

Top Quark Production and Decay (Theory)

CTEQ 2002

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I. Motivation

II. Top strong interactions

III. Top weak interactions

I. Motivation

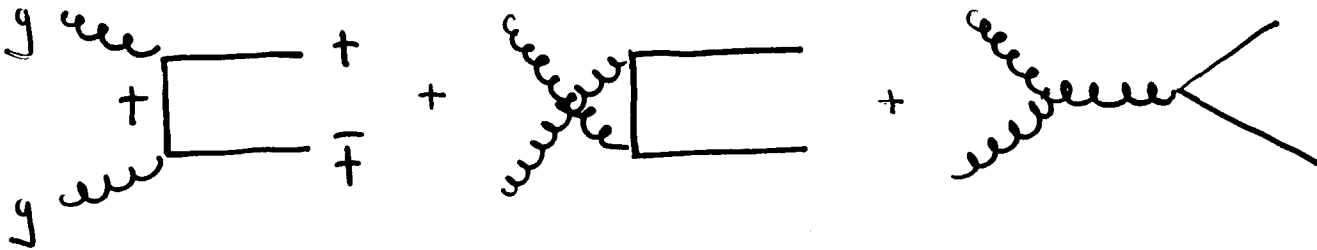
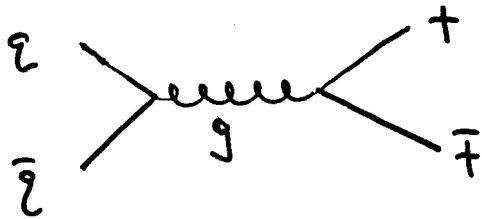
Why study the top quark?

1. Is the top quark just an ordinary quark, or is it exotic in some way?
2. Top may be useful to discover new particles (e.g., $p\bar{p} \rightarrow t\bar{t}h$)
3. Top is a background to almost everything else (e.g. $h \rightarrow w^+w^-$)

Strategy: get to know the top quark by measuring everything about it

II. Top strong interactions

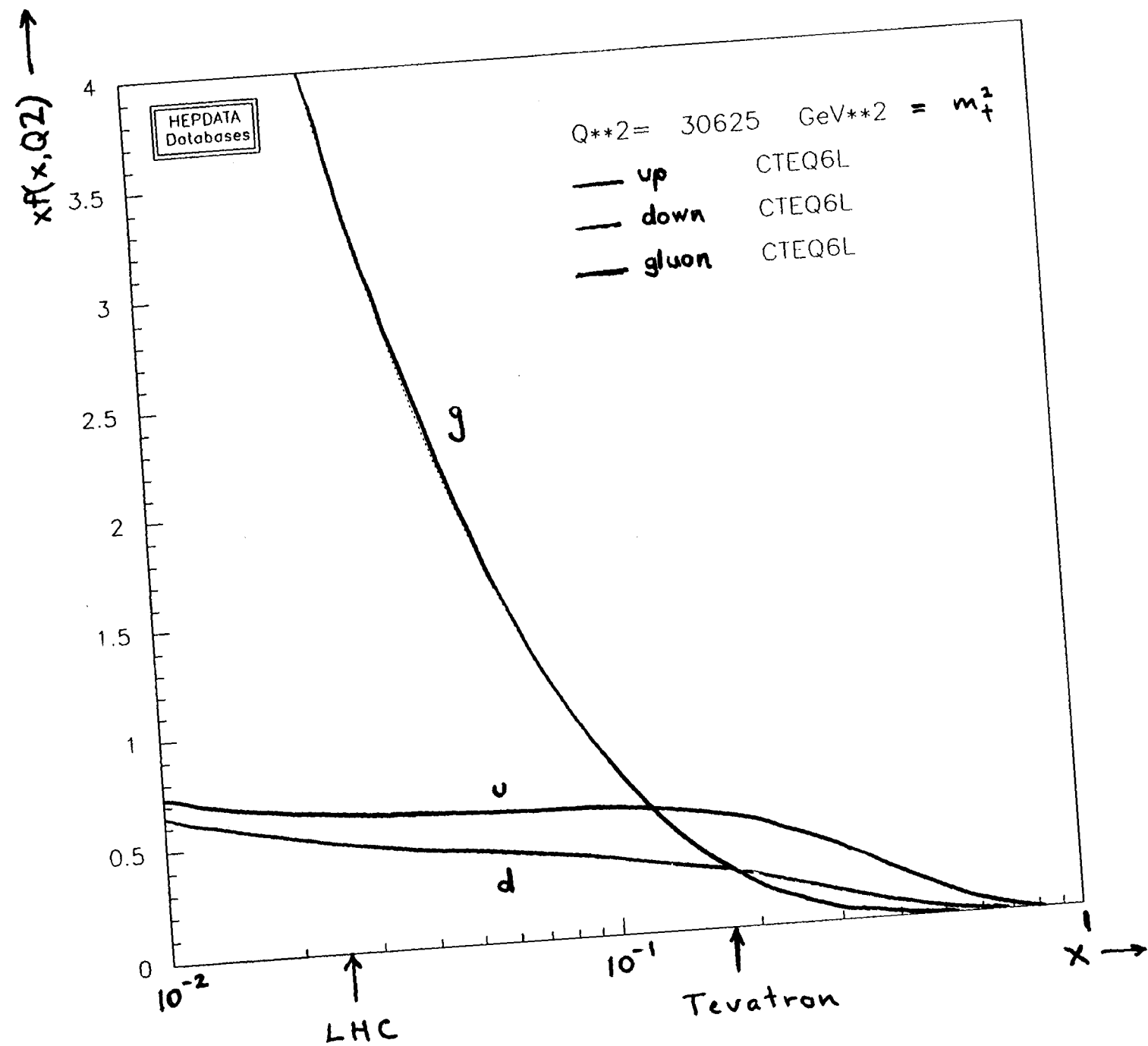
Top - quark production:



