

In Search of a Charged Higgs Boson

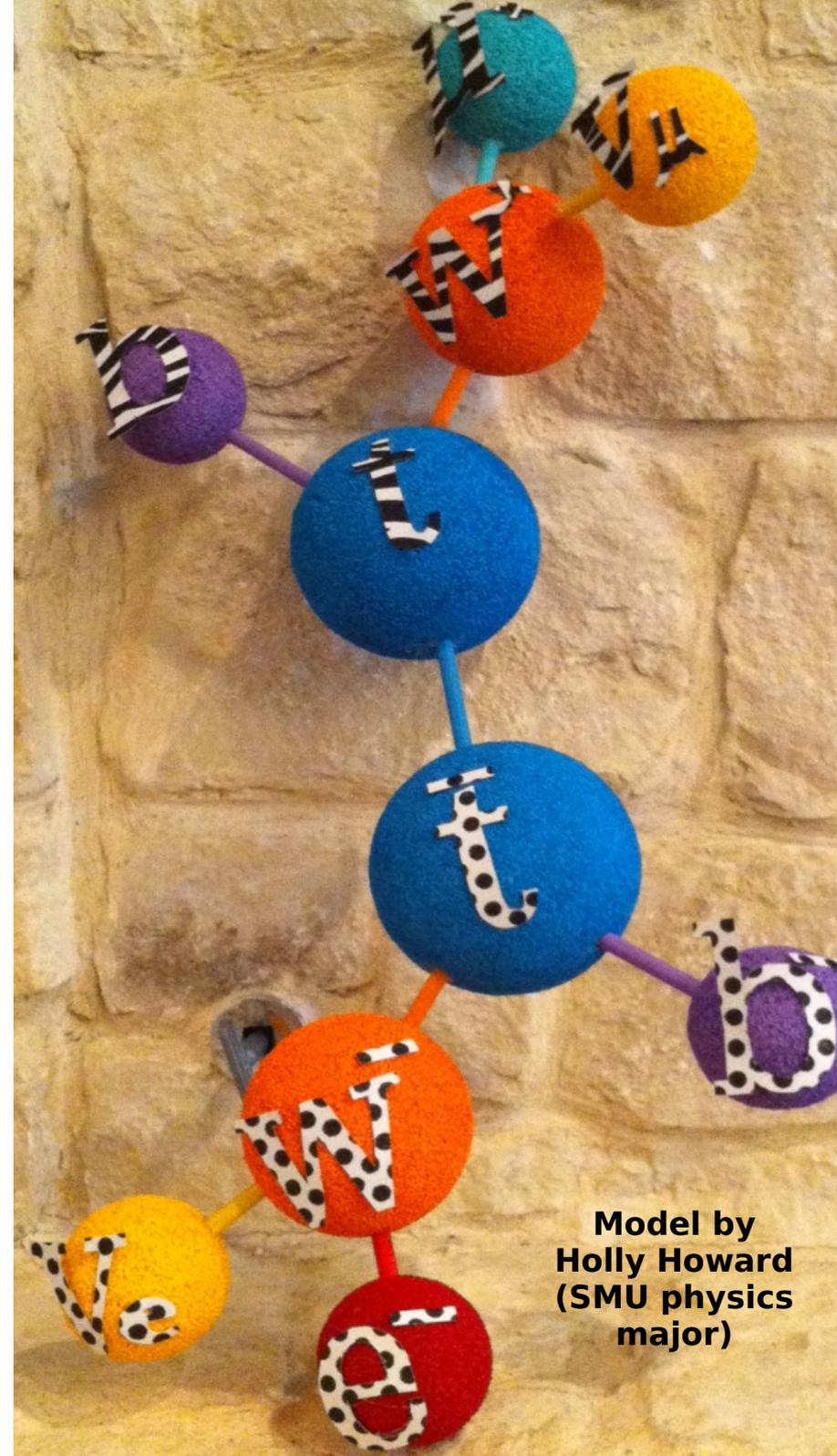
Stephen Sekula
SMU

Presented at the Physics Department Seminar
September 26, 2011



Programme

- Theoretical Motivation
 - Standard Model (SM)
 - Beyond the SM
 - Supersymmetry (SUSY) and the MSSM Higgs scenario(s)
- Low-energy implications
 - how does a charged Higgs affect precision physics at low energy?
- The direct test at high energy
 - The Large Hadron Collider and the ATLAS Experiment
 - Direct search for a “low mass” H^\pm



Model by
Holly Howard
(SMU physics
major)

The Standard Model Higgs

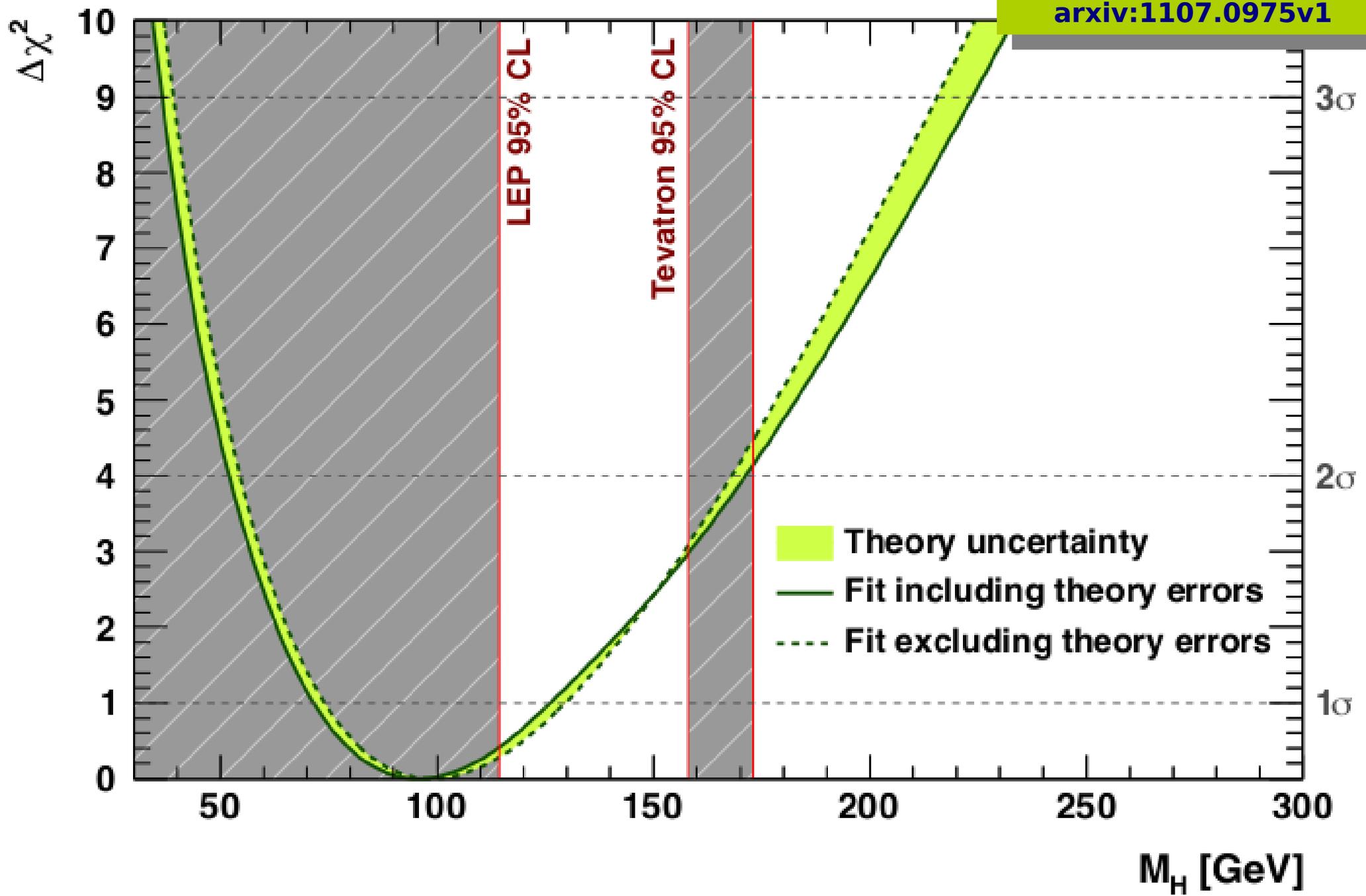
Mass is introduced in the Standard Model by introducing a complex isodoublet, Φ , whose self-interactions provide a mechanism for spontaneous symmetry breaking and gives rise to fermion and boson masses

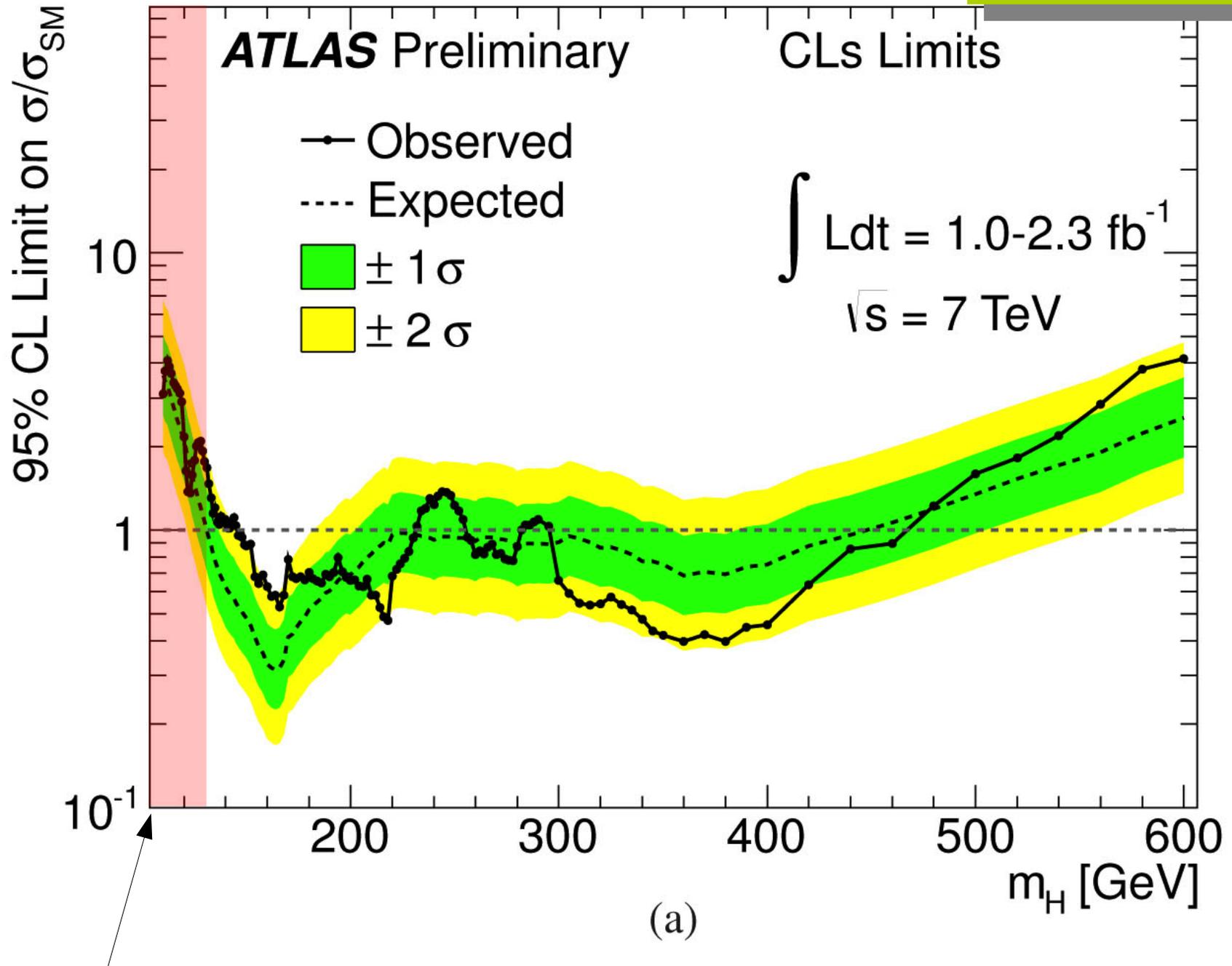
$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$$

→ Adds 4 degrees of freedom to the Standard Model, which by a clever gauge transformation can be made to appear at the three longitudinal polarizations of the W^+ , W^- , and Z^0 bosons and in the mass of a new scalar boson, the Higgs.

$$V(\Phi) = \mu^2 |\Phi|^2 + \lambda |\Phi|^4 \longrightarrow m_H = \sqrt{-2\mu^2}$$

$$m_H = \sqrt{-2\lambda v^2} \longrightarrow \text{Vacuum Expectation Value (VEV) of the Higgs Field, } v = 246 \text{ GeV}$$





Red region indicates central value + 1σ for indirect prediction of Higgs mass.

Extending the Standard Model

- We have good reason to believe the SM is incomplete
 - e.g. neutrino mass/mixing, divergence of corrections to Higgs mass, dark matter, dark energy, that pesky gravity problem . . .
- Extending the SM to solve some of these problems typically means adding more fermions and bosons
 - if the Higgs mechanism is the origin of mass, you need more Higgs bosons to give mass to all the new particles in any extension

Minimal Supersymmetric Standard Model

- Extend the SM with Supersymmetry
 - adds a fermion for every boson, a boson for every fermion
- Requires the addition of at least one more Higgs field doublet
 - adds 4 more degrees of freedom (read: more unknowns that must be measured)
 - leads to 5 physical Higgs boson states:
 - H^0, h^0, A^0, H^+, H^-
 - The MSSM requires that up-type fermions couple to one doublet and down-type fermions couple to the other (“Type-II two-Higgs-doublet Model” or “Type-II 2HDM”)

The Charged Higgs Boson

- Two of the new states are spin-0 and carry electric charge: H^\pm
- General statement:
 - there is no such “beast” in the Standard Model
 - observation of a charged scalar gauge boson would be unambiguous evidence of physics beyond the SM!
- Decay/Production Rates
 - depend on charged Higgs mass and the parameter $\tan \beta$:

$$\tan \beta \equiv \frac{v_2}{v_1}$$

VEV for doublet that couples to down-type fermions (d,s,b,e, μ , τ)

VEV for doublet that couples to up-type fermions (u,c,t, ν)

Charged Higgs Decay

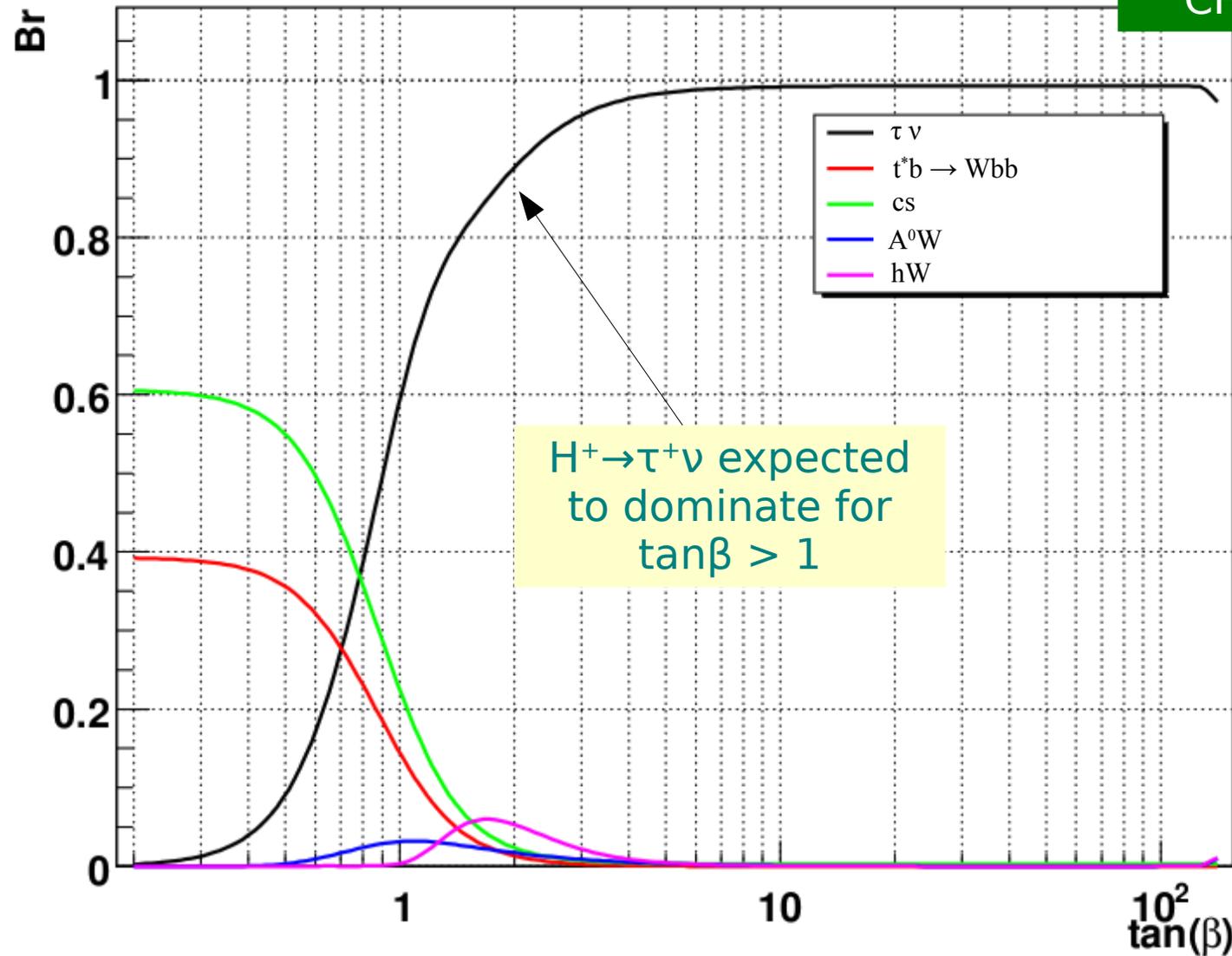
- Decay modes
 - $H^\pm \rightarrow \tau \nu$: dominates for large $\tan\beta$
 - this mode can be used to constrain $\tan\beta$ from above
 - $H^\pm \rightarrow c \bar{s}$: higher rate at low $\tan\beta$
 - can be used to constrain $\tan\beta$ from below
 - $H^\pm \rightarrow t^{(*)} \bar{b}$:
 - also useful at low $\tan\beta$
 - difficult due to off-shell top at low H mass and in general due to top decay to many jets/leptons + MET

“ m_h -max Scenario”

- One of the “benchmark points” in SUSY space for LHC searches
 - chosen to obtain the largest possible Higgs boson mass as a function of $\tan\beta$ (e.g. m_h up to $\sim 140\text{GeV}$)

$$m_t = 174.3 \text{ GeV}, \quad M_{SUSY} = 1 \text{ TeV}, \quad \mu = 200 \text{ GeV}, \quad M_2 = 200 \text{ GeV},$$
$$X_t^{\text{OS}} = 2 M_{SUSY} \text{ (FD calculation)}, \quad X_t^{\overline{\text{MS}}} = \sqrt{6} M_{SUSY} \text{ (RG calculation)}$$
$$A_b = A_t, \quad m_{\tilde{g}} = 0.8 M_{SUSY} .$$

arXiv:hep-ph/0202167
(Eur.Phys.J.C26:601-607,2003)
and arXiv:hep-ph/9812472
(Eur.Phys.J.C9:343-366,1999)

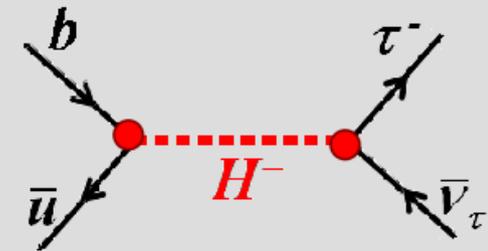
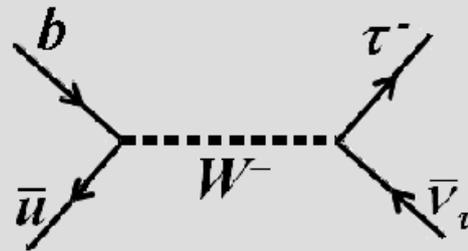


Trends in the decay branching fractions of H^+ for $m_{H^+} = 120 \text{ GeV}$
(computed using FeynHiggs in the m_h -max scenario of the
MSSM)

Low-Energy Implications

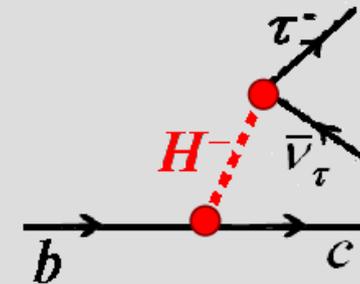
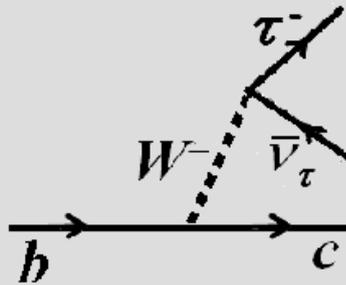
leptonic

