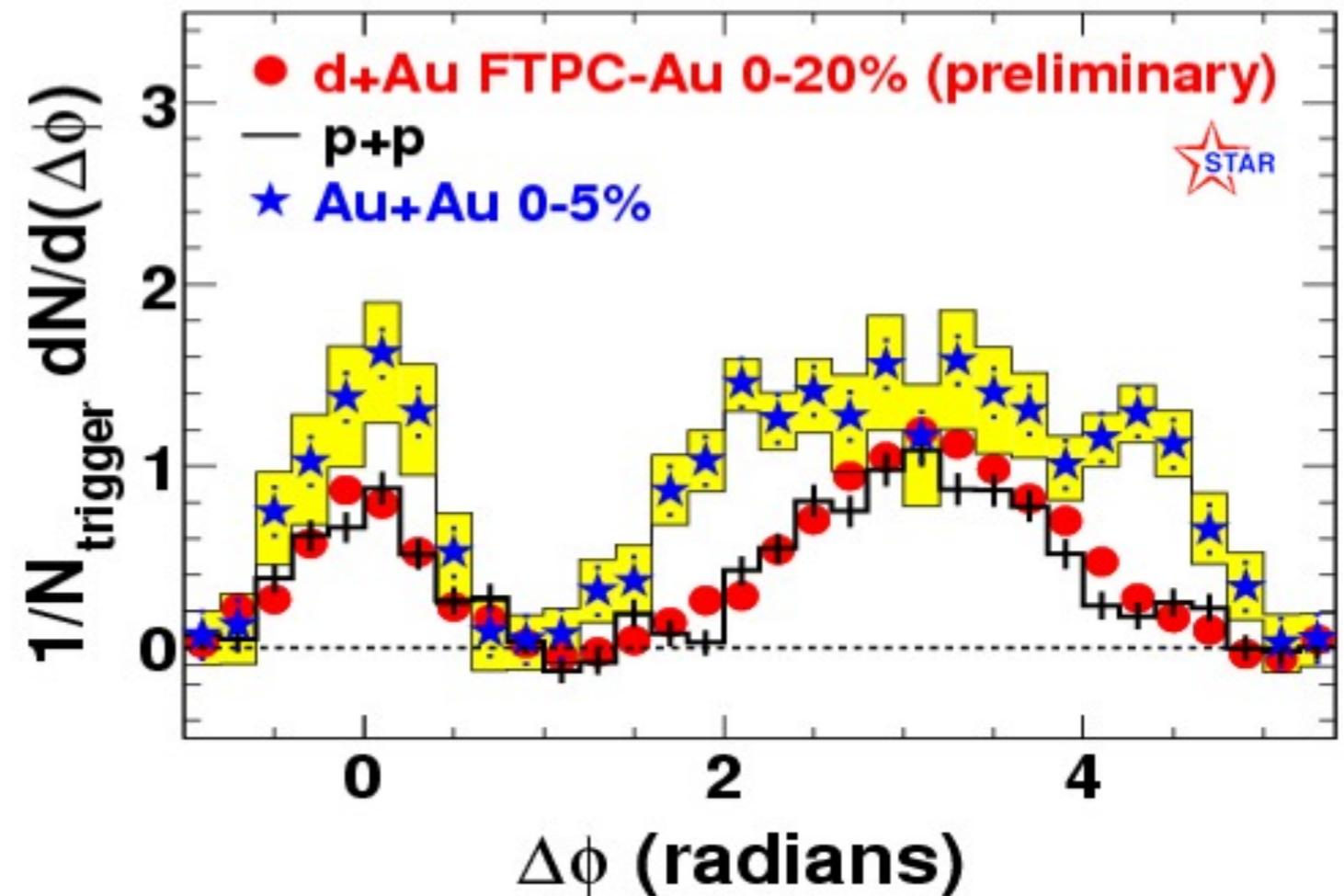
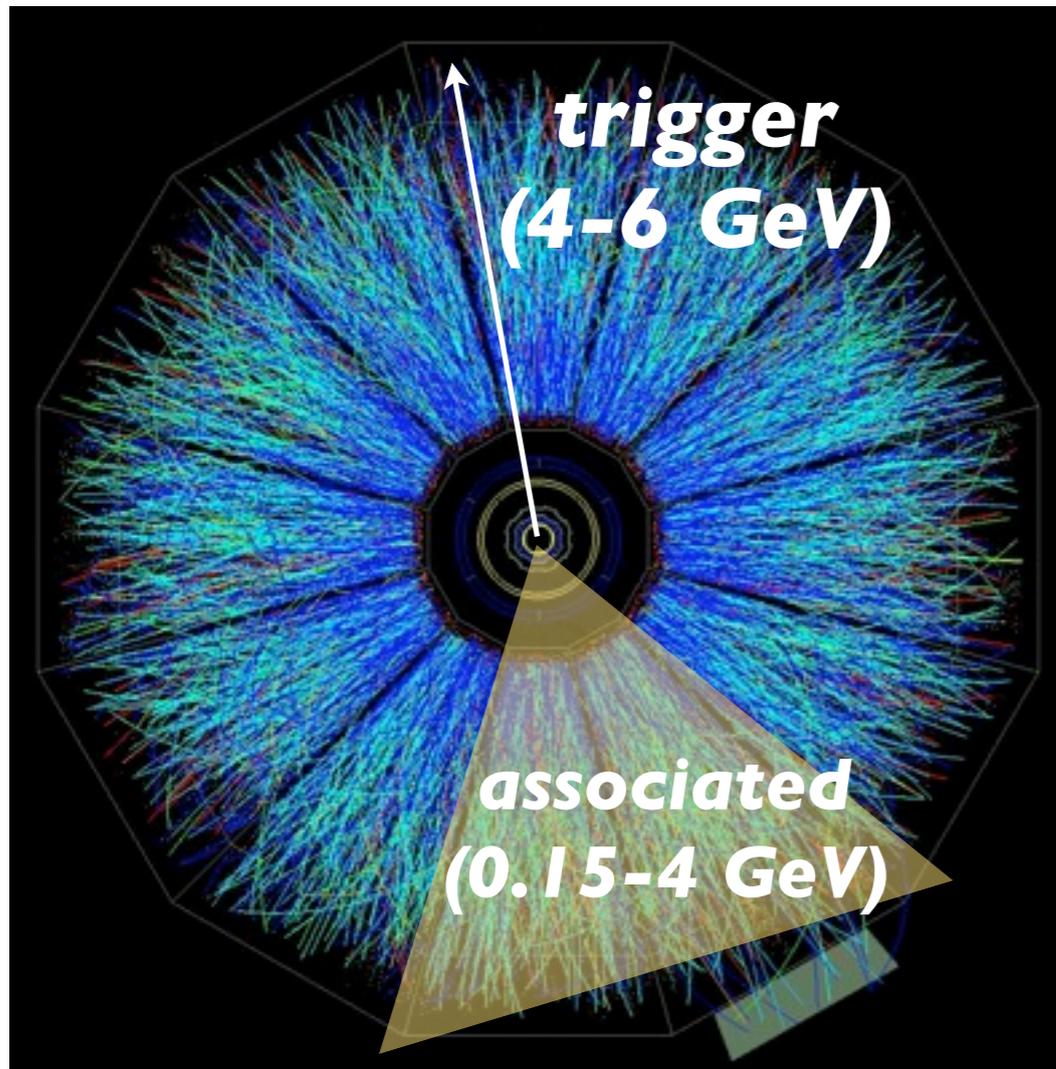
d+Au similar to p+p
Au+Au shows a disappearance
of the away side peak



Surface Bias

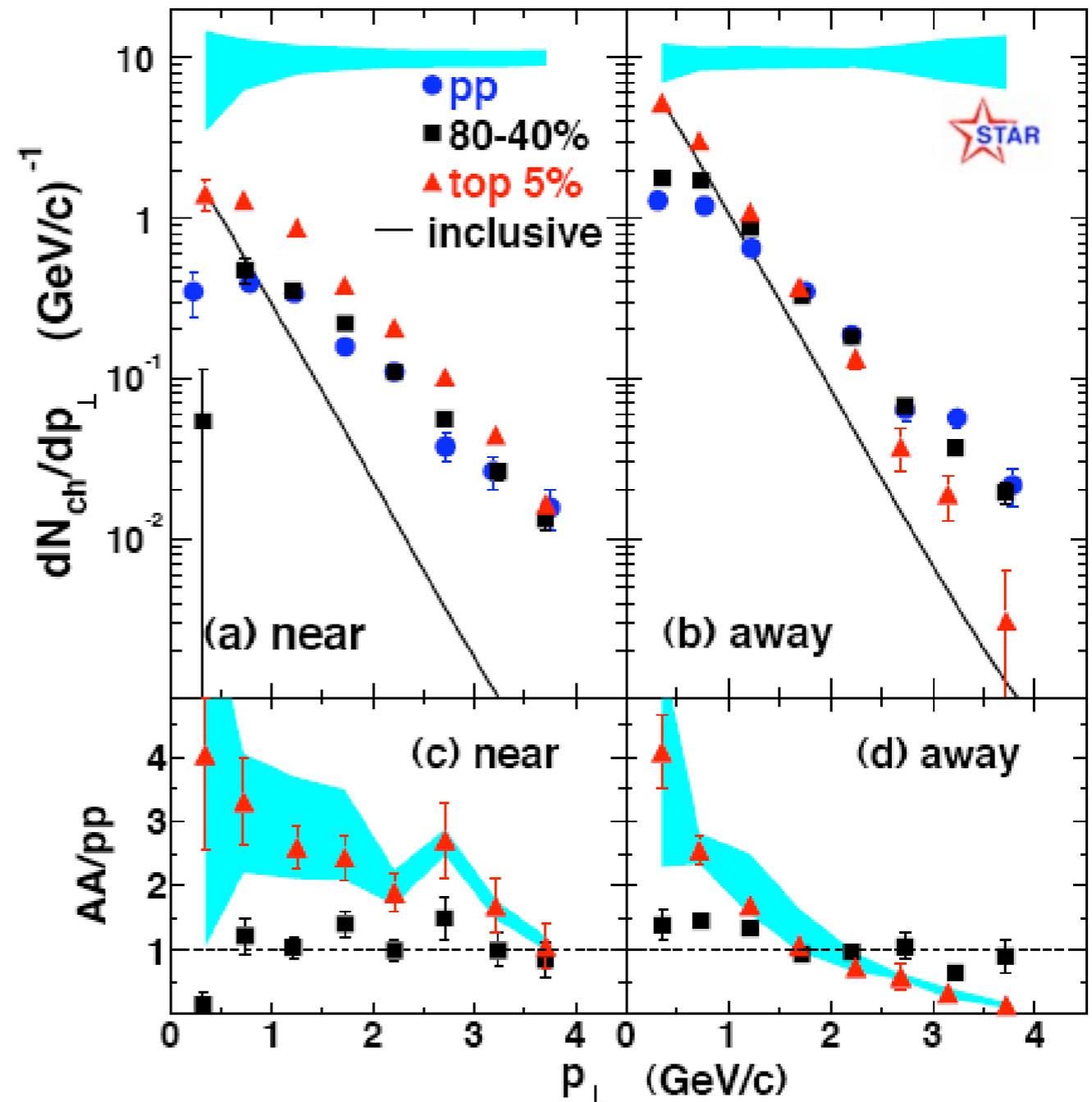
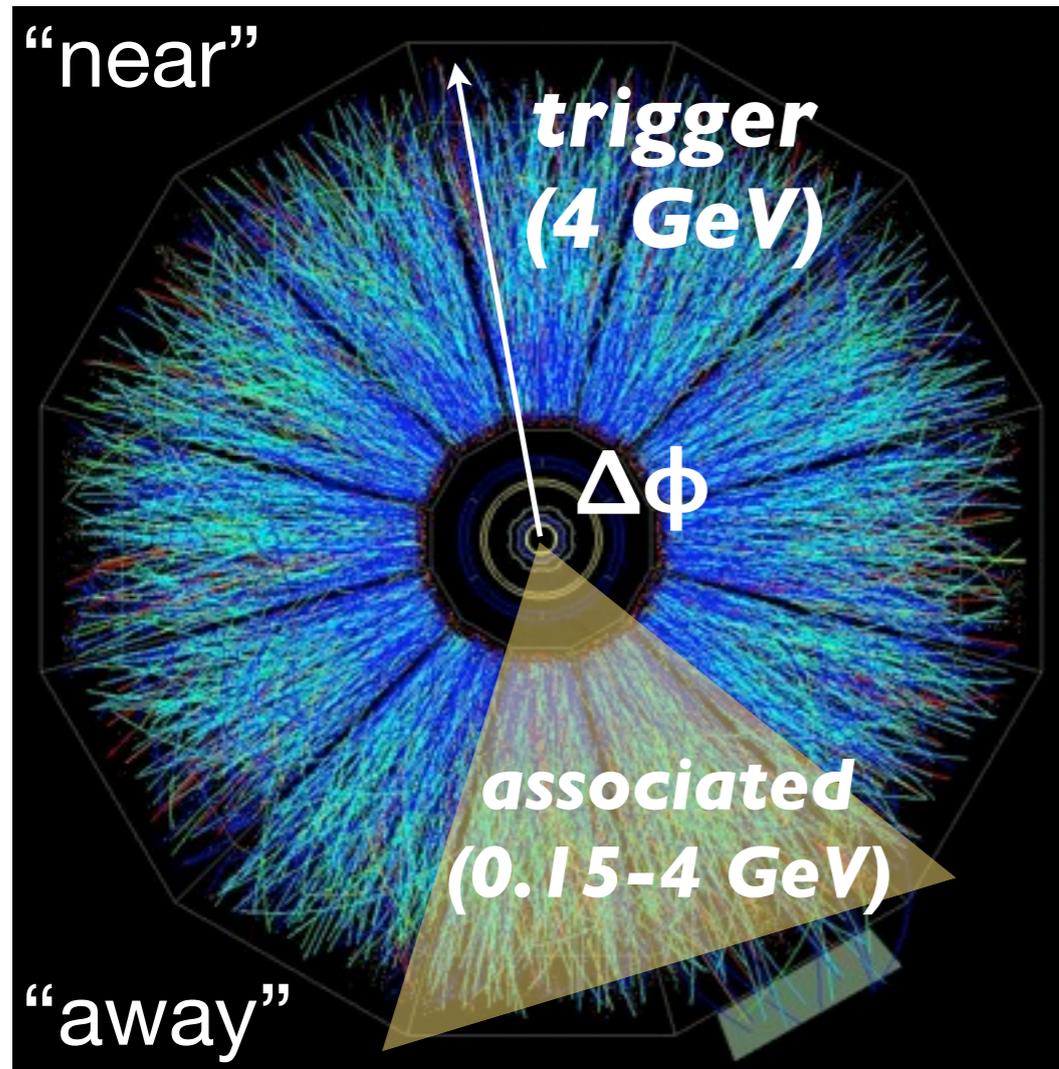


Return of the “Away Side”



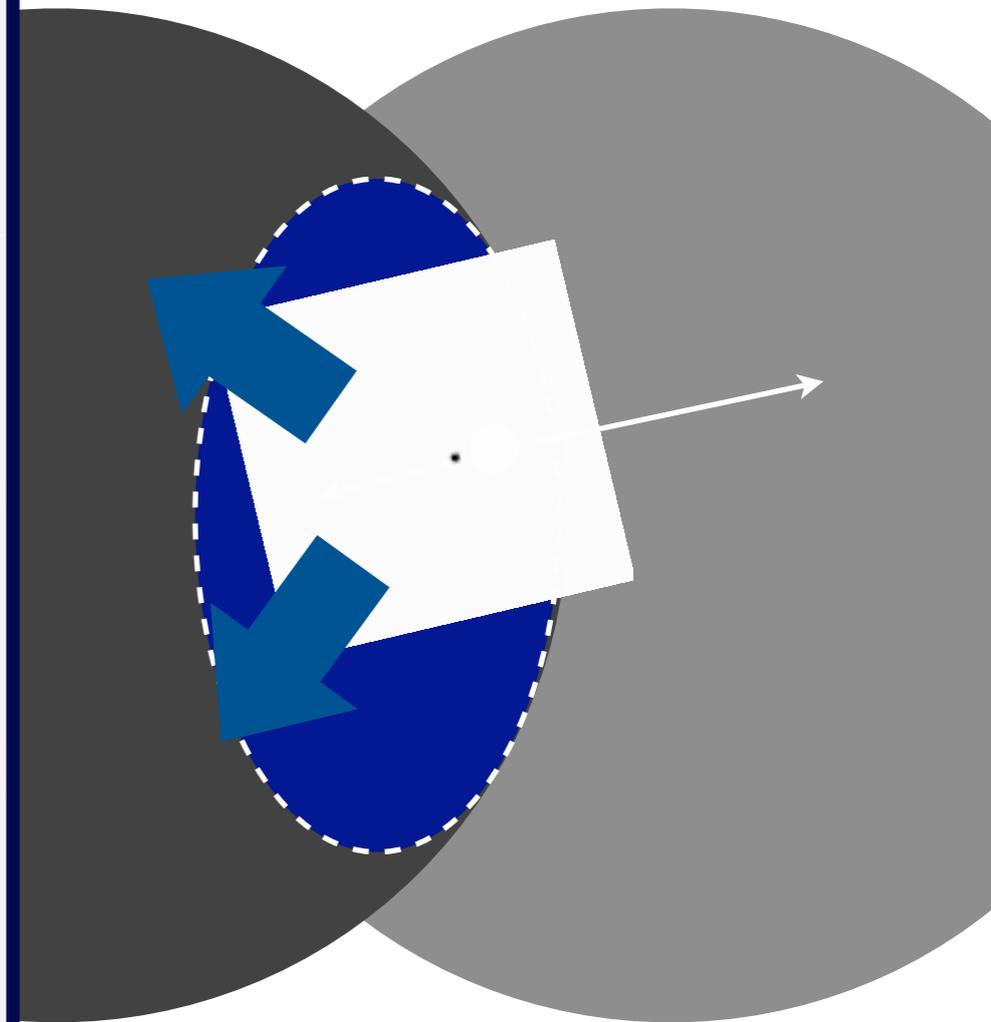
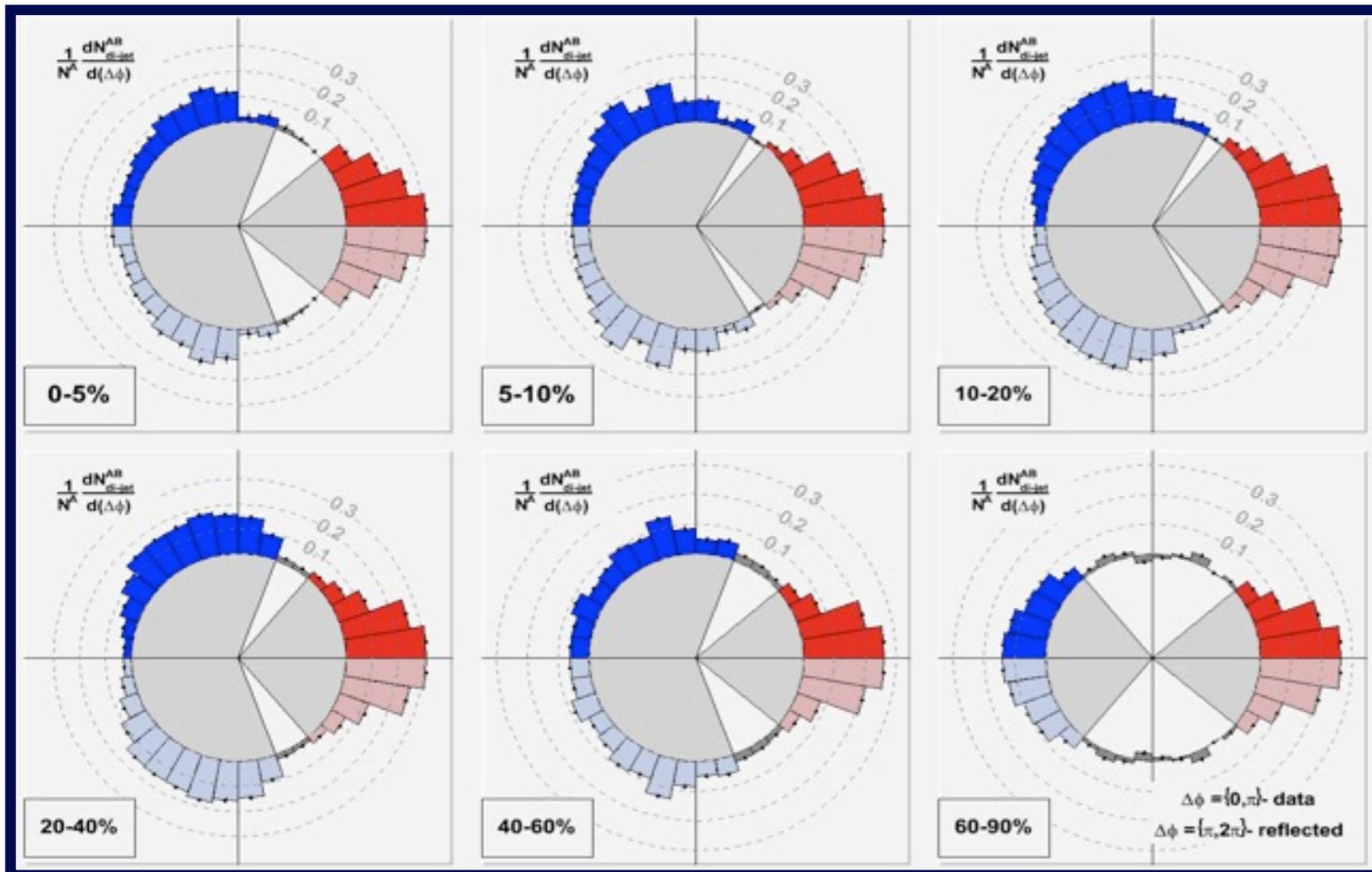
Including all particles (soft & hard) accounts for suppressed jet, but highly smeared-out in ϕ
Indicates non-trivial interaction with medium.

Spectral Modification (thermalization)



Away-side spectrum starts to resemble inclusive “thermal” one

Medium Effects on Jets



In central events, 2-particle correlations not back-to-back!