Predicted cross sections (NLO QCD) BCMN

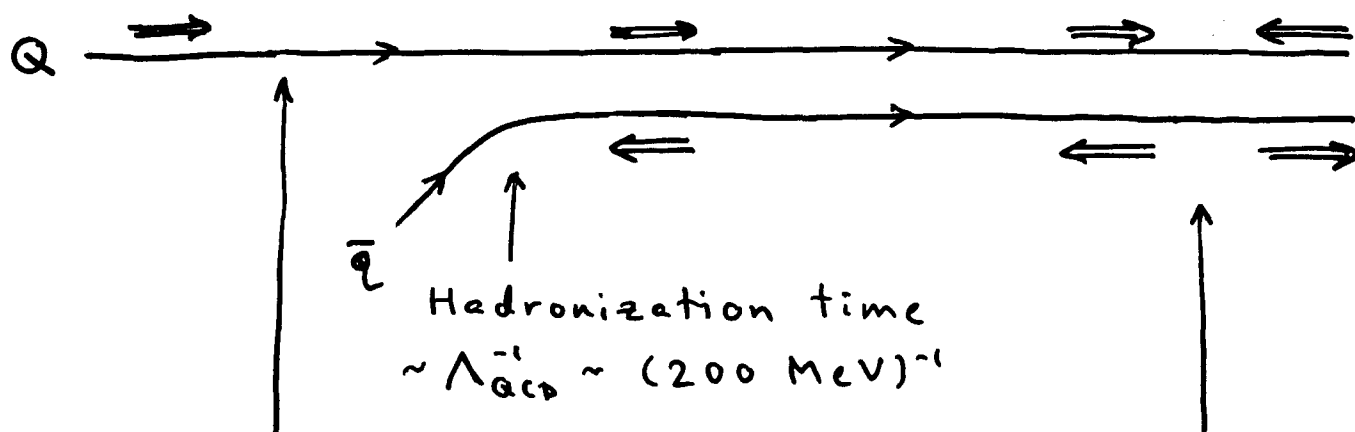
($m_t = 175 \text{ GeV}$)