- $B \rightarrow \tau \nu_\tau$



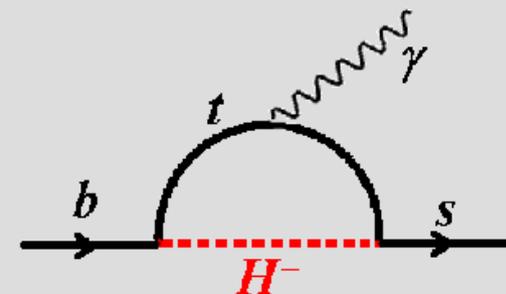
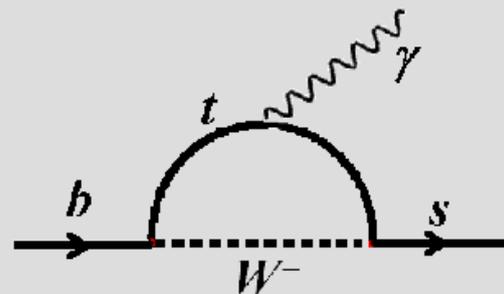
semileptonic

- $B \rightarrow D^{(*)} \tau \nu_\tau$

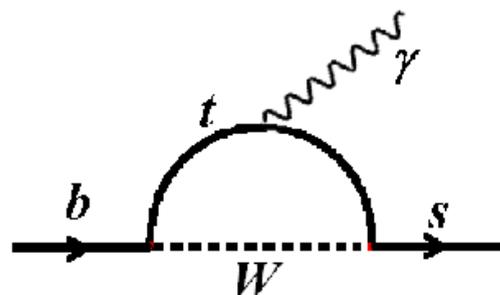


inclusive radiative

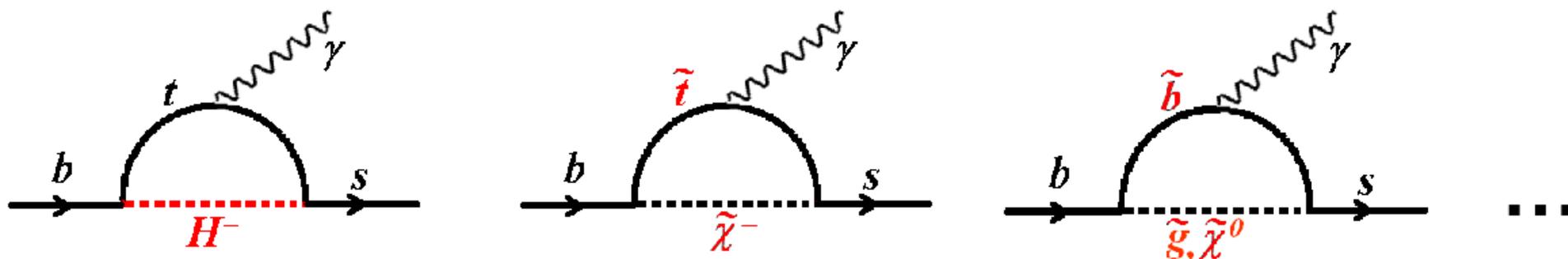
- $B \rightarrow X_S \gamma$



FCNC process in SM occurs via loop diagram

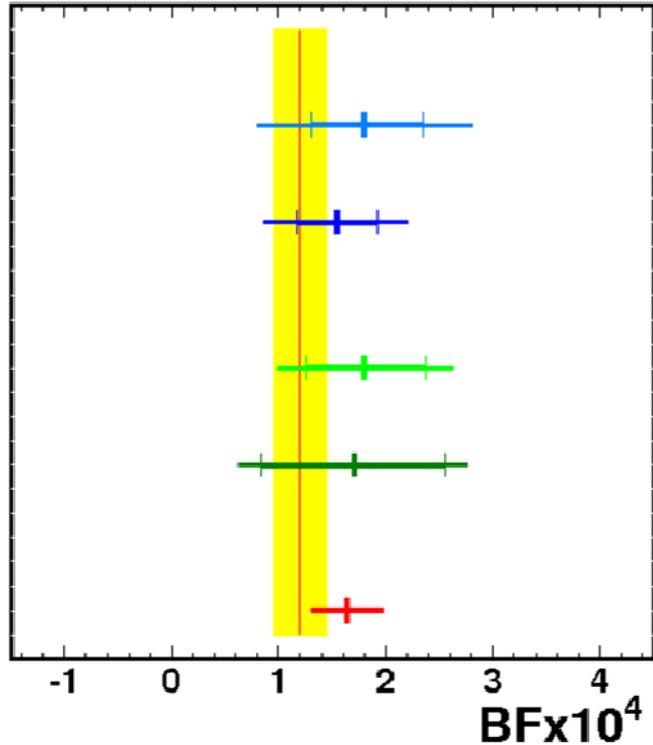


new physics can enter with size comparable to SM contributions



- ☺ BF -enhancement due to the amplitudes with H^\pm depends on m_{H^\pm} but is almost independent of $\tan\beta$
- ☹ more NP processes complicate the interpretation...

$$BF(B^+ \rightarrow \tau^+ \nu_\tau)$$



$$[1.79^{+0.56}_{-0.49}(\text{stat})^{+0.46}_{-0.51}(\text{syst})] \times 10^{-4}$$

$$[1.54^{+0.38}_{-0.37}(\text{stat})^{+0.35}_{-0.37}(\text{syst})] \times 10^{-4}$$

$$[1.80^{+0.57}_{-0.54}(\text{stat}) \pm 0.26(\text{syst})] \times 10^{-4}$$

$$[1.7 \pm 0.8(\text{stat}) \pm 0.2(\text{syst})] \times 10^{-4}$$

$$[1.64 \pm 0.34] \times 10^{-4}$$

HFAG
Aug. 2010



stat. syst.

hadronic tags

semileptonic tags



hadronic tags

semileptonic tags

average¹

Standard Model²

Results consistent within uncertainties,
but all above the SM prediction

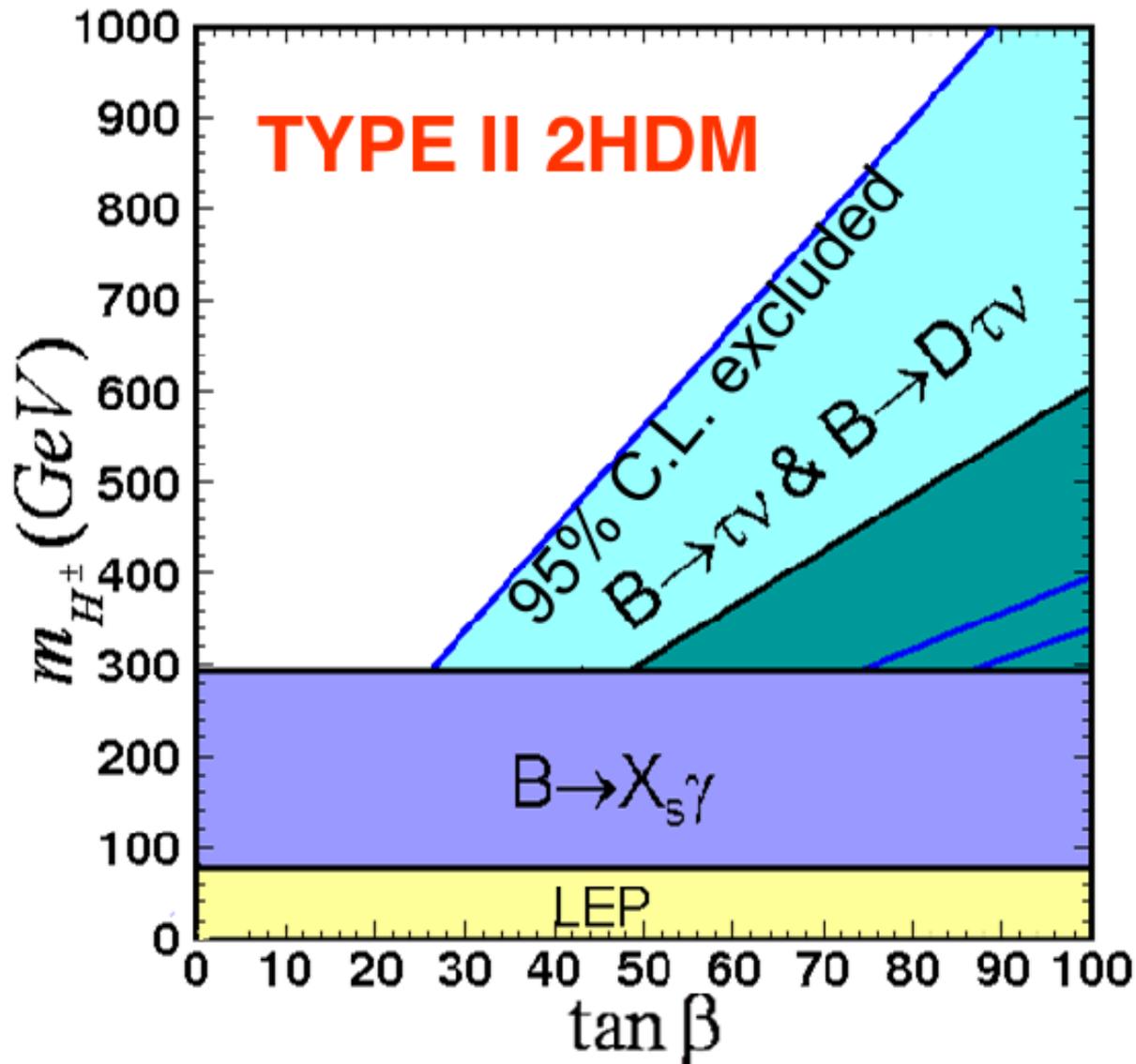
$$r_H = 1.37 \pm 0.39$$

¹ HFAG, <http://www.slac.stanford.edu/xorg/hfag>

² $|V_{ub}| = (4.32 \pm 0.16 \pm 0.29) \times 10^{-3}$ HFAG ICHEP08

$f_B = 190 \pm 13 \text{ MeV}$

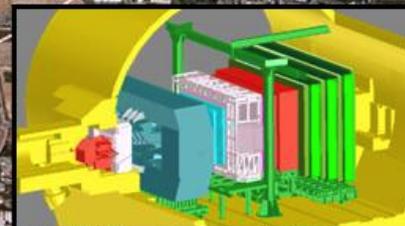
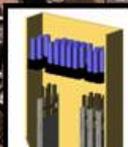
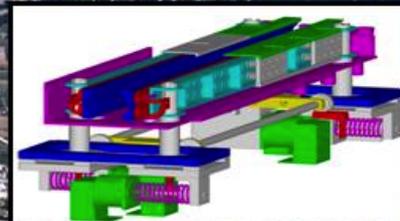
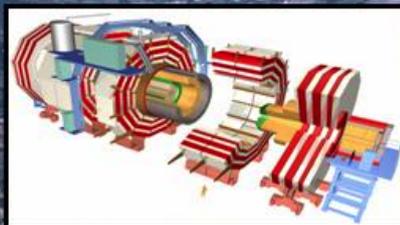
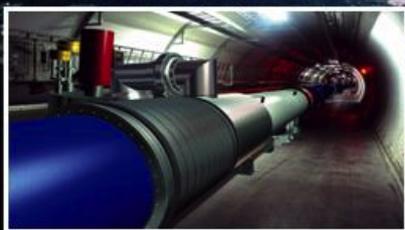
HPQCD arXiv:0902.1815



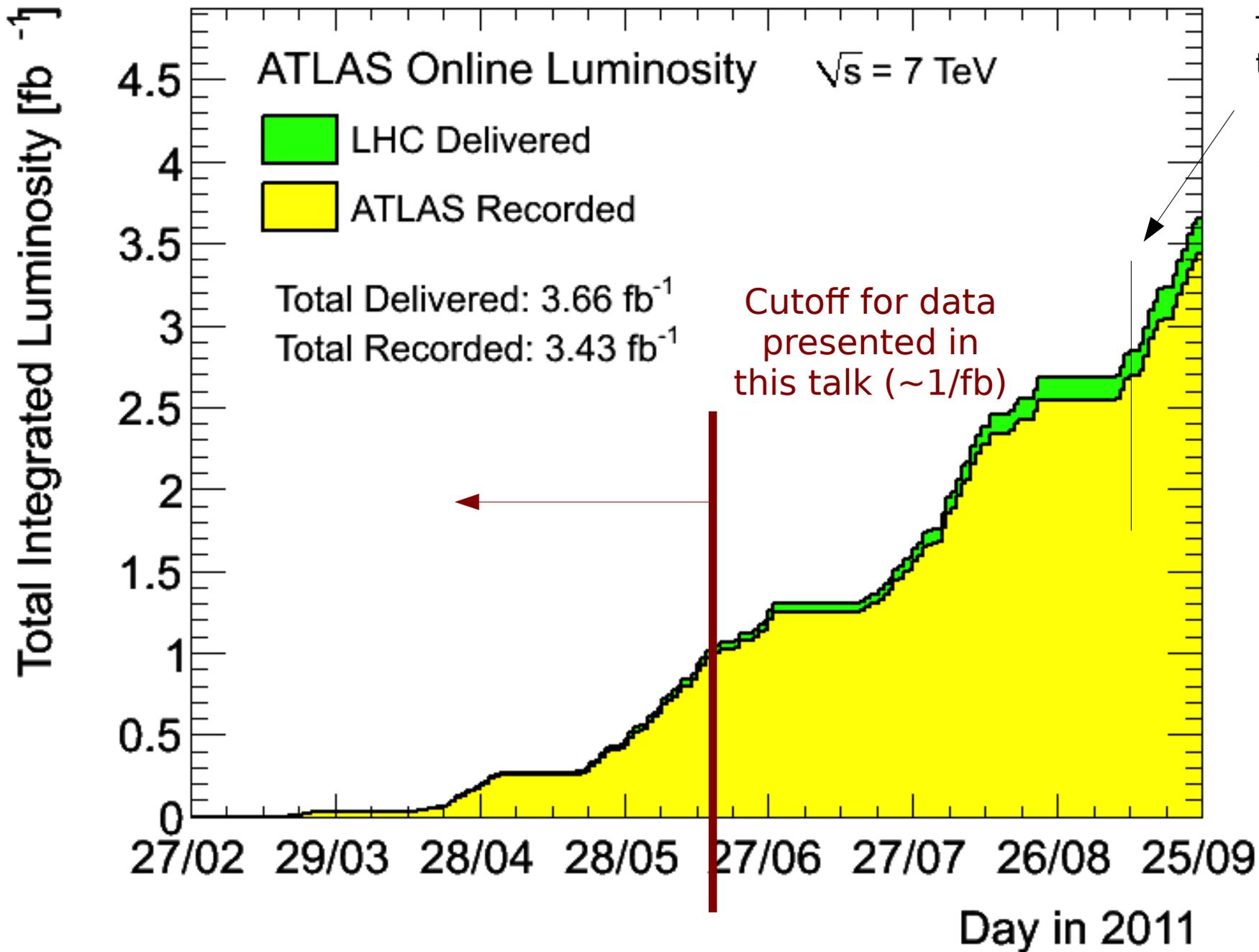
Indirect constraints would suggest that a charged Higgs is excluded whose mass is < 300 GeV ($\tan \beta < 25$) or whose mass is less than $[300, 1000]$ GeV for $\tan \beta > 25$. . .

But one must always be careful with indirect constraints. We have an unprecedented data sample for direct search, and we should use it!

The Large Hadron Collider @ CERN in Geneva, Switzerland



Currently accelerates 1380 proton bunches ($\sim 10^{11}$ per bunch), separated by 50ns, up to 7 TeV in energy (designed for 14 TeV)

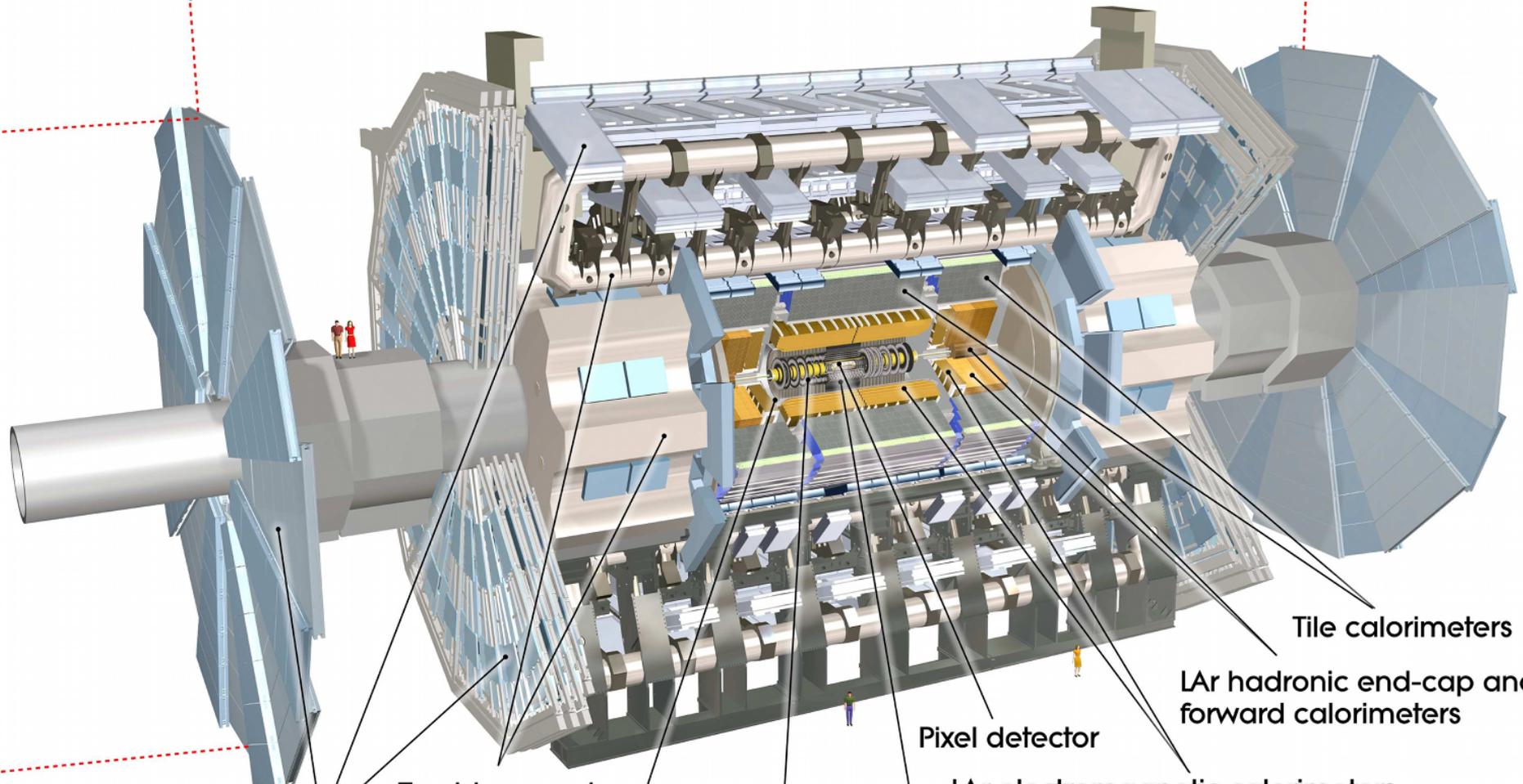


Remember: $L \cdot \sigma = N$

The ATLAS Experiment

44m

25m



Muon chambers

Toroid magnets

Solenoid magnet

Semiconductor tracker

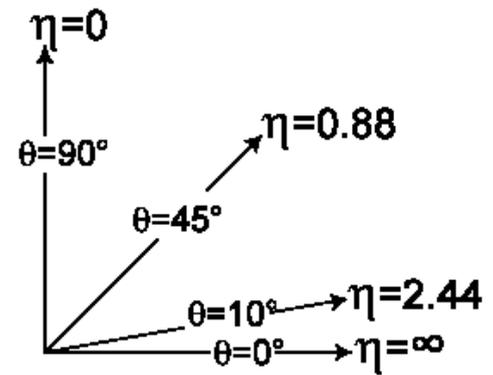
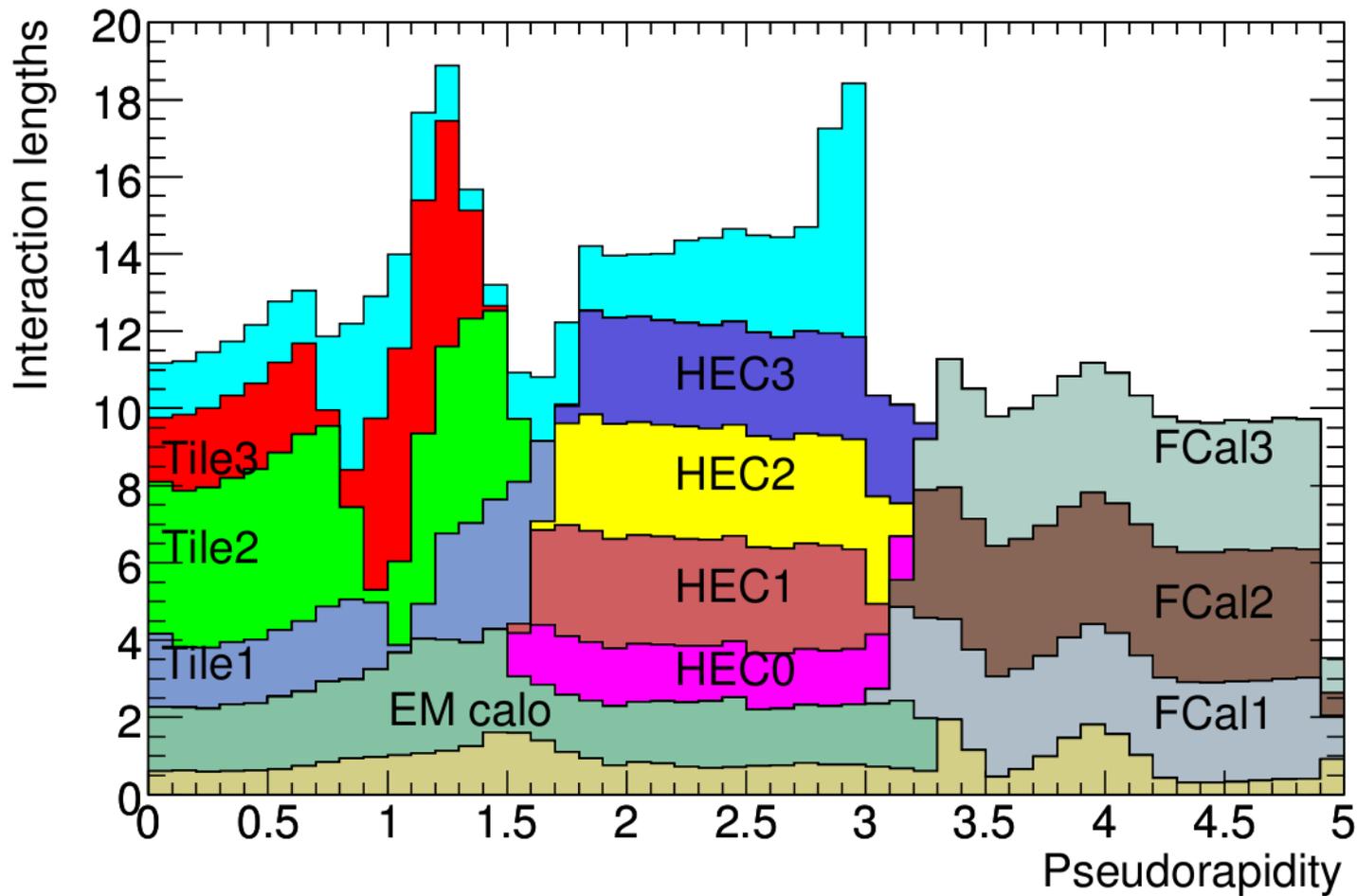
Transition radiation tracker