Suppression is a “redistribution” of energy/momentum.
Excitations couple strongly to medium, rapidly thermalize

May 2006

Drag force in AdS/CFT

Mach cones from AdS/CFT, Gubser et al hep-th/0607022

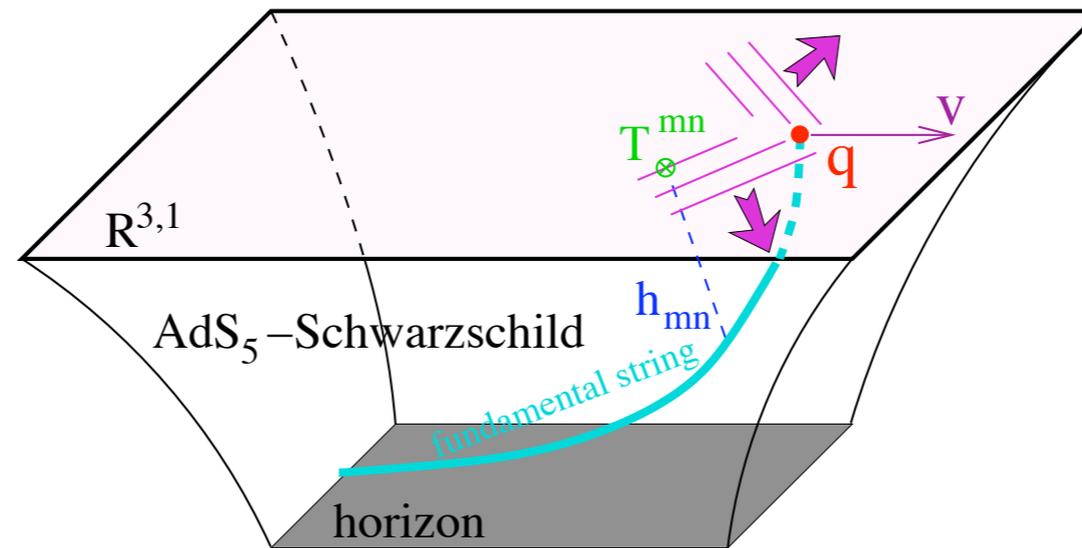
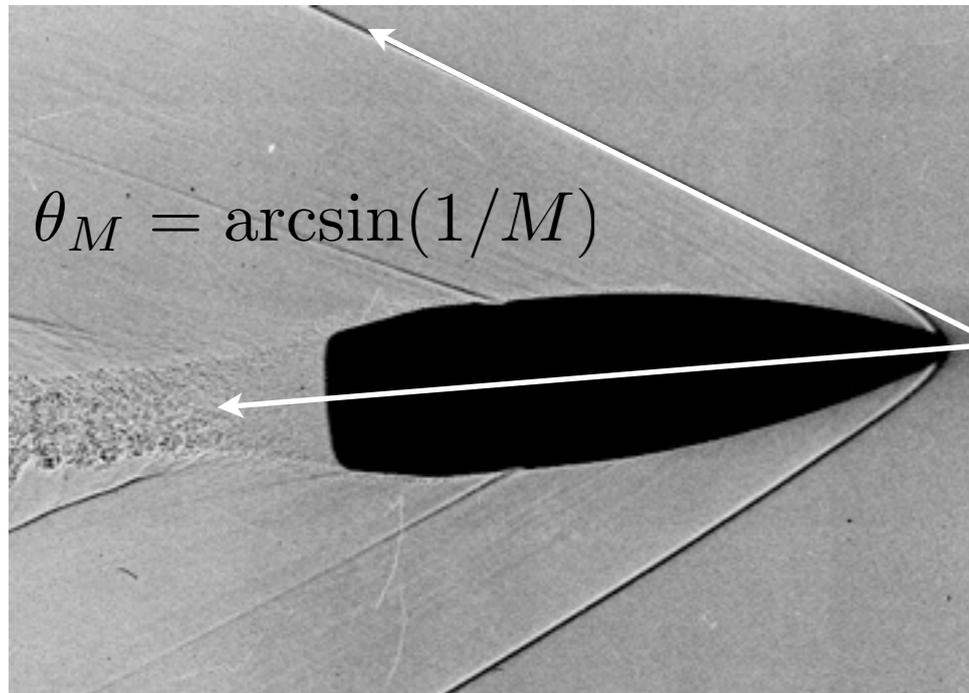
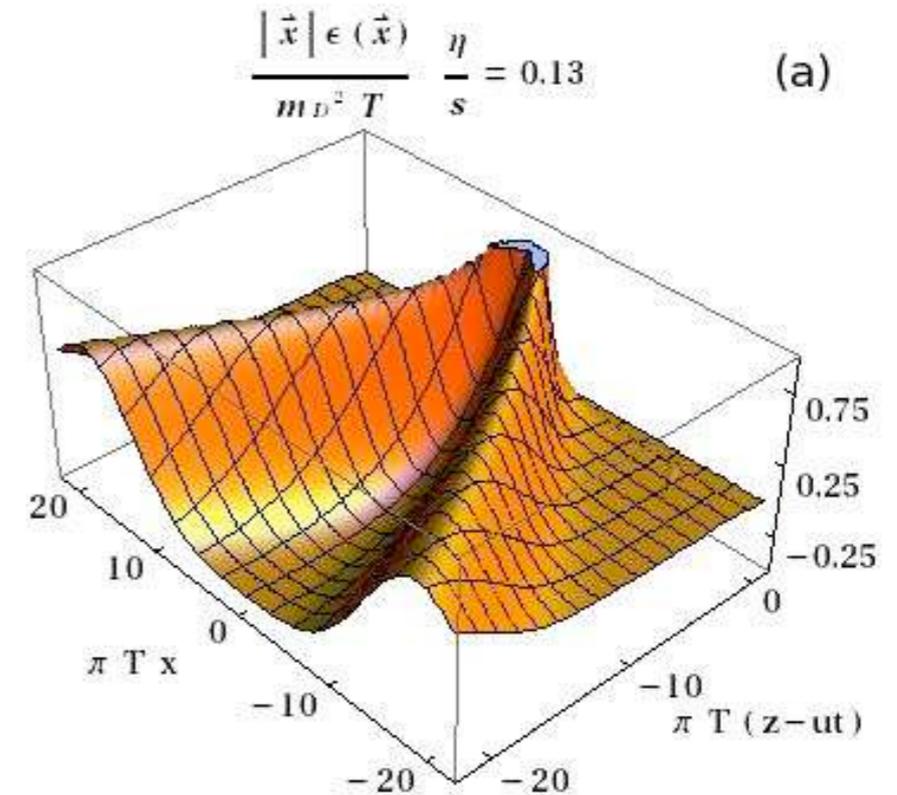


Figure 1: The AdS_5 -Schwarzschild background is part of the near-extremal D3-brane, which encodes a thermal state of $\mathcal{N} = 4$ supersymmetric gauge theory [24]. The external quark trails a string into the five-dimensional bulk, representing color fields sourced by its fundamental charge and interacting with the thermal medium.

QCD Mach Cones?



wQGP:
Neufeld,
Ruppert,
Mueller
(2008)

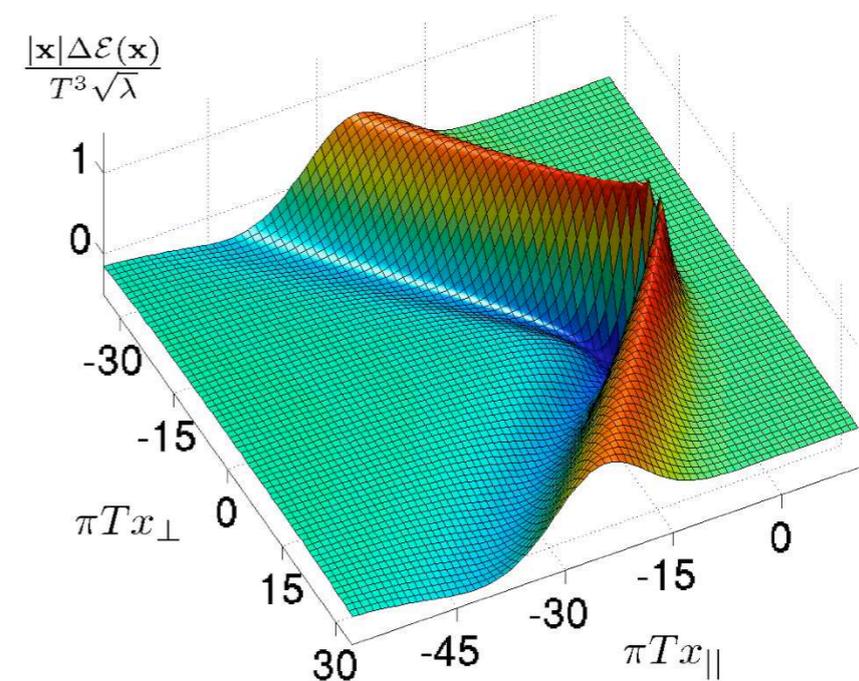


Speed of light: c
Speed of sound: $c/\sqrt{3}$

Expect shock wave ~ 1
radian from jet axis

Hot debate: 3p correlations

sQGP:
Chesler
Yaffe
(2007)



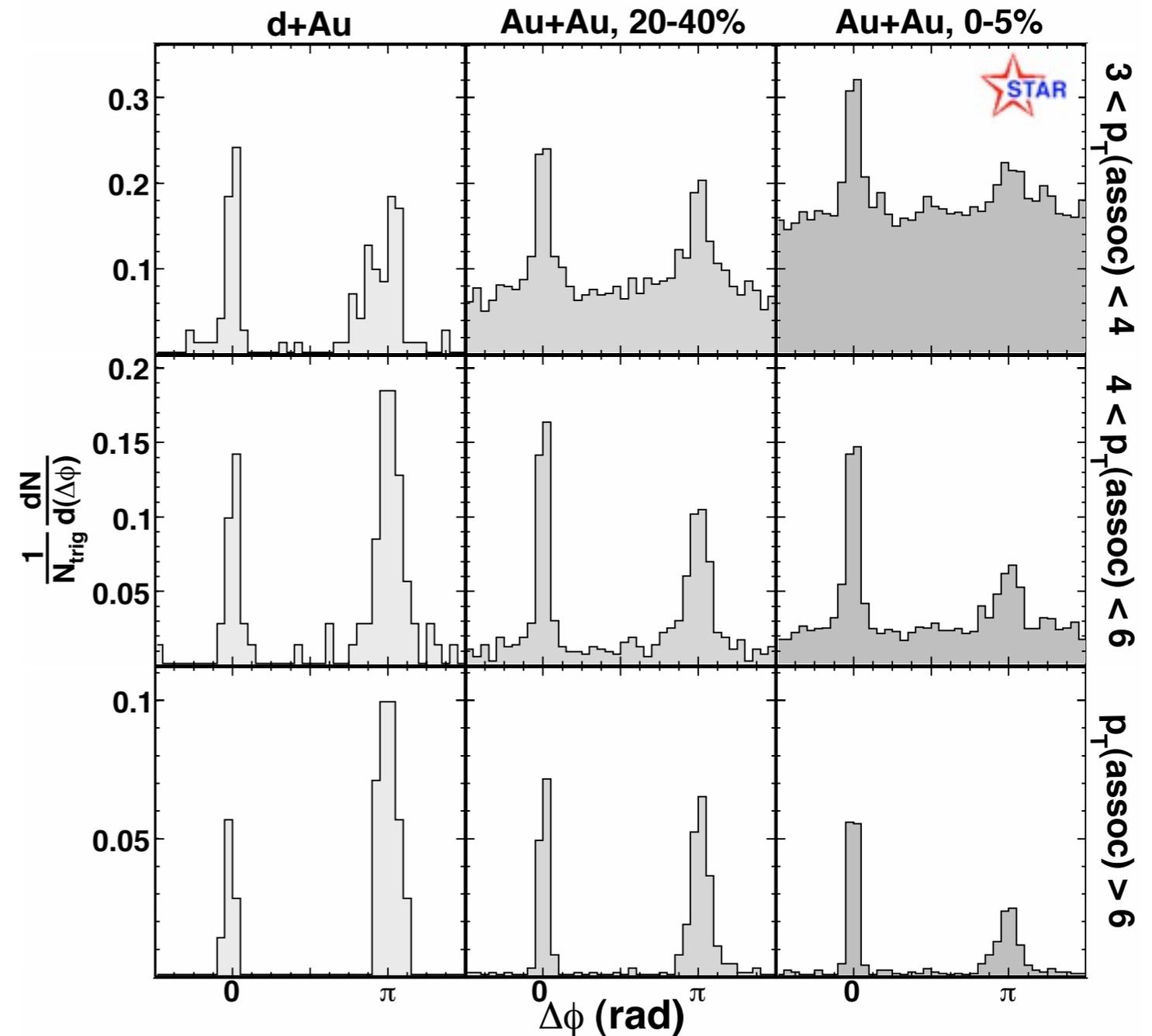
Back to Back comes back (STAR)

CTEQ

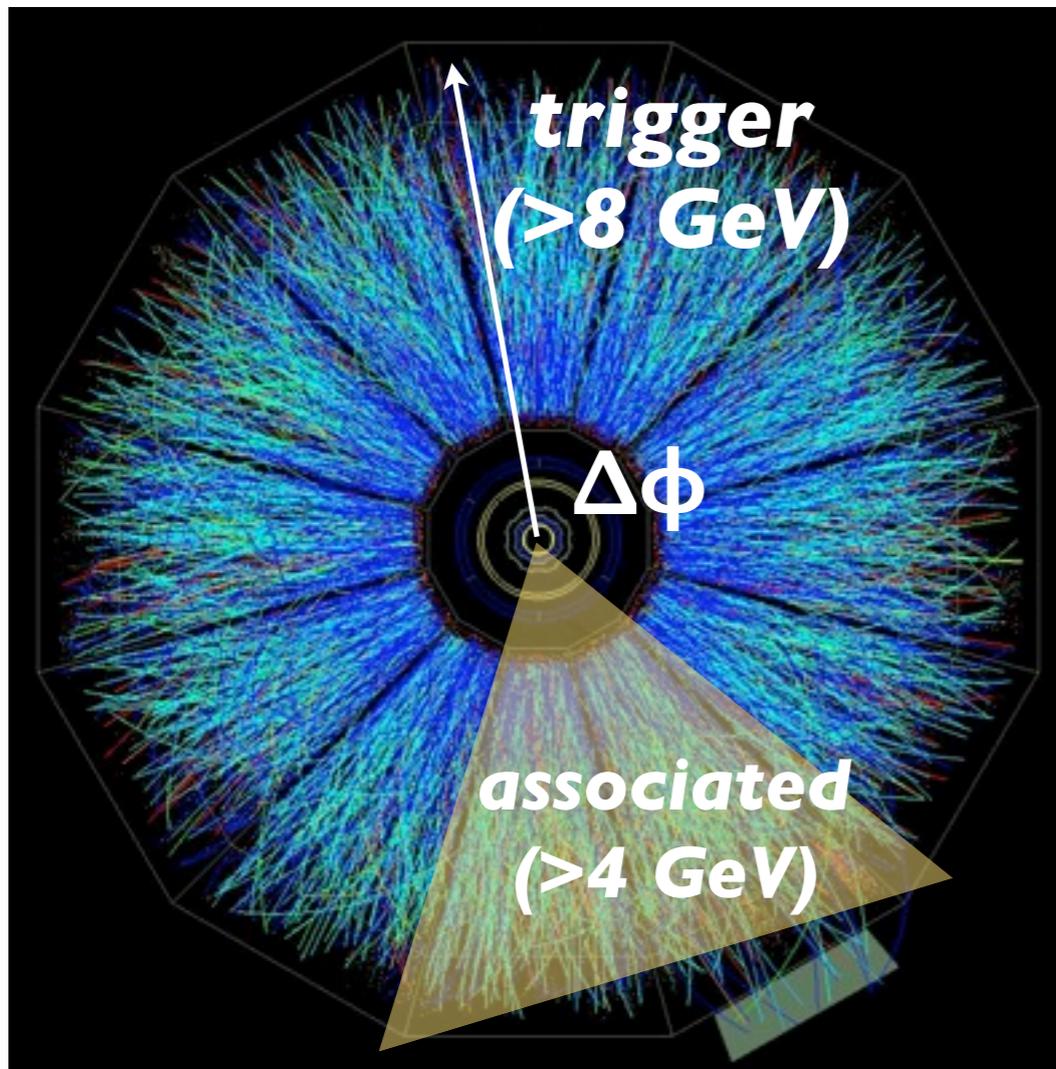
Previous results:
near-complete quenching in
strongly-coupled plasma.

Raising momentum of
trigger & associated particles
show back-to-back jets

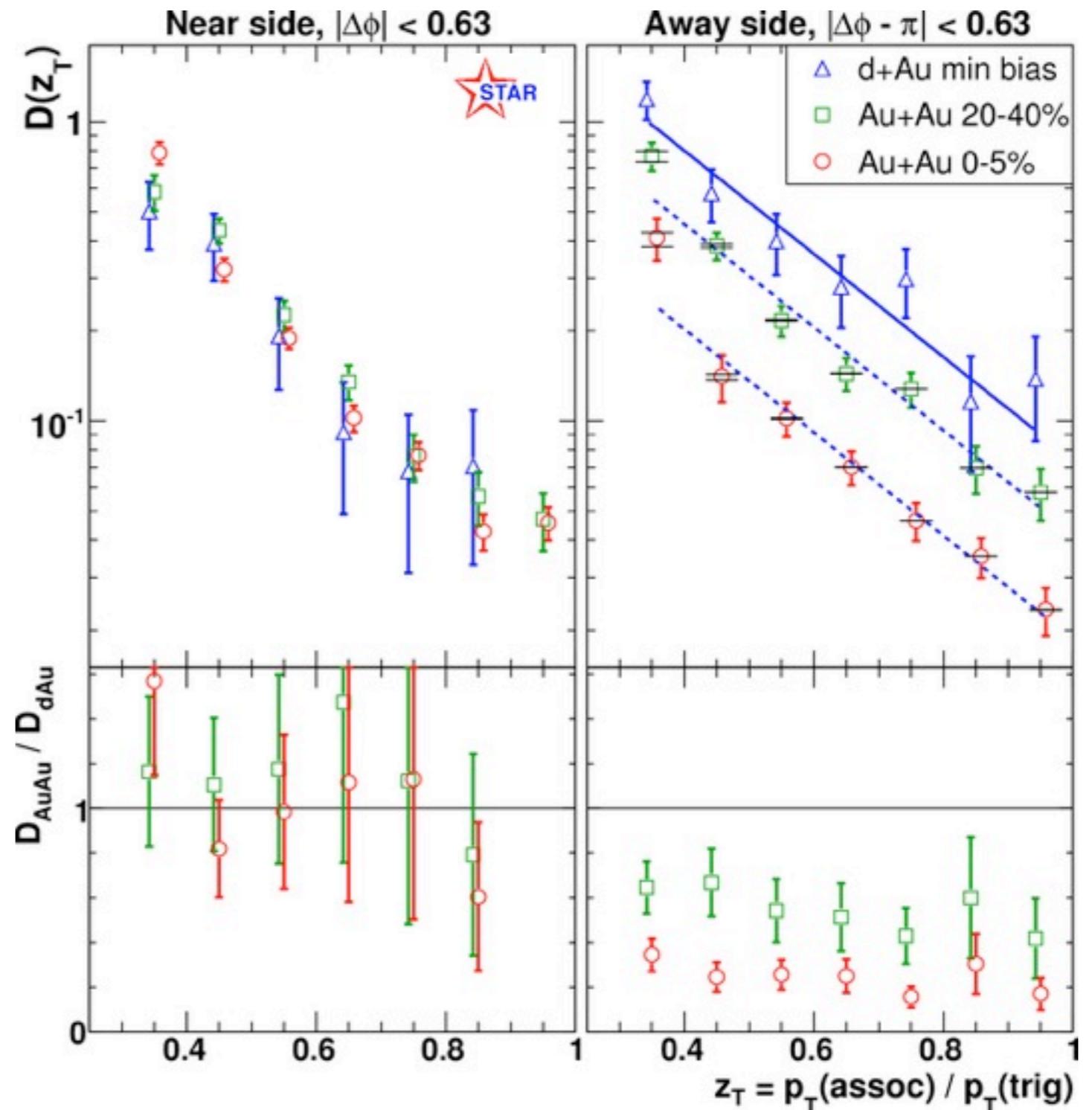
Observation of dijets
in heavy ion collisions



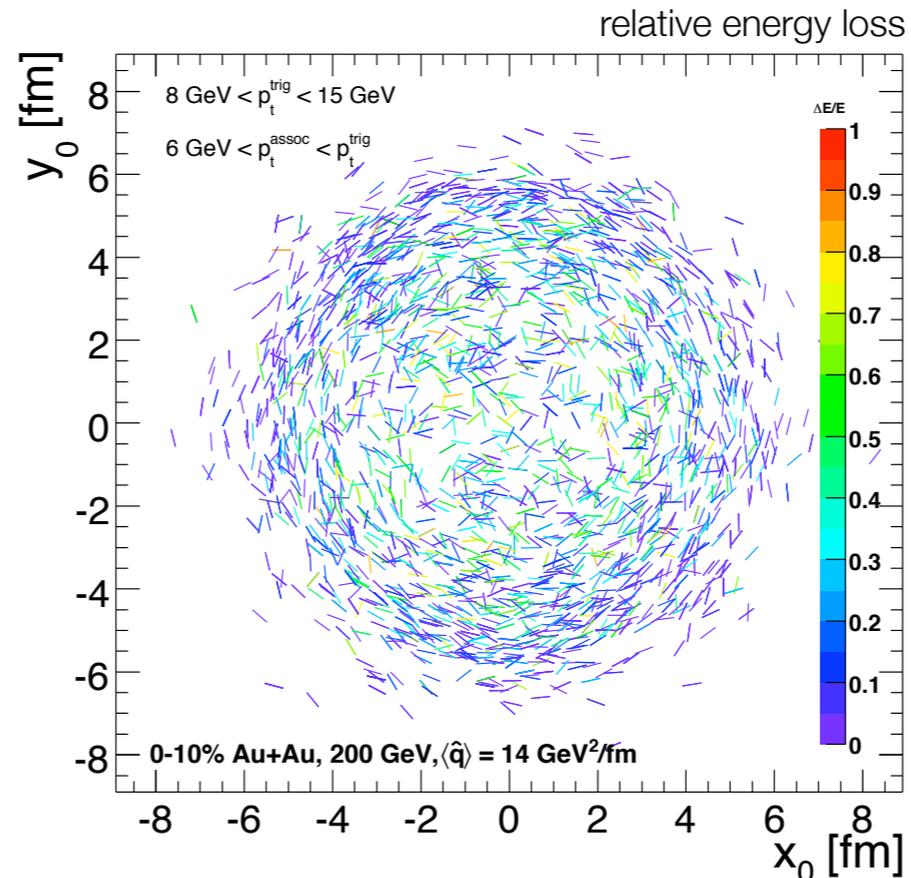
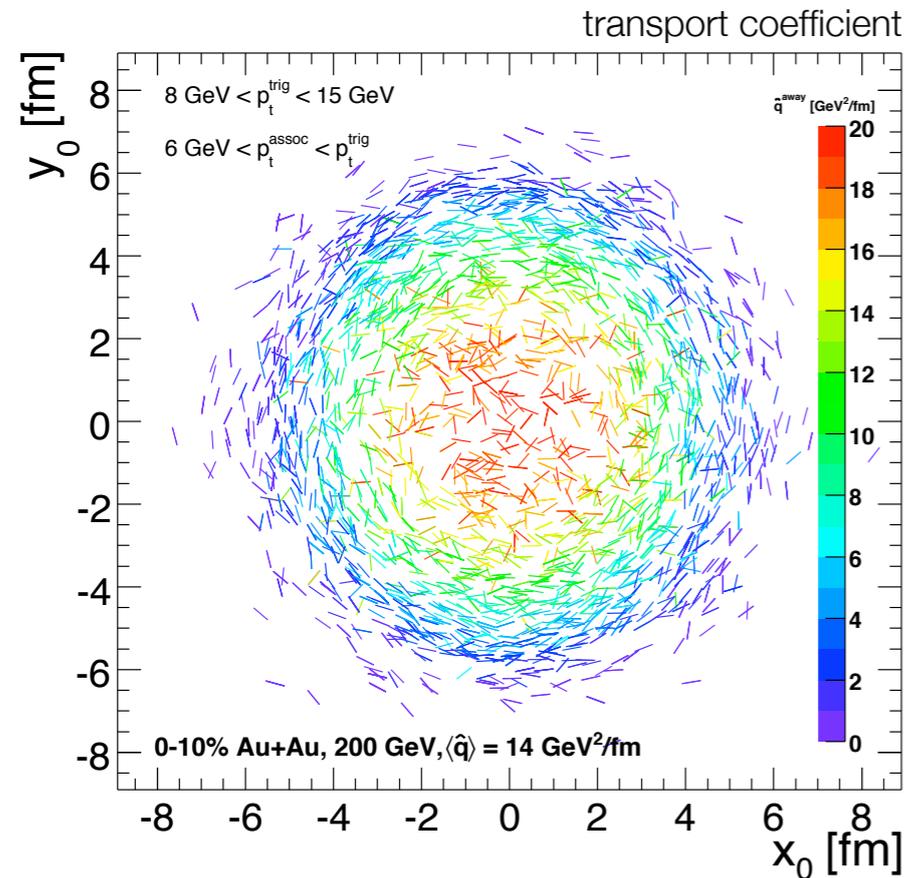
Spectral non-modification