	$\sigma_{\text{NLO}} (\text{pb})$	$q\bar{q} \rightarrow t\bar{t}$	$gg \rightarrow t\bar{t}$
Tevatron ($\sqrt{s} = 1.8 \text{ TeV } p\bar{p}$)	$4.87 \pm 10\%$	90%	10%
Tevatron ($\sqrt{s} = 2.0 \text{ TeV } p\bar{p}$)	$6.70 \pm 10\%$	85%	15%
LHC ($\sqrt{s} = 14 \text{ TeV } pp$)	$803 \pm 15\%$	10%	90%



typical $x \sim \frac{2 m_t}{\sqrt{s}}$

The top-quark spin is observable



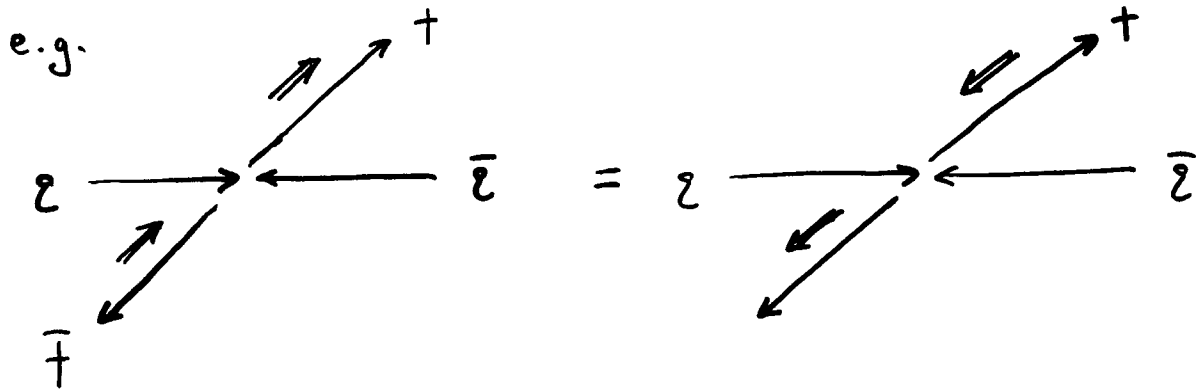
Top lifetime
 $= \Gamma_t^{-1} \approx (1.5 \text{ GeV})^{-1}$

Depolarization time
 $\sim \left(\frac{\Lambda_{QCD}^2}{m_t} \right)^{-1} \sim (1 \text{ MeV})^{-1}$

Top decays well before there is time
 to depolarize its spin.

BOKKZ

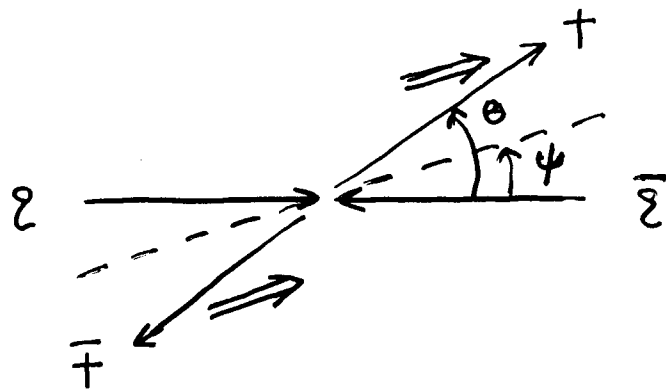
Unfortunately, QCD produces unpolarized top quarks (due to parity conservation):



Fortunately, t and $\bar{t}$ spins are correlated

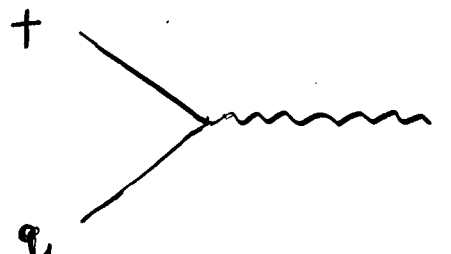
The correlation is 100% for $q\bar{q} \rightarrow t\bar{t}$

when a specific spin basis is chosen:



$$\tan \psi = \frac{\beta^2 \sin \theta \cos \theta}{1 - \beta^2 \sin^2 \theta}$$