Pixel detector

LAr electromagnetic calorimeters

LAr hadronic end-cap and forward calorimeters

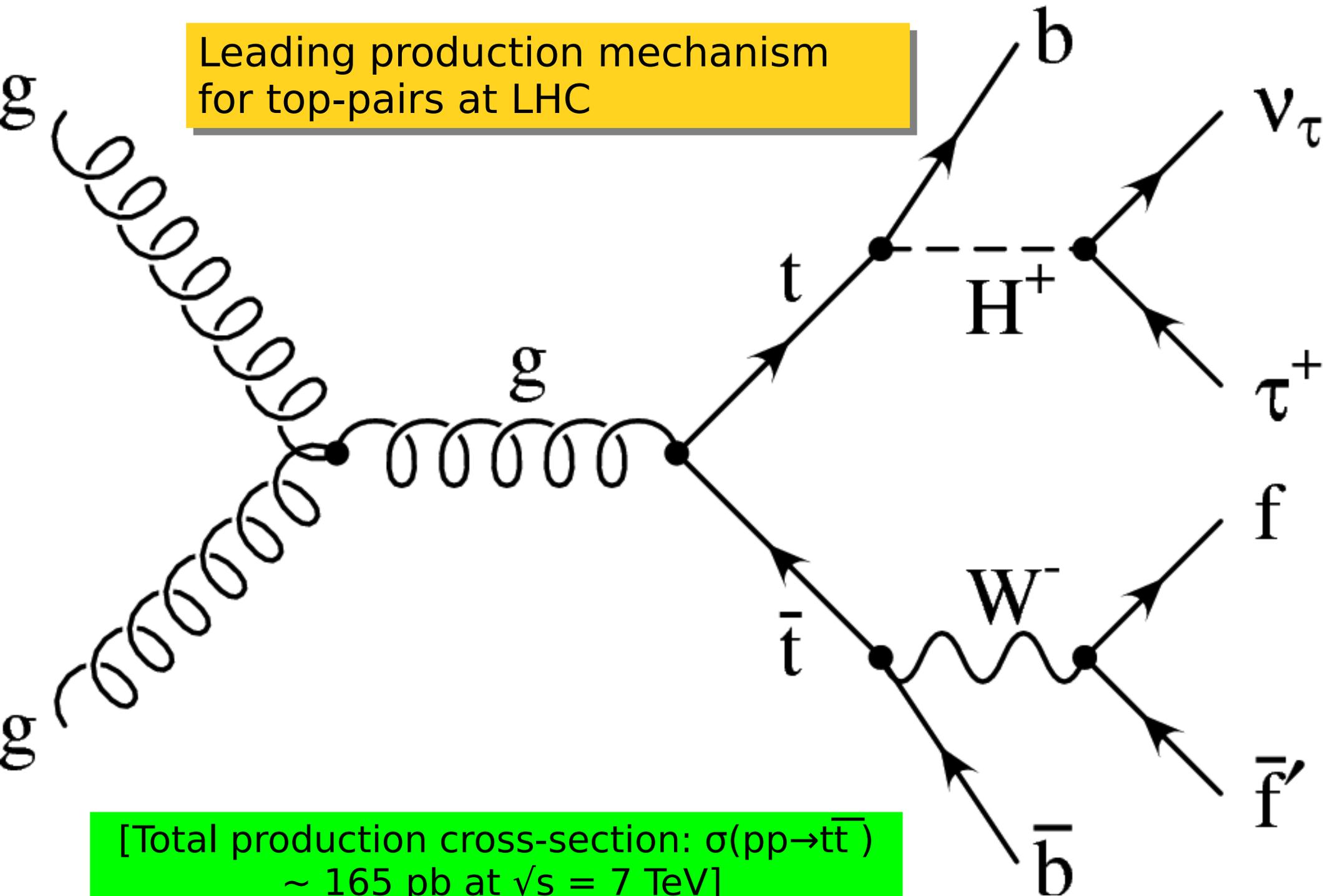
Tile calorimeters



$$\eta = -\log \left(\tan \left(\frac{\theta}{2} \right) \right)$$

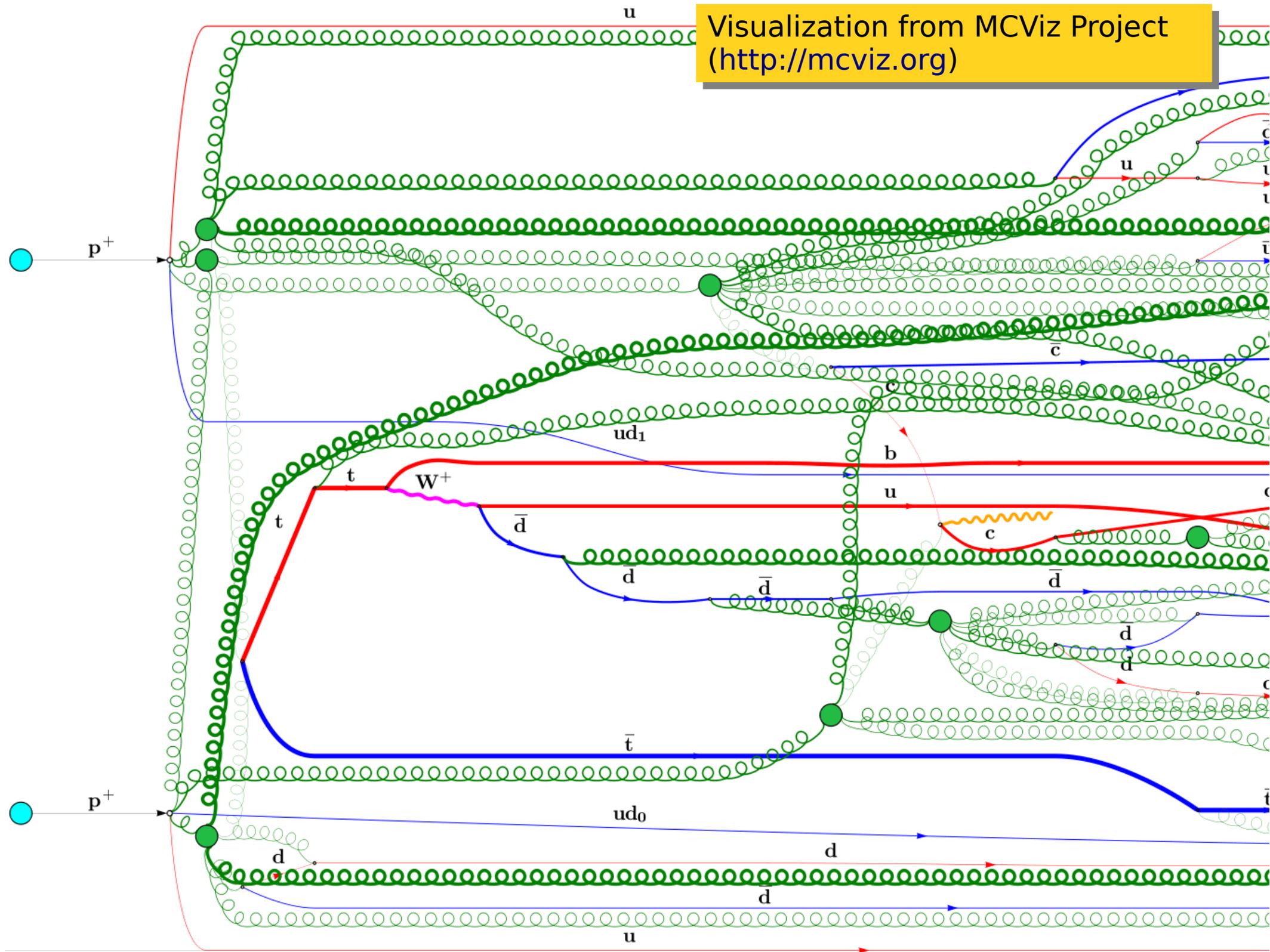
Since much of what I will be talking about today is based on jet reconstruction, it's useful to note that using jets with $|\eta| > 2.5$ means risking incomplete containment of the jets in the EM and Hadronic Calorimeters.

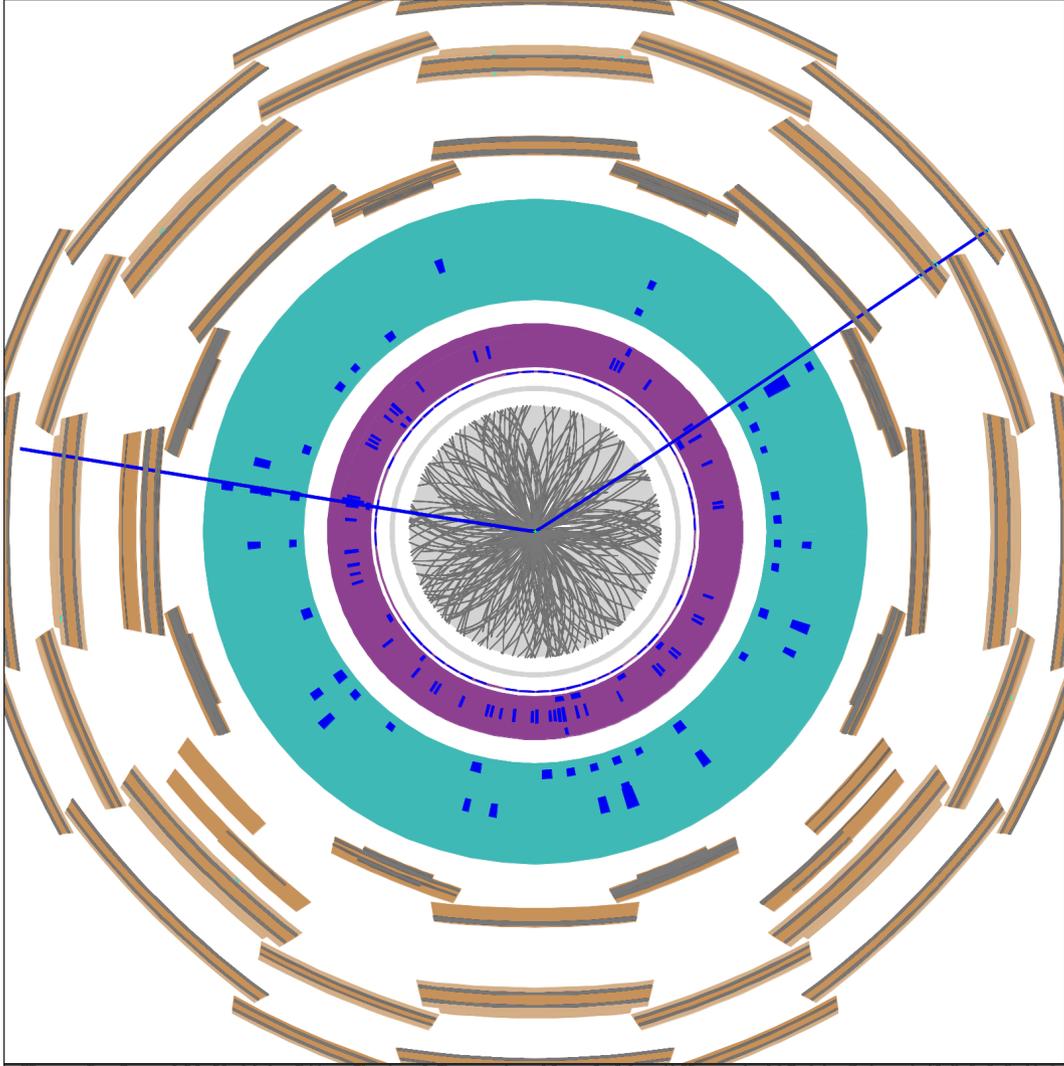
Leading production mechanism
for top-pairs at LHC



[Total production cross-section: $\sigma(pp \rightarrow t\bar{t})$
 ~ 165 pb at $\sqrt{s} = 7$ TeV]

Visualization from MCViz Project (<http://mcviz.org>)

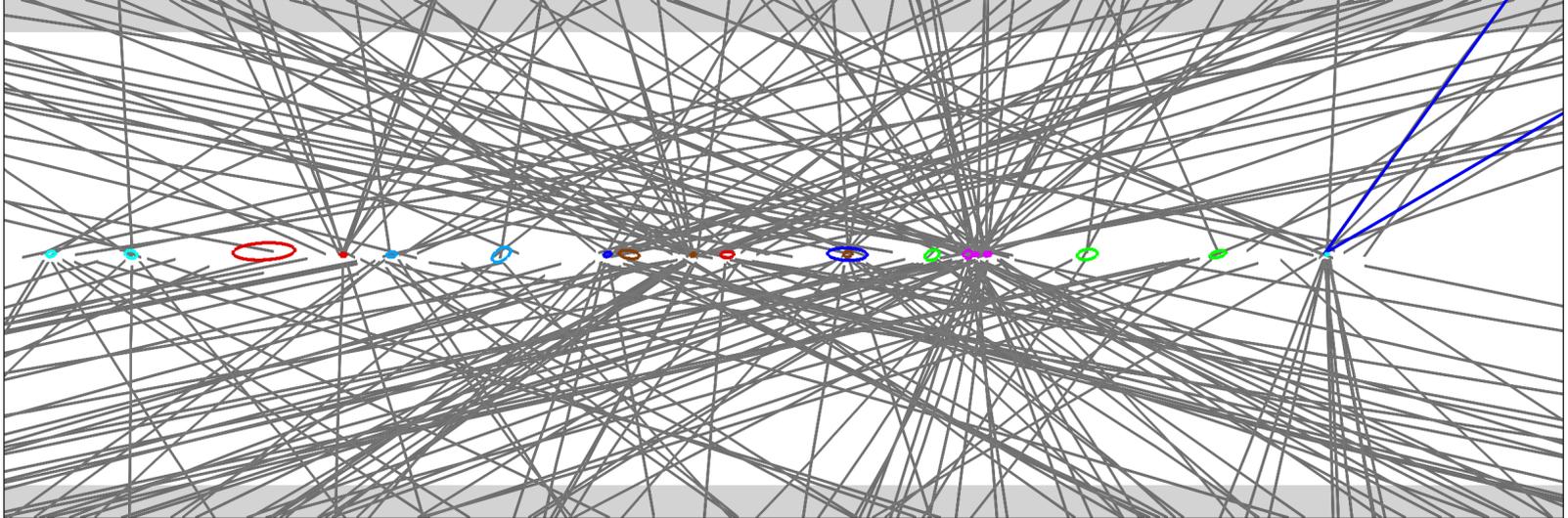
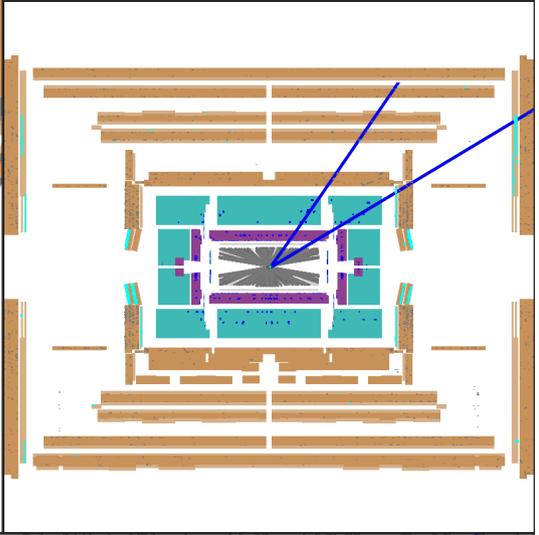