Dijet fragmentation on away side appears unmodified relative to d+Au

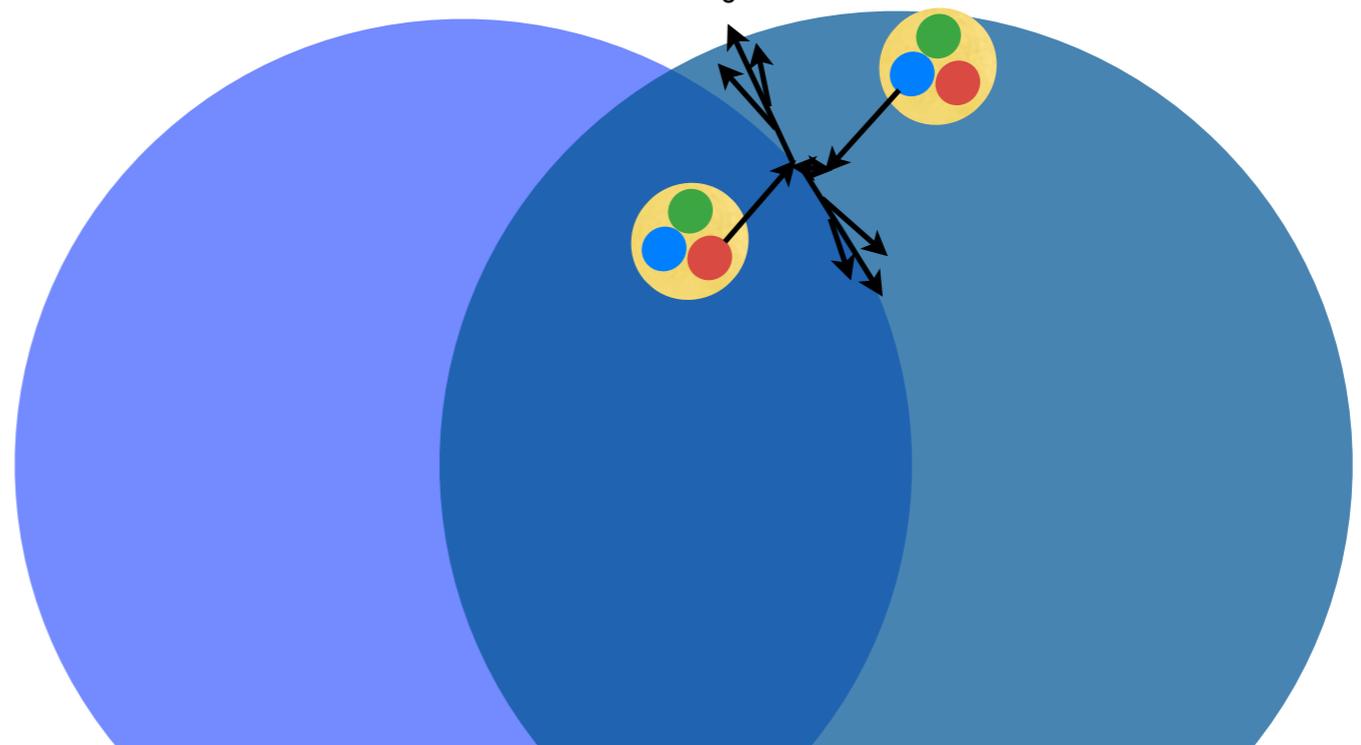


Tangential emission



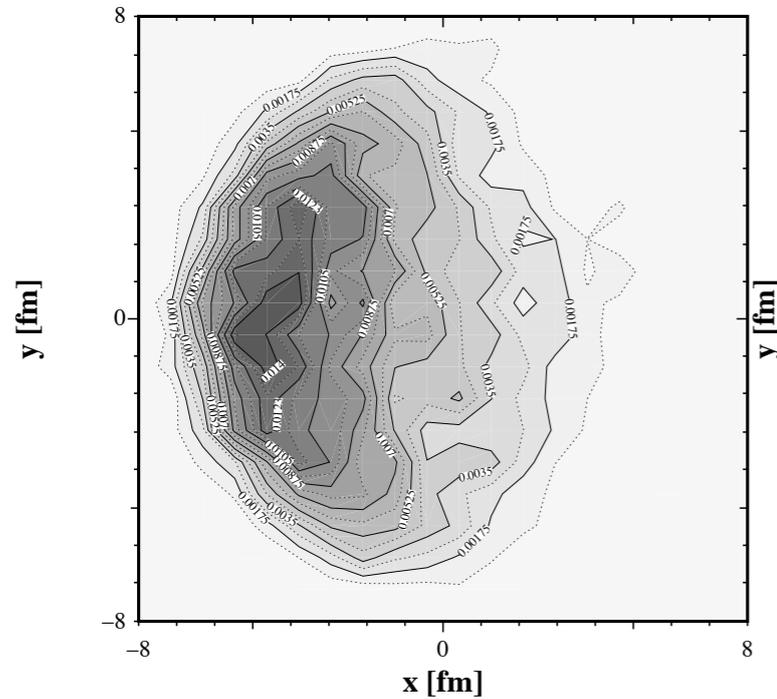
Loizides HQ06(2006)

Jet rates and no observable modification suggests that these jets “escaped”:
“tangential emission”

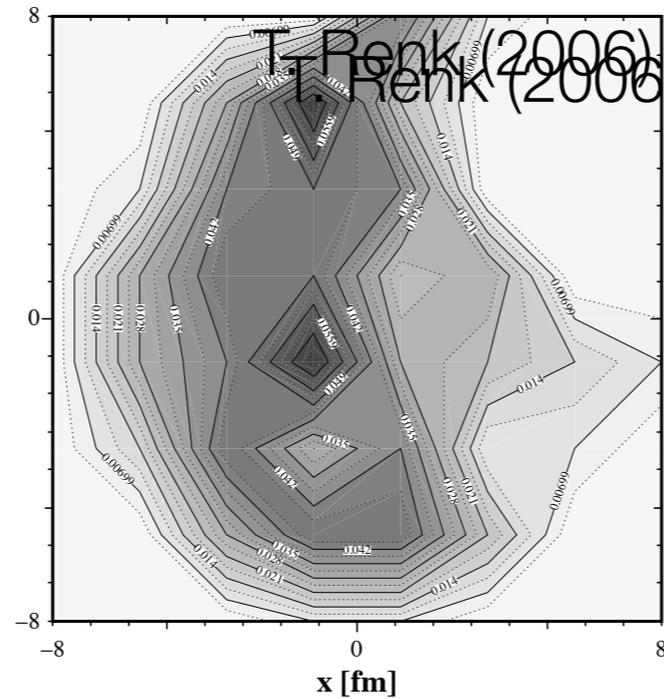


“Punch through”

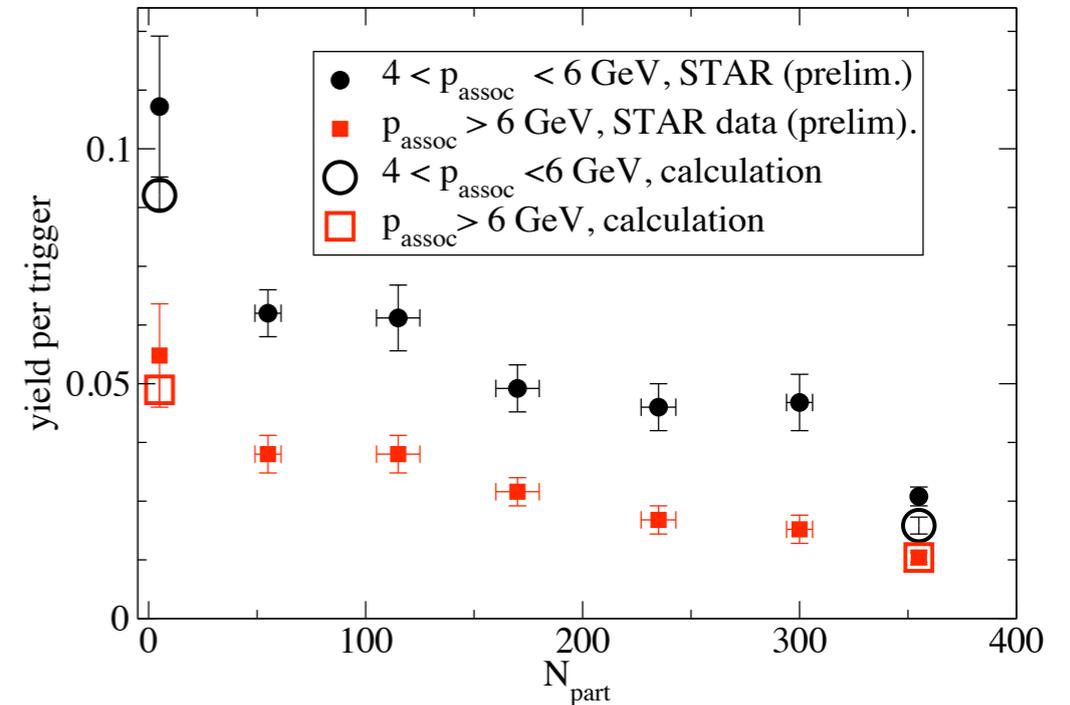
Triggered vertex distribution



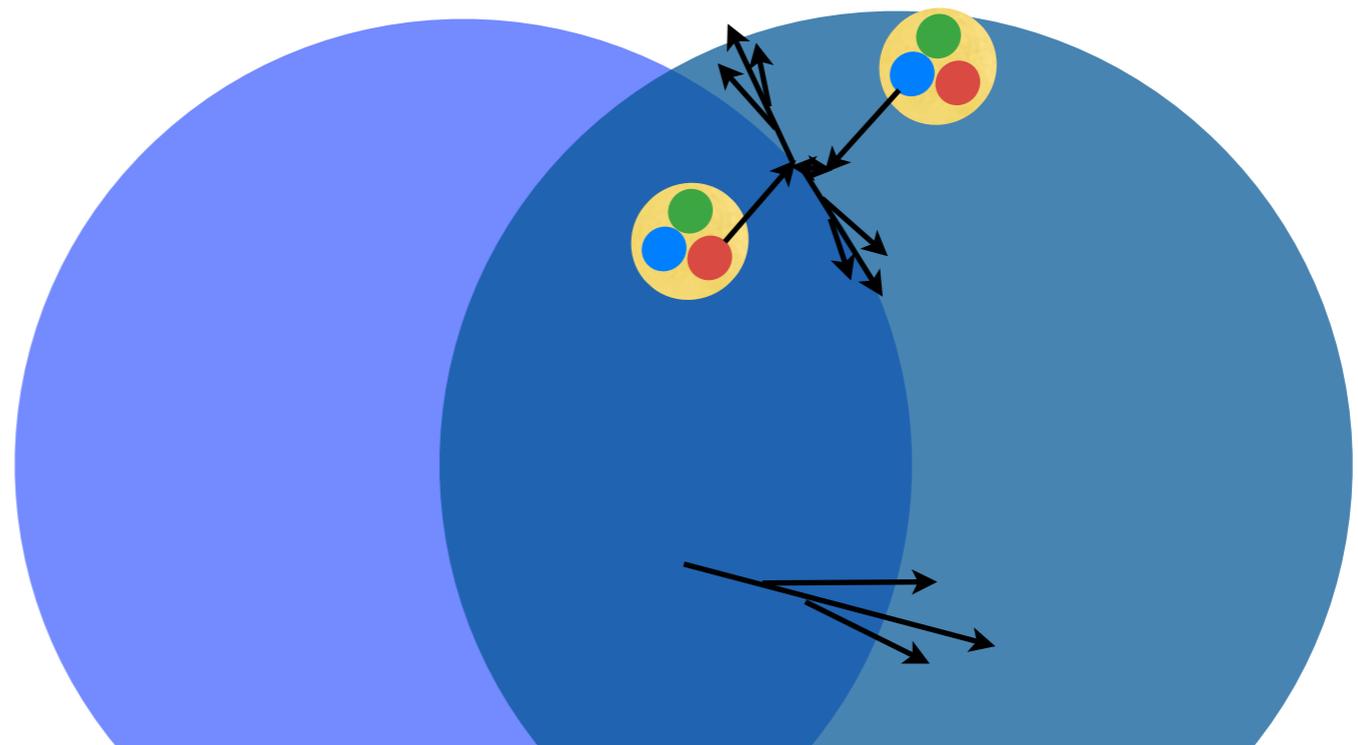
Punchthrough distribution



away side, trigger > 8 GeV



some jets will lose little energy even when passing through medium: “punch through” (25% in this case)



The Landscape of 2-particle $\Delta\Phi$ correlations

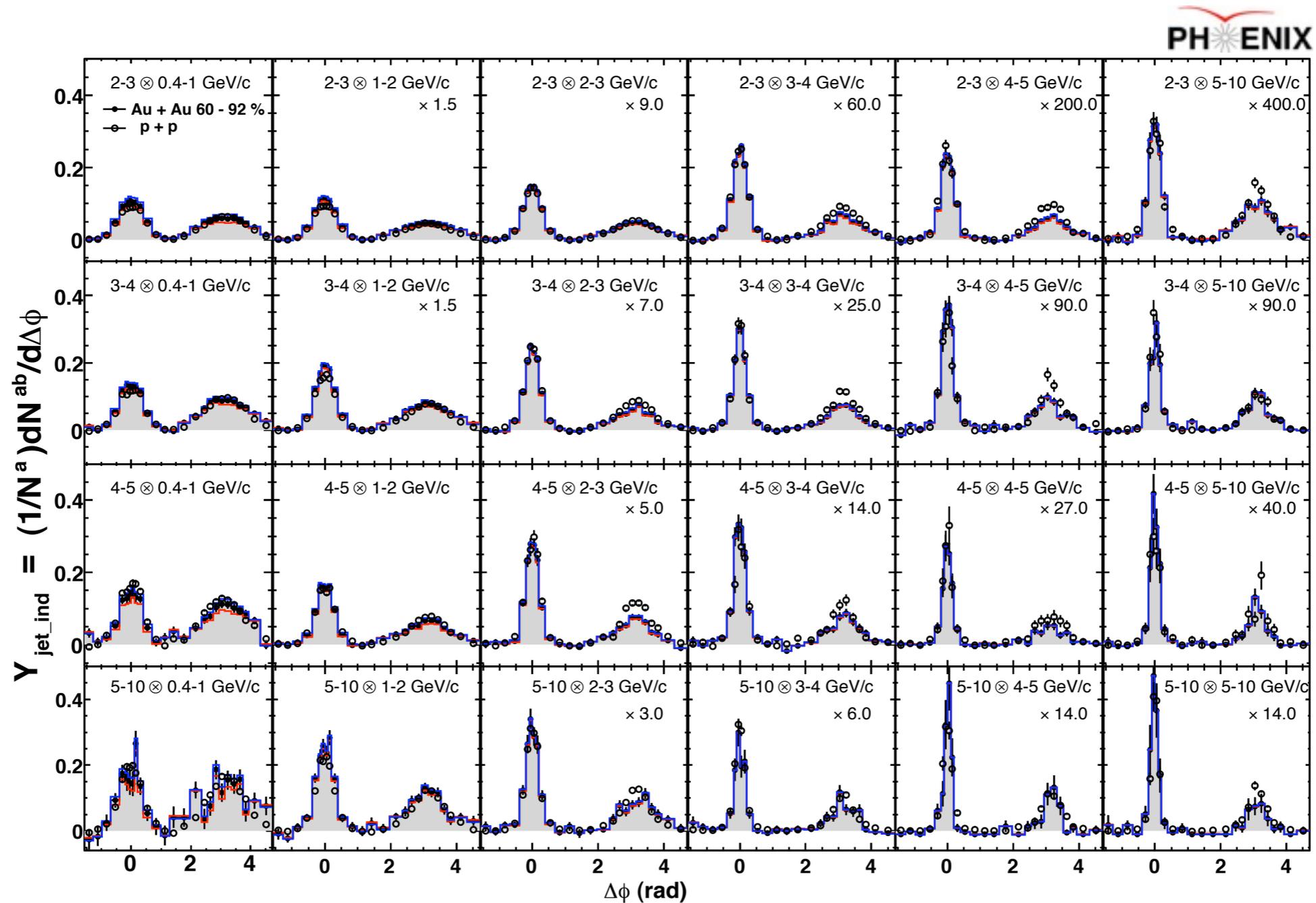


FIG. 38: (Color online) Per-trigger yield versus $\Delta\phi$ for successively increasing trigger and partner p_T ($p_T^a \otimes p_T^b$) in $p + p$ (open circles) and 60-92 % Au+Au (filled circles) collisions. Data are scaled to the vertical axes of the three left panels. Histograms indicate elliptic flow uncertainties for Au+Au collisions.

The Landscape of 2-particle $\Delta\Phi$ correlations

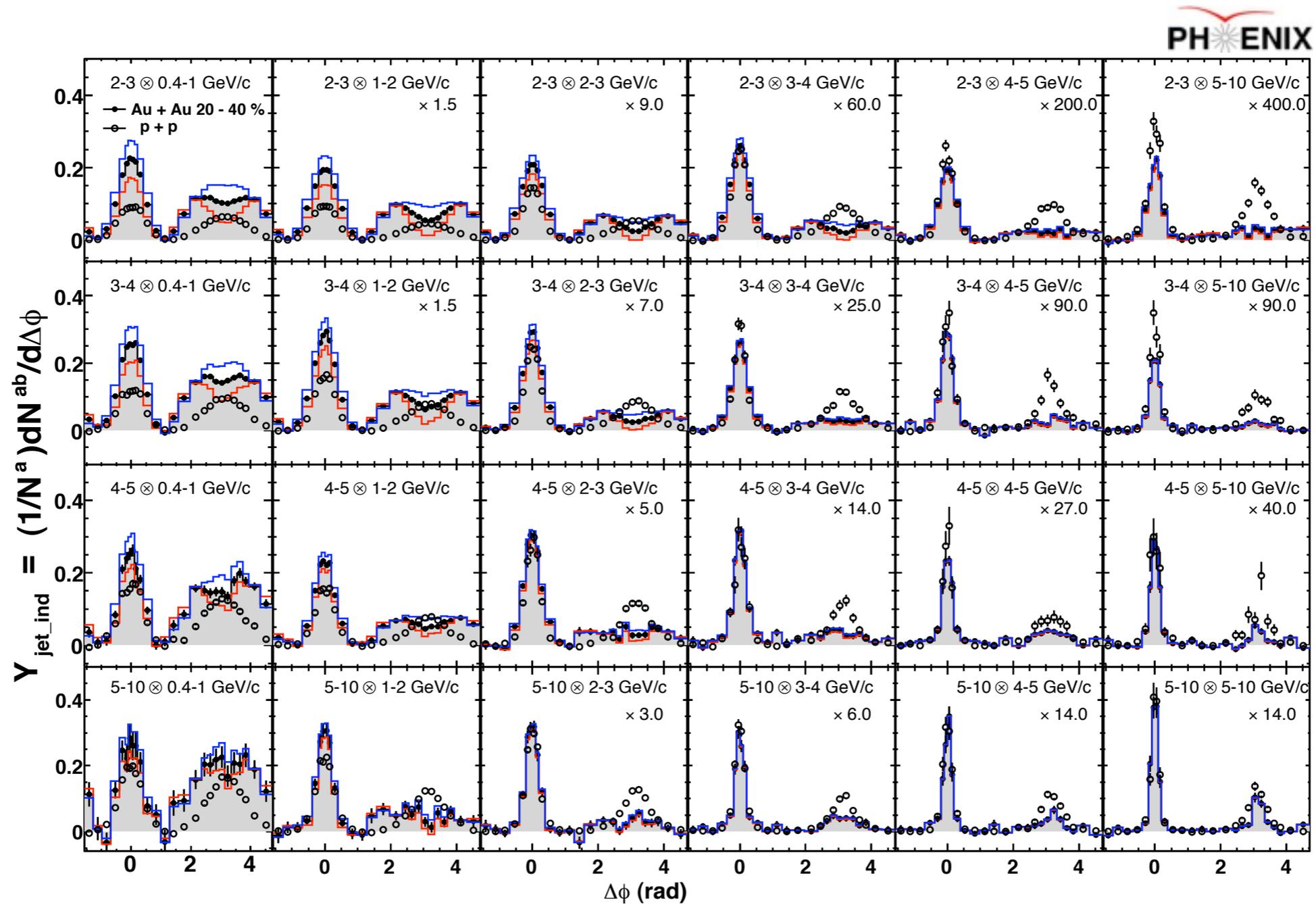
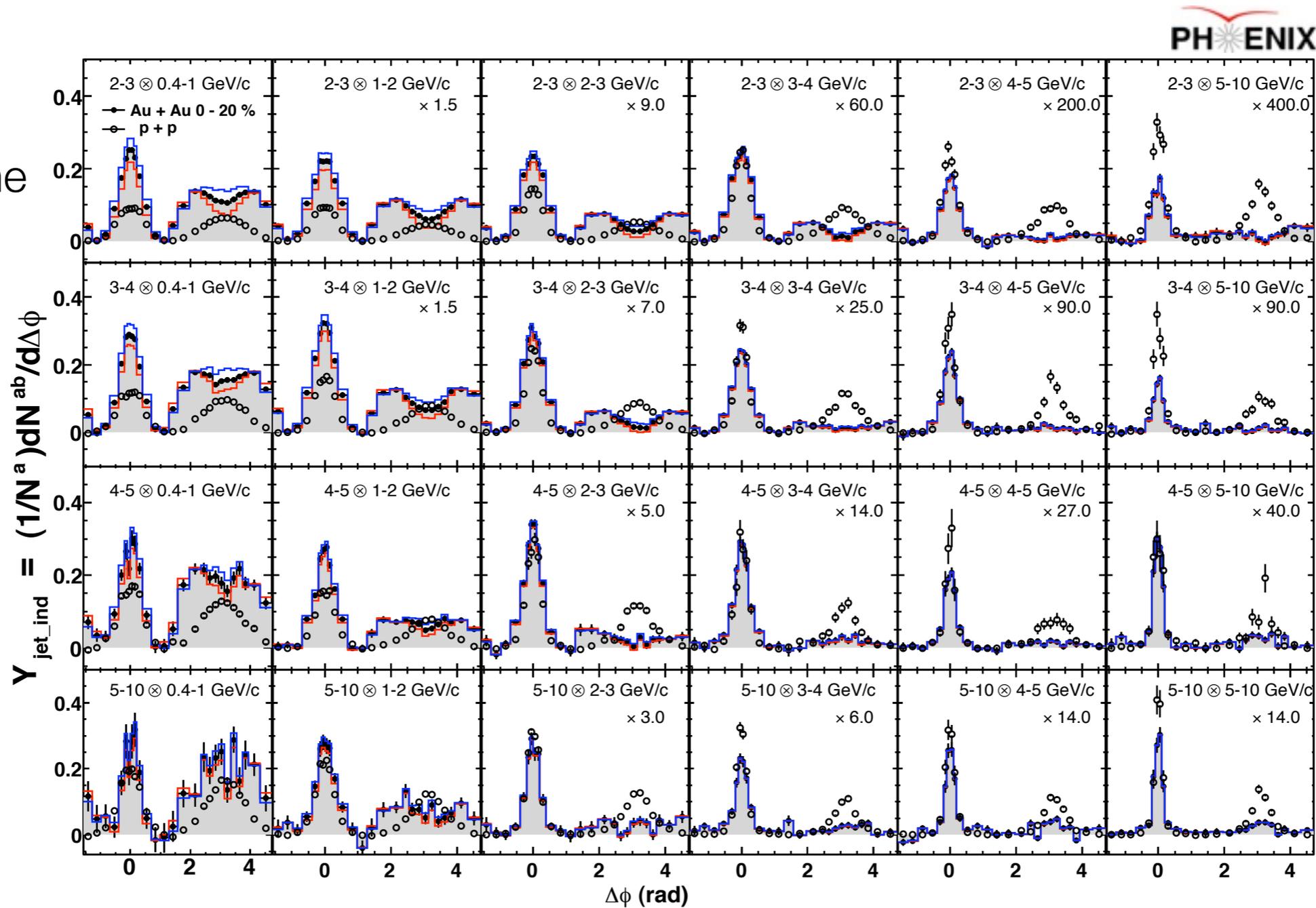


FIG. 37: (Color online) Per-trigger yield versus $\Delta\phi$ for successively increasing trigger and partner p_T ($p_T^a \otimes p_T^b$) in $p+p$ (open circles) and 20-40 % Au+Au (filled circles) collisions. Data are scaled to the vertical axes of the three left panels. Histograms indicate elliptic flow uncertainties for Au+Au collisions.

