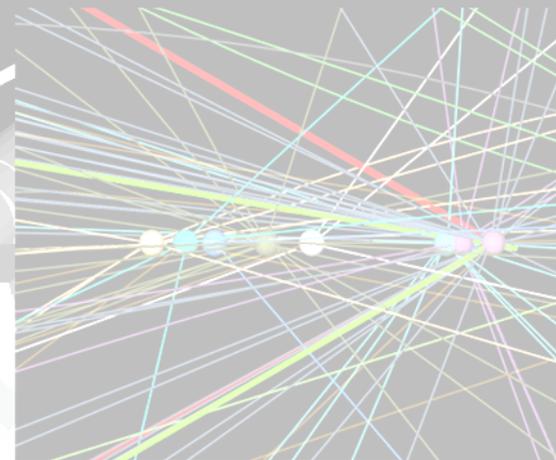
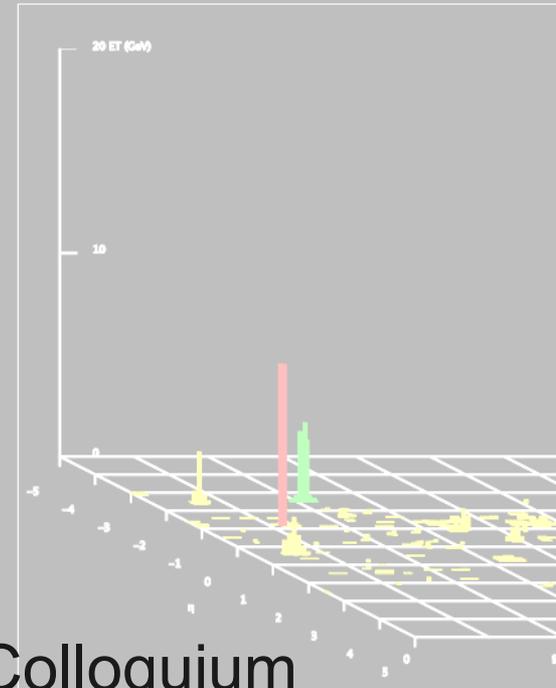


The Discovery of a New State with a Mass of $126 \text{ GeV}/c^2$

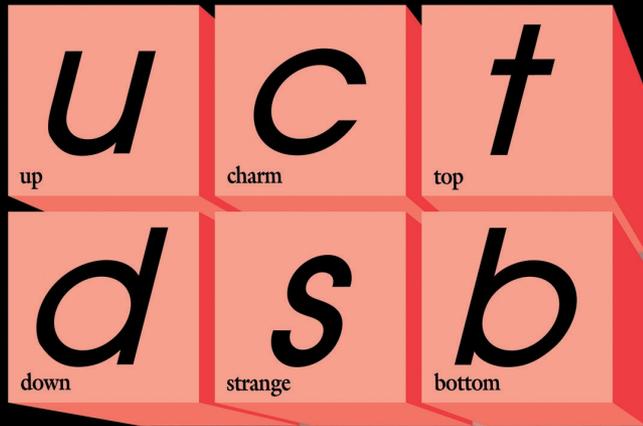
Stephen Sekula

Southern Methodist University
The ATLAS Collaboration

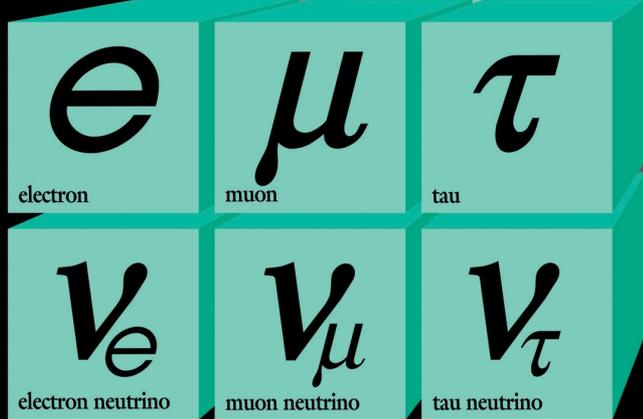
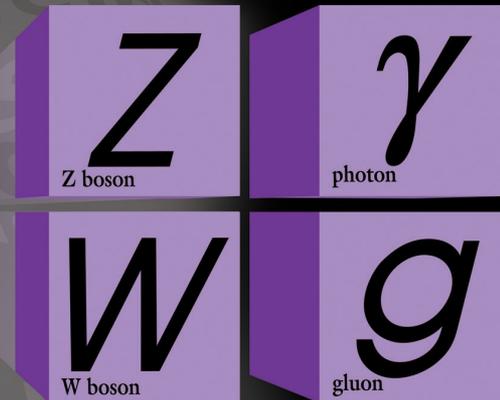
Presented at the SMU Physics Dept. Special Colloquium
September 13, 2012



Quarks

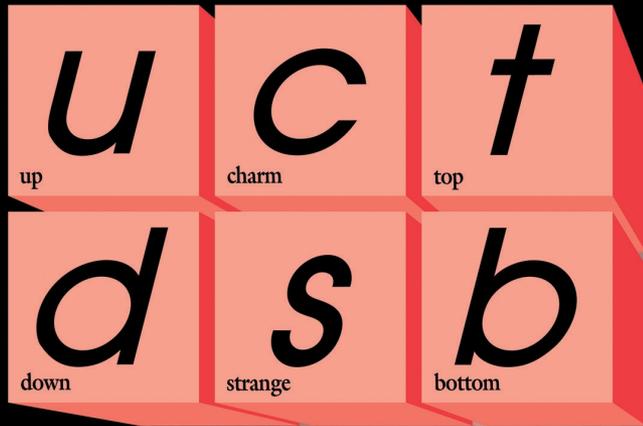


Forces

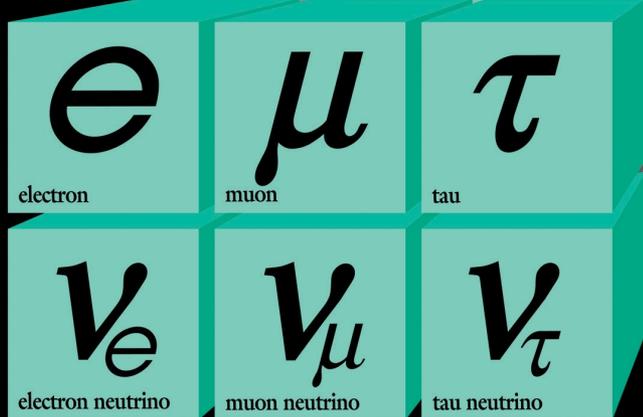
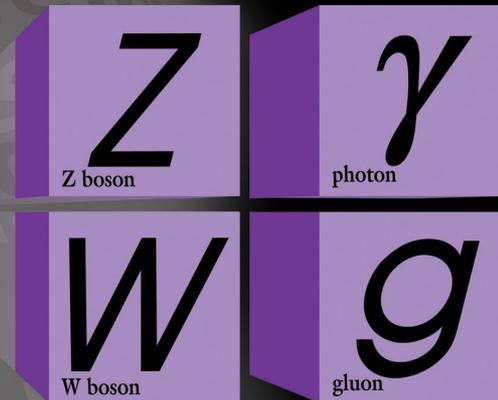


Leptons

Quarks



Forces



Leptons

According to the Standard Model (Electroweak symmetry breaking) this should be a spinless ($S=0$) particle, a boson (integer spin), whose couplings to other particles determines their mass.