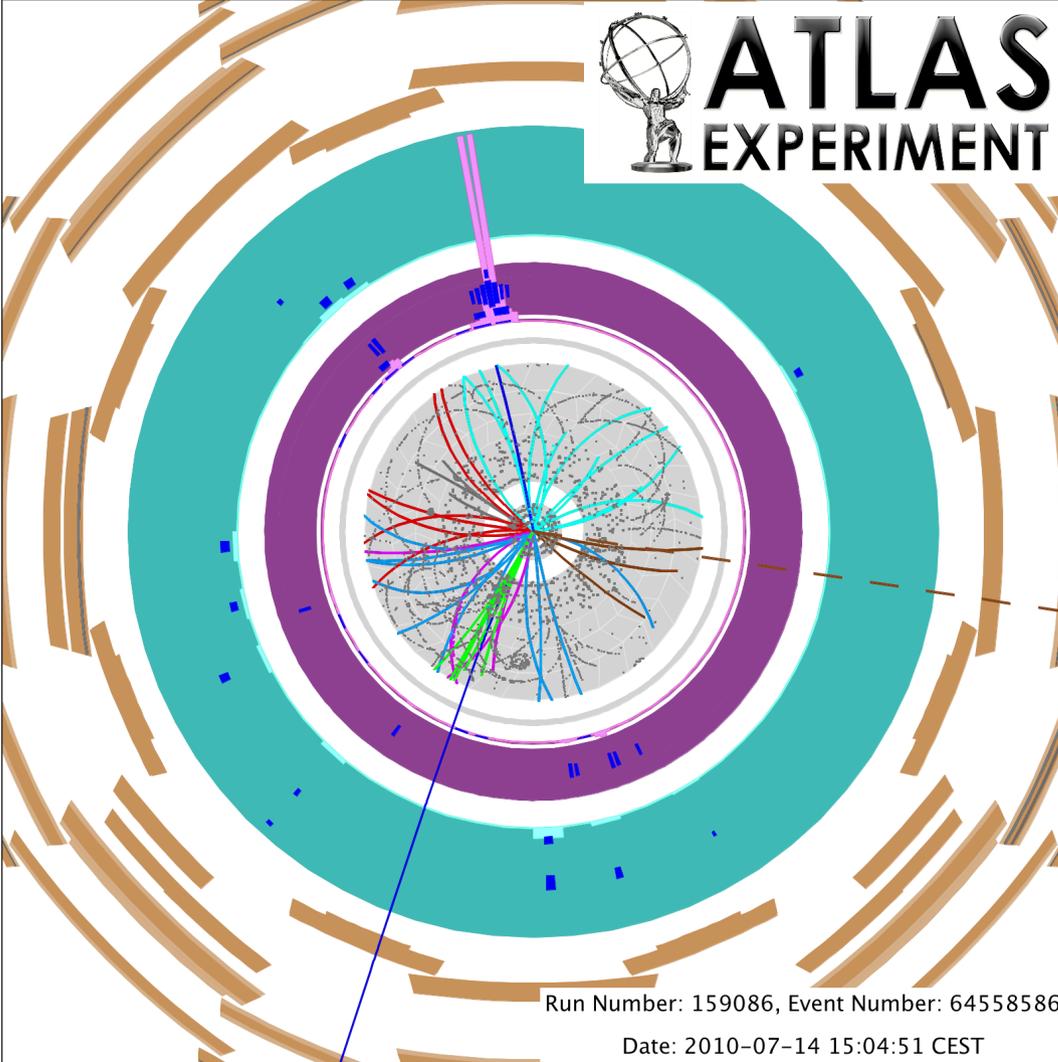


ATLAS EXPERIMENT

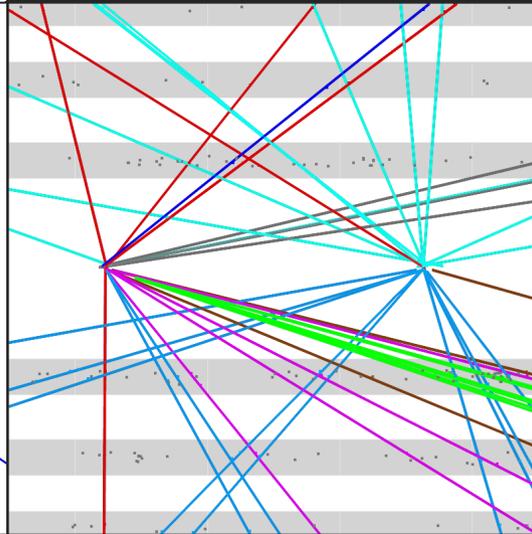
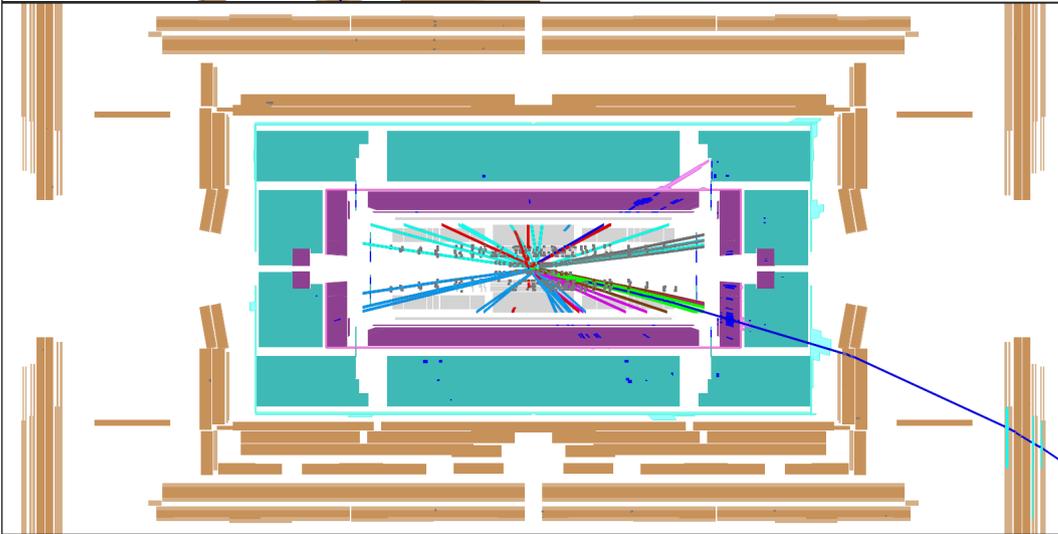
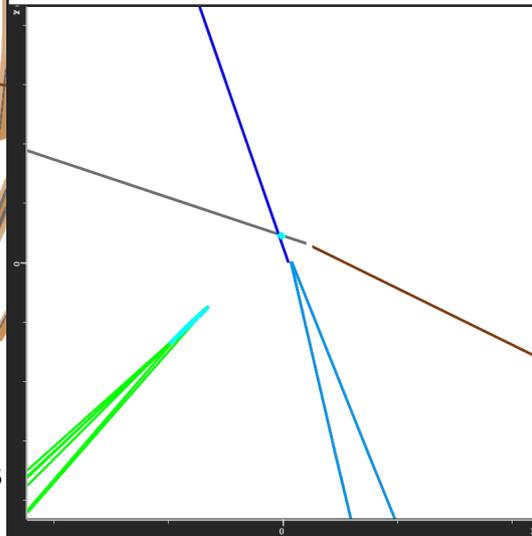
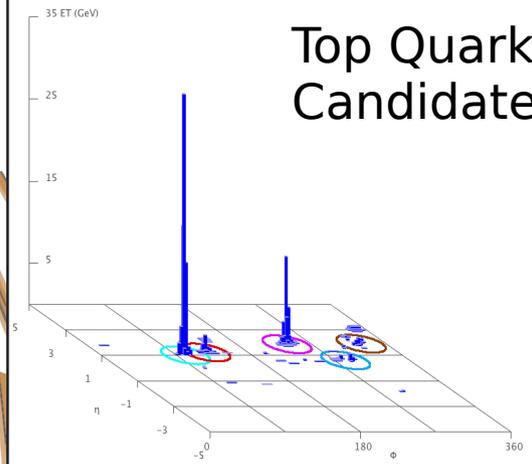
Run Number: 189280, Event Number: 1705325

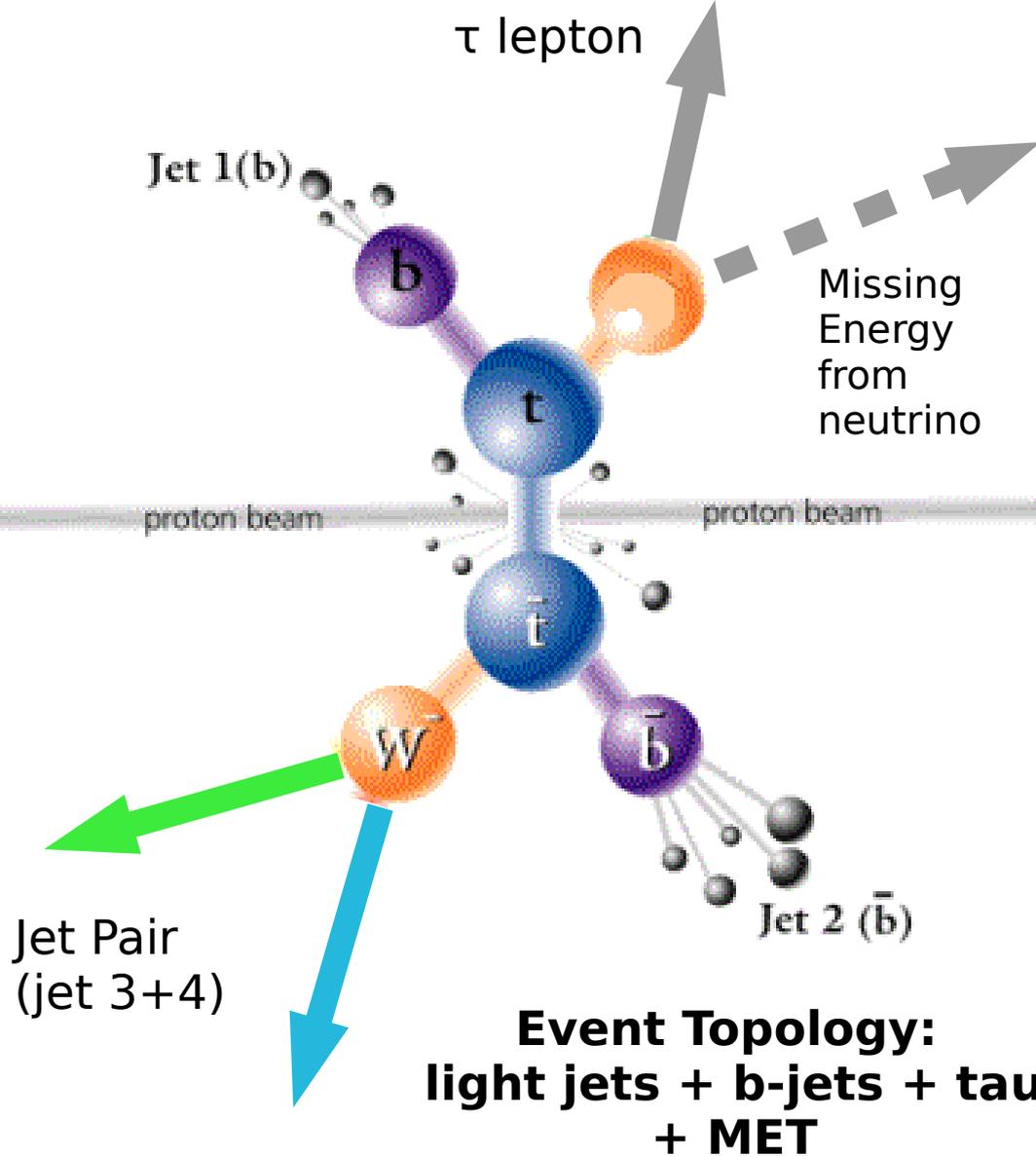
Date: 2011-09-14 02:47:14 CEST





Top Quark Pair Candidate Event





STRATEGY

- Require ≥ 4 jets
- Require ≥ 1 b-tagged jet
- Require a well-identified tau lepton
- Require significant missing transverse energy (MET)
- Combine (jj)b into a top quark candidate
- Study the tau+MET system

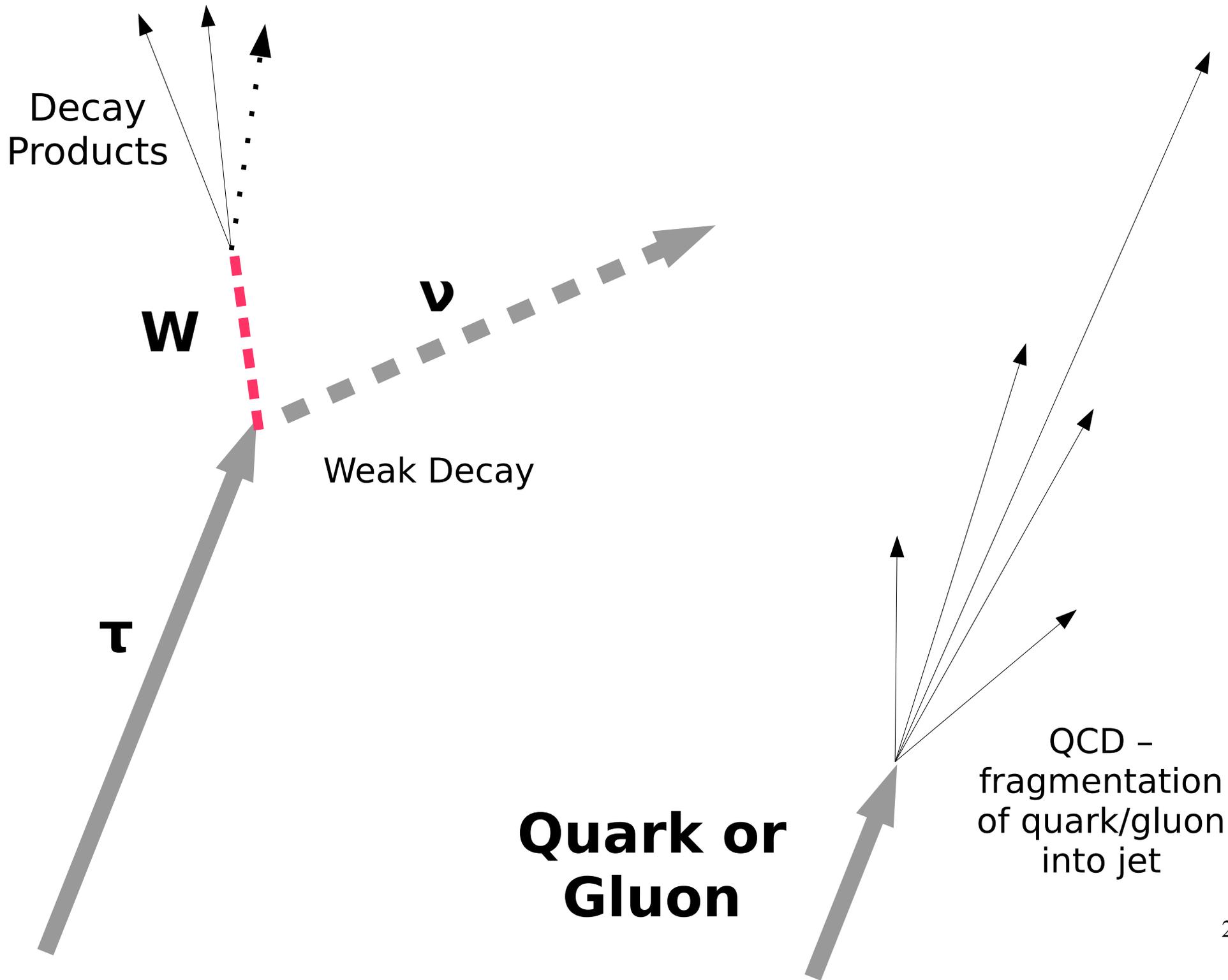
In the next part of the talk I'll focus in on a few key features of the reconstruction: tau identification and background rejection/estimation.

Simulations

Process	Generator	Cross section [pb]
$t\bar{t}$ with $\geq 1\ell$	MC@NLO	89.4
single-top (s, t, Wt channel)	MC@NLO	21.4, 1.41, 14.6
$W \rightarrow \ell\nu + \text{jets}$	ALPGEN	$3.1 \cdot 10^4$
$Z/\gamma^* \rightarrow \ell\ell + \text{jets}$	ALPGEN	$3.2 \cdot 10^3$
$t\bar{t} \rightarrow bH^\pm bW$ with $H^\pm \rightarrow \tau\nu$	PYTHIA	29.6

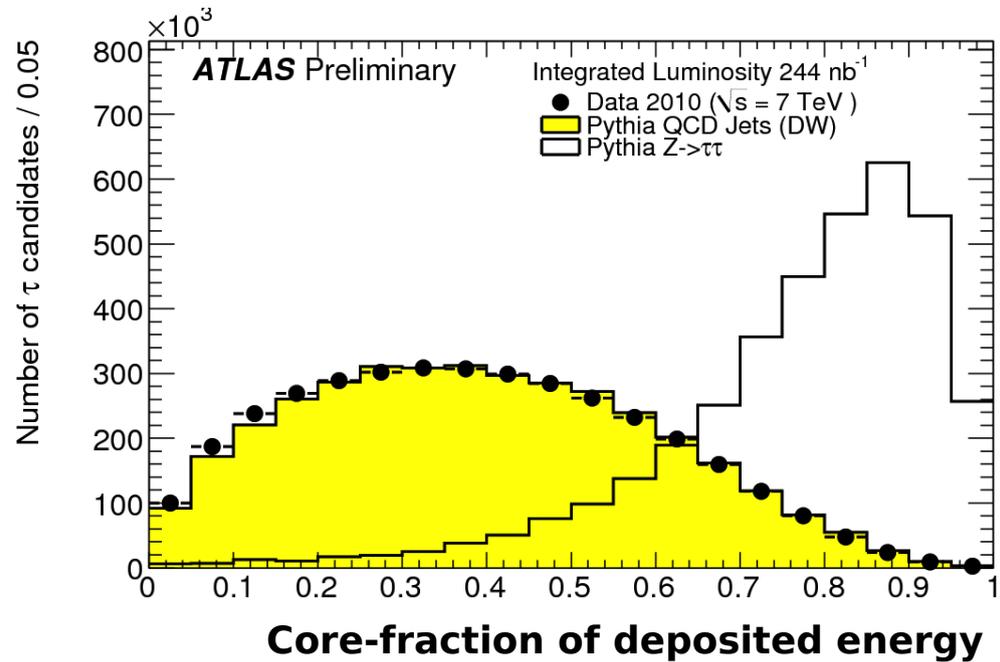
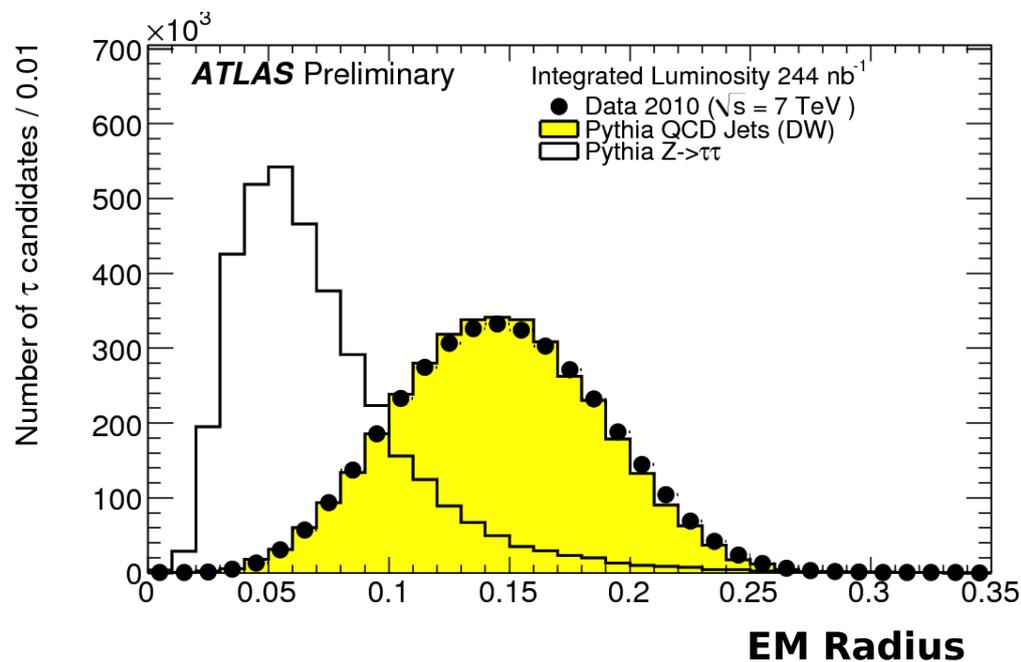
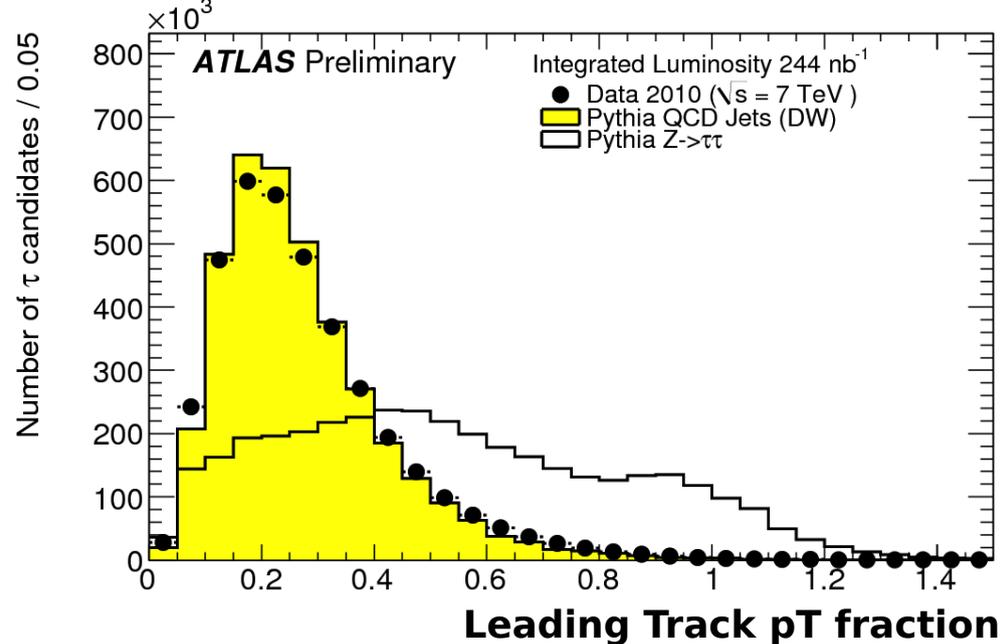
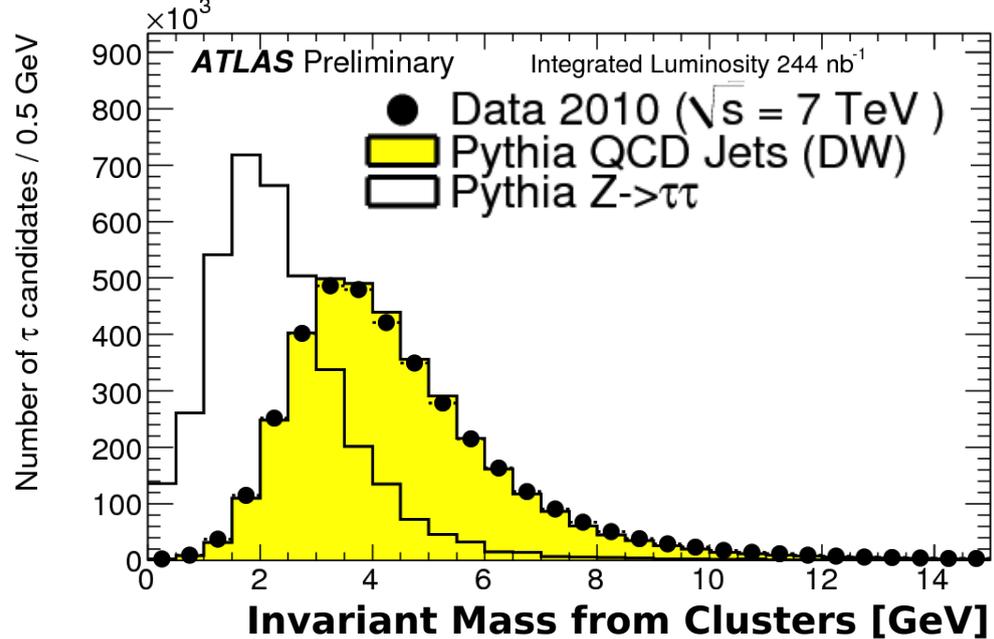
After all our selection criteria are applied (see next several slides), we find this dominates over all other backgrounds.