The Landscape of 2-particle $\Delta\Phi$ correlations

mach cone
+ "extra"



back-to-back
disappearance

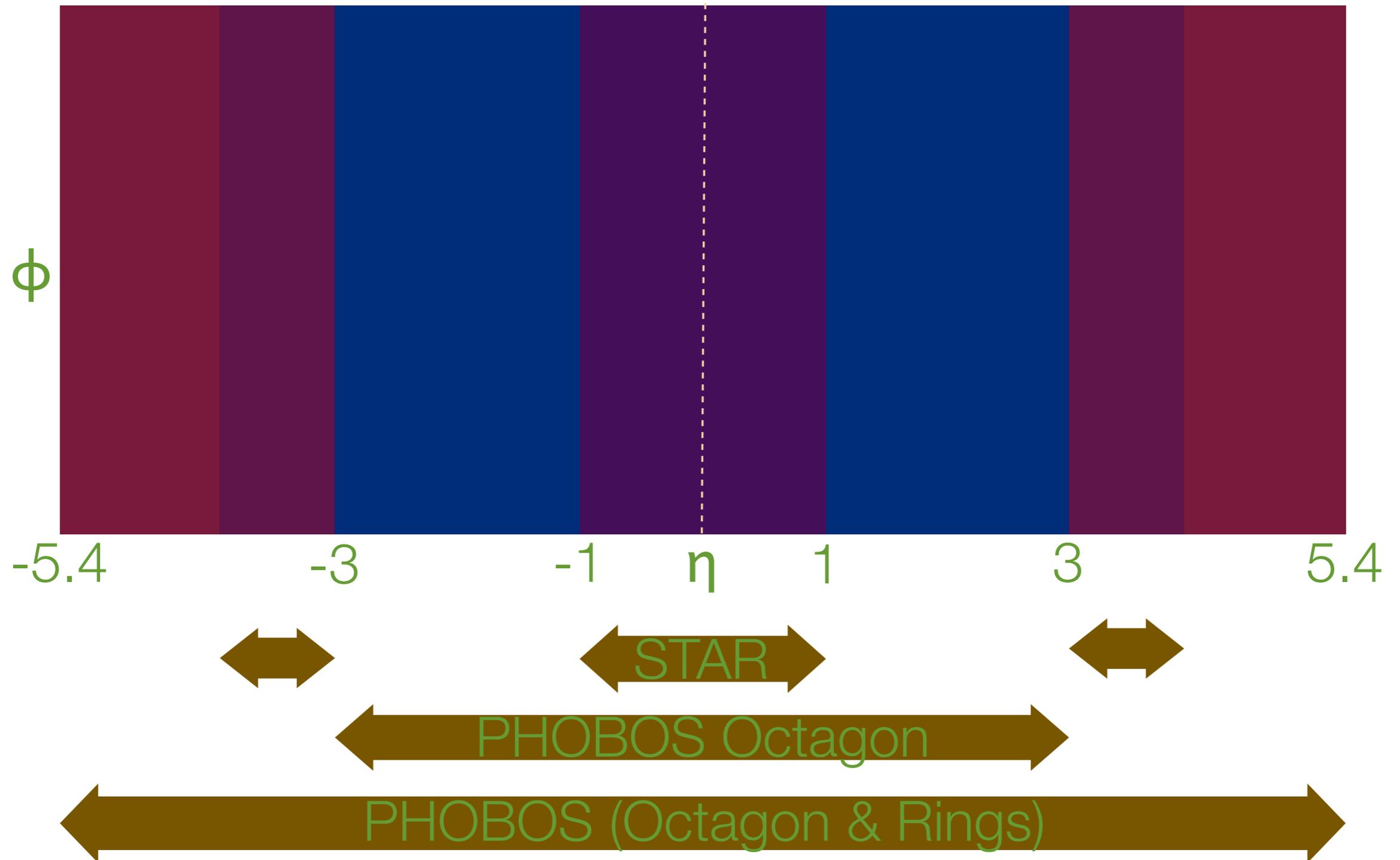
Au+Au
(central)

dijets

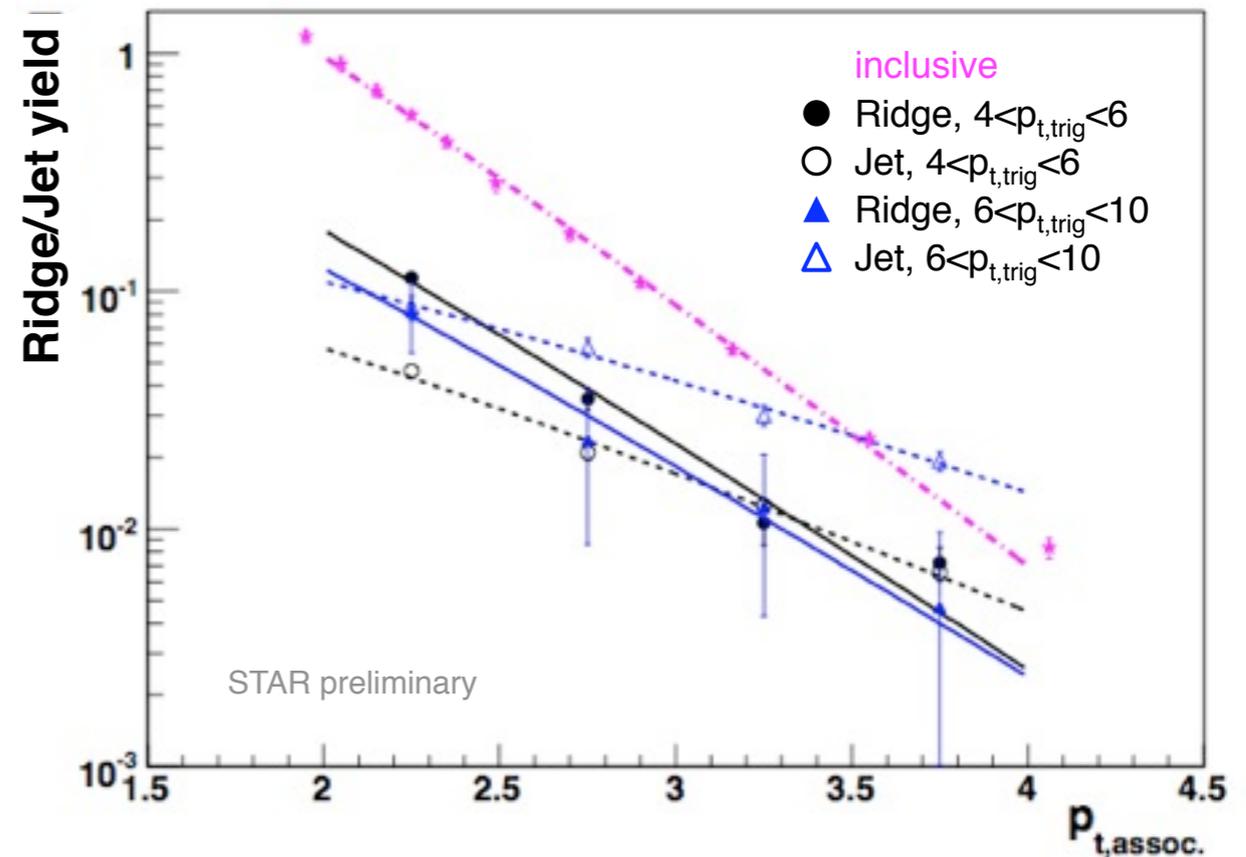
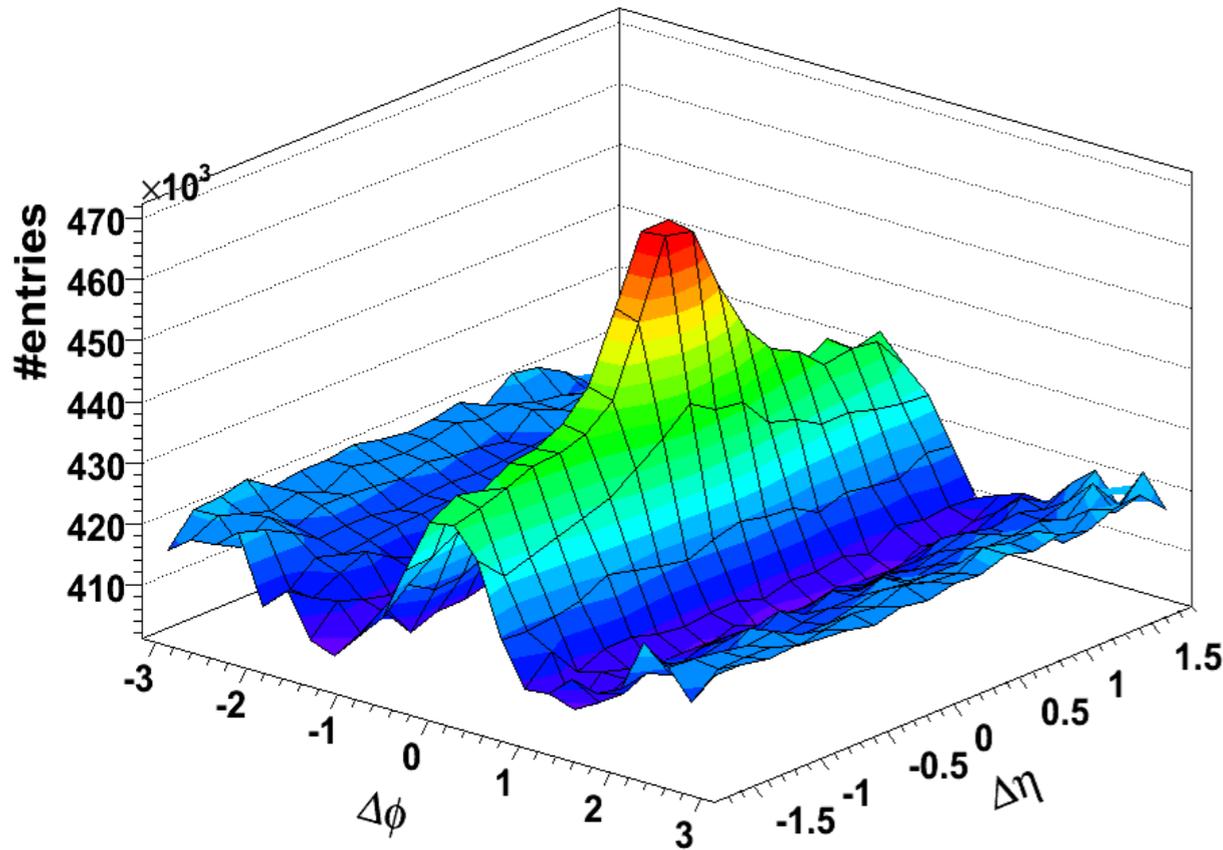
FIG. 36: (Color online) Per-trigger yield versus $\Delta\phi$ for successively increasing trigger and partner p_T ($p_T^a \otimes p_T^b$) in $p+p$ (open circles) and 0-20 % Au+Au (filled circles) collisions. Data are scaled to the vertical axes of the four left panels. Histograms indicate elliptic flow uncertainties for Au+Au collisions.

What about $\Delta\eta$?

CTEQ

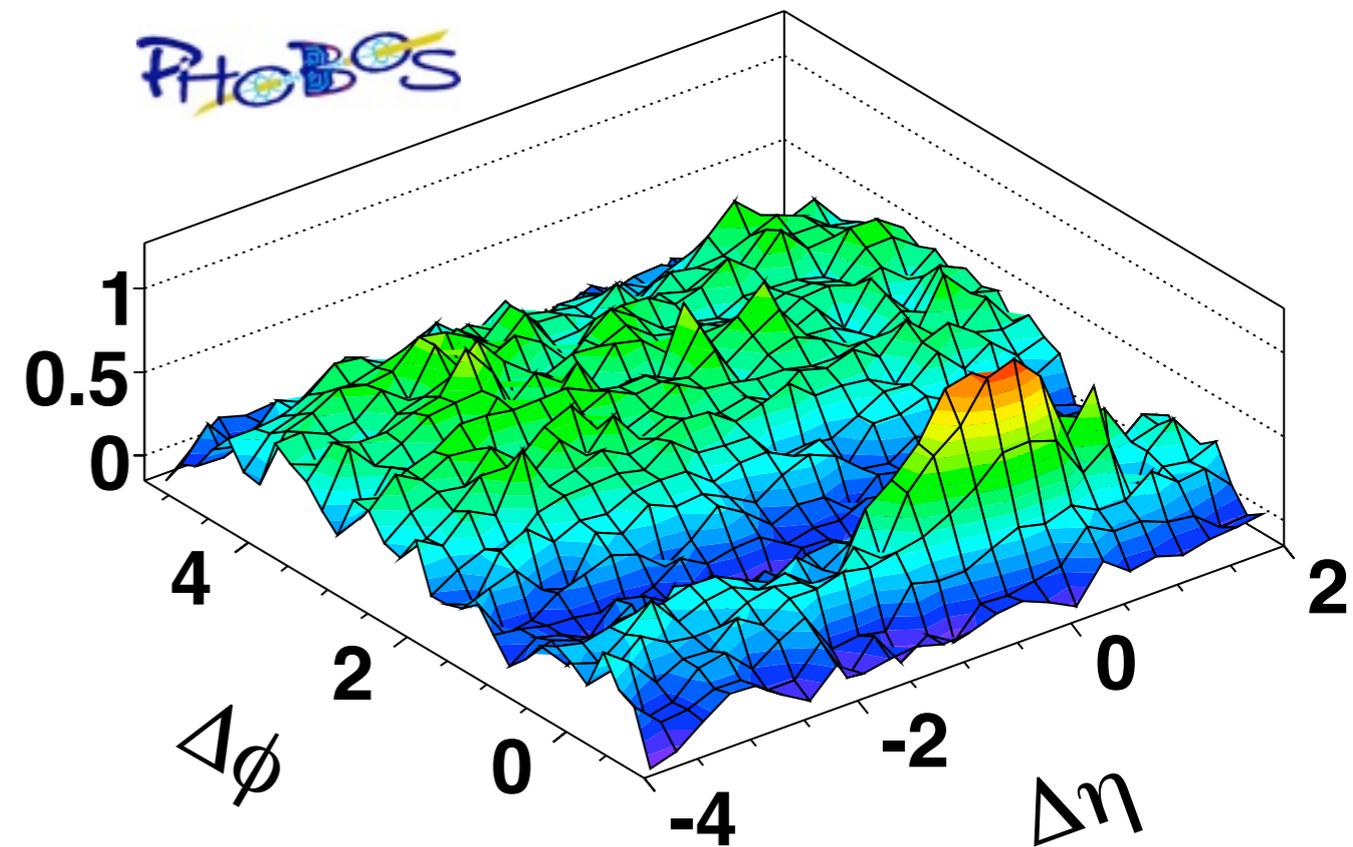
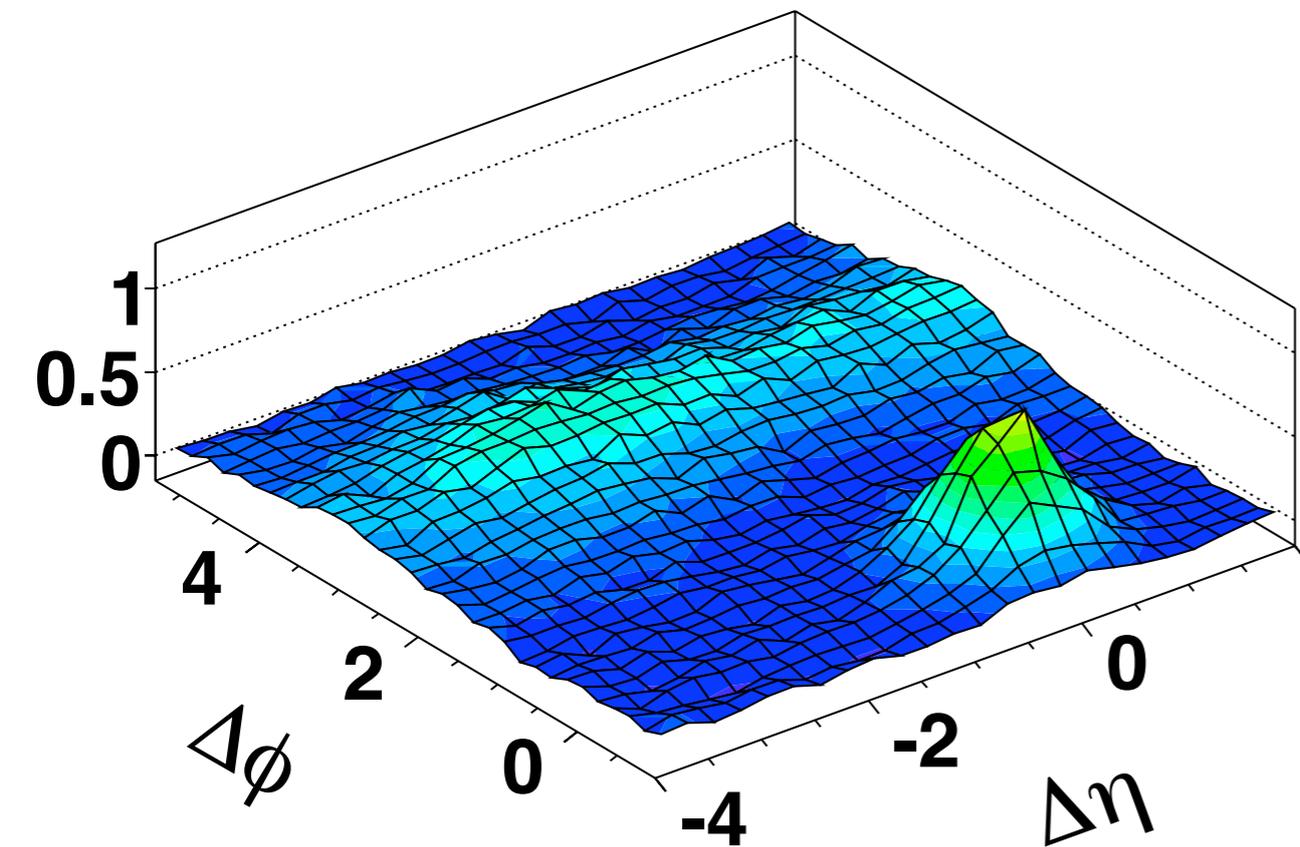


The STAR “ridge”