PARTICLE MASS



PARTICLE SPIN



PARTICLE CP



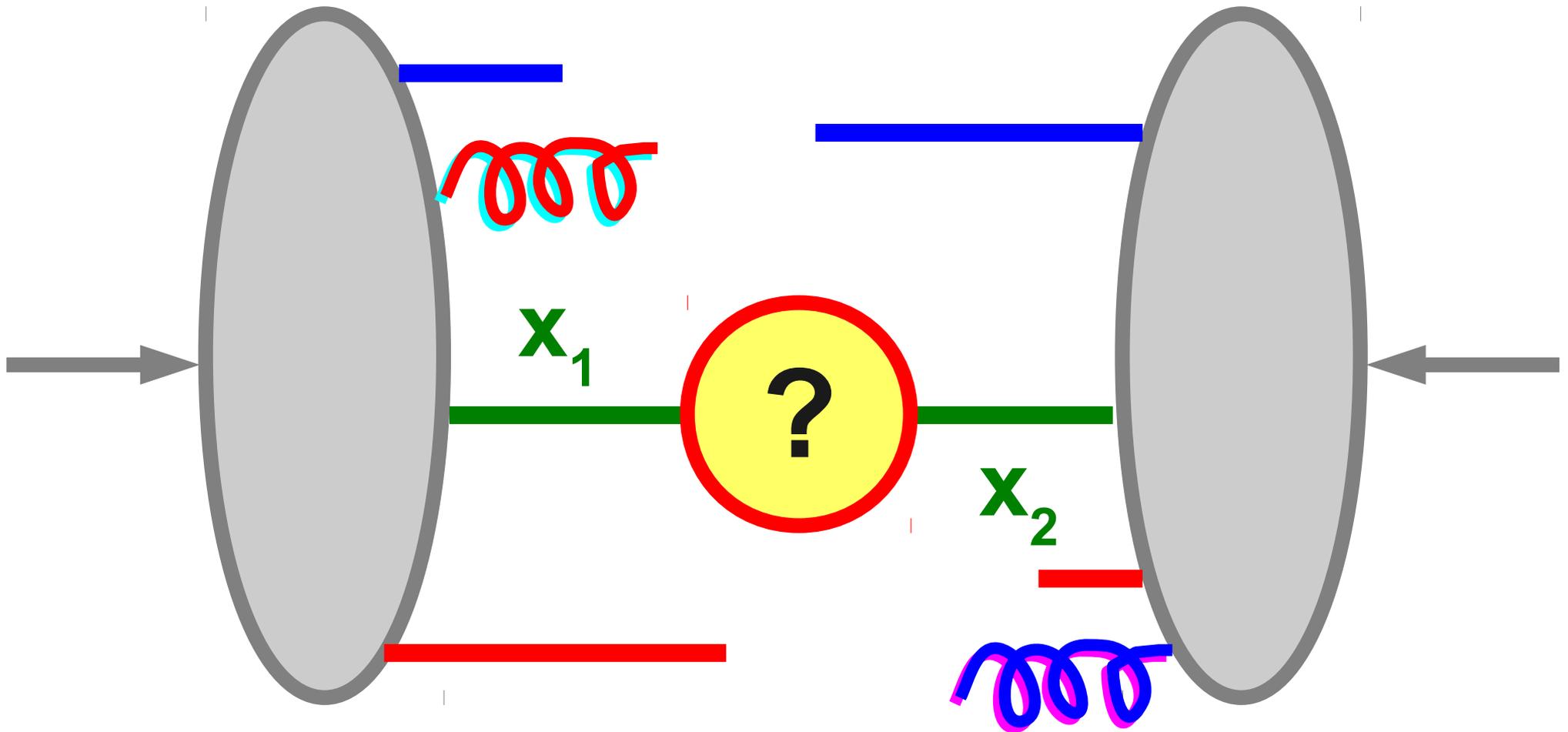
PARTICLE COUPLINGS

*[for instance:
rate of production
rate of decay]*



PROTON

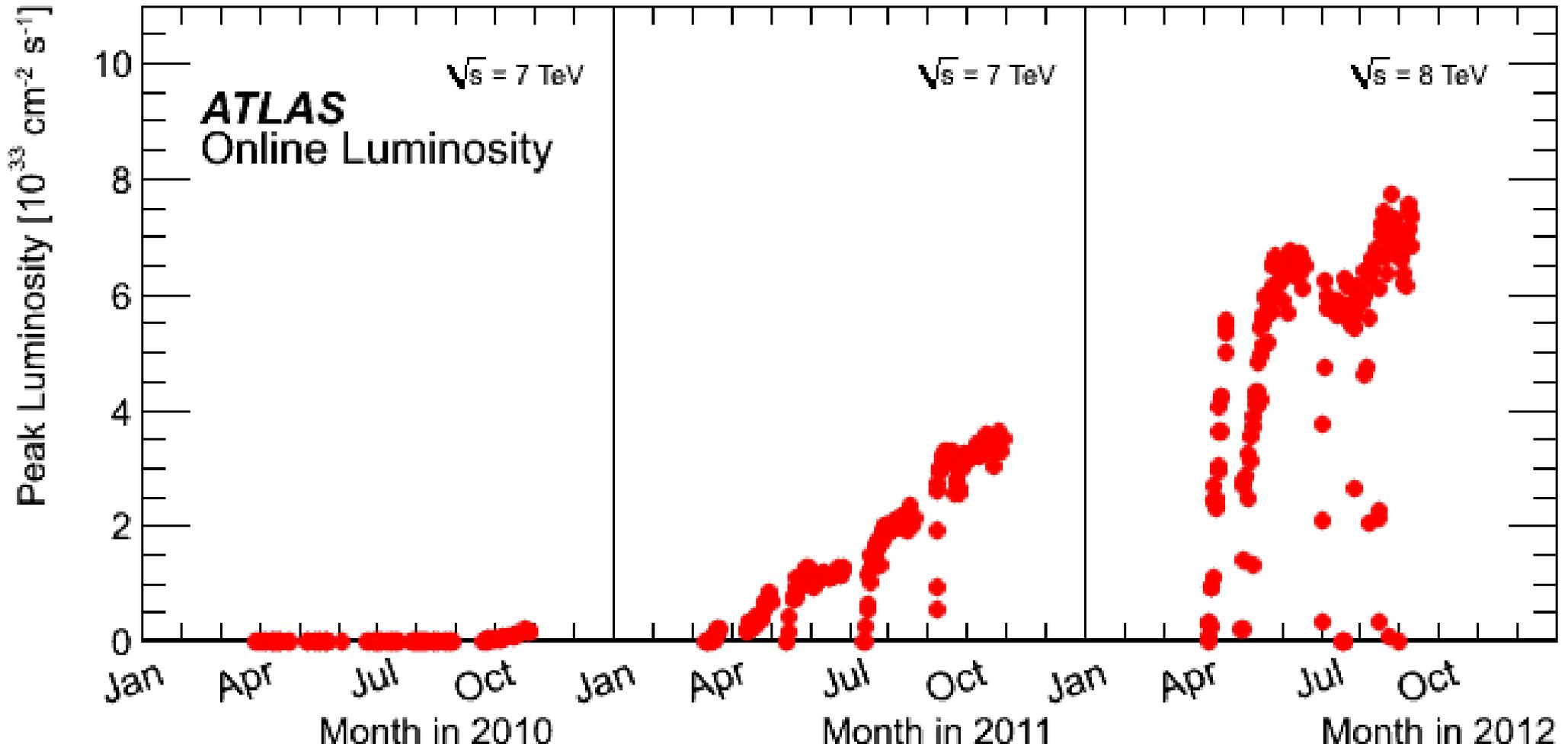
PROTON



The center-of-mass (CM) proton-proton energy is denoted by $E_{\text{CM}} = \sqrt{s}$, but the primary interaction only receives a fraction of the total energy of either proton. For a symmetric collider, the energy available for the interaction is $E_{\text{int}} = \sqrt{(x_1 x_2 s)}$





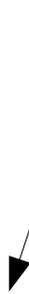


We routinely operate now at instantaneous luminosities of $(7.0\text{-}7.5)\times 10^{33}$ interactions per cm^2 per second ($\text{cm}^{-2}\text{s}^{-1}$).