One notable omission from this table: **generic multi-jet QCD!** More on that in a few slides . . .



Tau Reconstruction

- Tau Decay
 - mass: $1.777 \text{ GeV}/c^2$, with electric charge (either $\pm e$)
 - 35% leptonic final states ($e^+ \nu \nu$, $\mu^+ \nu \nu$),
65% hadronic final states (e.g. $\pi^+ \nu$, $\pi^+ \pi^0 \nu$, $\pi^+ \pi^- \pi^+ \nu$)
 - hadronic final states will superficially resemble jets \rightarrow jet background expected to be the largest problem
- Tau Reconstruction
 - seeded from either a track (trajectory left by a charged particle), a calorimeter jet candidate, or both
 - identification uses additional variables
 - e.g. mass (tracks, clusters), momentum of the leading (highest-pT) track, size and distribution of energy in calorimeters (EM and hadronic)



These and other variables are combined using multivariate algorithms (e.g. boosted decision trees or likelihood ratios) to define a final selector output.

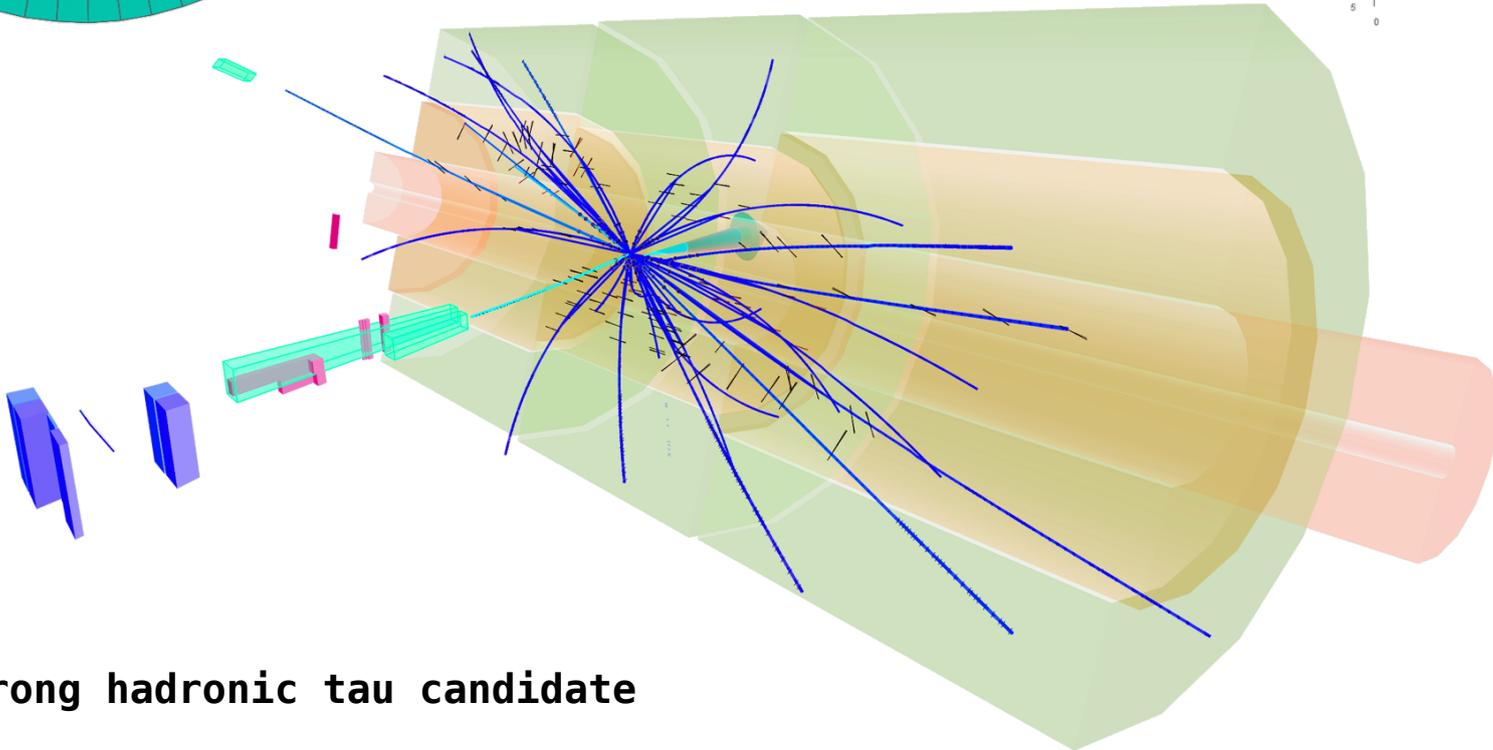
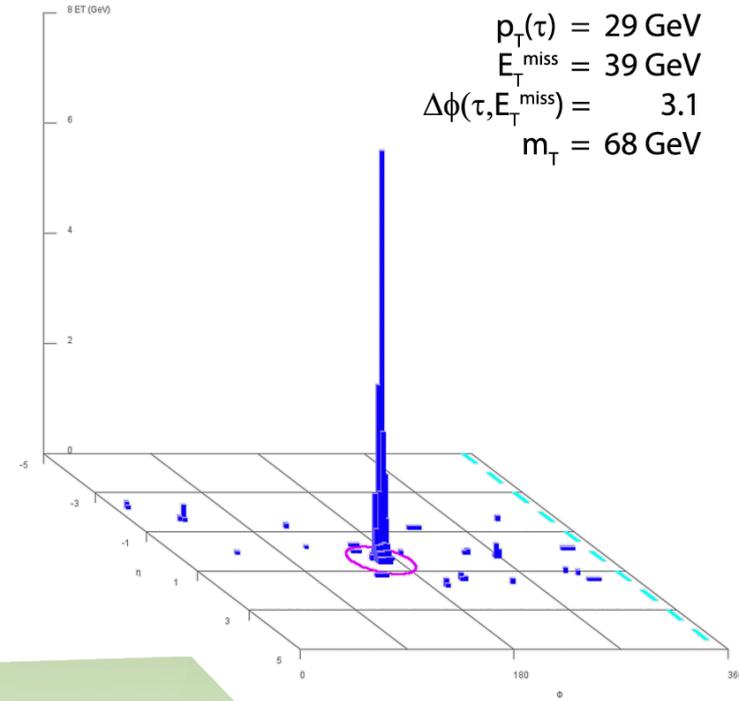
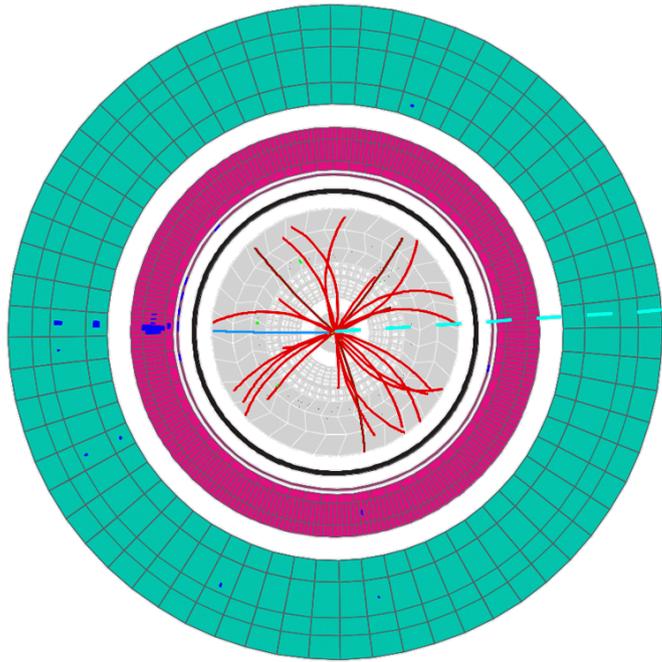


ATLAS EXPERIMENT

Run 155697, Event 6769403

Time 2010-05-24, 17:38 CEST

$W \rightarrow \tau \nu$ candidate in
7 TeV collisions



Single-prong hadronic tau candidate

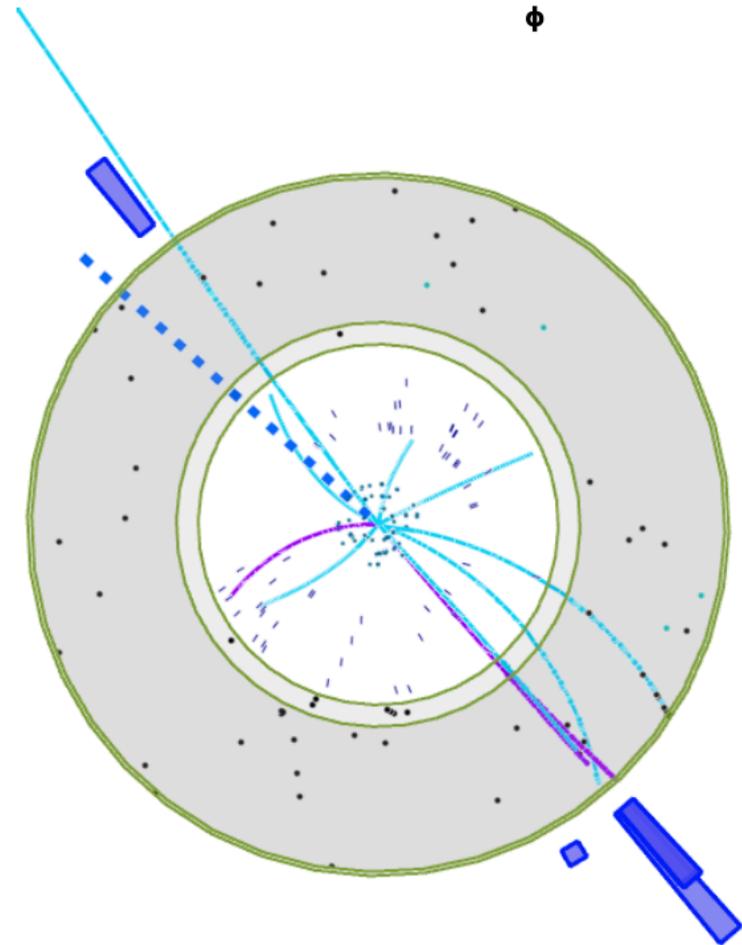
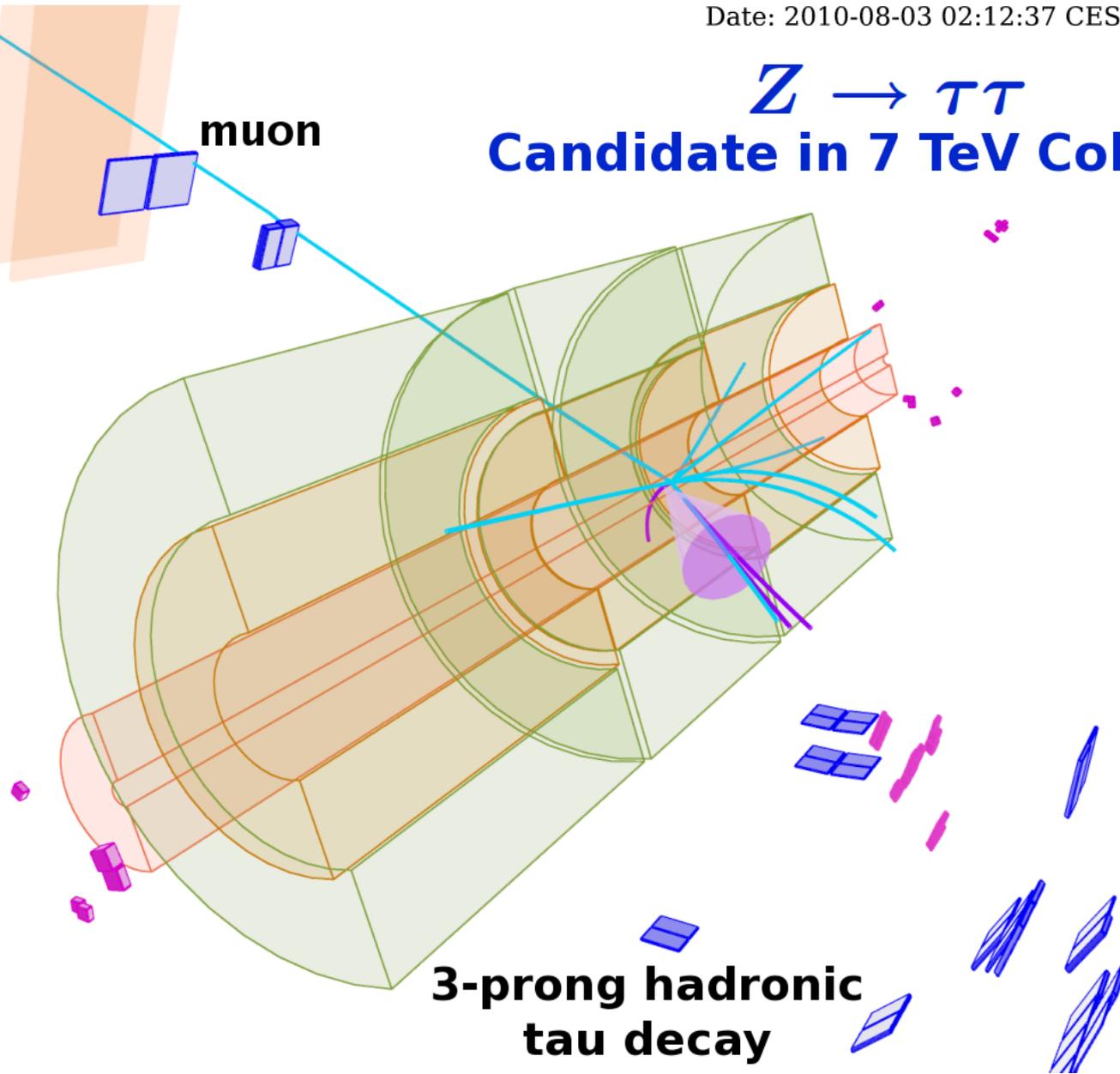
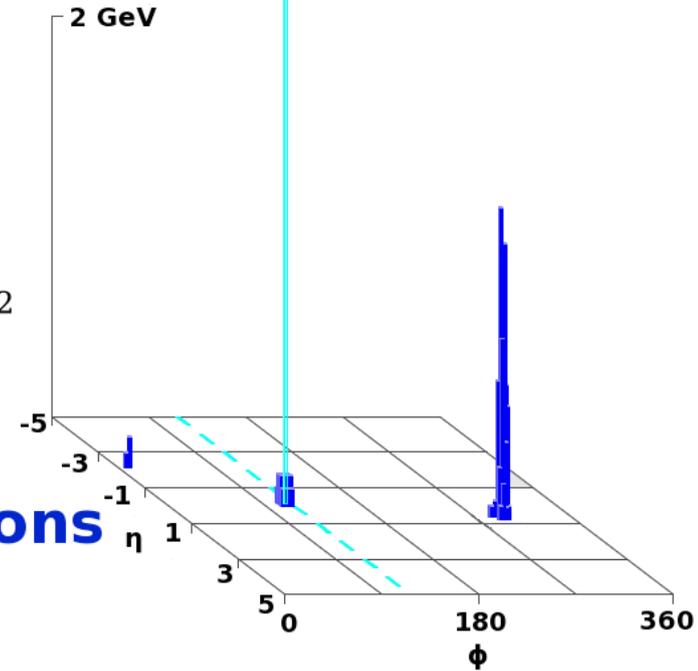
$p_T(\mu) = 18 \text{ GeV}$
 $p_T^{\text{vis}}(\tau_h) = 26 \text{ GeV}$
 $m_{\text{vis}}(\mu, \tau_h) = 47 \text{ GeV}$
 $m_T(\mu, E_T^{\text{miss}}) = 8 \text{ GeV}$
 $E_T^{\text{miss}} = 7 \text{ GeV}$



Run Number: 160613, Event Number: 9209492

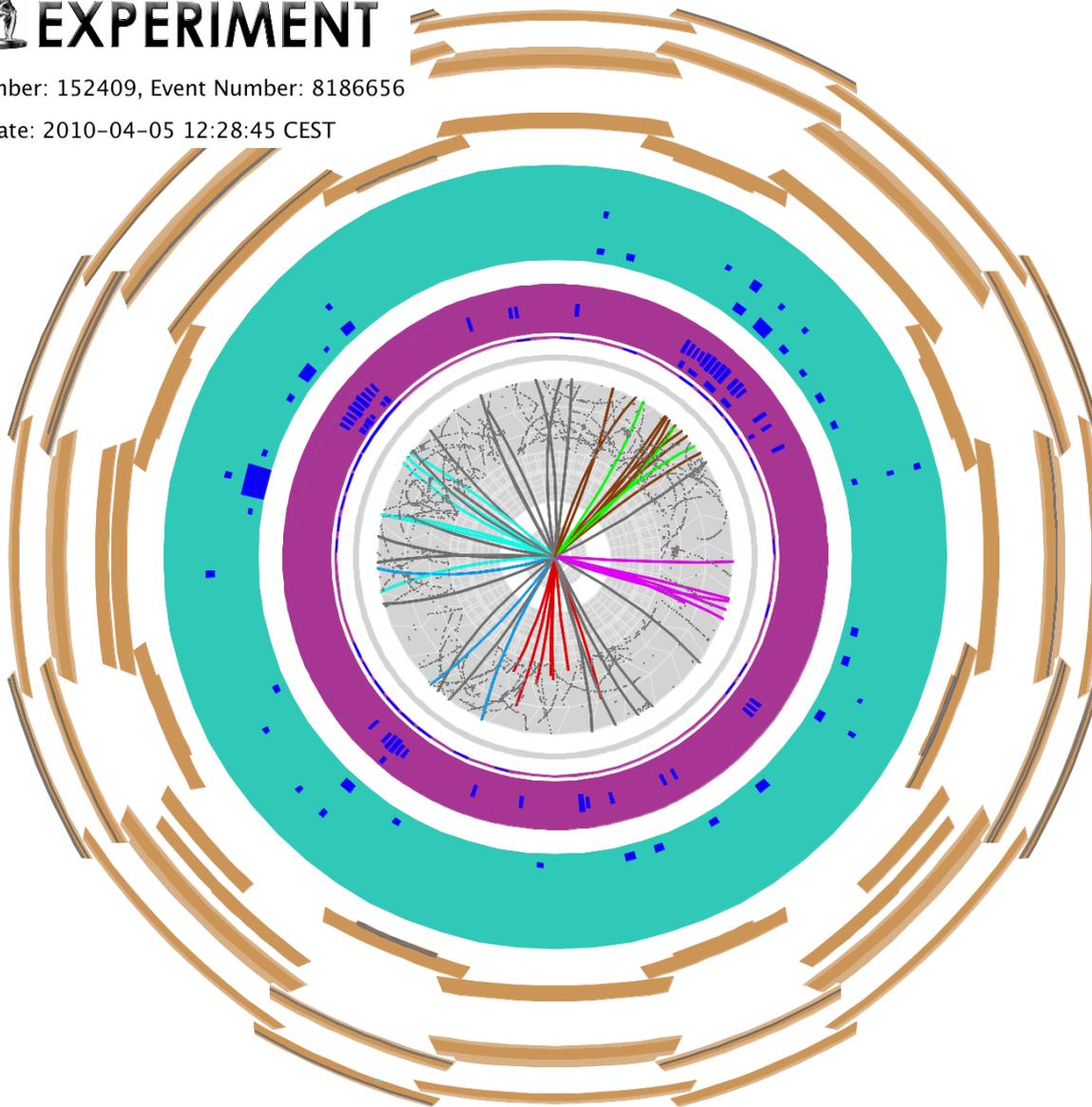
Date: 2010-08-03 02:12:37 CEST

$Z \rightarrow \tau\tau$ Candidate in 7 TeV Collisions

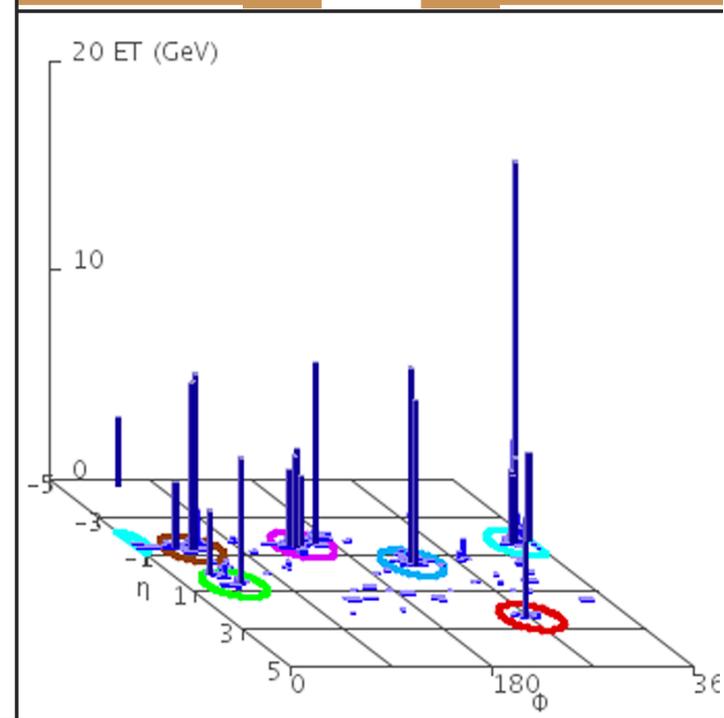
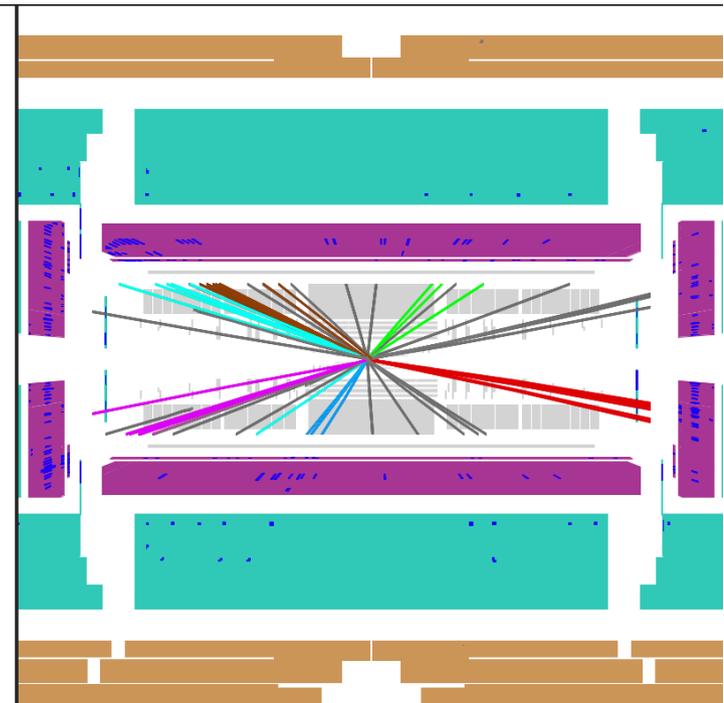