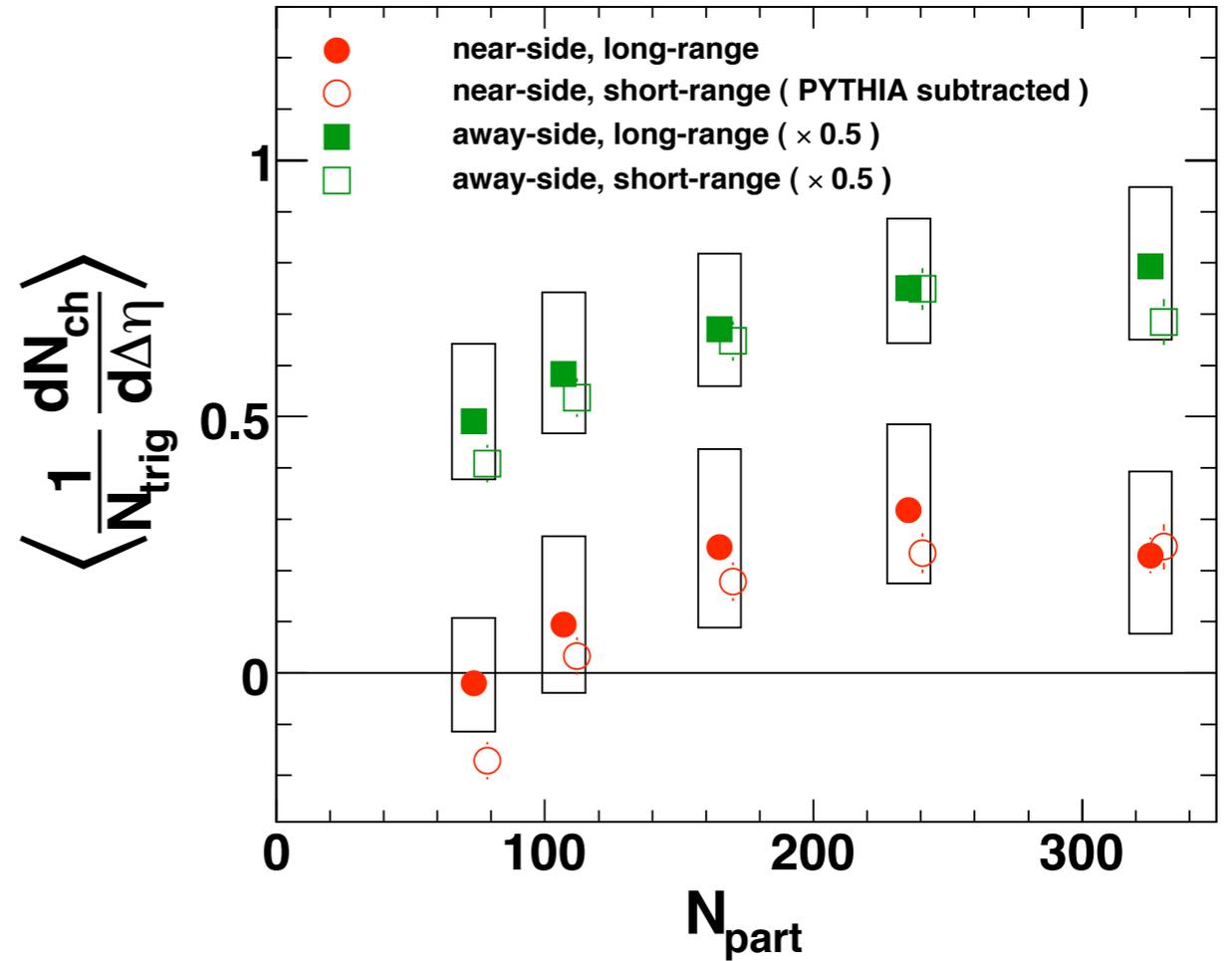
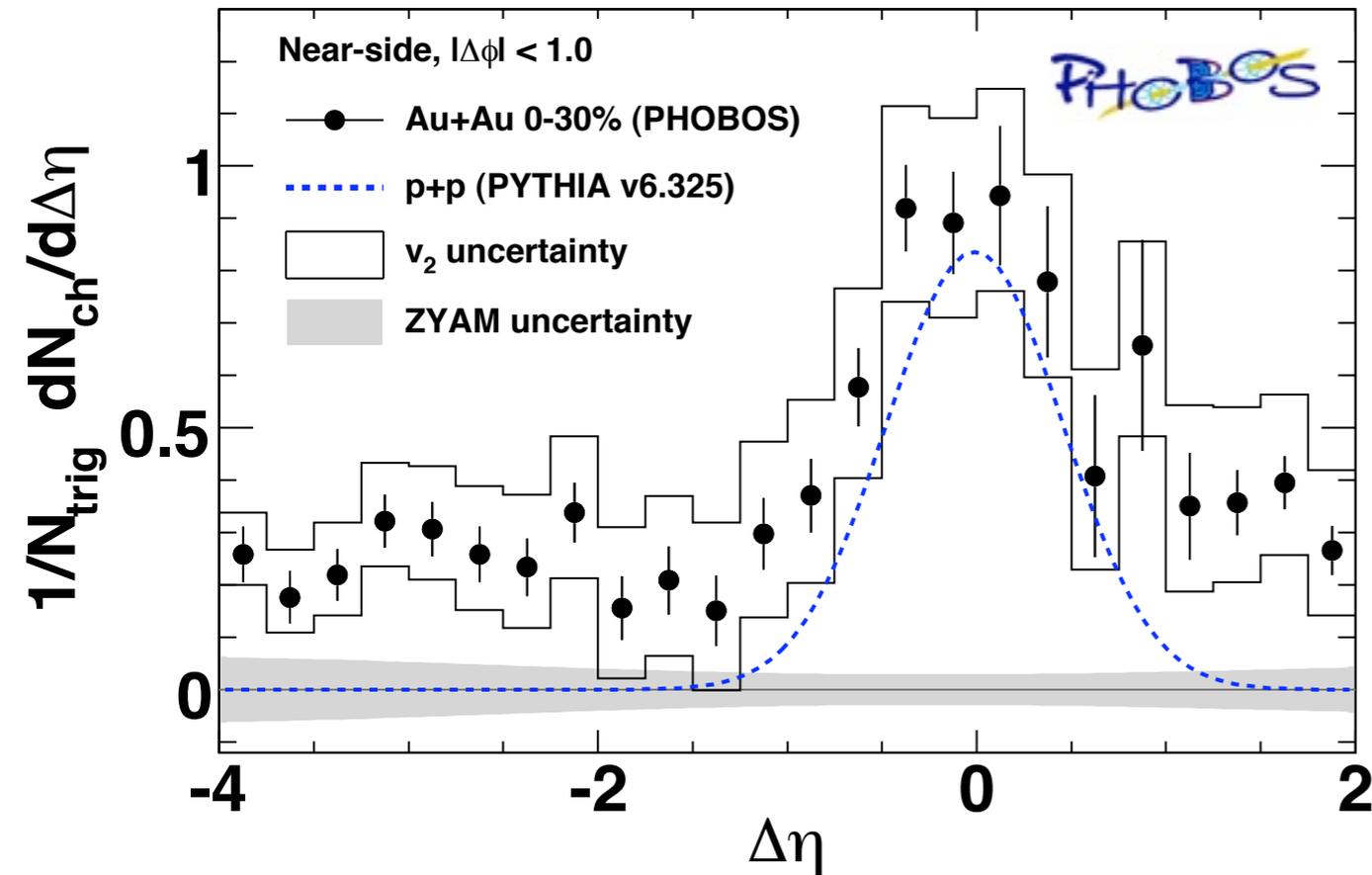
A surprising feature not predicted by anyone:
extended longitudinal correlations, with excess reflecting
properties of the “thermal” bulk

Long-range modifications



PHOBOS has measured a large-acceptance correlation function for $p_T(\text{trigger}) > 2.5 \text{ GeV}$ (PYTHIA used for p+p reference, confirmed by STAR data)

“The Ridge”

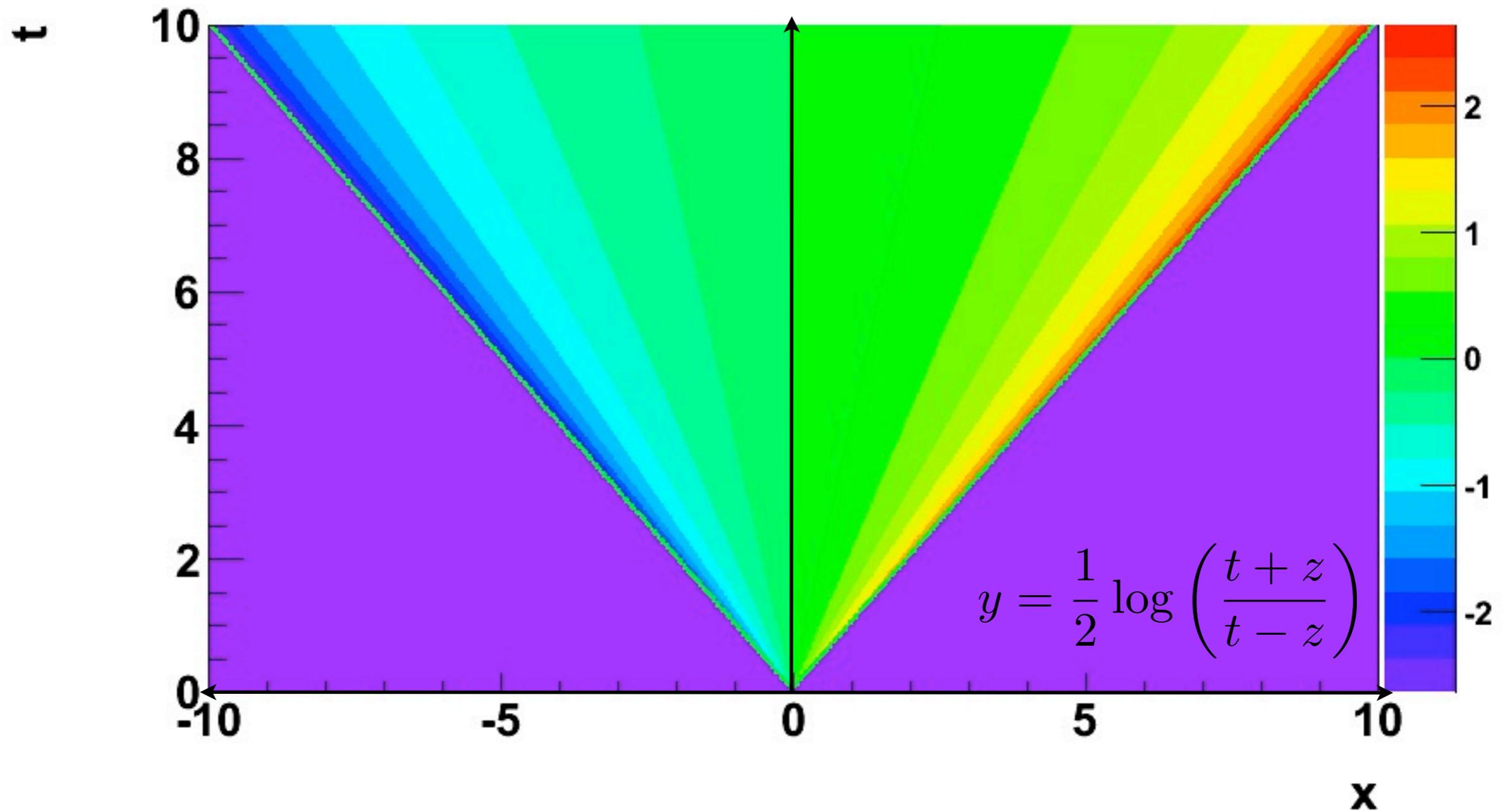


Substantial correlated yield out to $\Delta\eta=4$
Jet yield on near side seems to sit “on top” of the “ridge”

The Problem of the Ridge

- **Coupling of induced radiation to longitudinal flow**
Armesto et al., PRL 93, 242301
- **Recombination of shower + thermal partons**
Hwa, arXiv:nucl-th/0609017v1
- **Anisotropic plasma**
Romatschke, PRC 75, 014901
- **Turbulent color fields**
Shuryak, arXiv:0706.3531v1
- **Bremsstrahlung + transverse flow + jet-quenching**
Majumder, Muller, Bass, arXiv:hep-ph/0611135v2
- **Splashback from away-side shock**
Pantuev, arXiv:0710.1882v1
- **Momentum kick imparted on medium partons**
Wong, arXiv:0707.2385v2
- **Glasma Flux Tubes**
Dumitru, Gelis, McLerran, Venugopalan, arXiv:0804.3858; Gavin, McLerran, Moscelli, arXiv:0806.4718

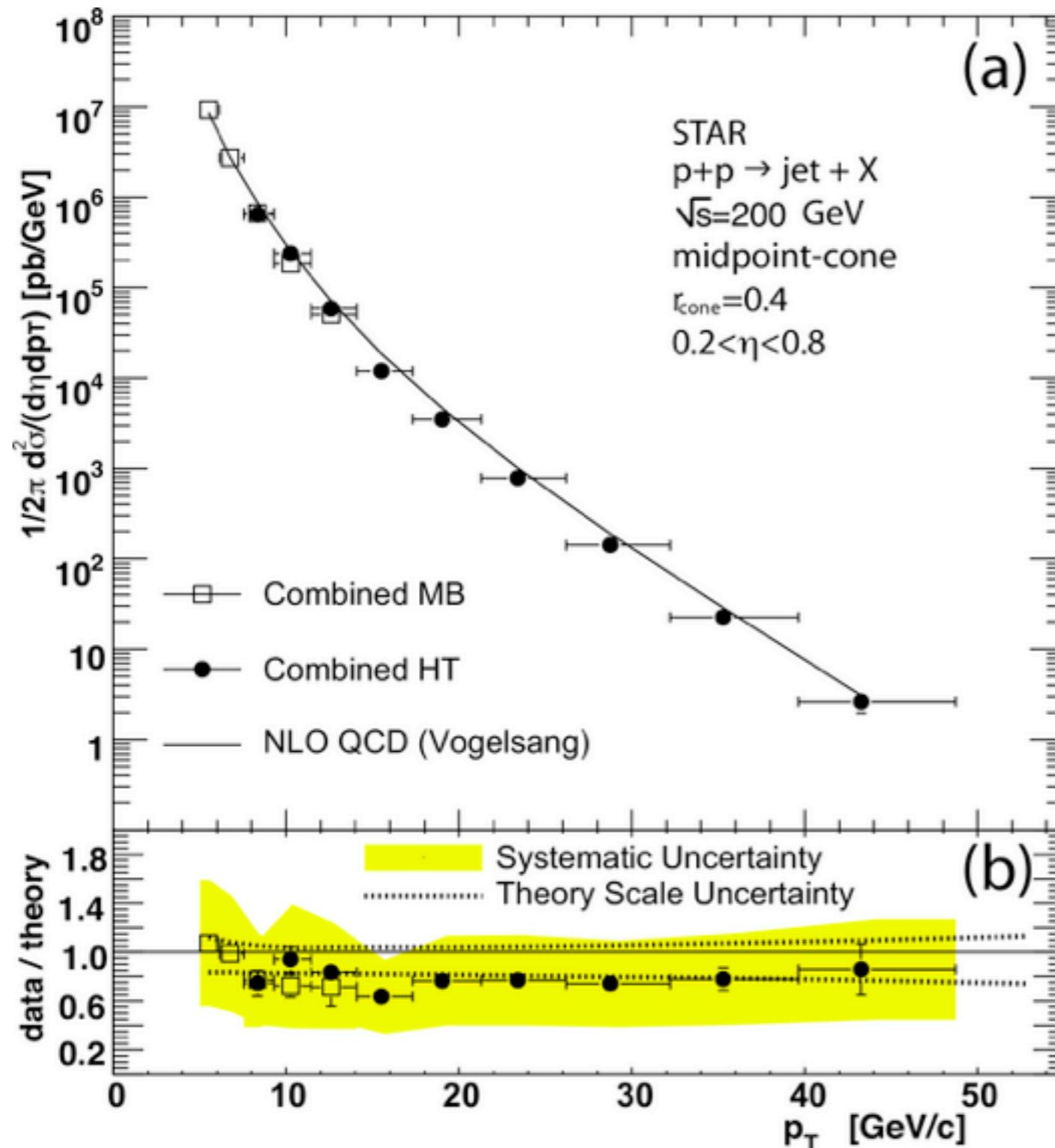
The Problem of the Ridge



Jumping rapidity bins requires large spatial correlation,
unless takes place at early time

Jets @ RHIC

CTEQ



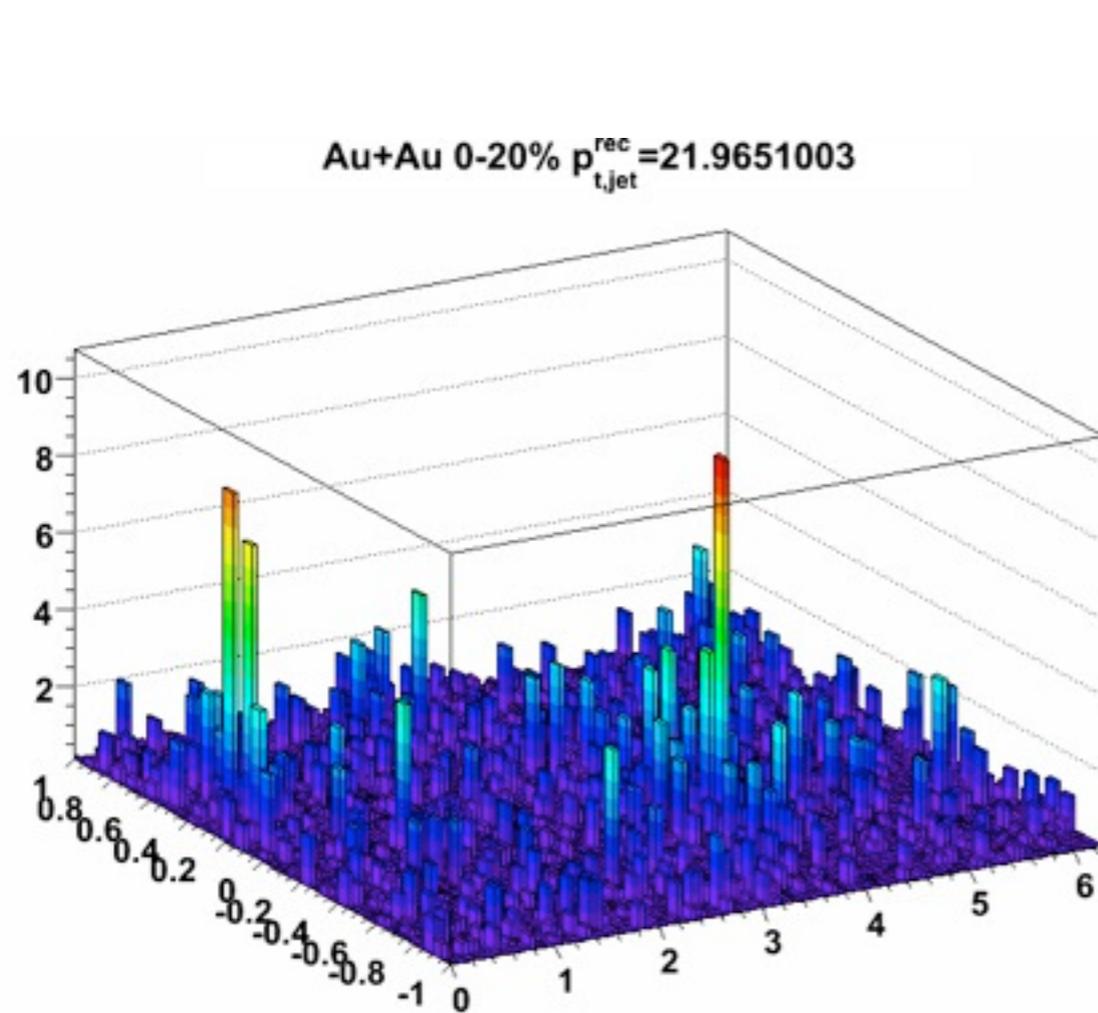
Midpoint cone $R=0.4$
- neutral energy from EMC
- charged particles from TPC

Good agreement with
NLO pQCD

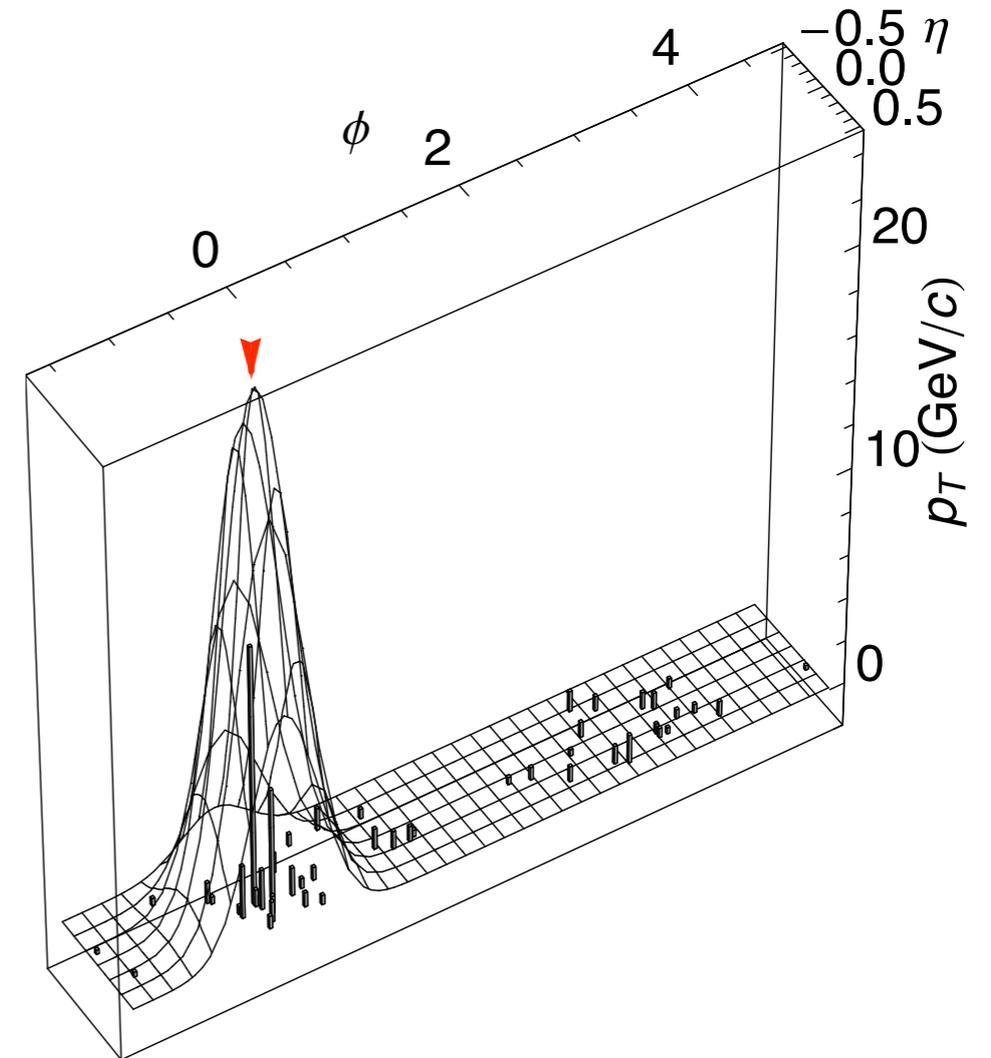
Experimental uncertainties
are large (50%)

Jets in Heavy Ions

CTEQ



STAR



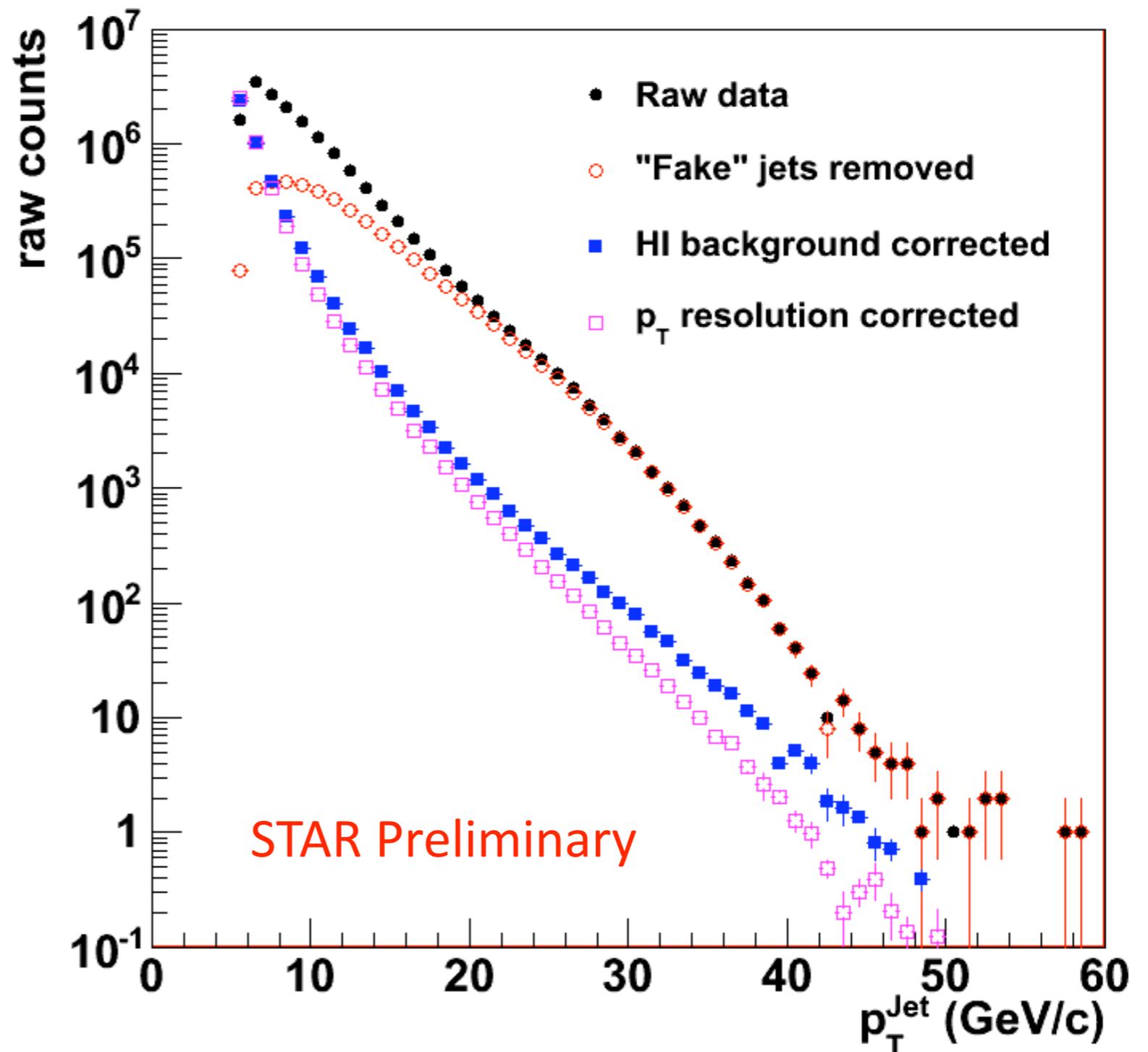
PHENIX

Despite enormous multiplicities, experiments @ RHIC are starting to reconstruct jets even in nuclear collisions