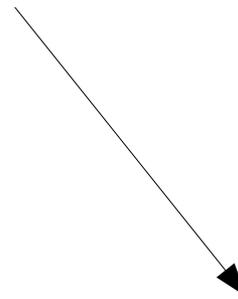
A convenient unit of area is a “barn” (as in, “You can’t hit the broad side of a barn.”) - it is the typical unit of cross-sectional area in nuclear and subnuclear physics.

1 barn = 10^{-24} cm^2 , so in these units the LHC achieves luminosities of $\sim 7.5 \times 10^9 \text{ b}^{-1} \text{ s}^{-1}$, or in units more appropriate to the scale of LHC physics, about $7.5 \times 10^{-6} \text{ fb}^{-1} \text{ s}^{-1}$

$$\mathcal{L} = \int L dt$$

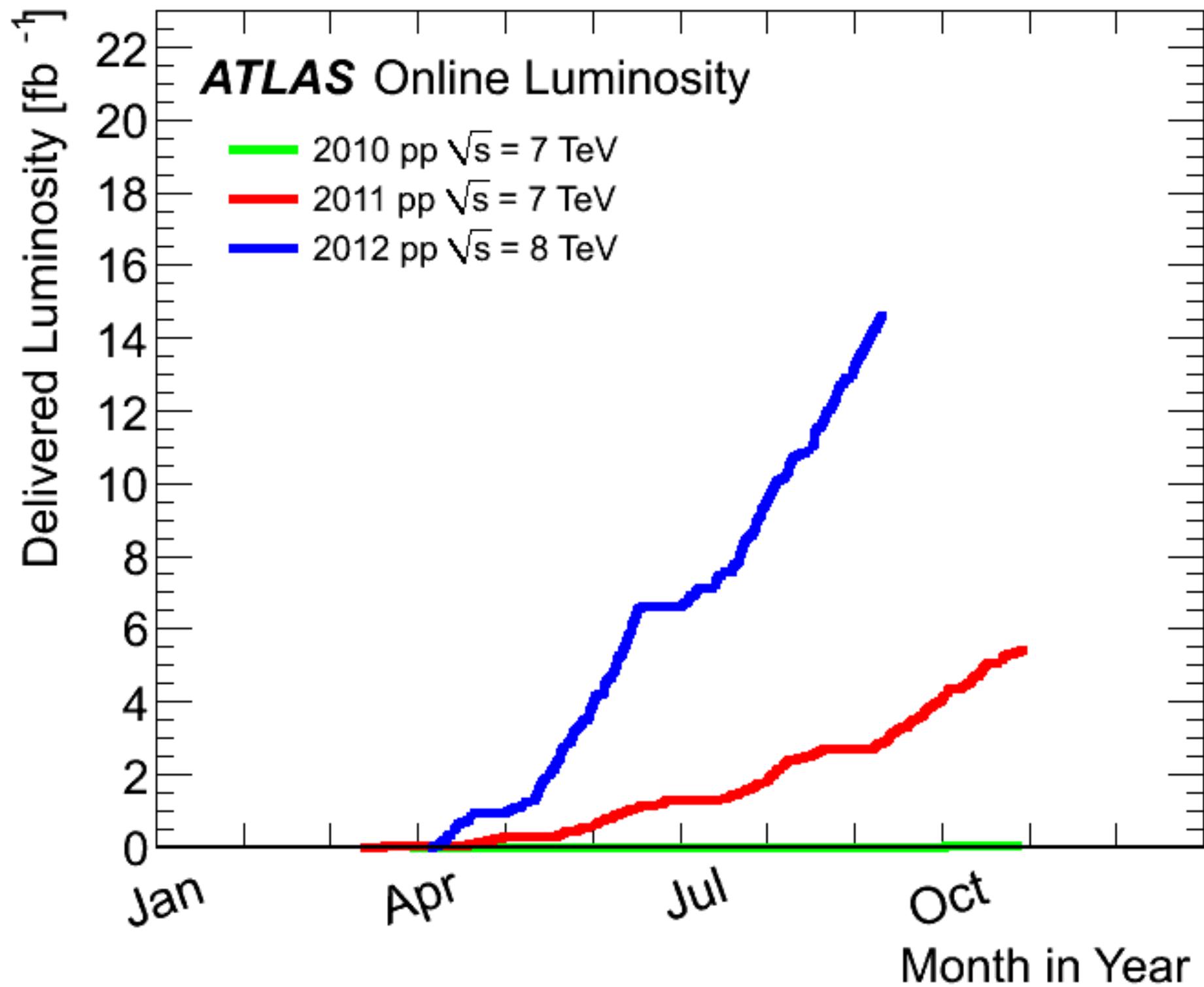


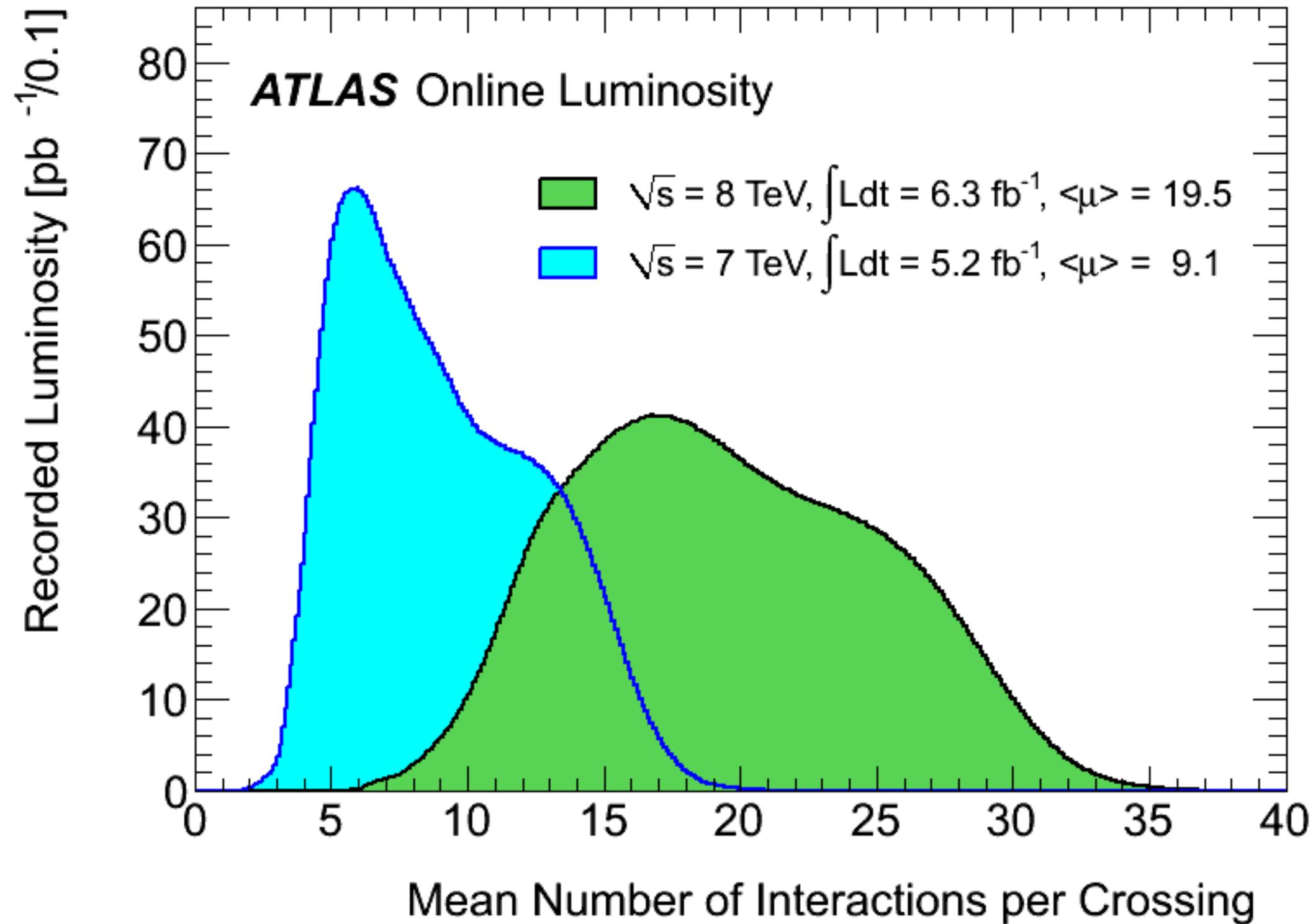
“Integrated Luminosity”
measured in fb^{-1}



“Instantaneous Luminosity”
measured in $\text{fb}^{-1} \text{ s}^{-1}$

LHC data sample sizes are measured in units of **fb^{-1}** because once you know this number, and you know the cross-section (probability of producing) a specific process, you can estimate how often that process will appear in your collider.