III. Top weak interaction

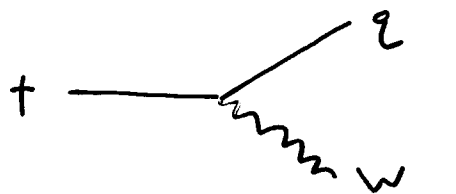


$$= -i \frac{g}{2\sqrt{2}} \gamma^\mu (1 - \gamma_5) V_{tq}$$

$\uparrow$
 CKM matrix
 element

$q = d, s, b$

Top-quark decay:



$$\propto |V_{tq}|^2$$

CDF has measured

$$\frac{\text{BR}(t \rightarrow Wb)}{\text{BR}(t \rightarrow Wq)} = 0.99 \pm 0.29$$

$$\parallel$$

$$\frac{|V_{tb}|^2}{|V_{td}|^2 + |V_{ts}|^2 + |V_{tb}|^2}$$

This implies $|V_{tb}| \gg |V_{td}|, |V_{ts}|$

If you assume 3 generations, then

CKM unitarity implies

$$|V_{td}|^2 + |V_{ts}|^2 + |V_{tb}|^2 = 1$$

Thus

$$|V_{tb}| = 0.99 \pm 0.15 \quad (3 \text{ gen})$$

Consistent with

$$V_{ij} = \begin{pmatrix} V_{ud} & V_{us} & 0.002 - 0.005 \\ V_{cd} & V_{cs} & 0.037 - 0.043 \\ V_{td} & V_{ts} & 0.9990 - 0.9993 \end{pmatrix} \begin{matrix} \leftarrow V_{ub} \\ \leftarrow V_{cb} \\ \leftarrow V_{tb} \end{matrix}$$

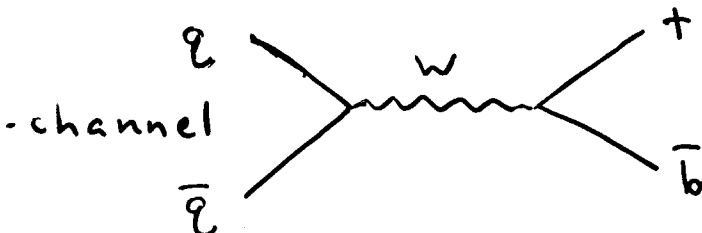
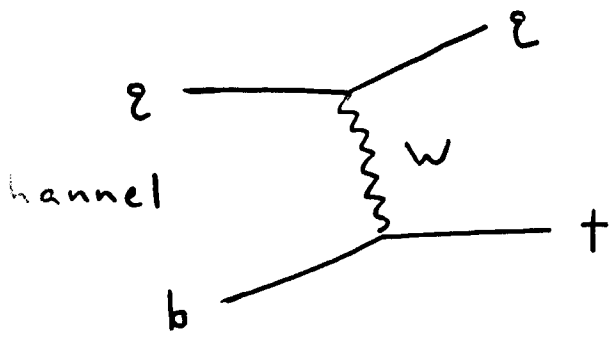
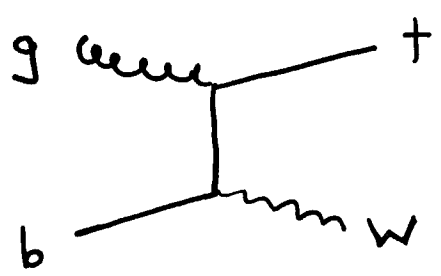
(3 gen)



third column

Top-quark production via the weak interaction (Single top production)

measures $|V_{tb}|$ directly: $\sigma \propto |V_{tb}|^2$

	<u>Tevatron</u>	<u>LHC</u>
 s-channel	0.90	10.6
 t-channel	2.1	250
 gluon fusion	0.1	75

These are predicted cross sections (pb)

(NLO QCD)