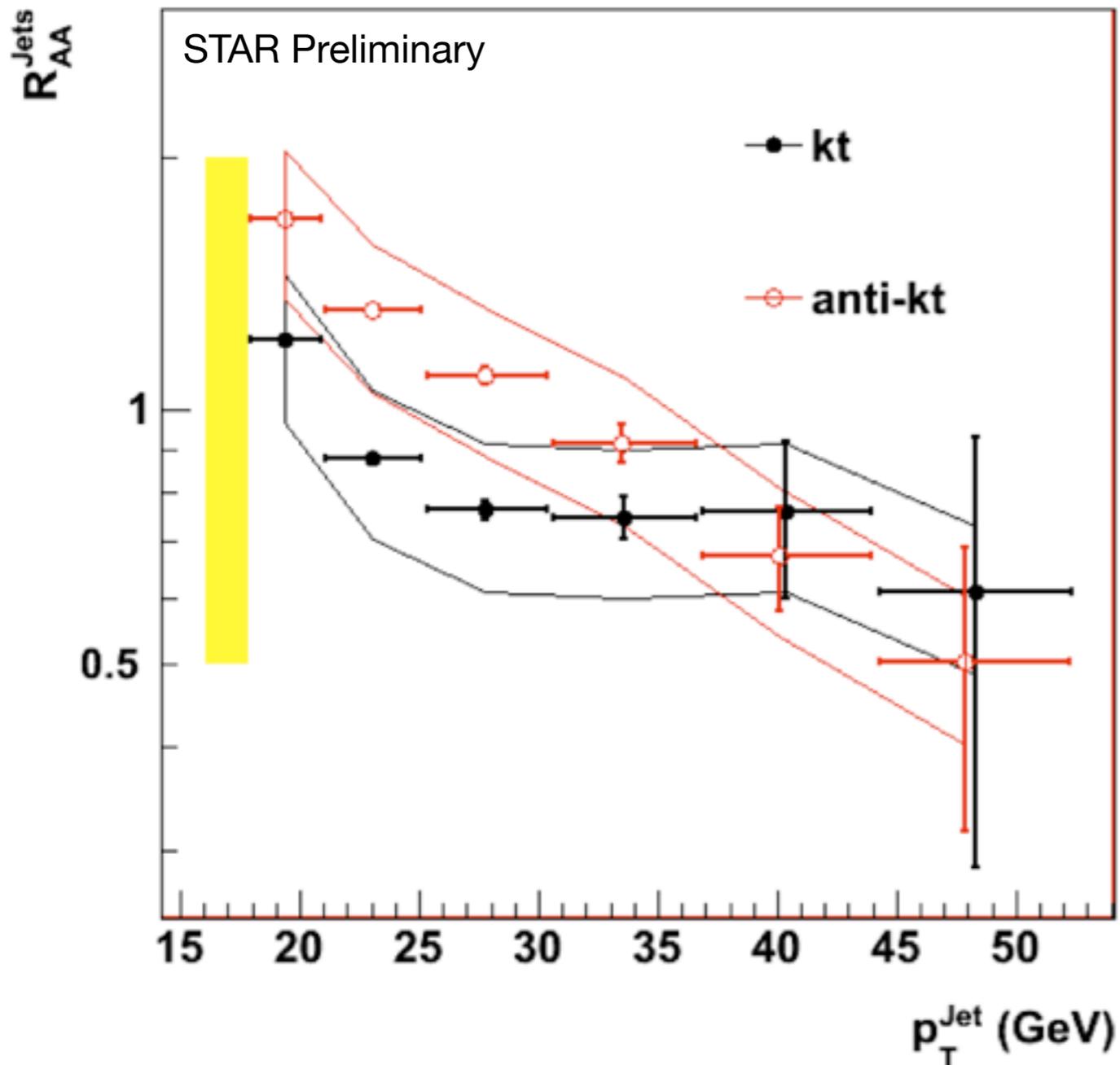
STAR jet spectrum extraction

Running k_T and anti- k_T algorithms, but with small cone sizes of $R=0.2, 0.4$

Need to reject fakes (look at “jets” at 90°), remove background (pedestal subtract) and unfold spectrum (correct for resolution)



Jet suppression at RHIC

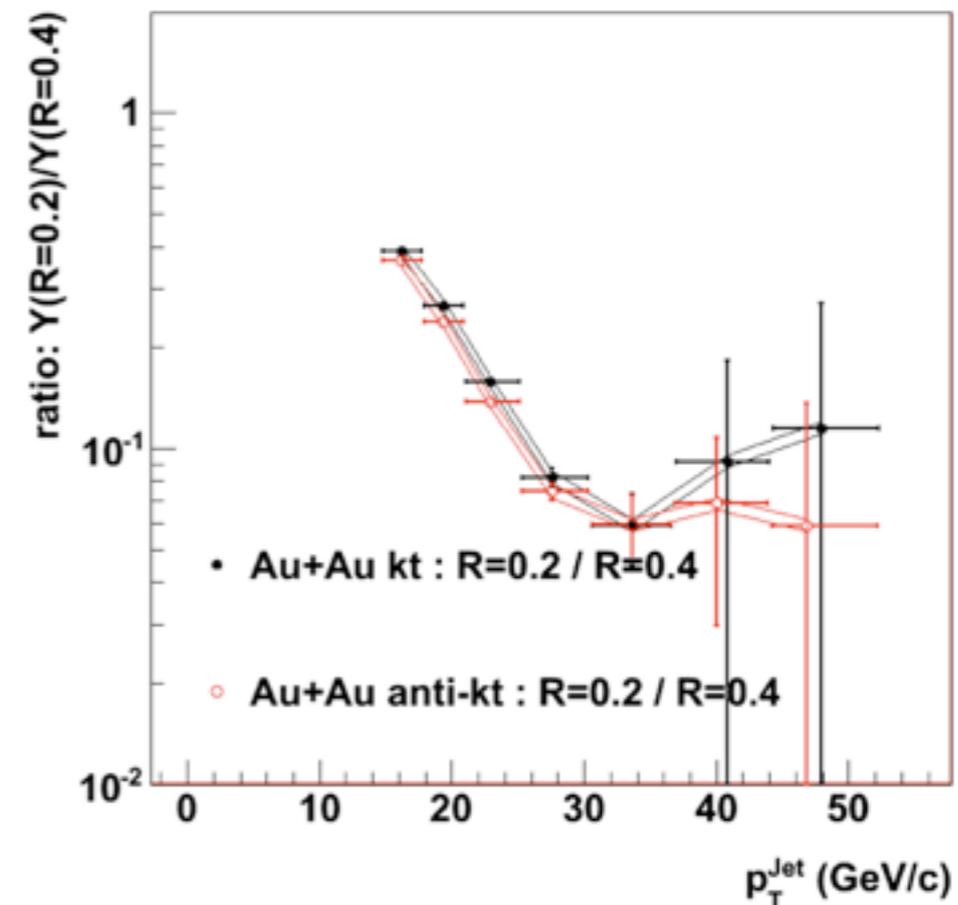
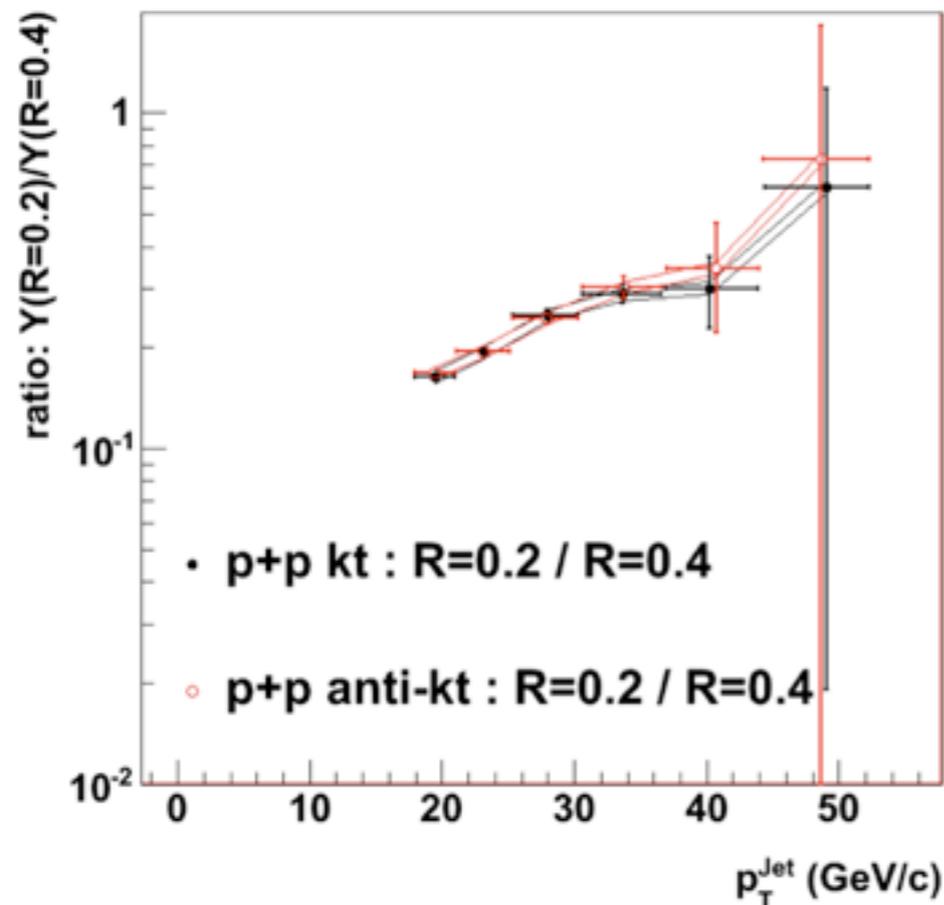


Strong algorithm dependence

No substantial suppression of jets themselves

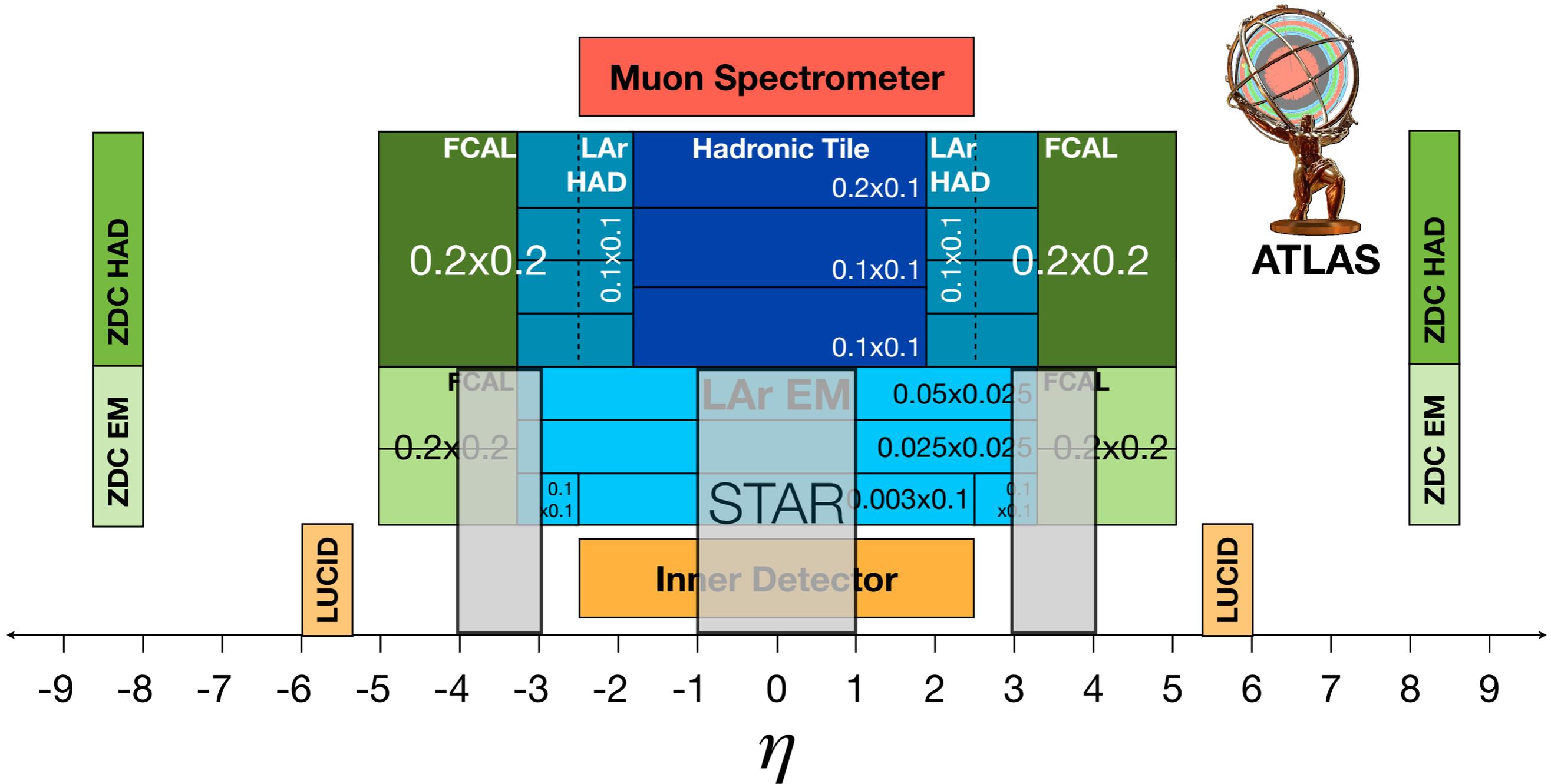
Suggestive that it is mainly the fragmentation that is modified

First measurements of jet shape



Comparison of jets with $R=0.2$ & 0.4 suggest that jets in Au+Au get broader with p_T (cf. p+p, where they apparently narrow)

The Future of Jets in HI (ATLAS case)



The LHC detectors provide unprecedented capability for high p_T

Z as a probe of nuclear PDFs

PDFs show “shadowing” at low- x

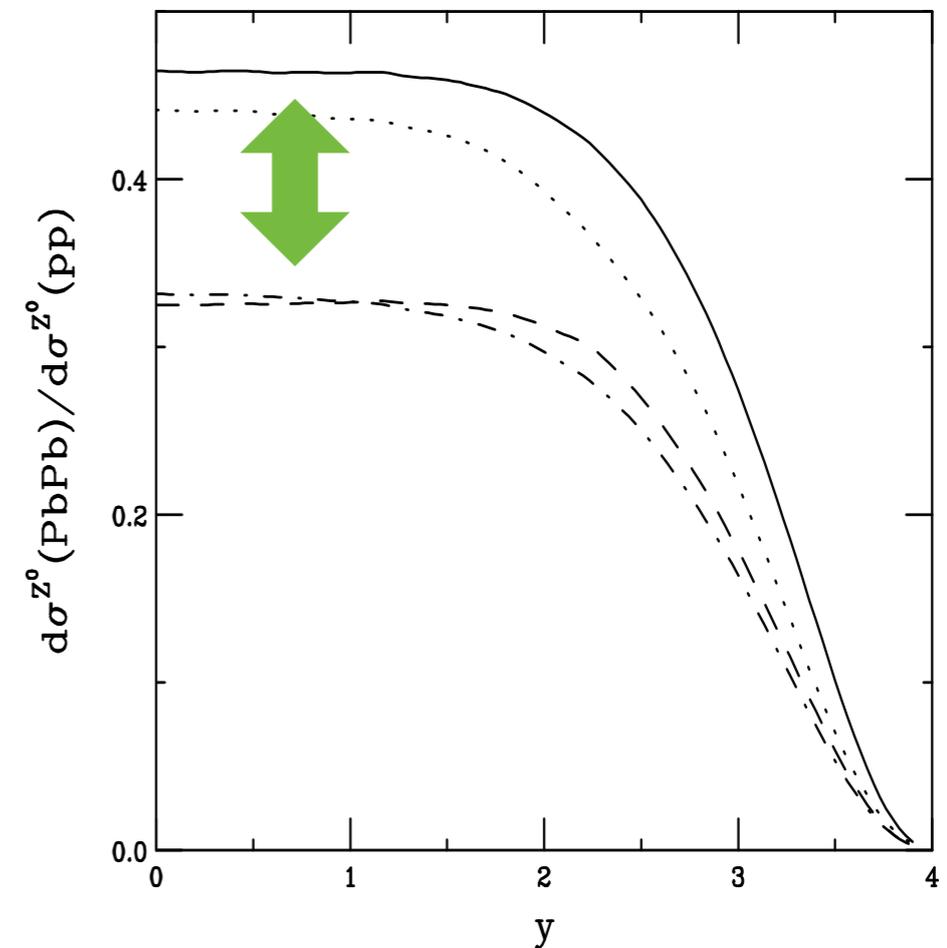
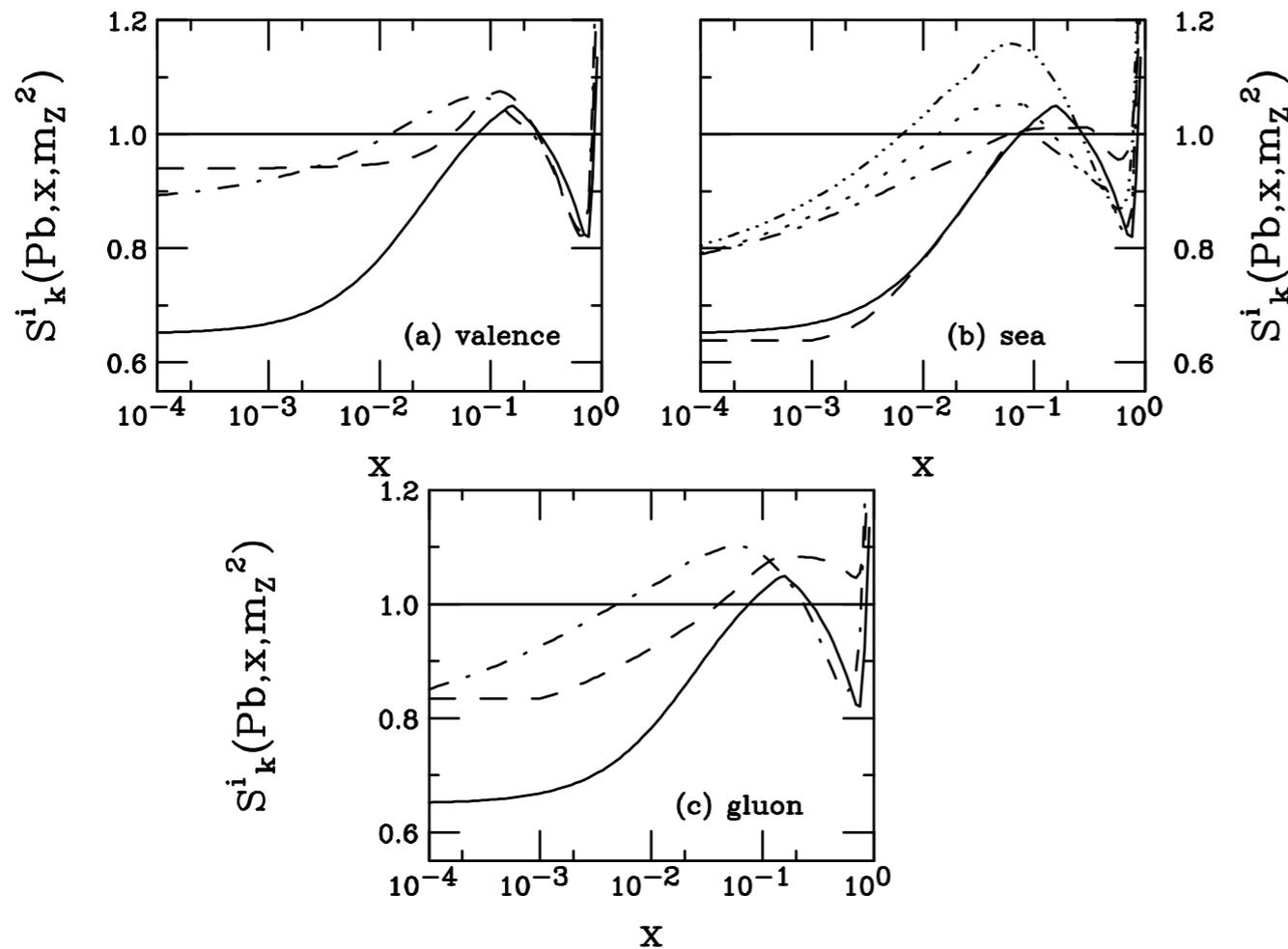


FIG. 10. The ratios of the Z^0 rapidity distributions in Pb+Pb collisions relative to pp collisions, calculated with the MRST HO distributions. The solid curve is the ratio without shadowing. The homogeneous shadowing results are given in the dashed, S_1 , dot-dashed, S_2 , and dotted, S_3 , lines.