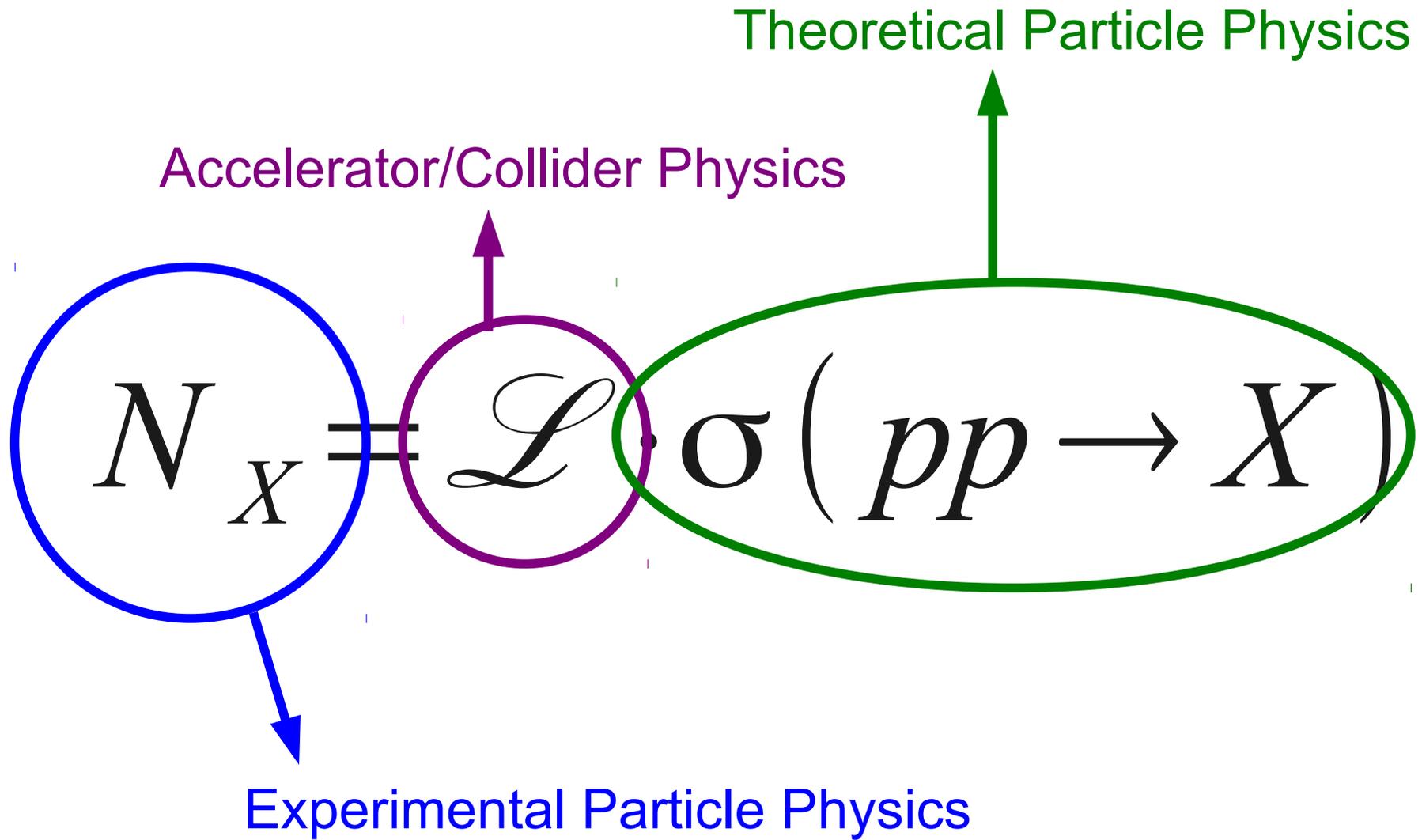




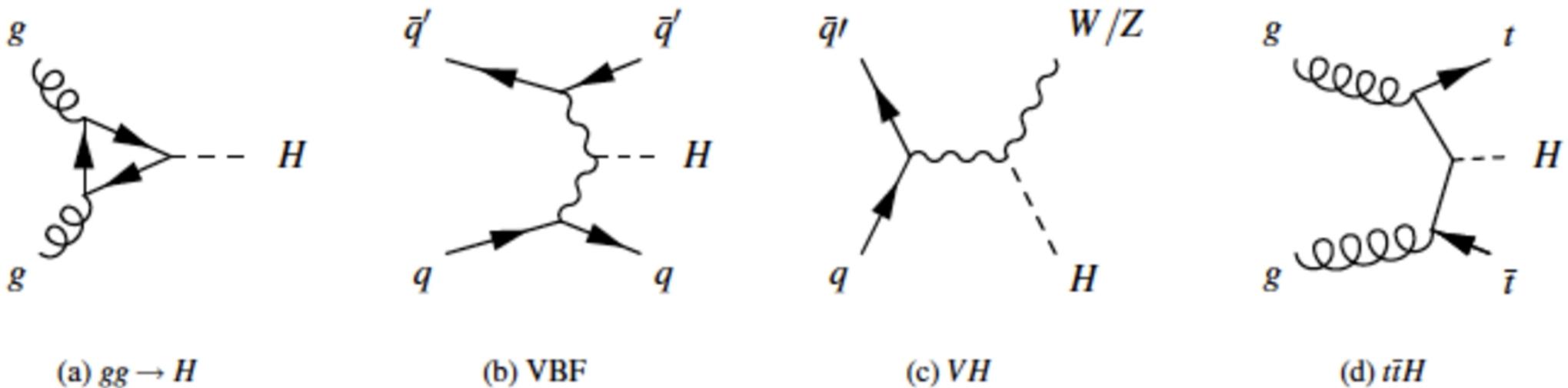
$$N_X = \mathcal{L} \cdot \sigma (pp \rightarrow X)$$

Example:

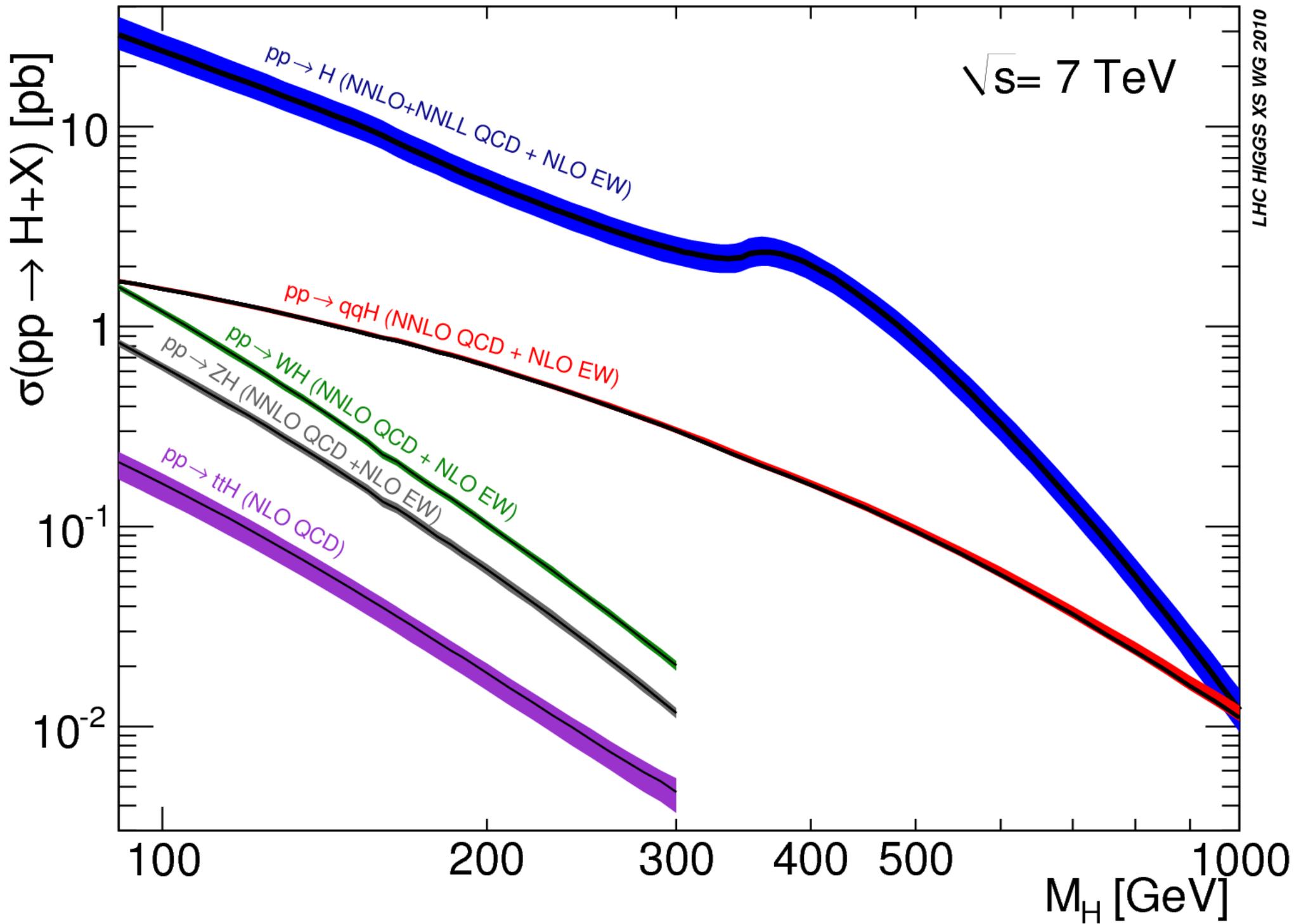
- Let's say that the cross-section for producing the Higgs with a certain mass at a certain center-of-mass collision energy is **1 pb**
- Let's say that the total integrated luminosity of the collider is **5 fb⁻¹ = 5,000 pb⁻¹**
- Then the total number of collisions that produce a Higgs is **N = 5,000**
- This DOESN'T take into account detector acceptance, trigger efficiency, selection efficiency, etc.

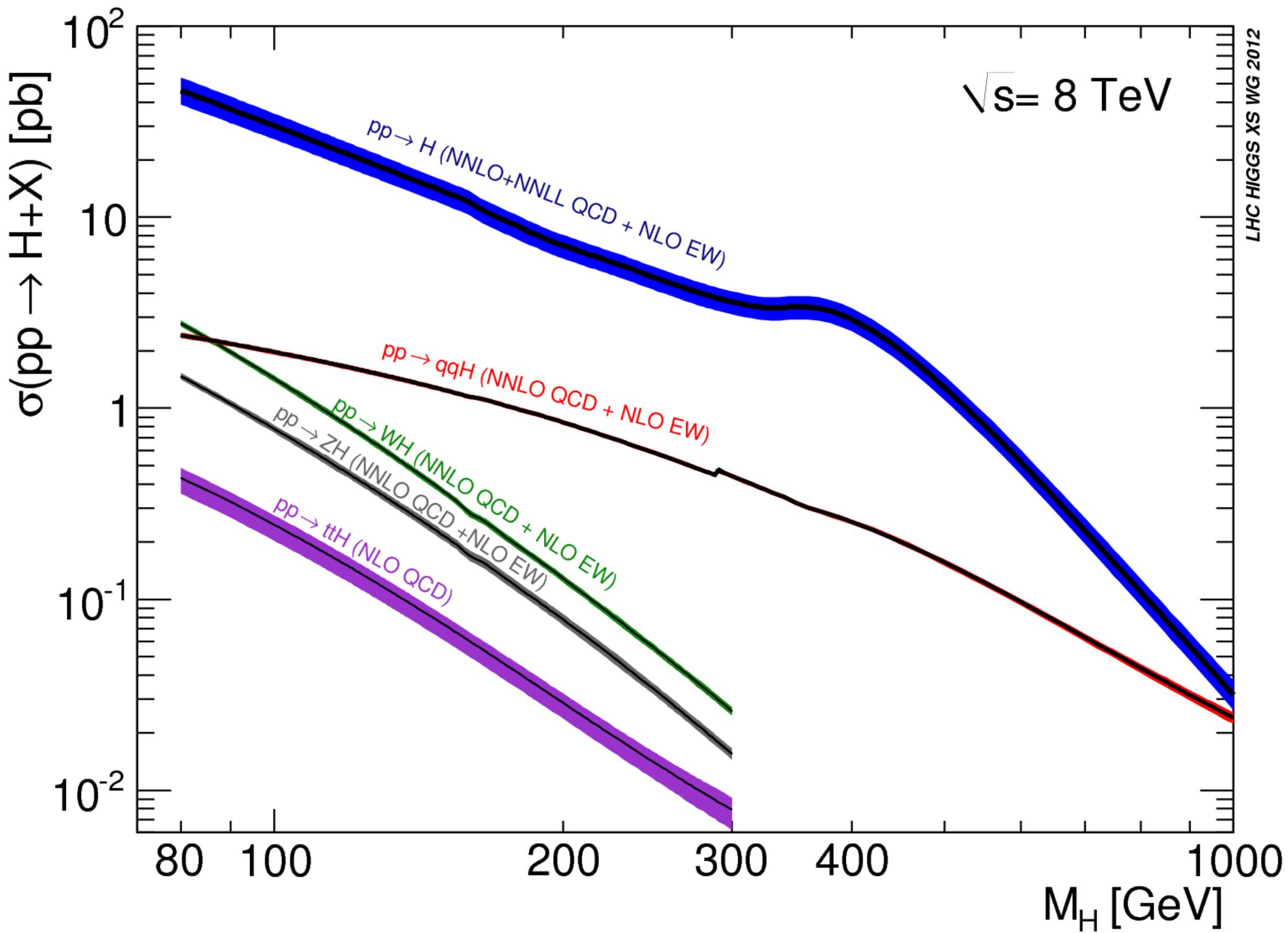


How should the Higgs be
produced?



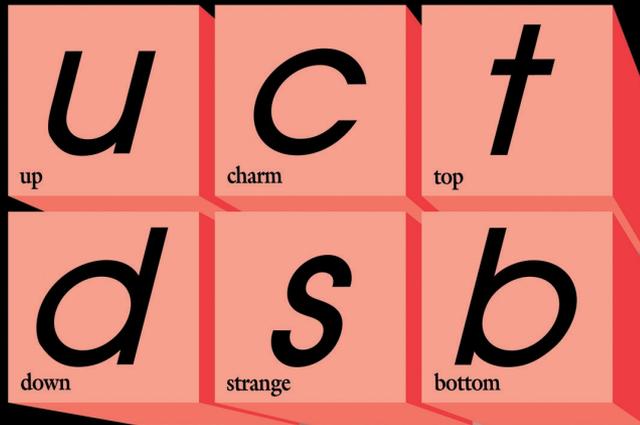
These are “leading order” diagrams – there are many higher-order diagrams that can also contribute for each of these processes.