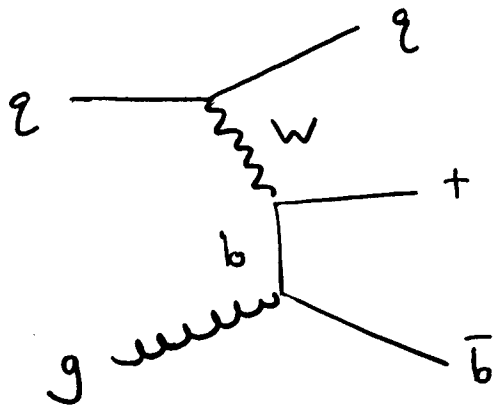
SSW

Compare $\sigma(t\bar{t})$

6.7

800

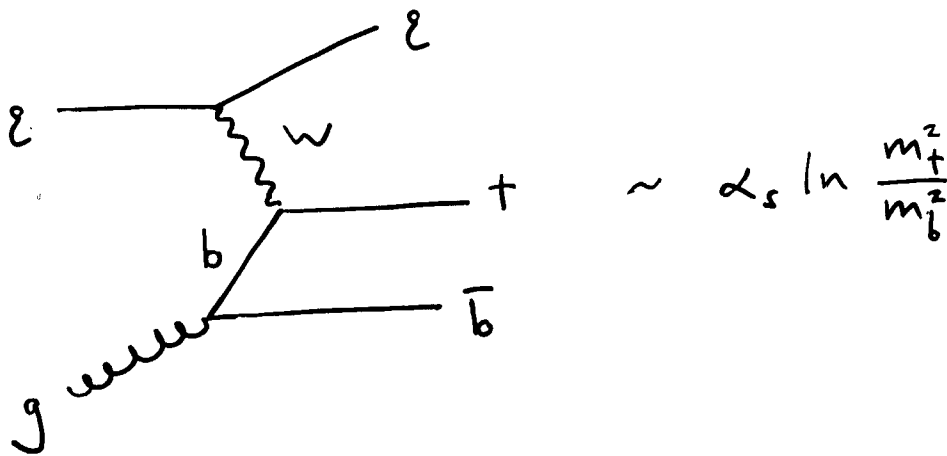
More on t -channel single top:



"W-gluon fusion"

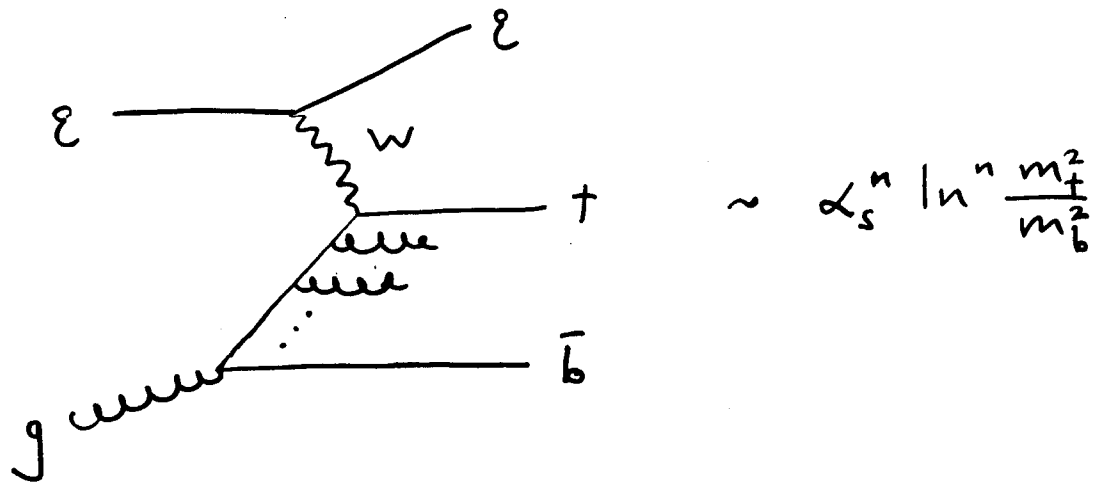
This is correct LO subprocess if $\bar{b}$ is at high p_T .

If we integrate over all p_T of the $\bar{b}$ we obtain a collinear enhancement:

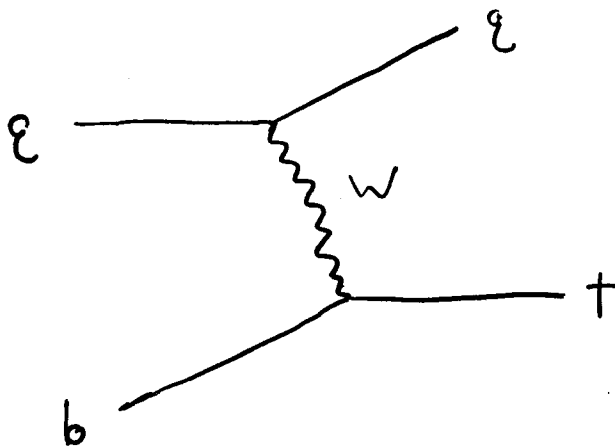


This is desirable but...

... it also makes perturbation theory less convergent:



Solution: Use DGLAP equations to sum enhanced terms to all orders into a b distribution function:



↑ Like any other PDF except it is

perturbatively calculable.