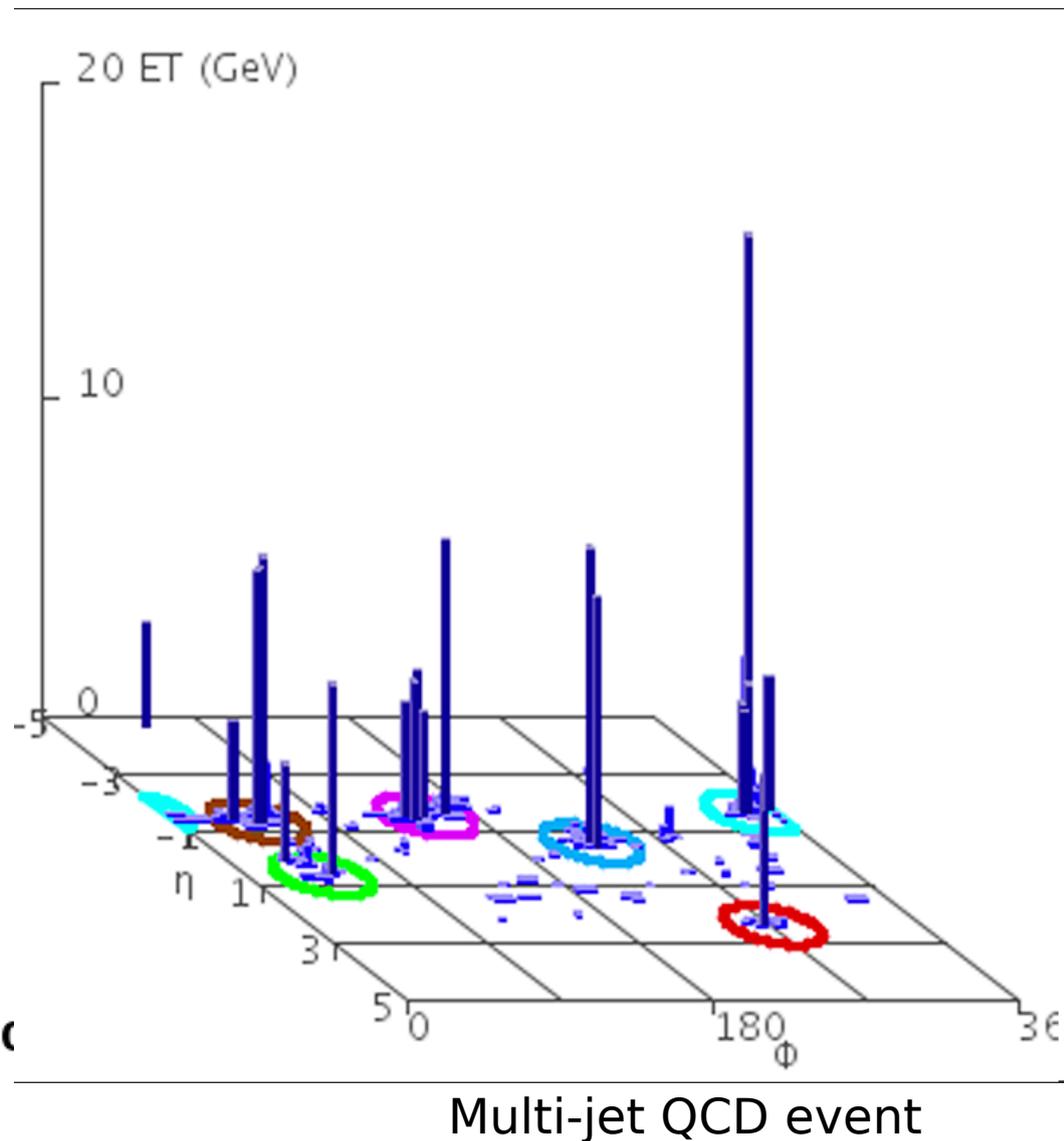
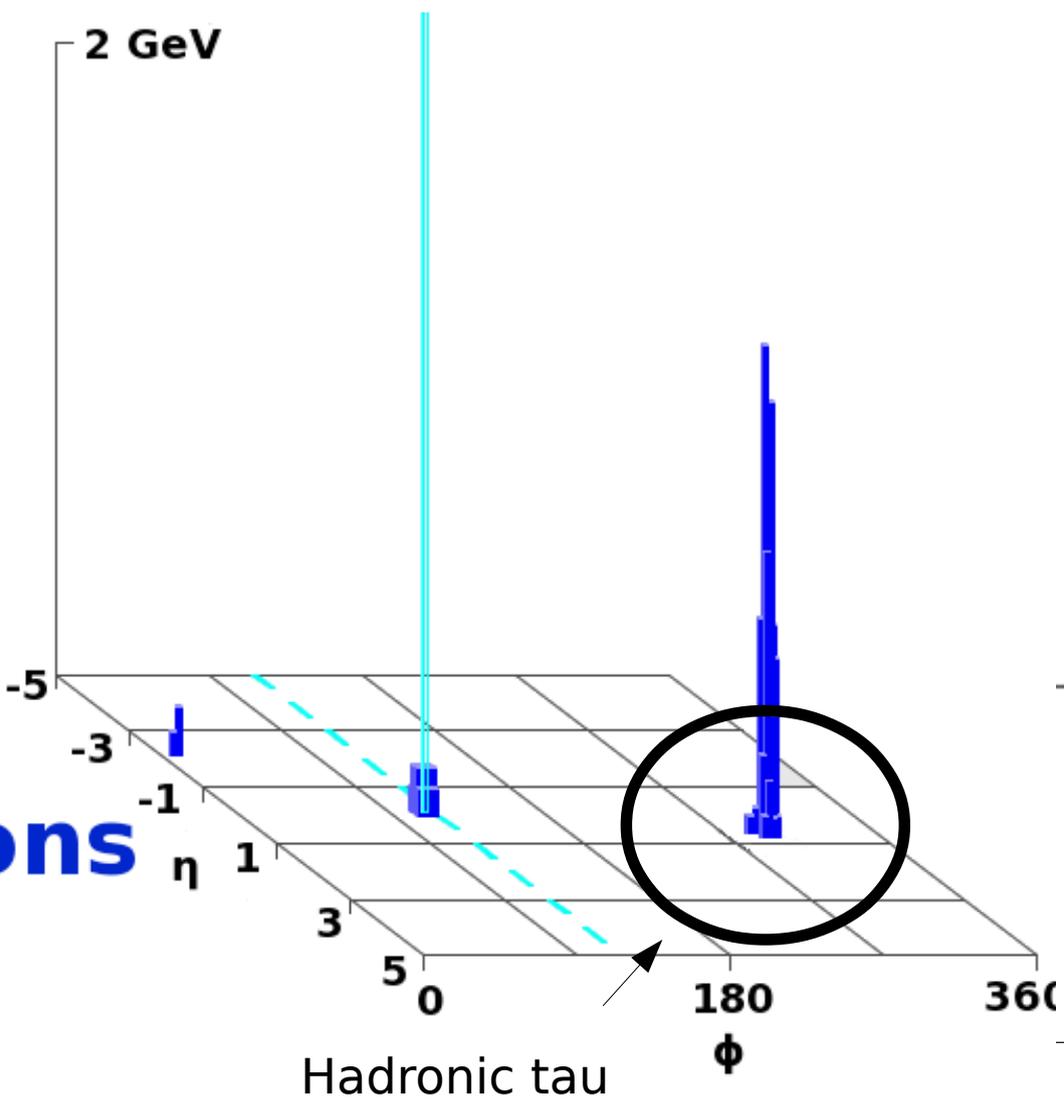
Run Number: 152409, Event Number: 8186656

Date: 2010-04-05 12:28:45 CEST

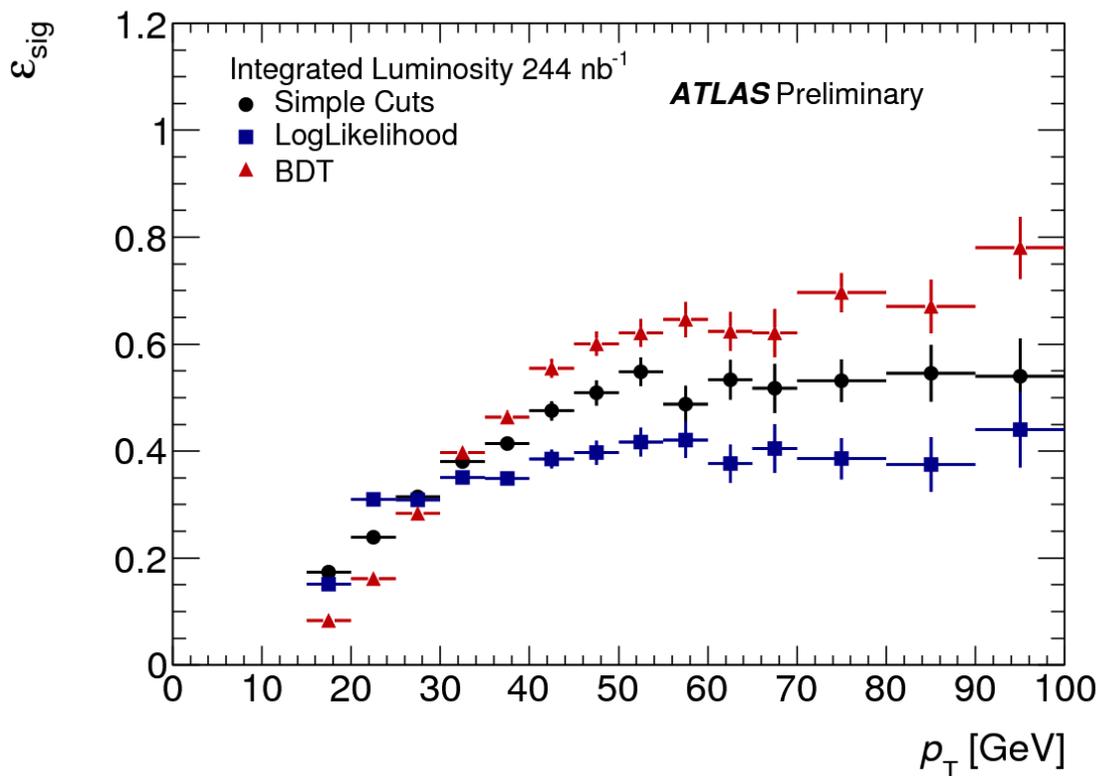
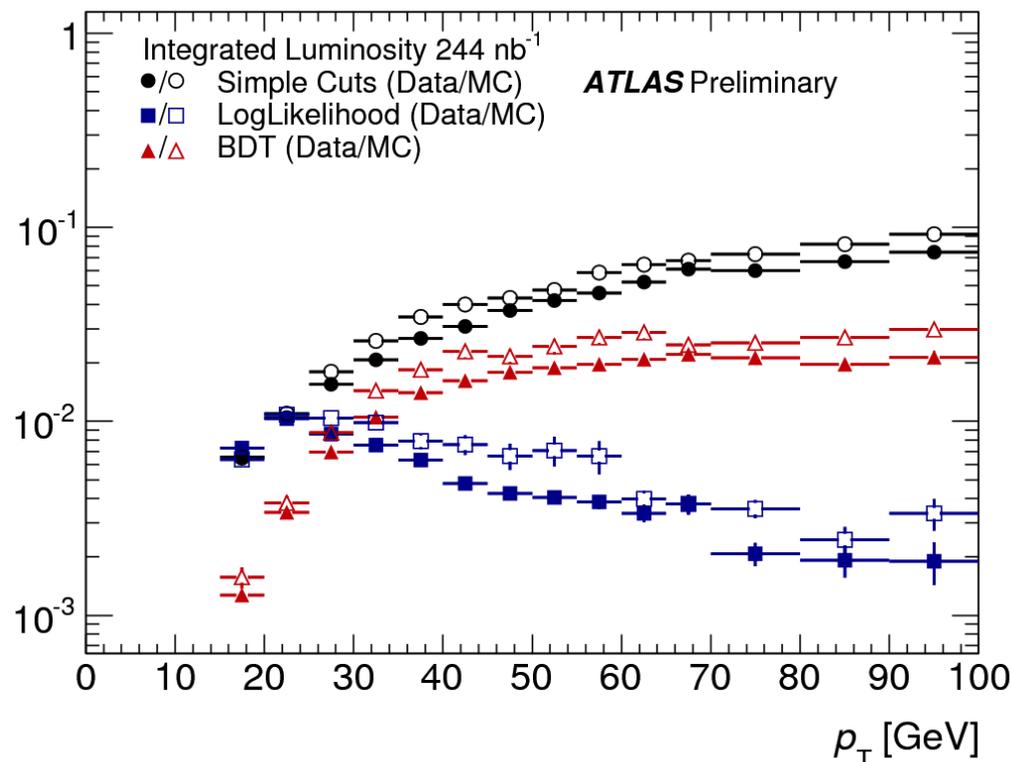


6 Jet Event in 7 TeV Collisions





Tau Selection for $H^+ \rightarrow \tau^+ \nu$



- We use the likelihood-based selector for the charged Higgs search
- Improves on background efficiency by ~order of magnitude while maintaining an acceptable signal efficiency
- We measure our τ signal efficiency using $Z \rightarrow \tau\tau$ events in data and find ~30% efficiency

Baseline Event Selection

- Trigger:
 - require at least 1 tau object ($p_T > 29$ GeV) and $MET > 35$ GeV
- Event refinement
 - ≥ 4 jets (excluding tau jets) with $p_T > 20$ GeV and $|\eta| < 2.5$
 - 1 tight-identified tau jet, 1 or 3 prongs, $p_T > 35$ GeV and $|\eta| < 2.3$
 - veto events containing identified electrons or muons
 - $MET > 40$ GeV
 - Large fluctuations in the absolute measured MET are possible, event-by-event, due to the limited energy resolution on jets, etc. We reject large METs due to this effect by requiring

$$\frac{MET}{(1/2)\sqrt{\sum E_T^{\text{visible}}}} > 8 \text{ GeV}^{1/2}$$

- Top mass, the mass of the highest- p_T (jj)b system, must satisfy $[120, 240]$ GeV
- Use transverse mass of the tau+MET system to finalize the analysis

Transverse Mass

From Phys. Rev. Lett. 50, 1738-1740 (1983)

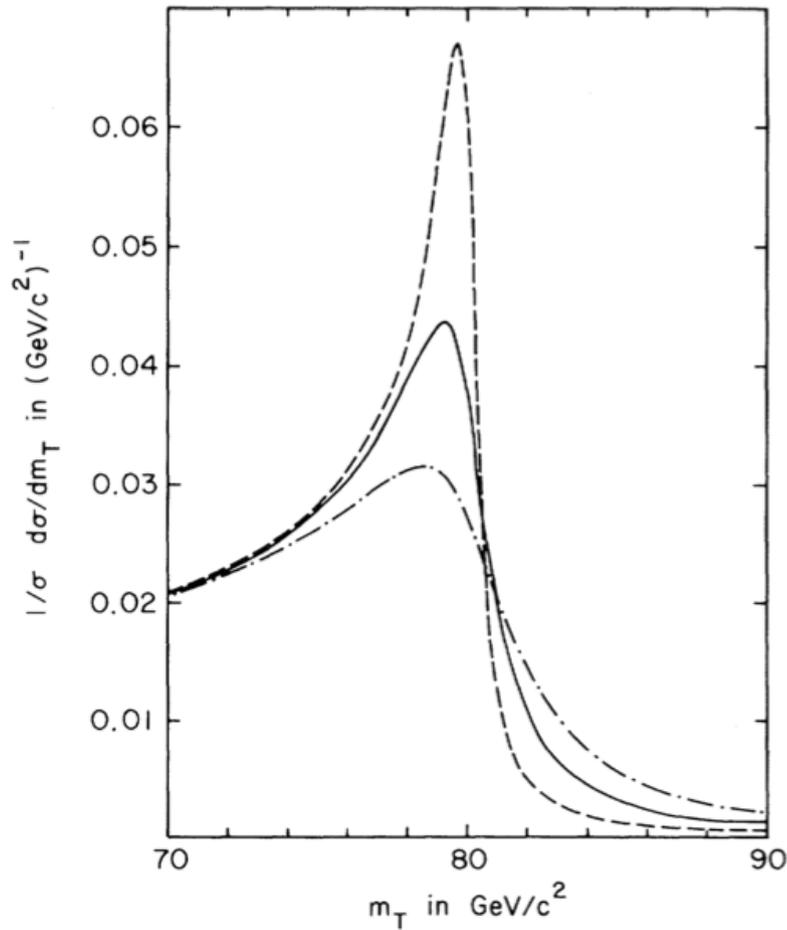


FIG. 2. $\sigma^{-1}d\sigma/m_T$ for $M=80 \text{ GeV}/c^2$ and $p_T^W = 0 \text{ GeV}/c^2$. The dashed, solid, and dot-dashed lines refer to $\Gamma = 1, 2.5, \text{ and } 5 \text{ GeV}/c^2$, respectively.