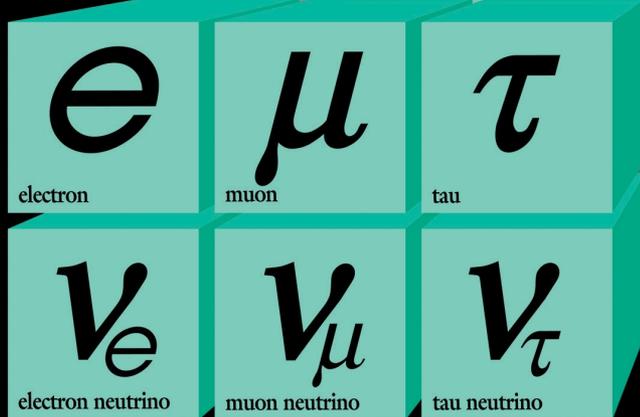
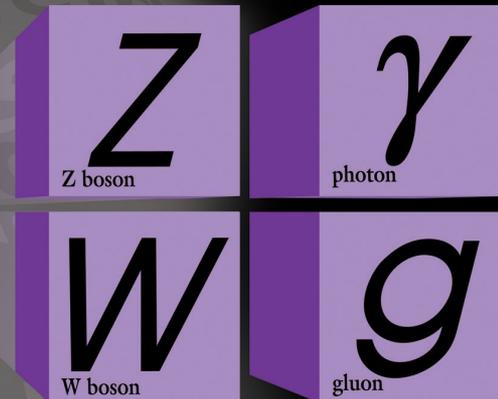


How should the Higgs
boson decay?

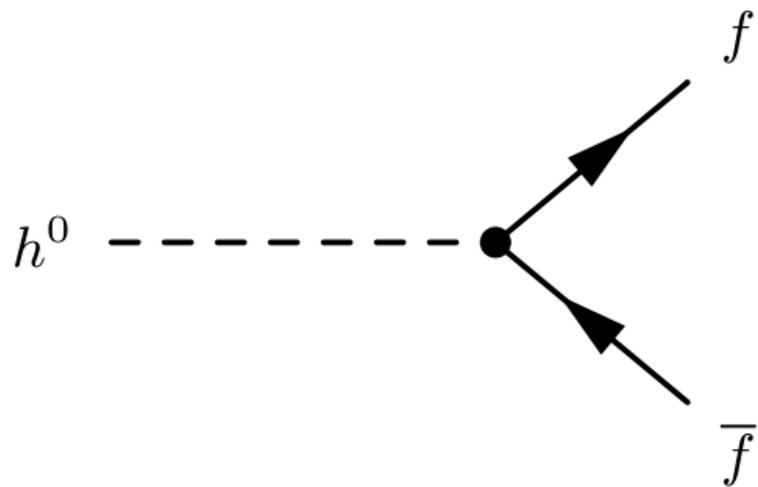
Quarks



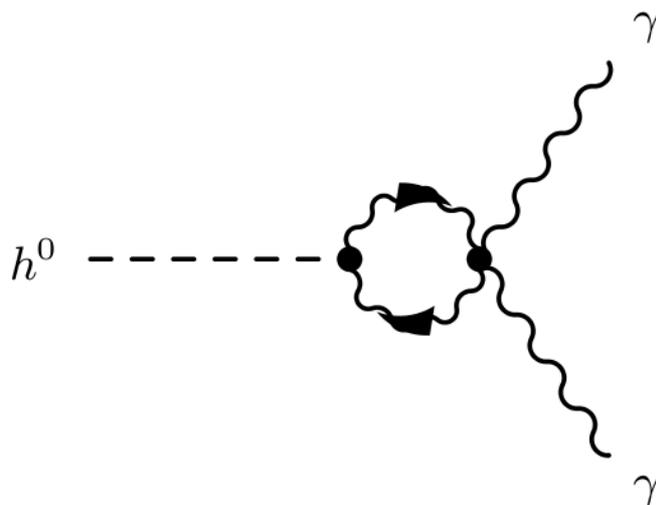
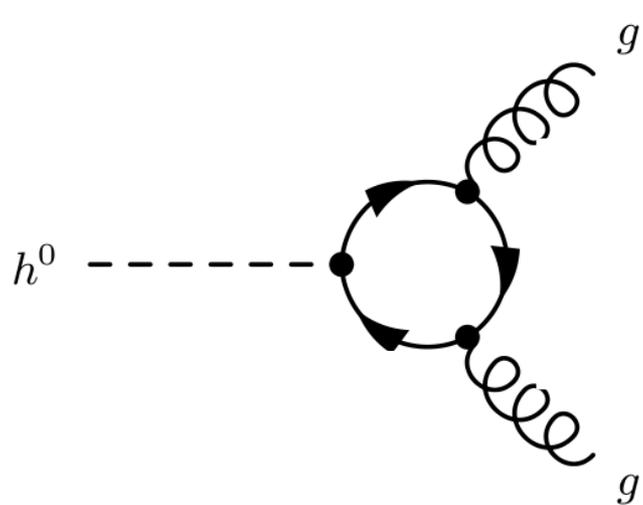
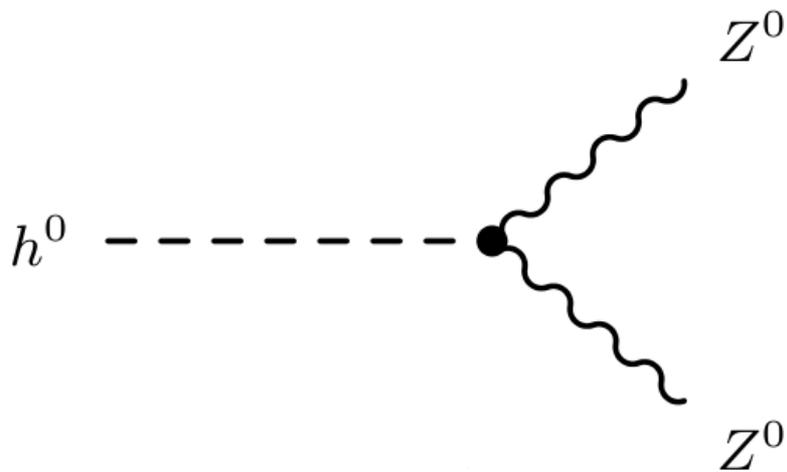
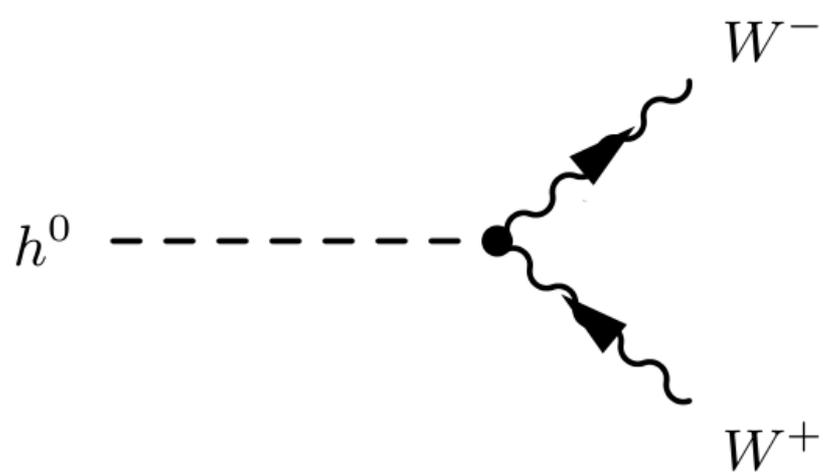
Forces