ACOT

Initial condition: $b(x, \mu^2 = m_b^2) = 0$ ($\overline{MS}$)

Use DGLAP equations to calculate

$b(x, \mu^2)$ for $\mu^2 > m_b^2$.

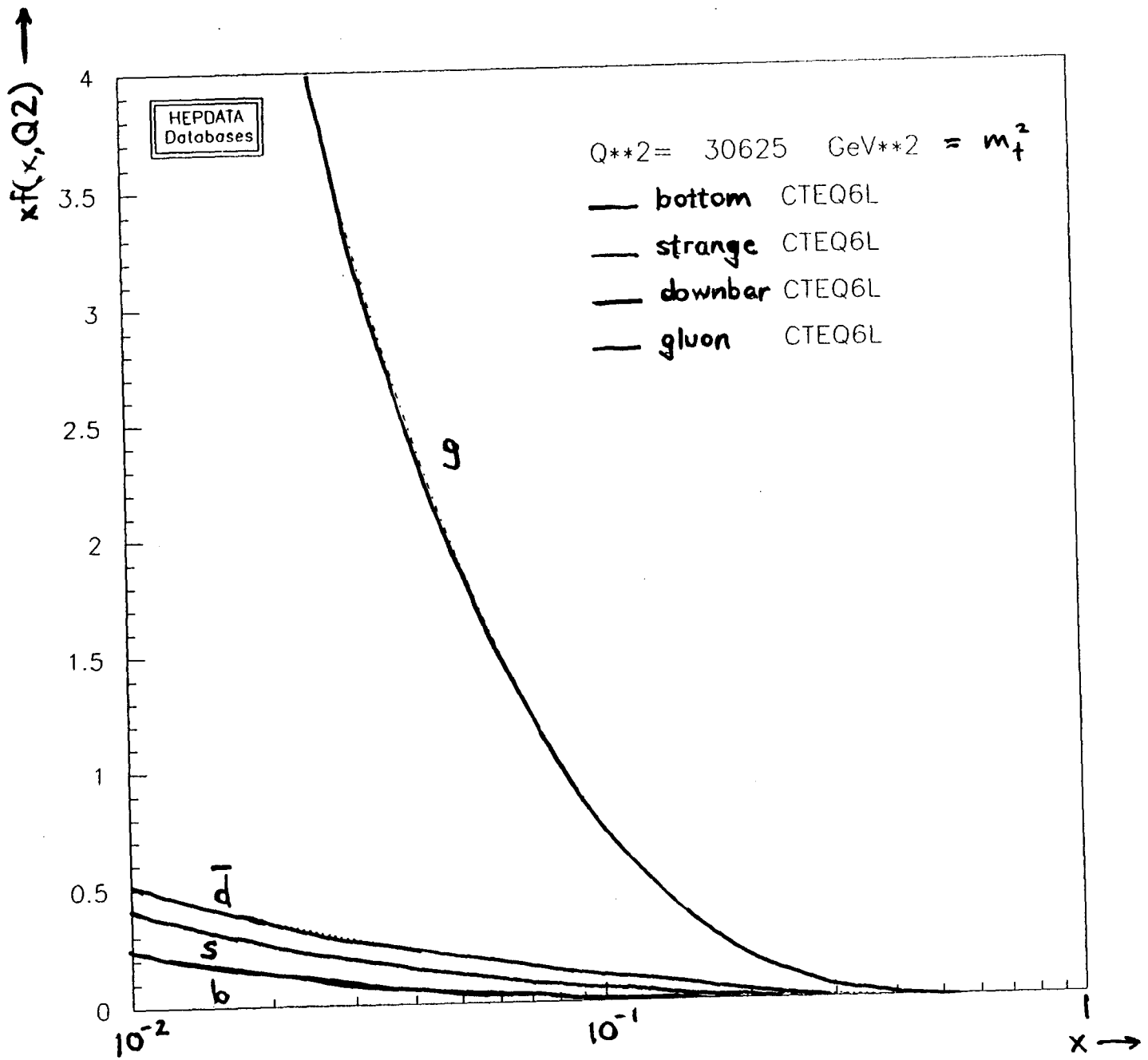