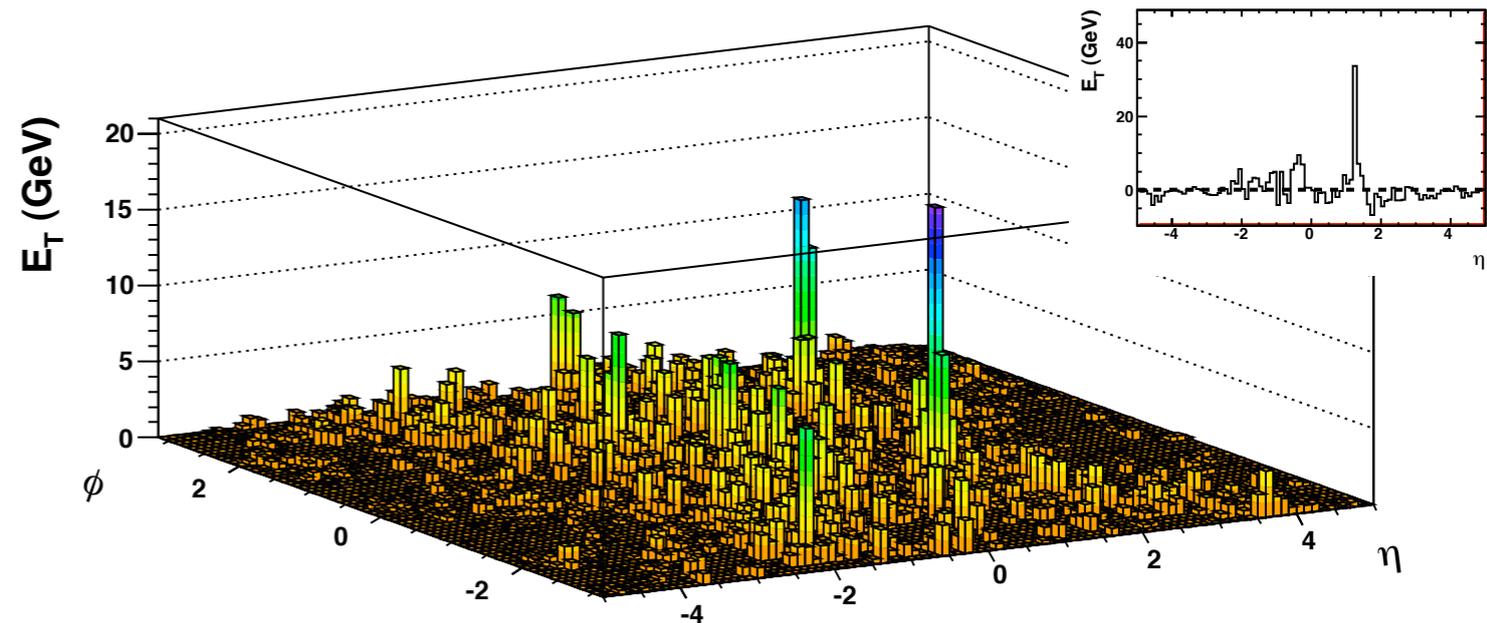
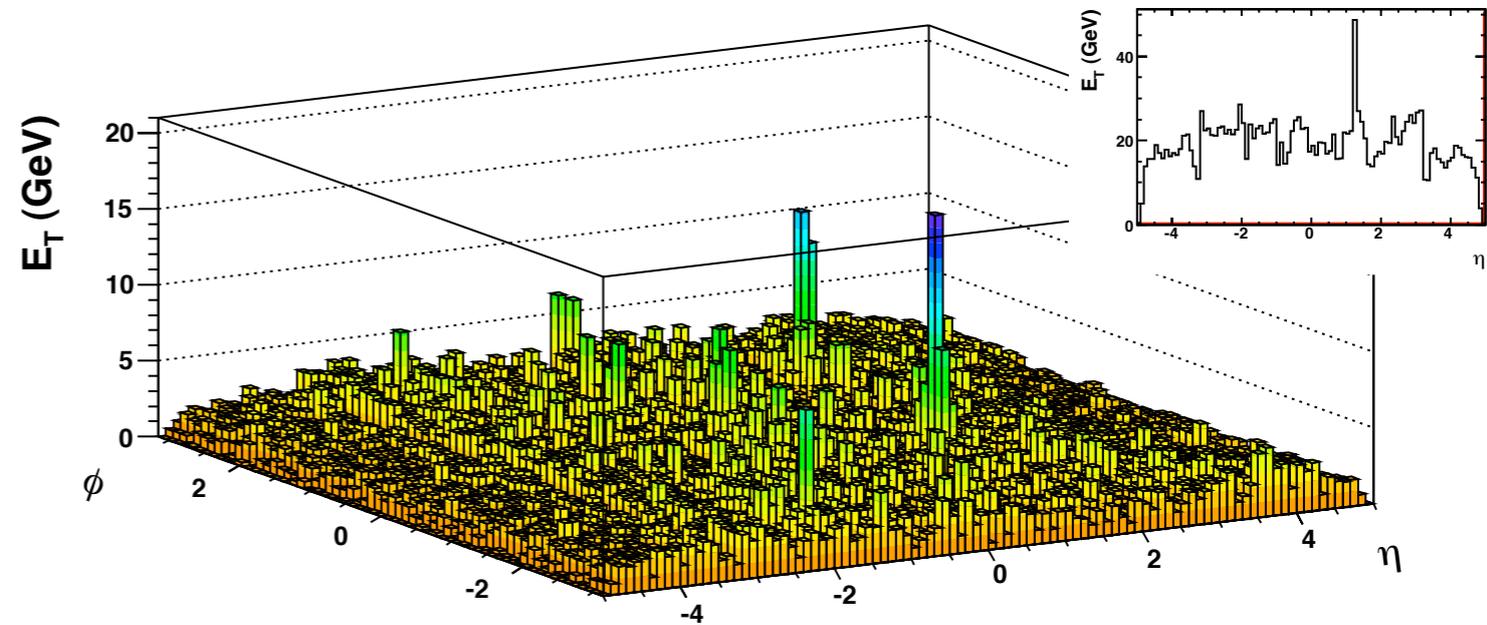
At the LHC, can use the production of Z bosons as a background-free probe of the very initial PDFs.

Expect 1000's of Z's each LHC year

Jet Finding in HI

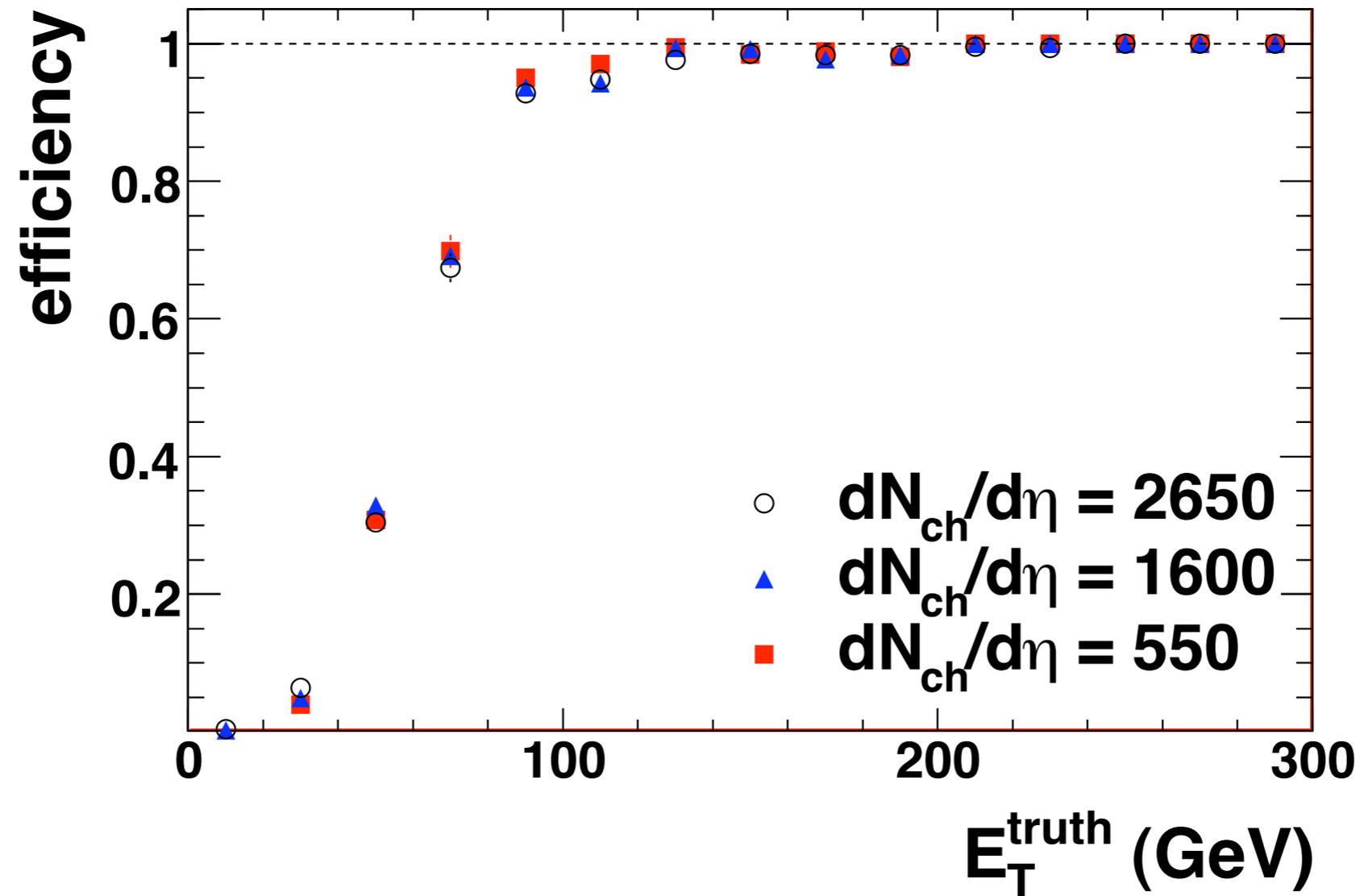
CTEQ

- Cells collected in to 0.1×0.1 Towers
- Seed finding used to exclude candidate jets
- Background calculated in Eta bins, subtracted layer by layer
- Seeded midpoint cone algorithm to find jets
- Calibrations applied after jet finding



Jet Efficiency

CTEQ

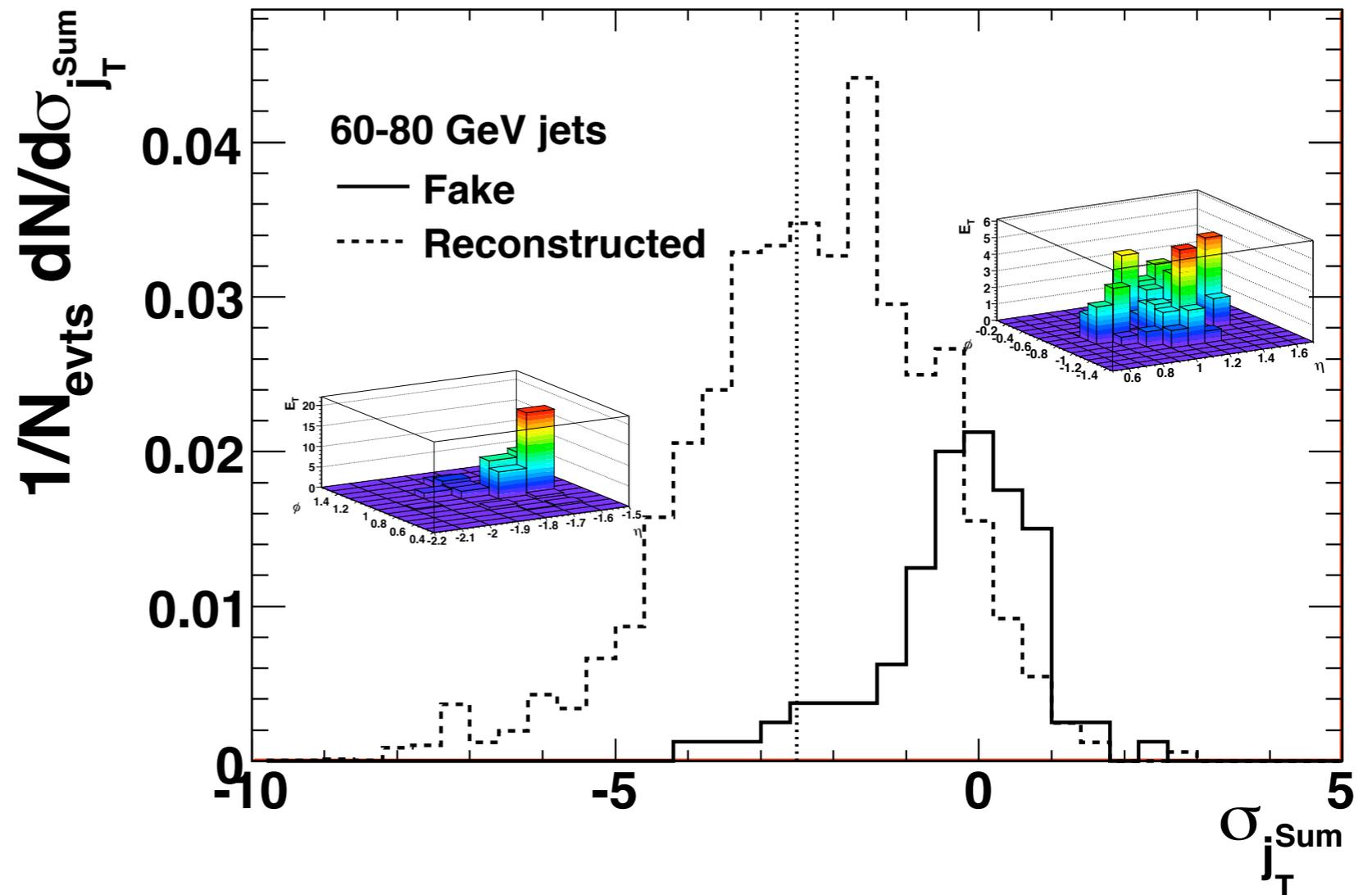


Efficiency turn-on curve reaches 100% at ~ 100 GeV
(it's "easy" above that energy)

Fake Jets

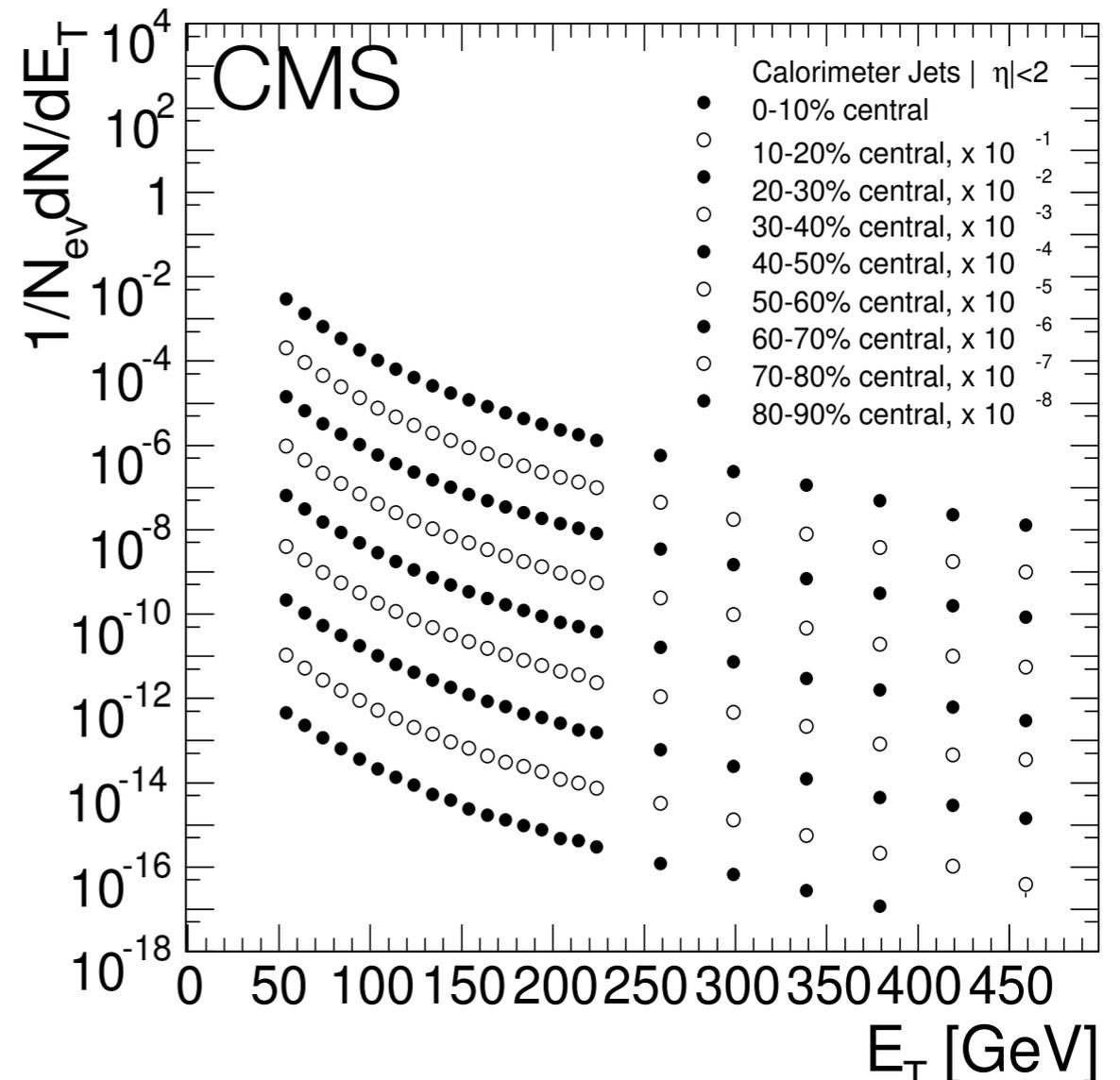
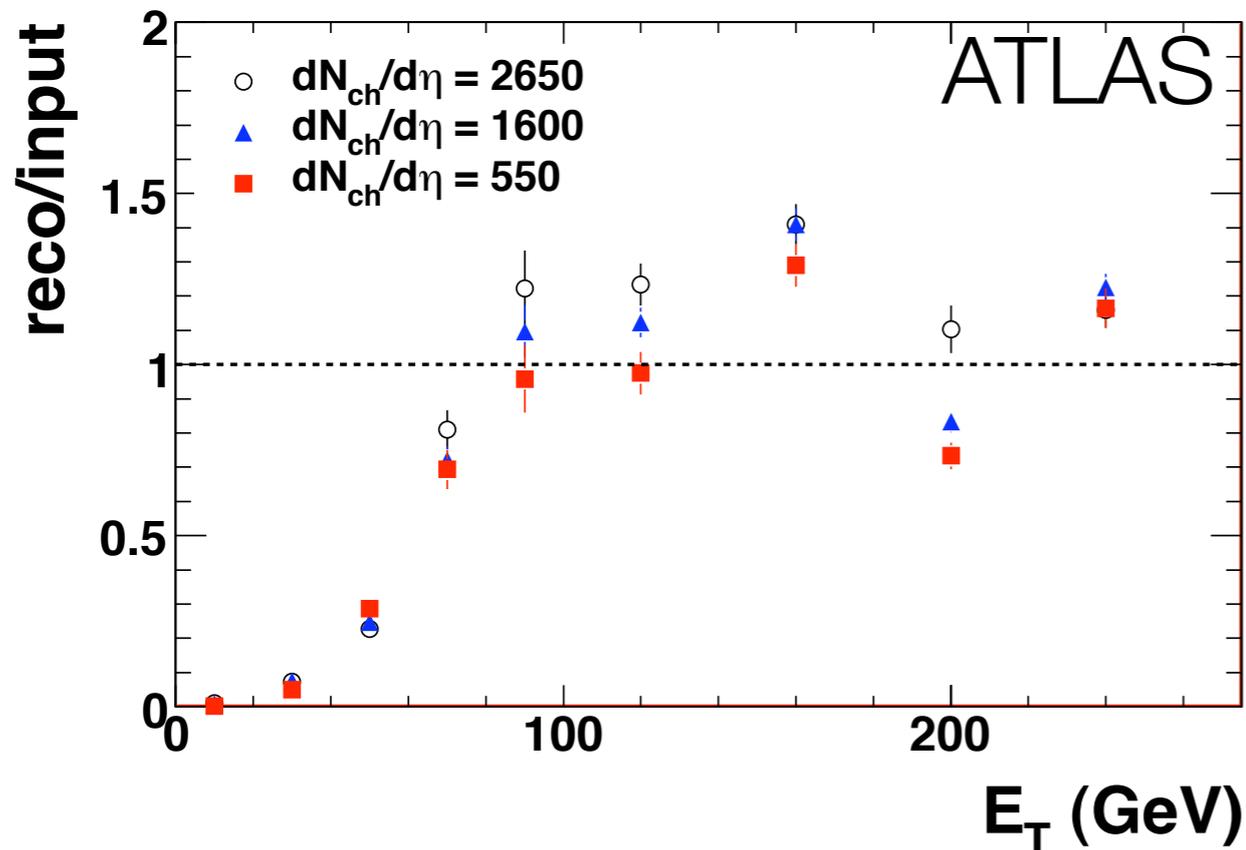
“Jets” triggered by a seed and supported by background have a different shape than normal jets

Experimented with a variety of estimators to cut against fakes, “SumJt” most effective



$$j_T^{Sum} = \sum_{cell} E_T^{cell} \sin R_{cell} \quad \sigma_{j_T^{Sum}} = \frac{j_T^{Sum}(E_T) - \langle j_T^{Sum} \rangle(E_T)}{\sigma(E_T)}$$