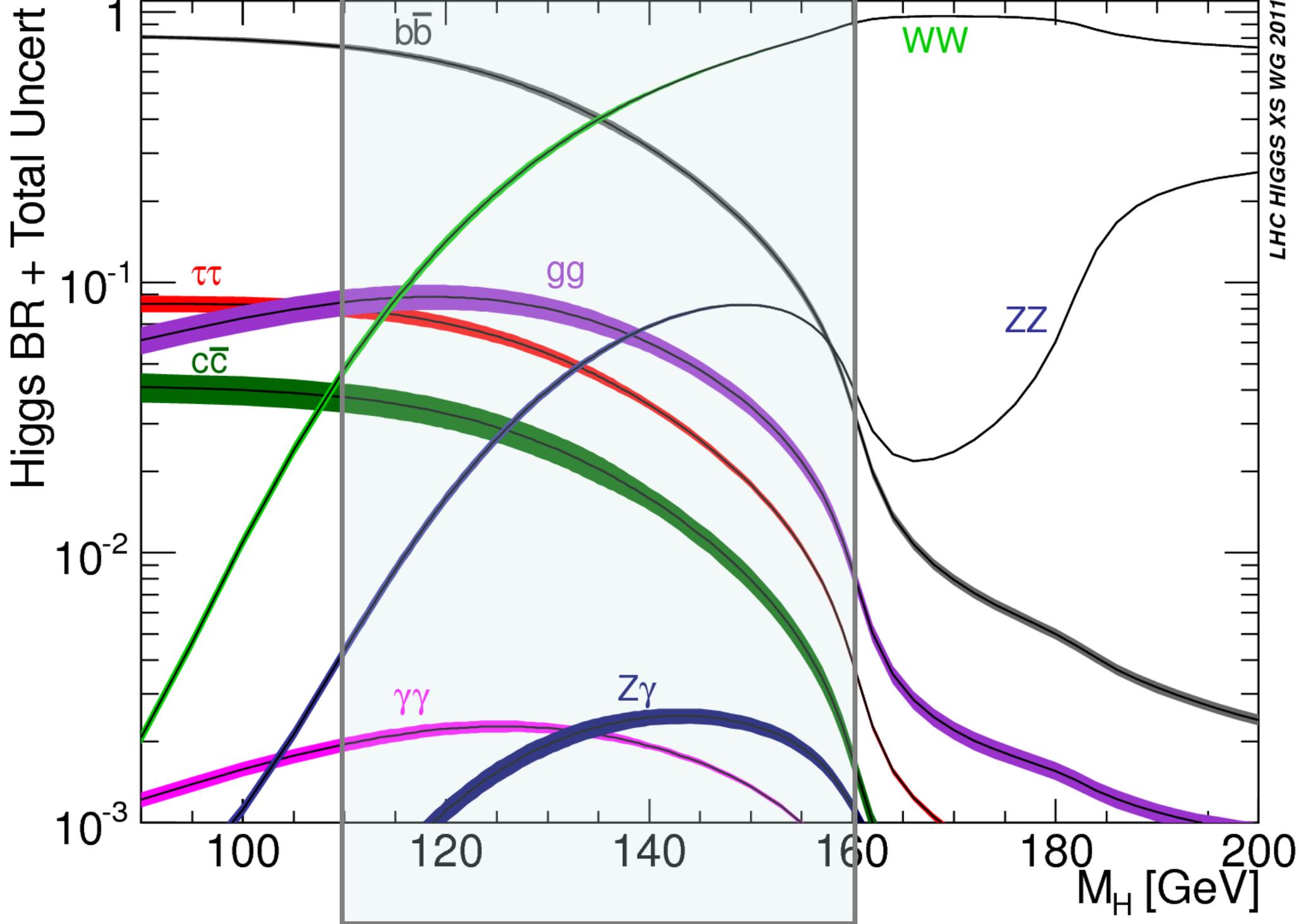


Leptons



Most likely fermions (based on mass):
 τ lepton and b -quark



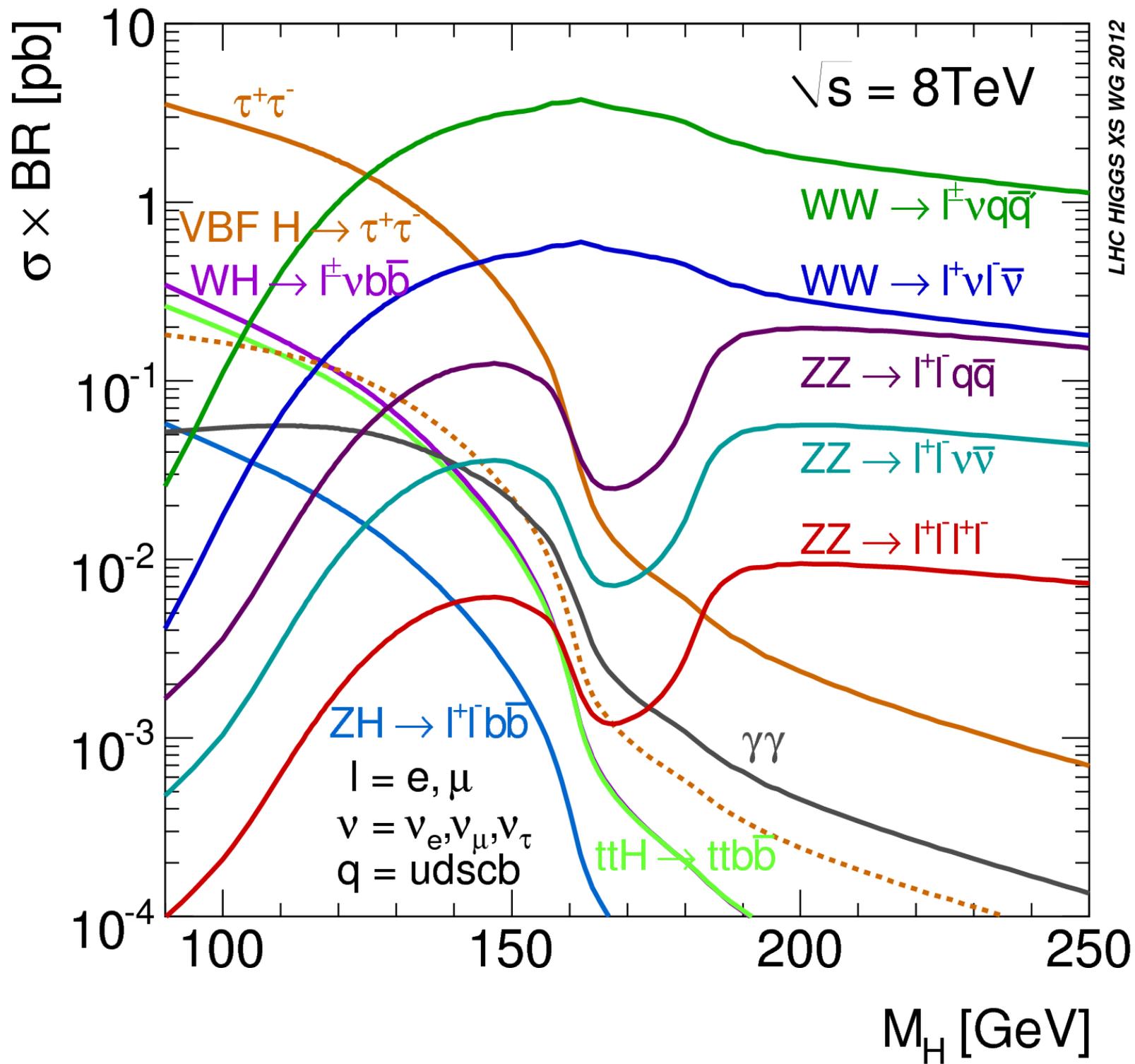


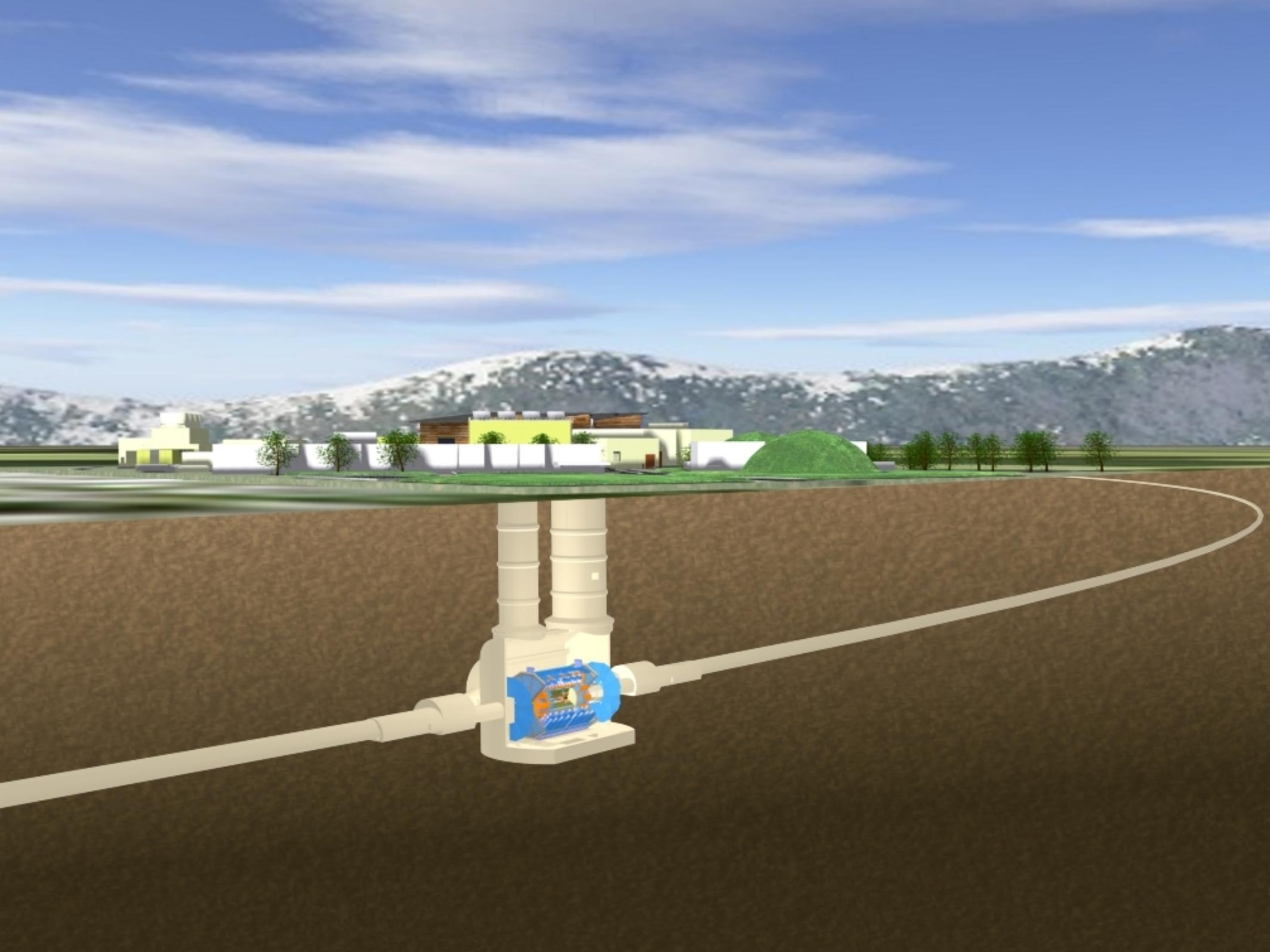
What, then, are the best ways to look for the Higgs boson?

Final states involving only quarks or gluons are experimentally extremely challenging
[in a hadron collider, nearly everything produced is gluons and quarks \rightarrow JETS]