General Definition:

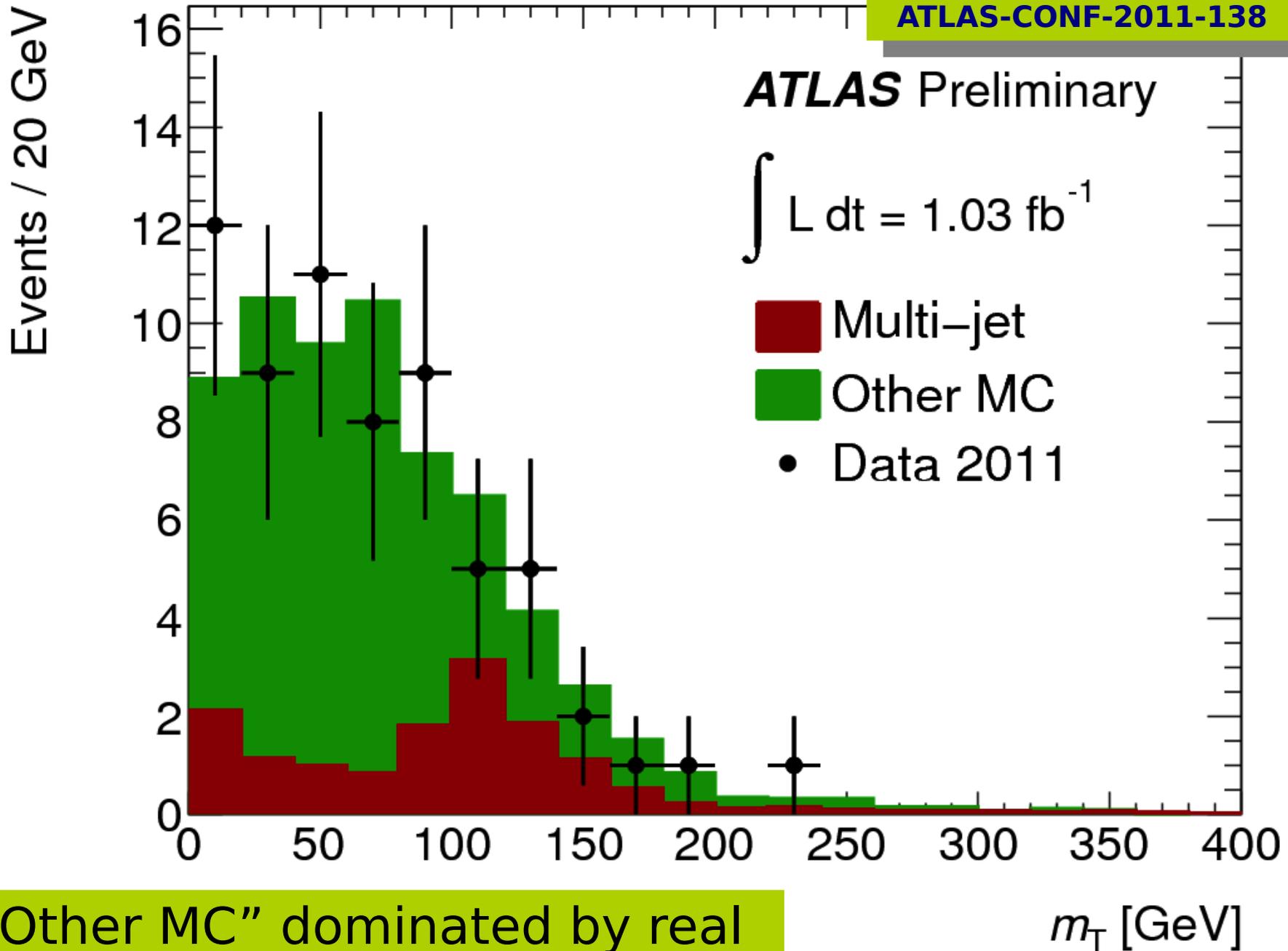
$$M_T^2 \equiv (E_T^l + E_T^{\nu})^2 - (\vec{p}_T^l + \vec{p}_T^{\nu})^2$$

Specific to $W \rightarrow l \nu$ (where final-state particles are essentially “massless” on these energy scales):

$$M_T^2 \equiv 2 p_T^l E_T^{\text{miss}} (1 - \cos \Delta \phi)$$

Transverse mass of the tau+MET system in tau + jets + MET events

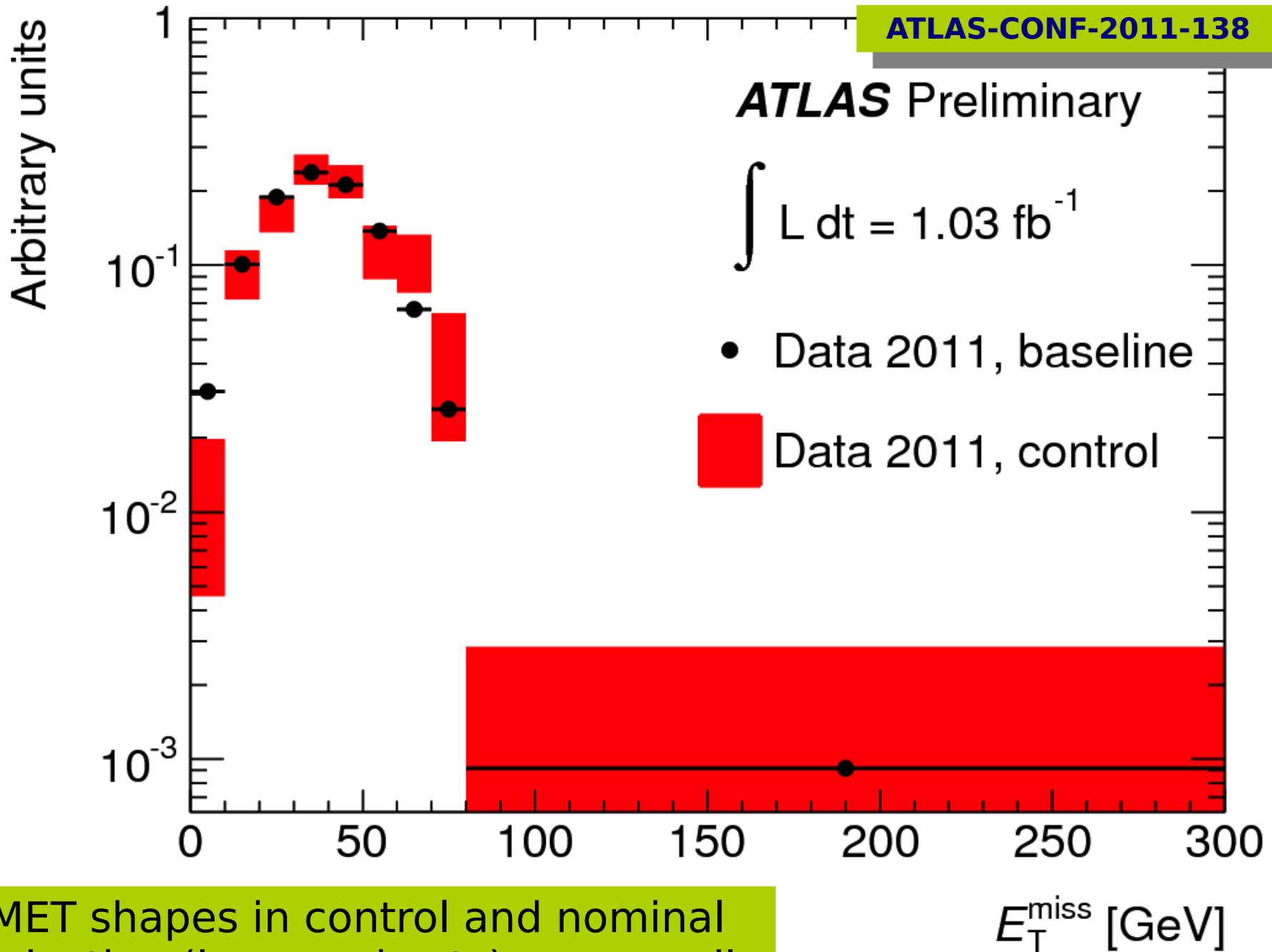
ATLAS-CONF-2011-138



“Other MC” dominated by real Standard Model $t\bar{t}$ events

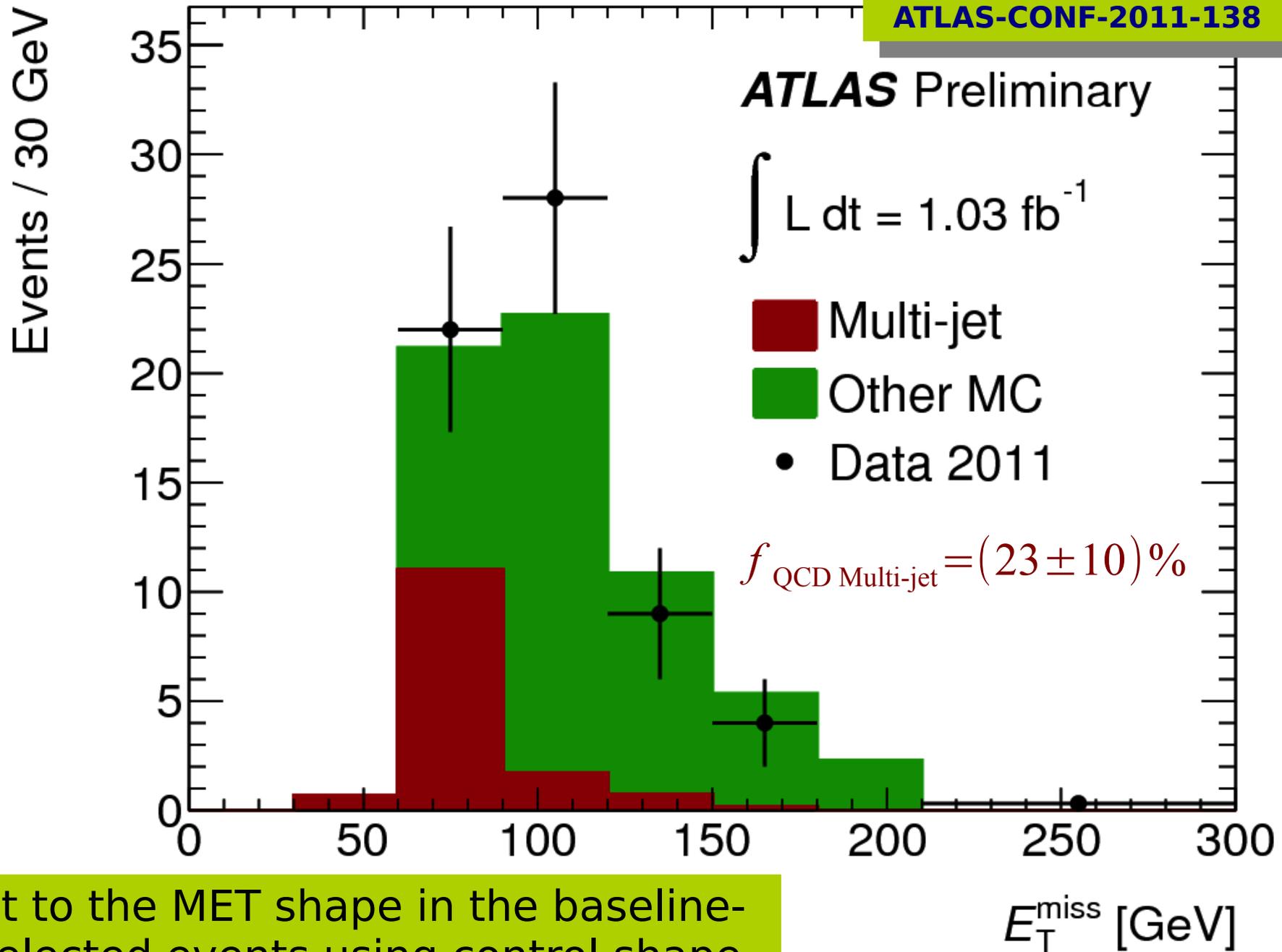
QCD Background Estimate

- MC Simulations of the QCD background are insufficient to understand this background
 - too few events are generated for a precision study – few/none survive all of our selection criteria
 - variable shapes not in great agreement at early stages of the selection when there is more QCD MC available
 - use data-driven method to study this background
- Method – define a “control sample” [A. Randle-Conde]
 - invert the b-tagging and tau selection cuts
 - No b-tagged jet AND no tau candidate passing the nominal criteria
 - subtract $t\bar{t}$ and other well-modeled MC shapes from MET distribution
 - Look at MET in these control events and develop a template to fit the MET shape in baseline-selection events



MET shapes in control and nominal selection (loosened cuts) agree well

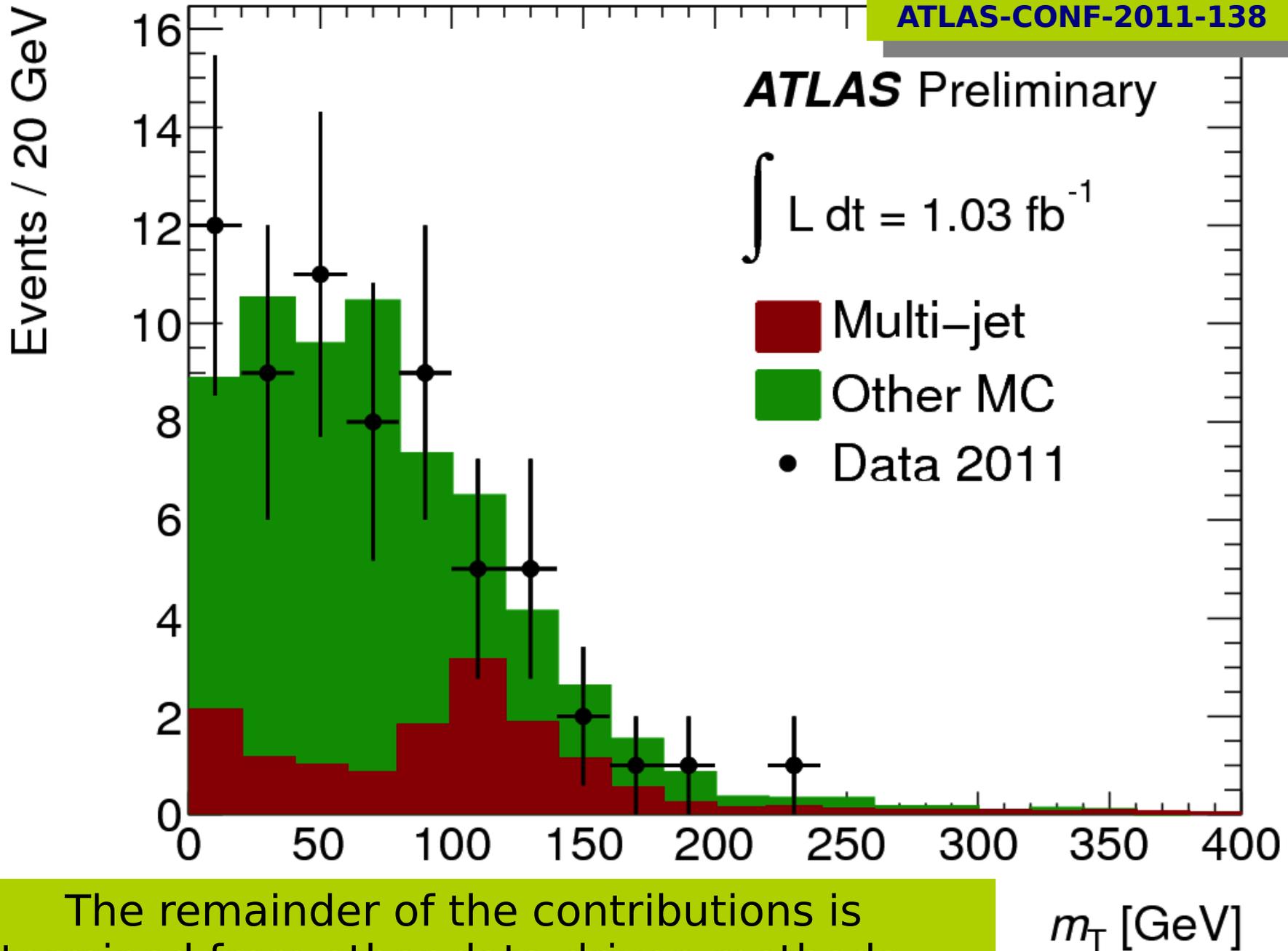
E_T^{miss} [GeV]



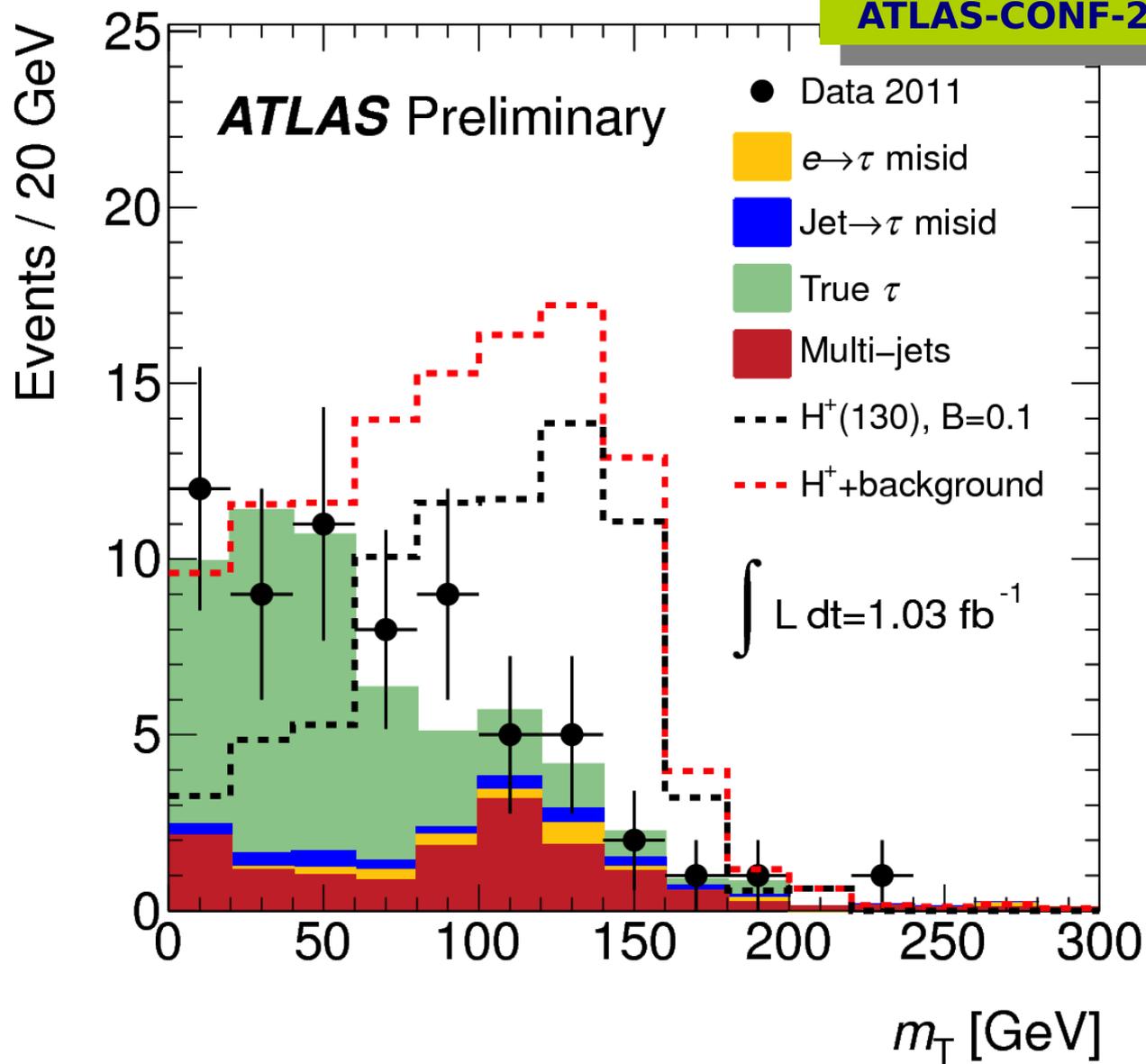
Fit to the MET shape in the baseline-selected events using control shape and MC (e.g. SM $t\bar{t}$).

Transverse mass of the tau+MET system in tau + jets + MET events

ATLAS-CONF-2011-138



The remainder of the contributions is determined from other data-driven methods . . .