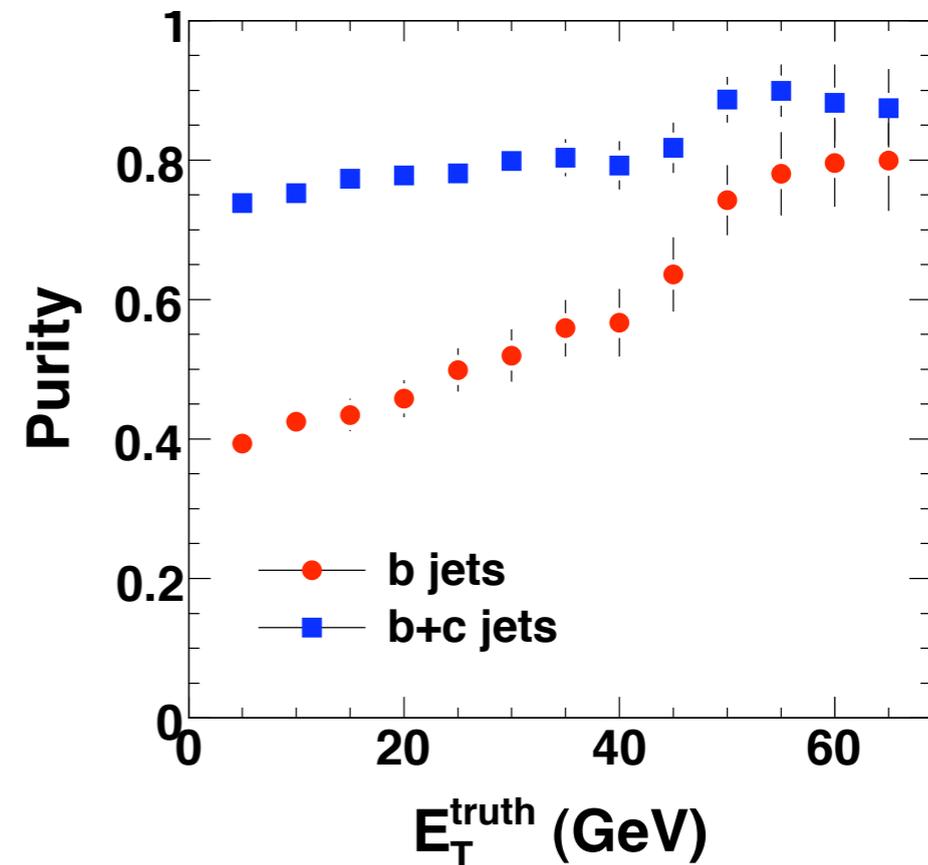
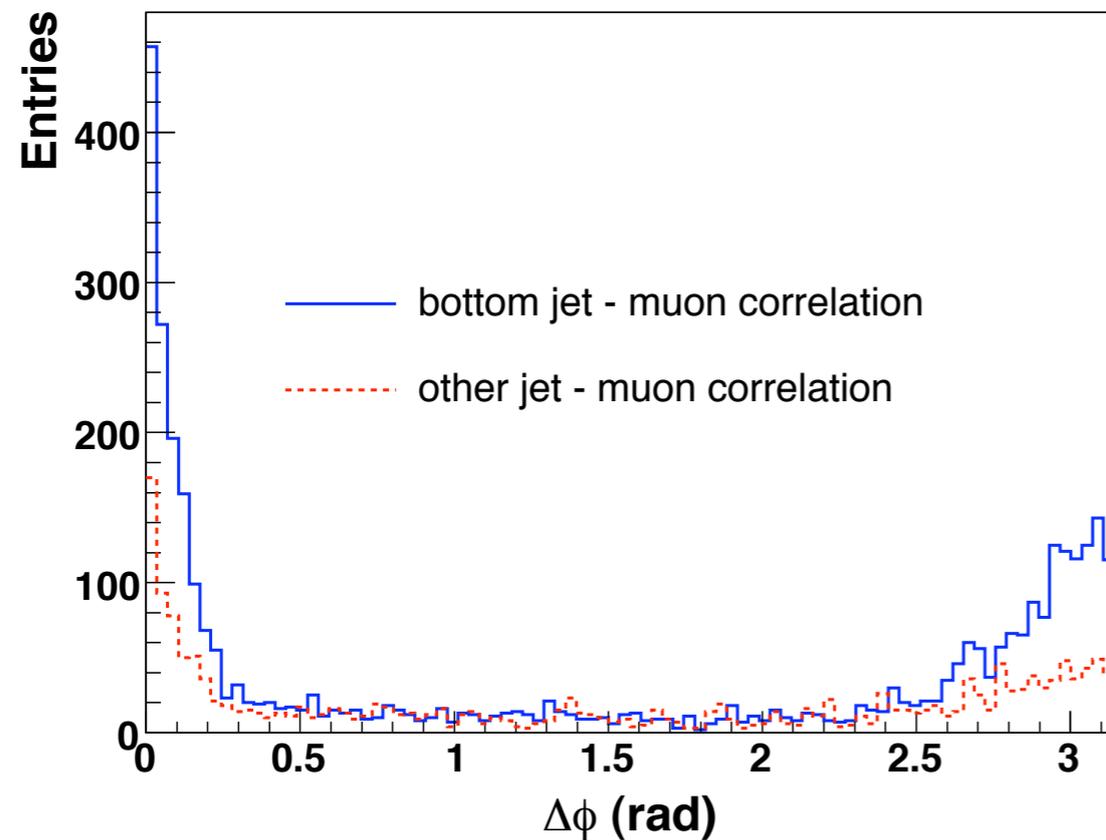
Jet reconstruction & expected rates



Straghtforward to achieve 20% reconstruction of jet spectrum.

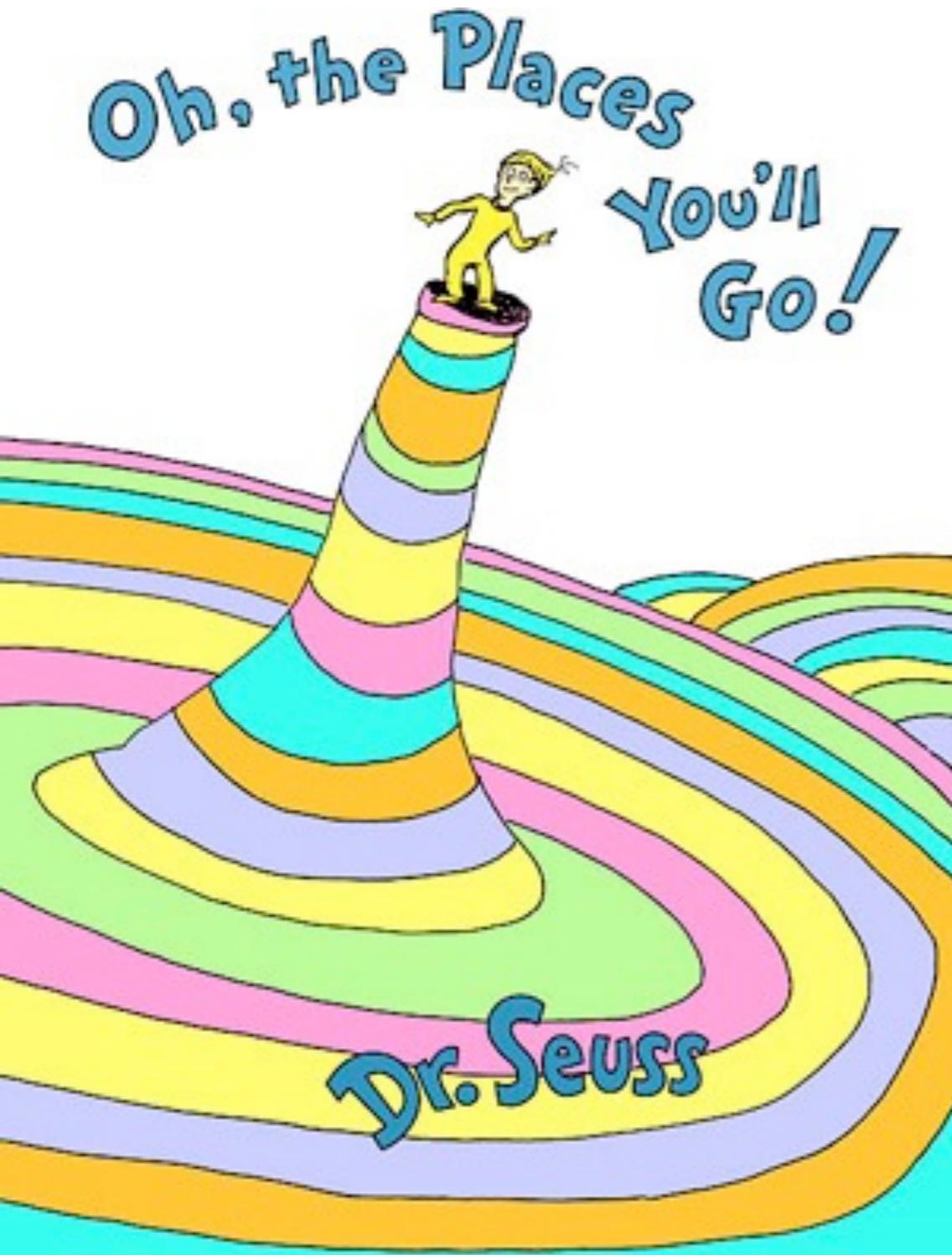
LHC Luminosity should provide out to 400 GeV (2xRHIC \sqrt{s} !)

Heavy Flavor Jets

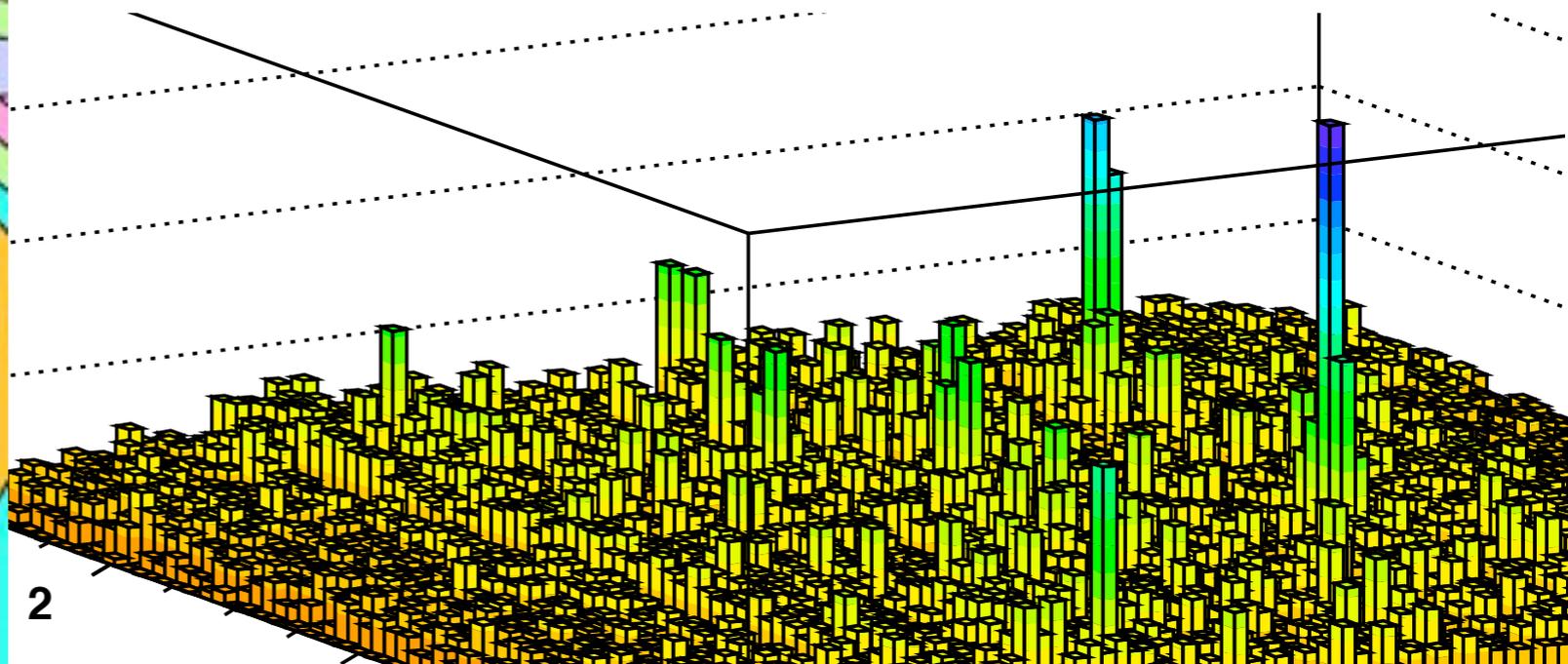


Requiring a muon to be correlated with a jet enhances the contribution from b-quarks: can look for HQ modifications to jet energy loss!

Oh, the places we'll go with jets @ LHC



- **Jet rates**
 - Singles (RAA)
 - Correlations (dijet broadening)
- **Fragmentation functions**
 - Direct measurement of modified FF's
- **Flavor tagging**
 - s,c,b can be tagged via leading hadrons
- **γ +Jet, Z+jet**



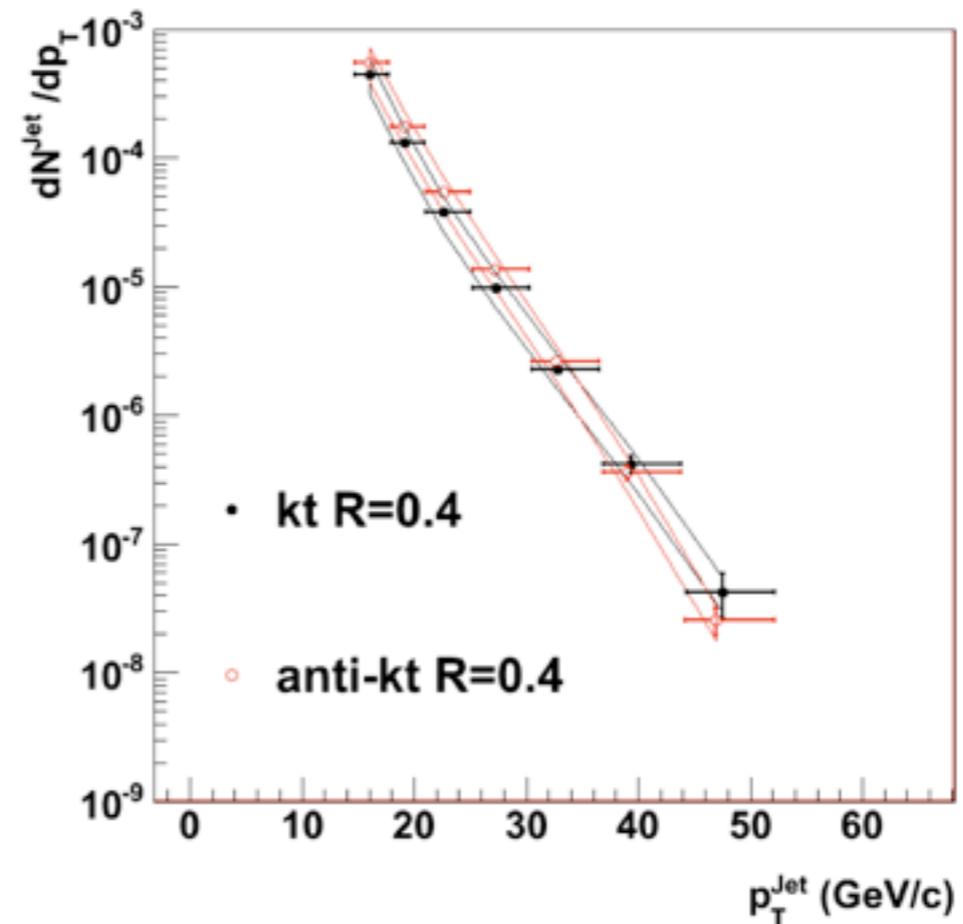
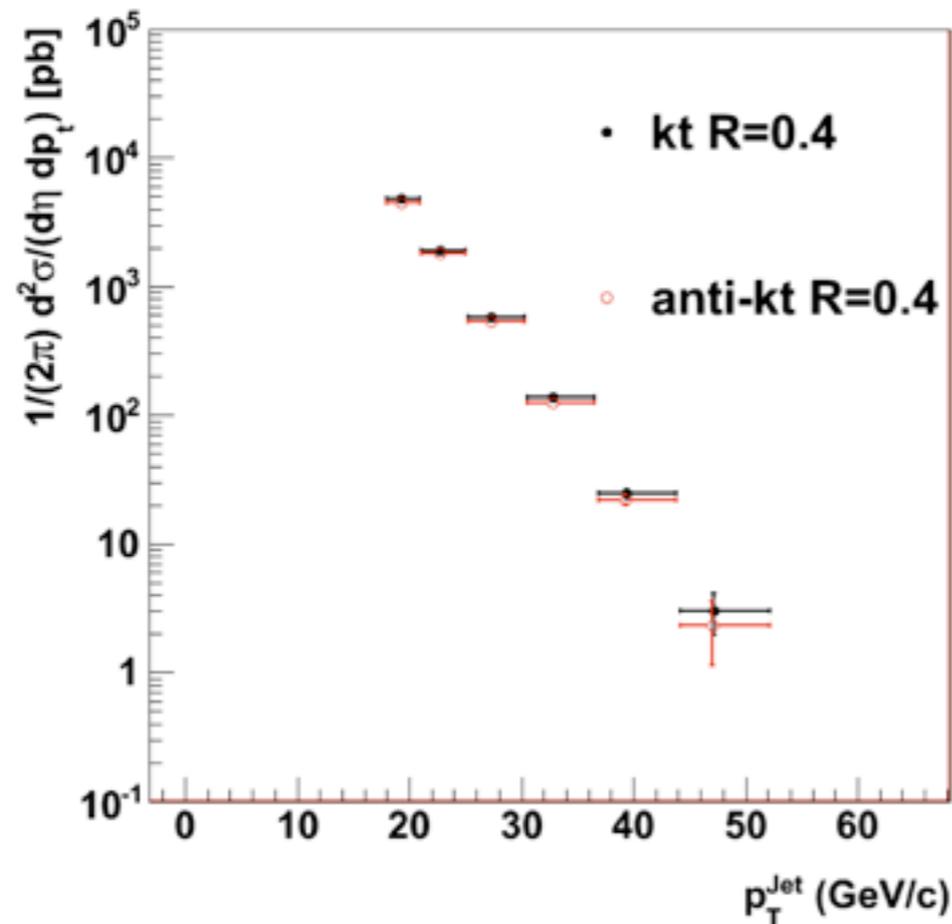
“Conclusions” (or, a new beginning!)

- **RHIC program has been incredibly productive**
 - Discovery of what seems to be a new strongly-coupled state of matter!
- **The future is particularly bright, especially for high p_T**
 - RHIC II is on the way: increased luminosity
 - LHC is on the way: increased energy of hard probes
 - pQCD technology will be essential for progress here
- **Opportunities in “soft” QCD**
 - Interesting phenomena not necessarily treatable by pQCD
 - Entropy production, cone, ridge, etc. require understanding of both soft sector and hard sector
 - Even jet quenching requires detailed knowledge of initial conditions etc. & viscous hydro evolution...

Nuclear modification factor

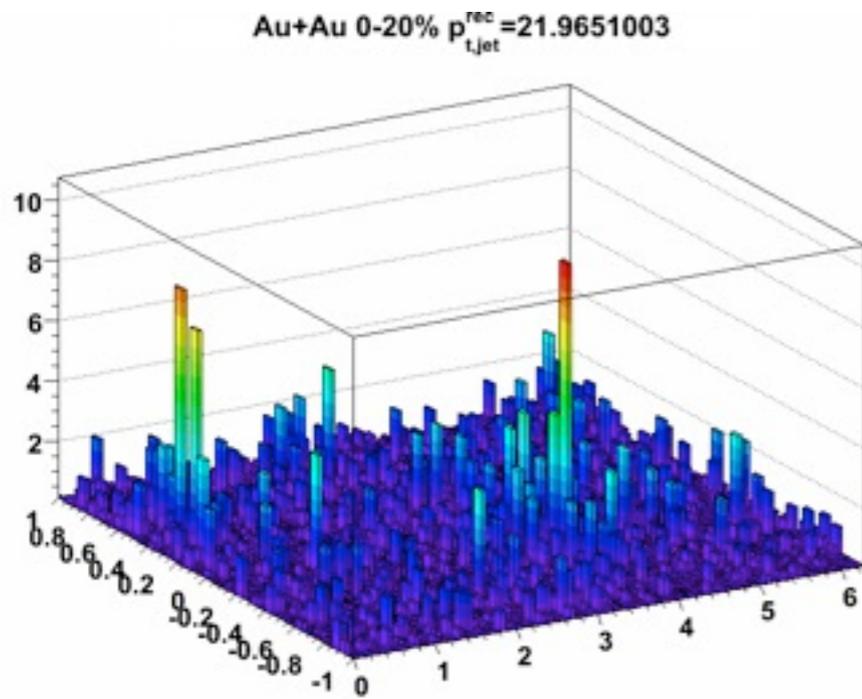
$$\begin{aligned} R_{A+A} &= \frac{N_{A+A}}{\langle T_{A+A} \rangle \sigma_{p+p}} \\ &= \frac{d^2 N_{A+A} / dp_T dy}{\langle T_{A+A} \rangle d^2 \sigma_{p+p} / dp_T dy} \\ &= \frac{1}{\langle T_{A+A} \rangle \sigma_{p+p}} \frac{d^2 N_{A+A} / dp_T dy}{d^2 N_{p+p} / dp_T dy} \\ &= \frac{1}{N_{coll}} \frac{d^2 N_{A+A} / dp_T dy}{d^2 N_{p+p} / dp_T dy} \end{aligned}$$

Results for p+p & Au+Au

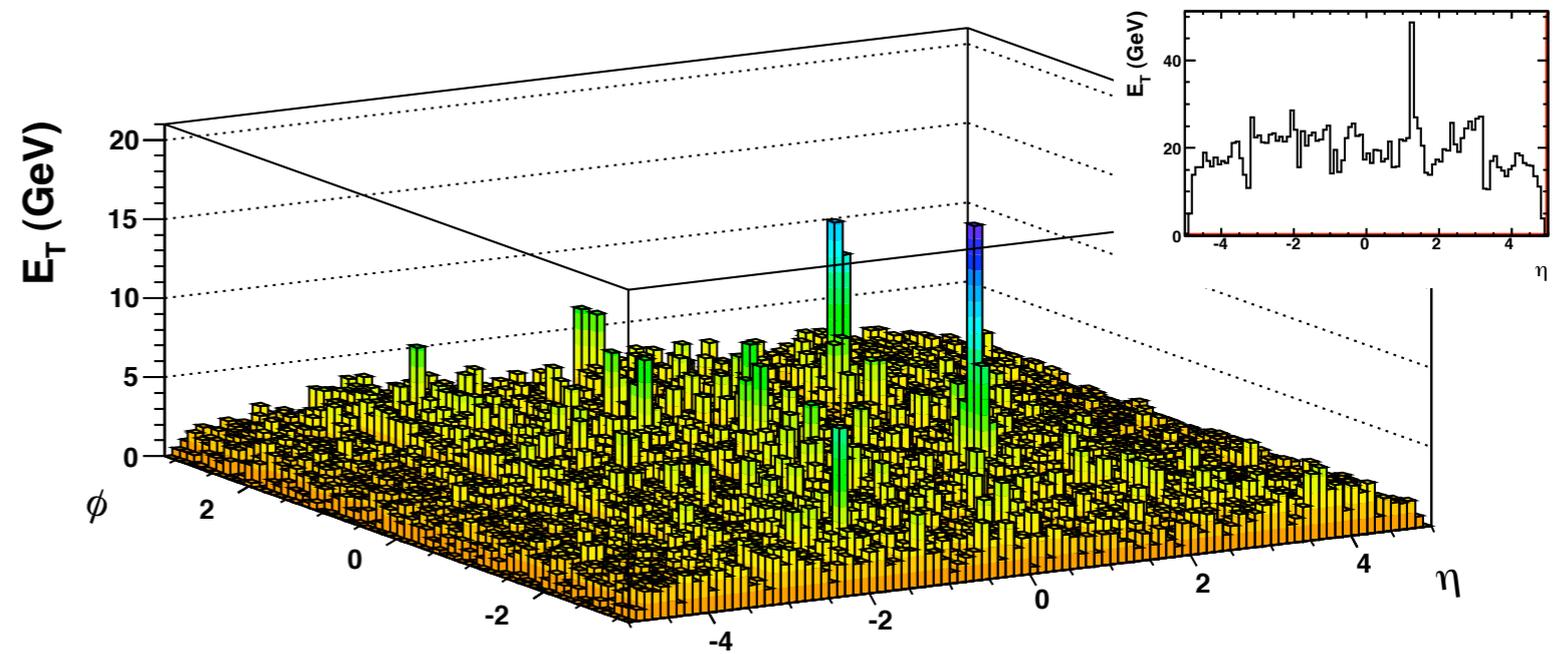


Spectral shape apparently quite similar between p+p & Au+Au for k_T and anti- k_T

Increased acceptance at the LHC



$$|\eta| < 1$$



$$|\eta| < 5$$