Power counting: Crude solution to

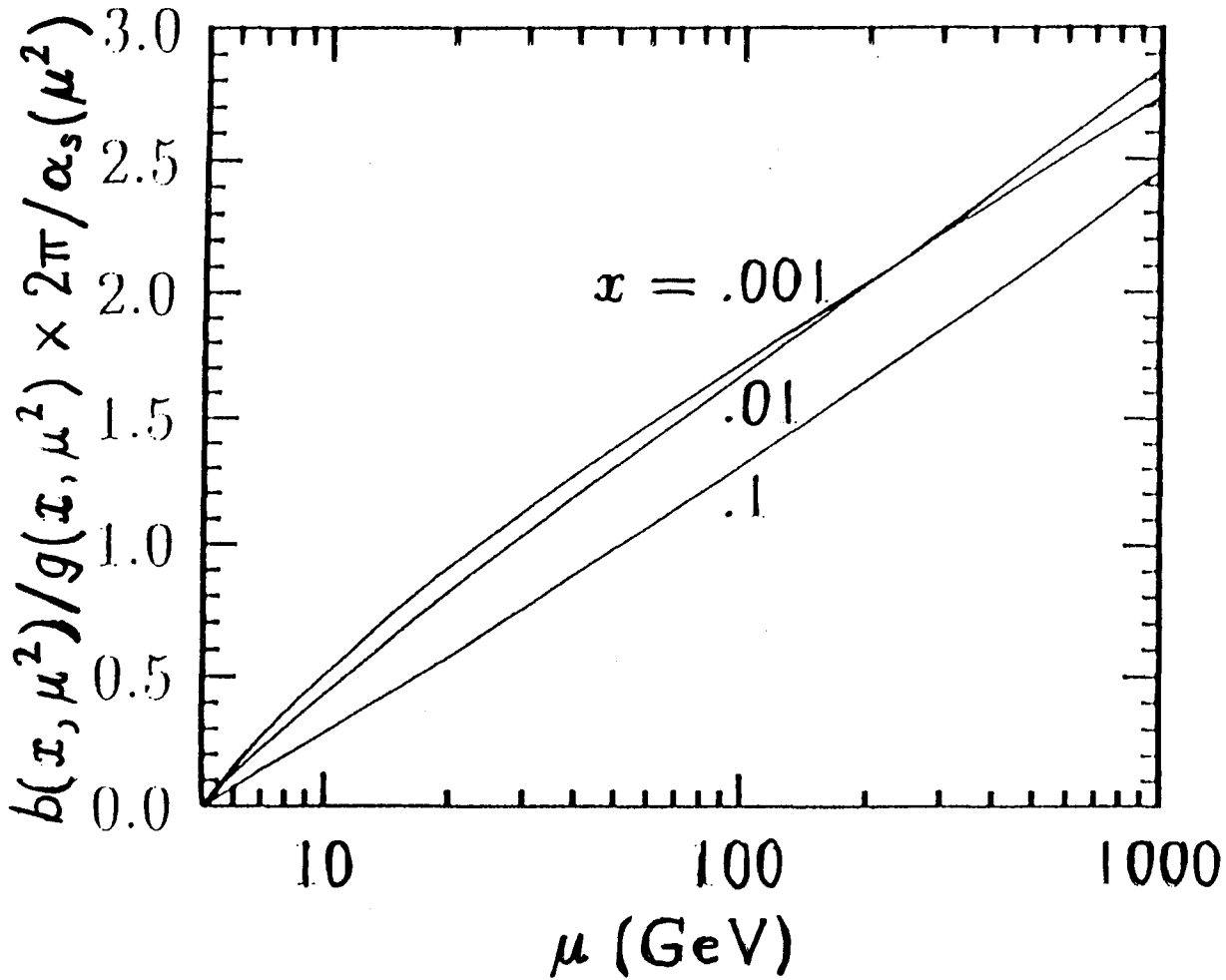
DGLAP equations is

$$b(x, \mu^2) = \frac{\alpha_s(\mu^2)}{2\pi} \ln \frac{\mu^2}{m_b^2} g(x, \mu^2)$$