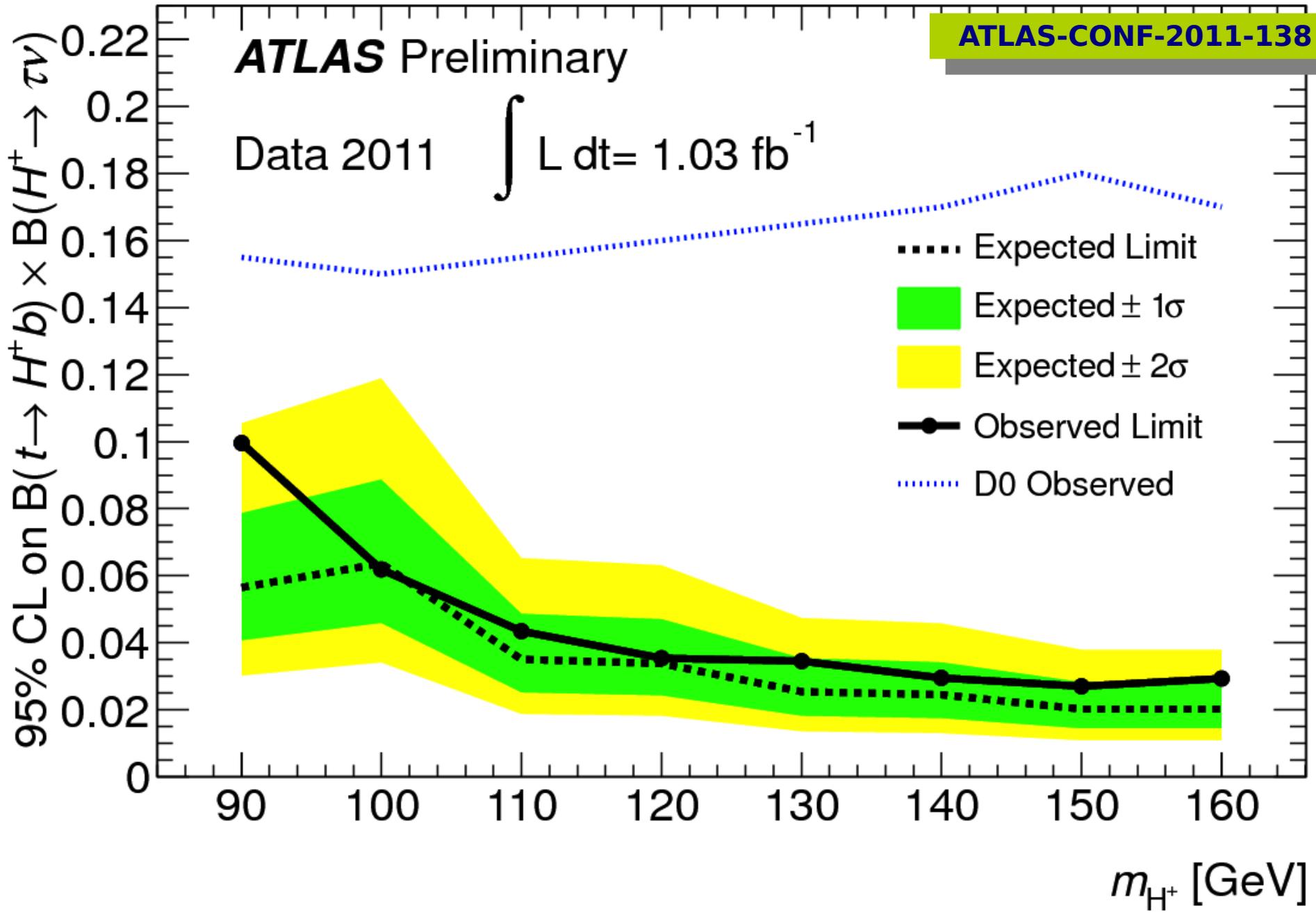
	Events with/from				expected (sum)	data
	true τ jets	jet $\rightarrow \tau$ mis-id	$e \rightarrow \tau$ mis-id	multi-jet		
$m_T > 40 \text{ GeV}$	21 ± 5	2.4 ± 0.7	1.9 ± 0.2	12 ± 5	37 ± 7	43

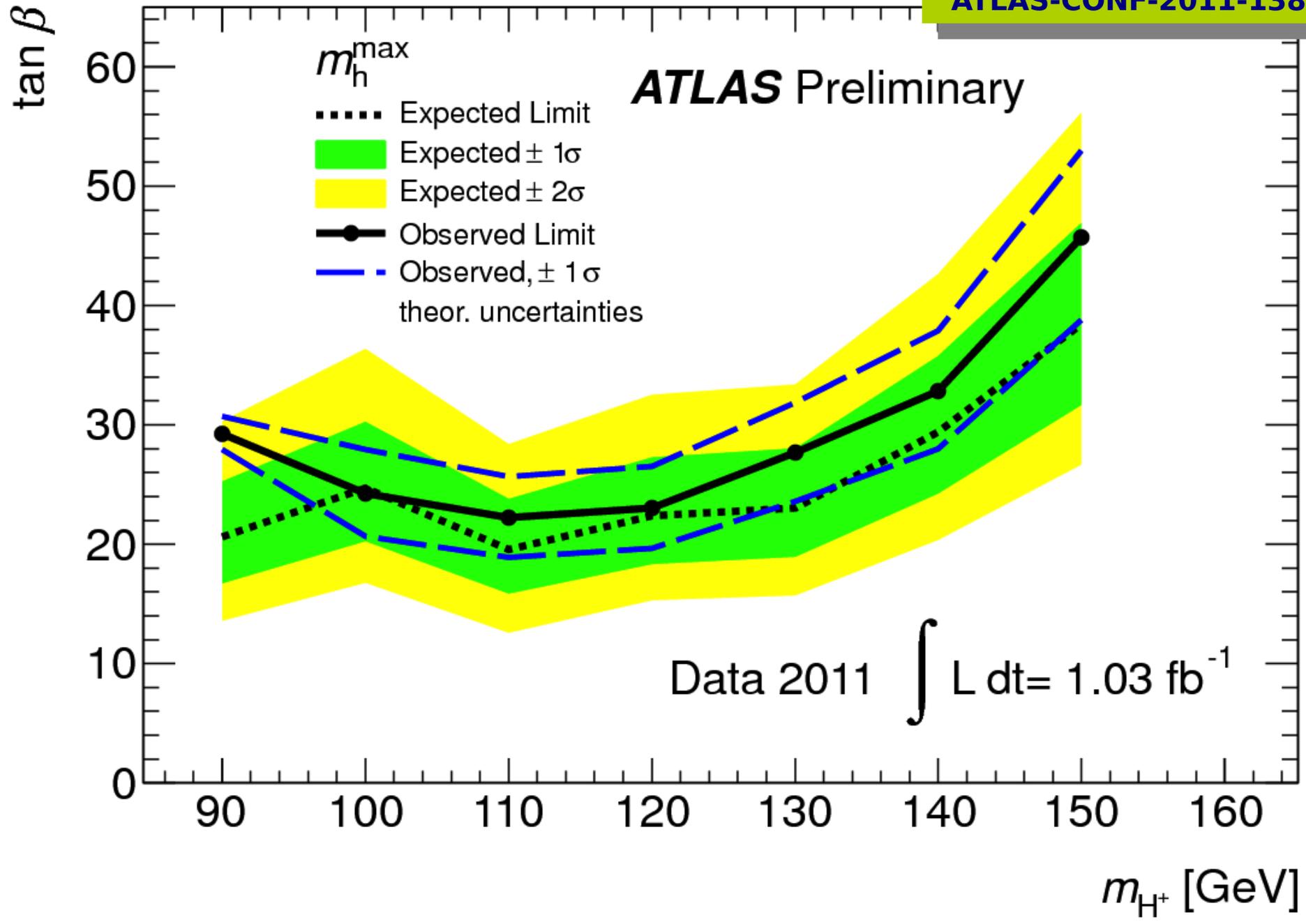
Systematics

Quantity	Uncertainty
Luminosity [36]	$\pm 3.7\%$
Jet energy resolution (JER)	$\pm(10 - 30)\%$, depending on p_T and η
Jet energy scale (JES)	$\pm(2.5 - 14)\%$, depending on p_T and η
E_T^{miss}	Uncertainty due to scale/resolution uncertainties (e.g. JES); additional 10% of pile-up-related uncertainty
b -tagging efficiency SF unc.	$\pm(0.05 - 0.15)$, depending on p_T and η
b -tagging mistag rate	$\pm(0.16 - 0.39)$, depending on p_T and η
b jets JES uncertainty	an additional $\pm 2.5\%$ on top of the standard JES
τ identification efficiency	$\pm(8.5 - 9.9)\%$, depending on p_T
τ energy scale	$\pm(4.5 - 6.5)\%$, depending on p_T , η , number of associated tracks
τ electron mis-id correction factors	$\pm(23 - 100)\%$, depending on η ; for one-prong only
$\tau + E_T^{\text{miss}}$ trigger	$\pm 9\%$
e reco. efficiency SF	$\pm(0.7 - 1.8)\%$, depending on η
e identification efficiency SF	$\pm(2.2 - 3.8)\%$, depending on E_T and η
e energy scale	$\pm(0.3 - 1.8)\%$, depending on p_T and η
e energy resolution	$\pm(0.5 - 2.4)\%$ (additional constant term), depending on p_T and η
μ reco. efficiency SF	$\pm(0.25 - 0.55)\%$, depending on the data-taking period
μ momentum scale and resolution	$\pm(0.4 - 0.7)\%$, depending on η
Initial/final state radiation modelling	$-16\% / +19\%$ ($t\bar{t}$ signal and background)
Acceptance	$\pm 4\%$ (background), $\pm 10\%$ (signal)
$t\bar{t}$ cross section	$165_{-9}^{+4}(\text{scale})_{-7}^{+7}(\text{pdf}) \text{ pb}$

ATLAS Preliminary

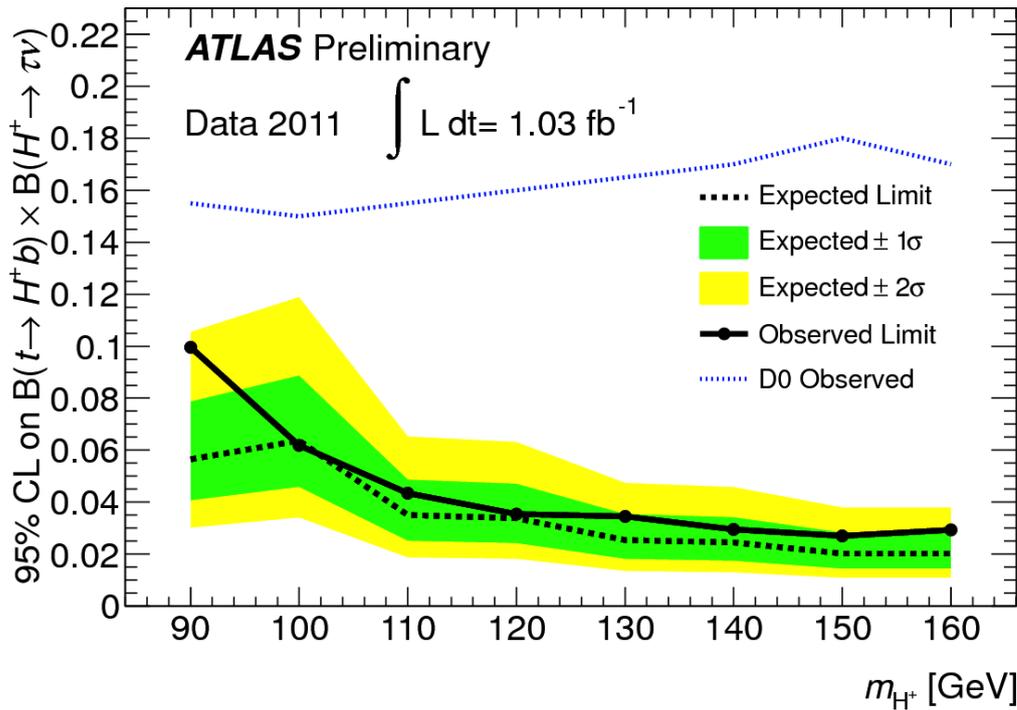
Data 2011 $\int L dt = 1.03 \text{ fb}^{-1}$



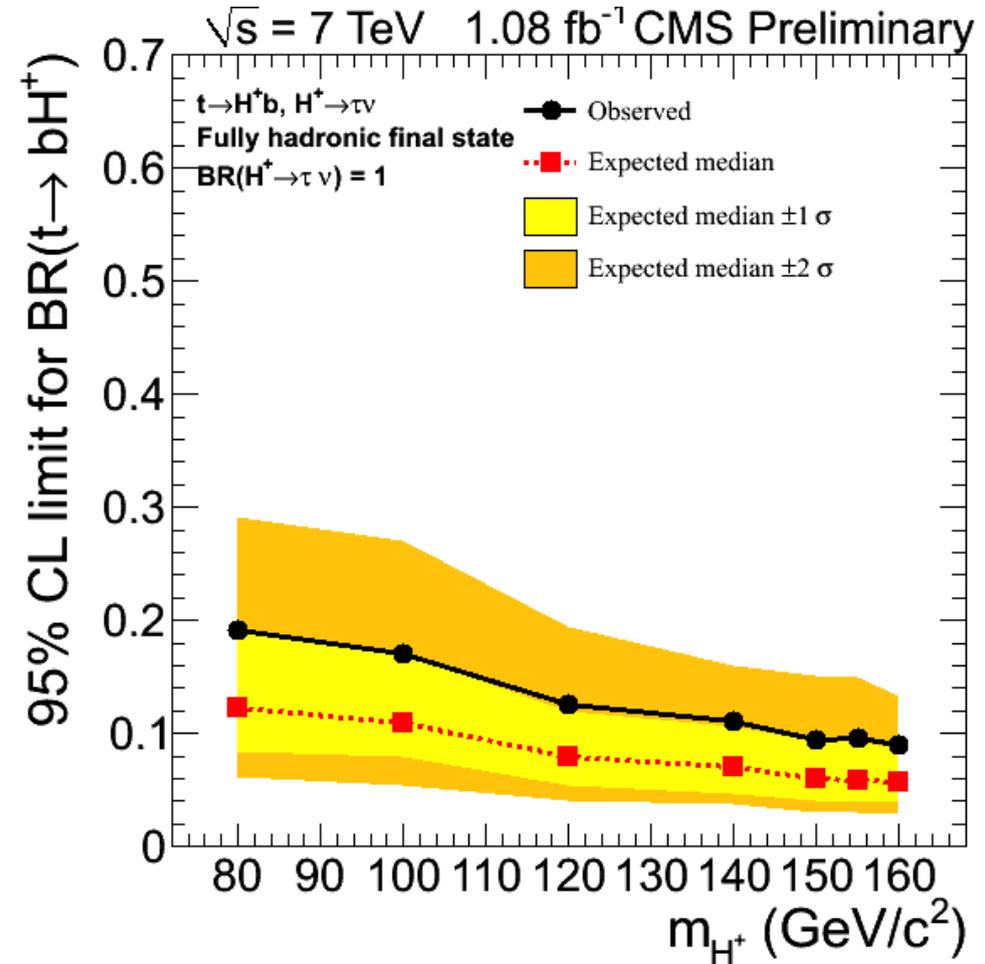


Comparison to other measurements

hadronic tau + jets

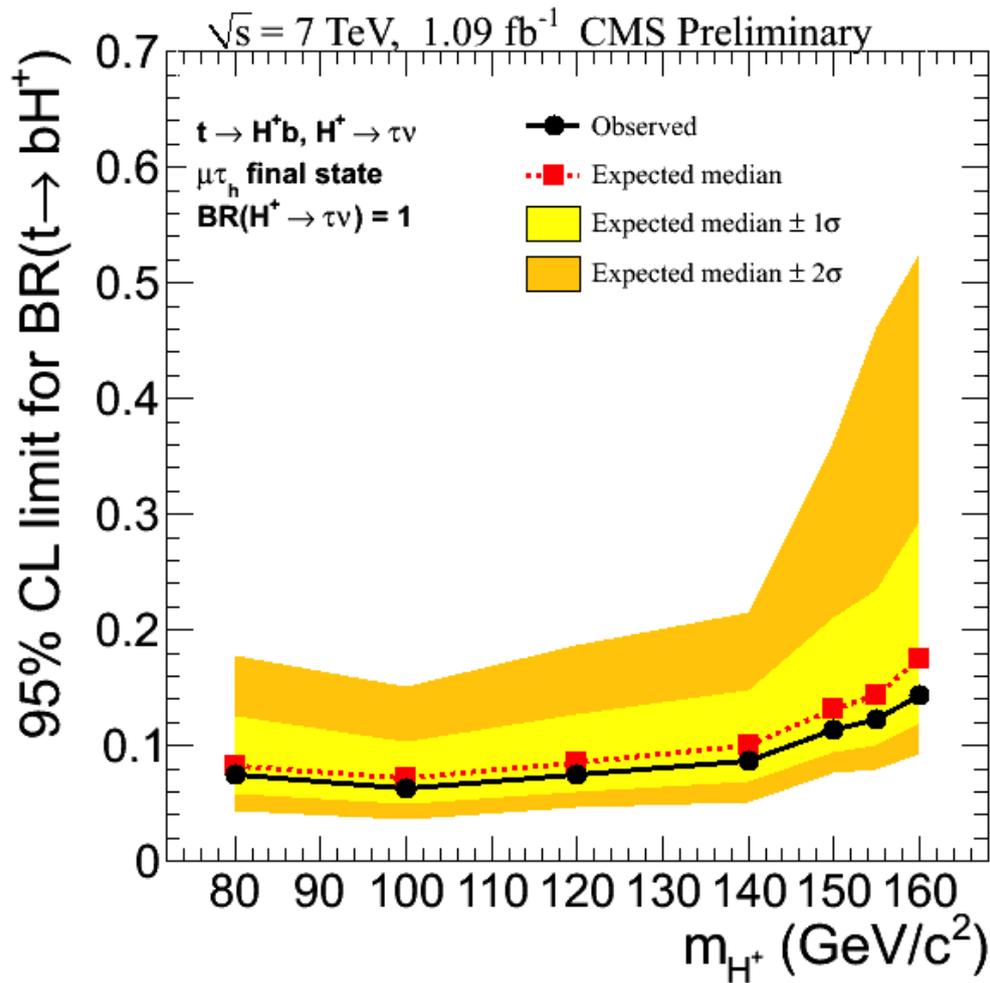


ATLAS-CONF-2011-138



CMS-PAS-HIG-11-008

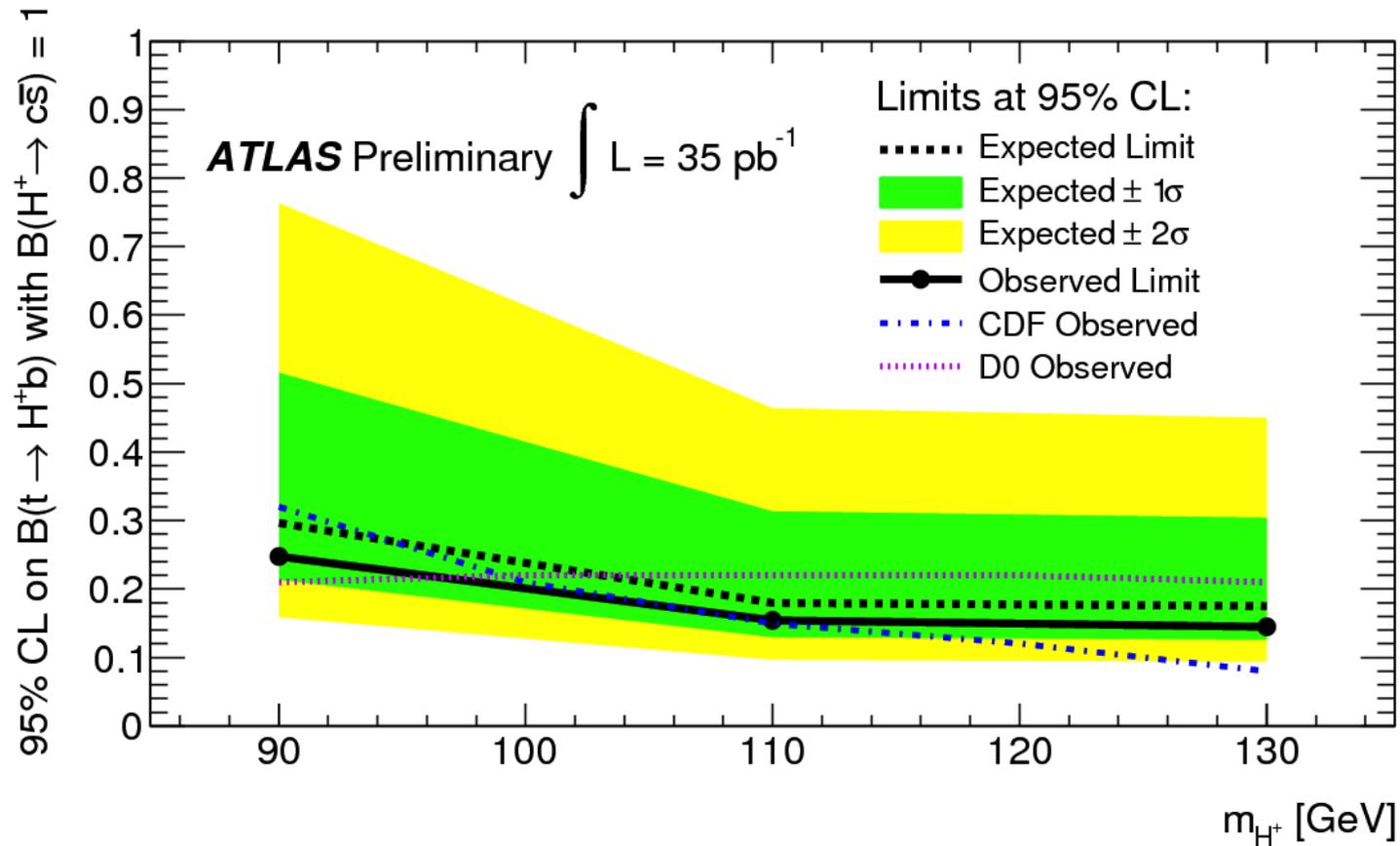
Comparison (continued)



ATLAS tau+lepton search was not ready in time for summer conferences but will be ready for our publication on this topic.

Comparison (cont.)

ATLAS-CONF-2011-094



ATLAS is pursuing other final states besides $\tau\nu$

Future Directions/Issues

- Increased Pileup Collisions
 - we are already at $L = 3.3 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$
 - corresponds to $\langle \mu \rangle \sim 11$ per crossing
 - at $L = 2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ (~ 2016), $\langle \mu \rangle \sim 46$
 - Impact on MET, tau, and jet triggers
 - Impact on offline reconstruction of MET, etc.
 - We are involved in trigger rates and object triggering studies for these conditions
- New methods for searching
 - $H^+ \rightarrow W^+ Z^0$
 - $H^+ \rightarrow tb$
 - $H^+ \rightarrow \text{SUSY}???$

Conclusions

- An exciting time for discovery-oriented physics
 - when keeping up with the data becomes your biggest problem, you're living in a golden age
- Direct tests of SUSY/MSSM hypothesis are possible at the target energy scales
 - SUSY is supposed to be interesting once you get to around 1 TeV
 - Here we are – where are you, SUSY?
- Testing the Higgs Mechanism
 - The race to discover a SM Higgs is hotter than ever
 - If the Higgs is the mechanism for the attainment of mass, and there are more particles out there (SUSY or no), then we need more Higgses
- Unprecedented sensitivity to $t \rightarrow Hb$
 - We're approaching the 1% level in sensitivity – precision tests are coming
 - Need to open up the space for higher-mass charged Higgs searches