so b distribution function is intrinsically

$$\mathcal{O}\left(\alpha_s \ln \frac{\mu^2}{m_b^2}\right)$$



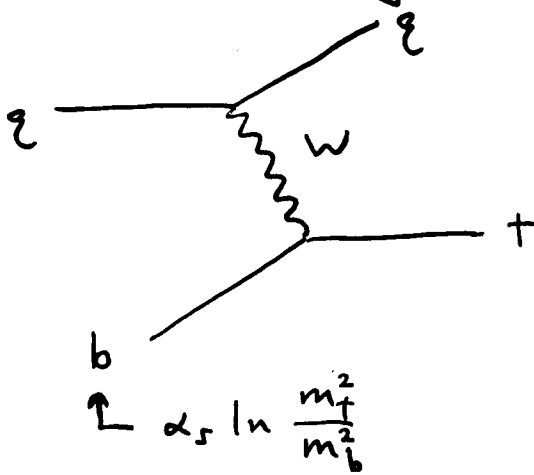


The ratio of the b distribution function to the gluon distribution function, times $2\pi/\alpha_s(\mu^2)$, versus the factorization scale μ , for various fixed values of x . The curves are approximately linear when μ is plotted on a logarithmic scale, indicating that $b(x, \mu^2) \propto [\alpha_s(\mu^2)/2\pi] \ln(\mu^2/m_b^2)g(x, \mu^2)$, as suggested by the approximation of Eq. (1).

SSW

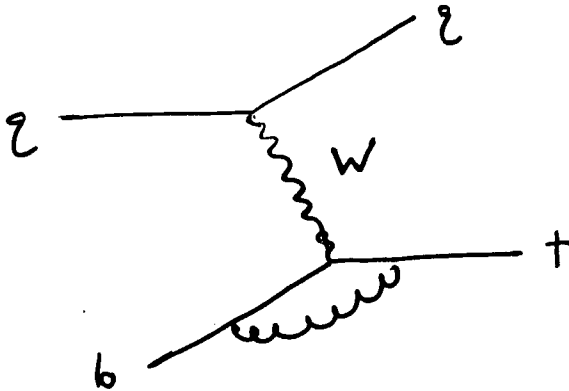
Power counting:

LO:



$$\sim \alpha_s \ln \frac{m_t^2}{m_b^2}$$

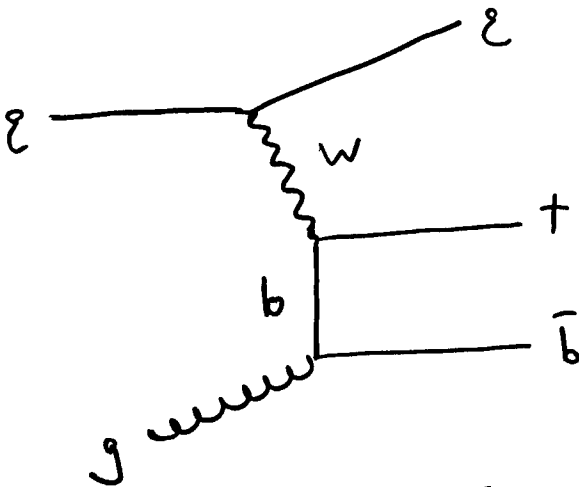
NLO:



$$\sim \alpha_s \text{ correction}$$

$$\alpha_s(m_t^2) \sim \frac{1}{\ln \frac{m_t^2}{\Lambda_{QCD}^2}}$$

NLO:



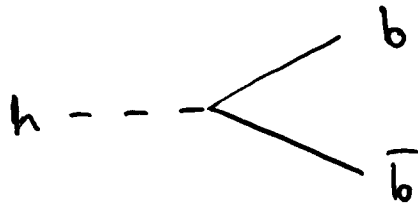
minus collinear enhancement

$$\sim \alpha_s \ln \frac{m_t^2}{m_b^2} + \alpha_s$$

$$\uparrow \frac{1}{\ln \frac{m_t^2}{m_b^2}} \text{ relative to LO}$$

SSW

Exercise: The Higgs boson couples to bottom quarks:



Draw the LO Feynman diagrams for these three processes:

1. $p\bar{p} \rightarrow h + X$
2. $p\bar{p} \rightarrow hb + X$
3. $p\bar{p} \rightarrow hb\bar{b} + X$

where the b quarks are at high p_T .

Estimate the size of processes 2. and 3. relative to 1. (power counting)

CEMW

Selected references

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