Final states involving tau leptons are challenging because $\tau \rightarrow X + \nu$ ($X = e, \mu, \text{quarks}$)

Final states involving only photons, electrons, or muons are experimentally very accessible, but happen at lower rates overall.

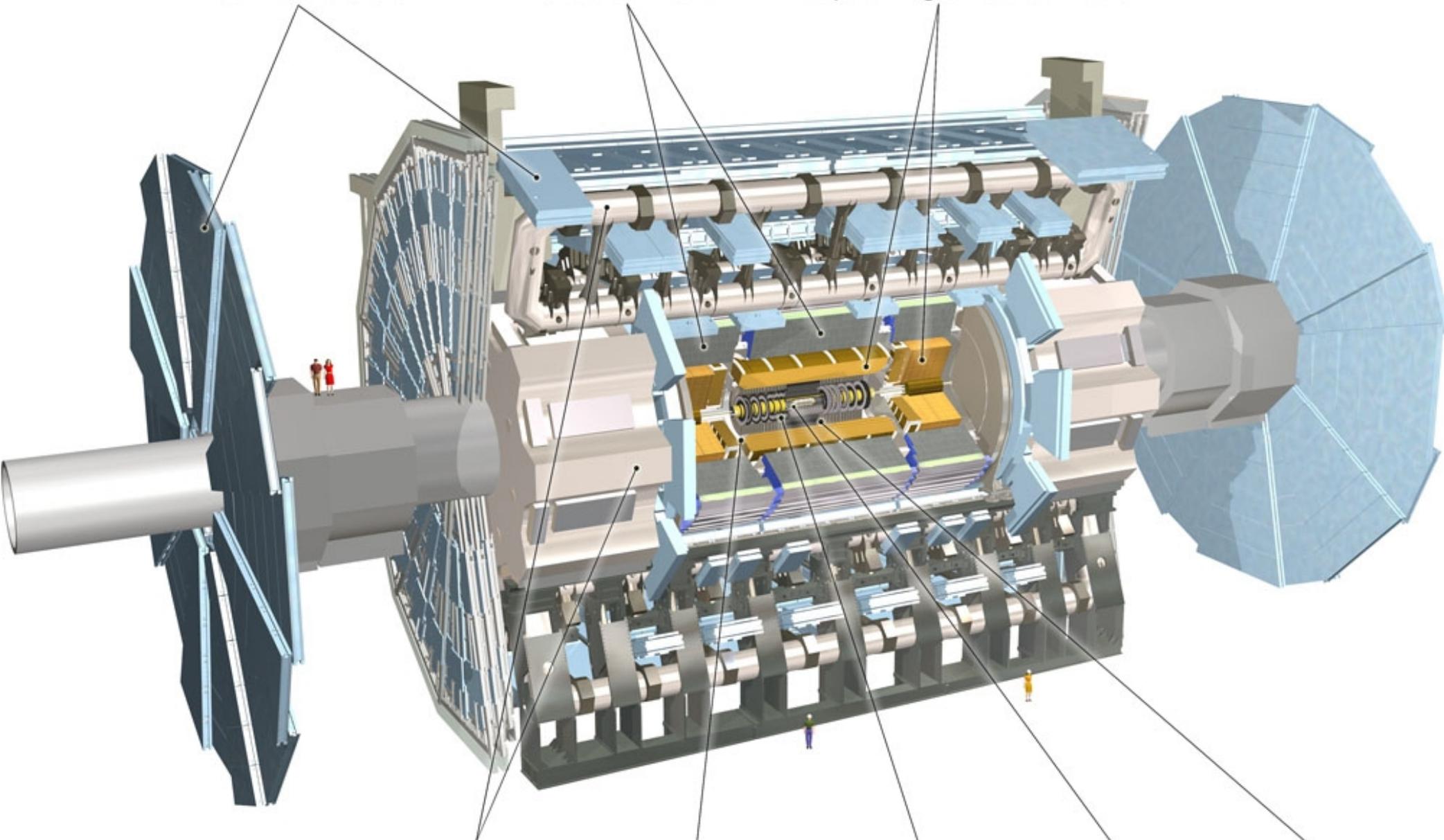




Muon Detectors

Tile Calorimeter

Liquid Argon Calorimeter



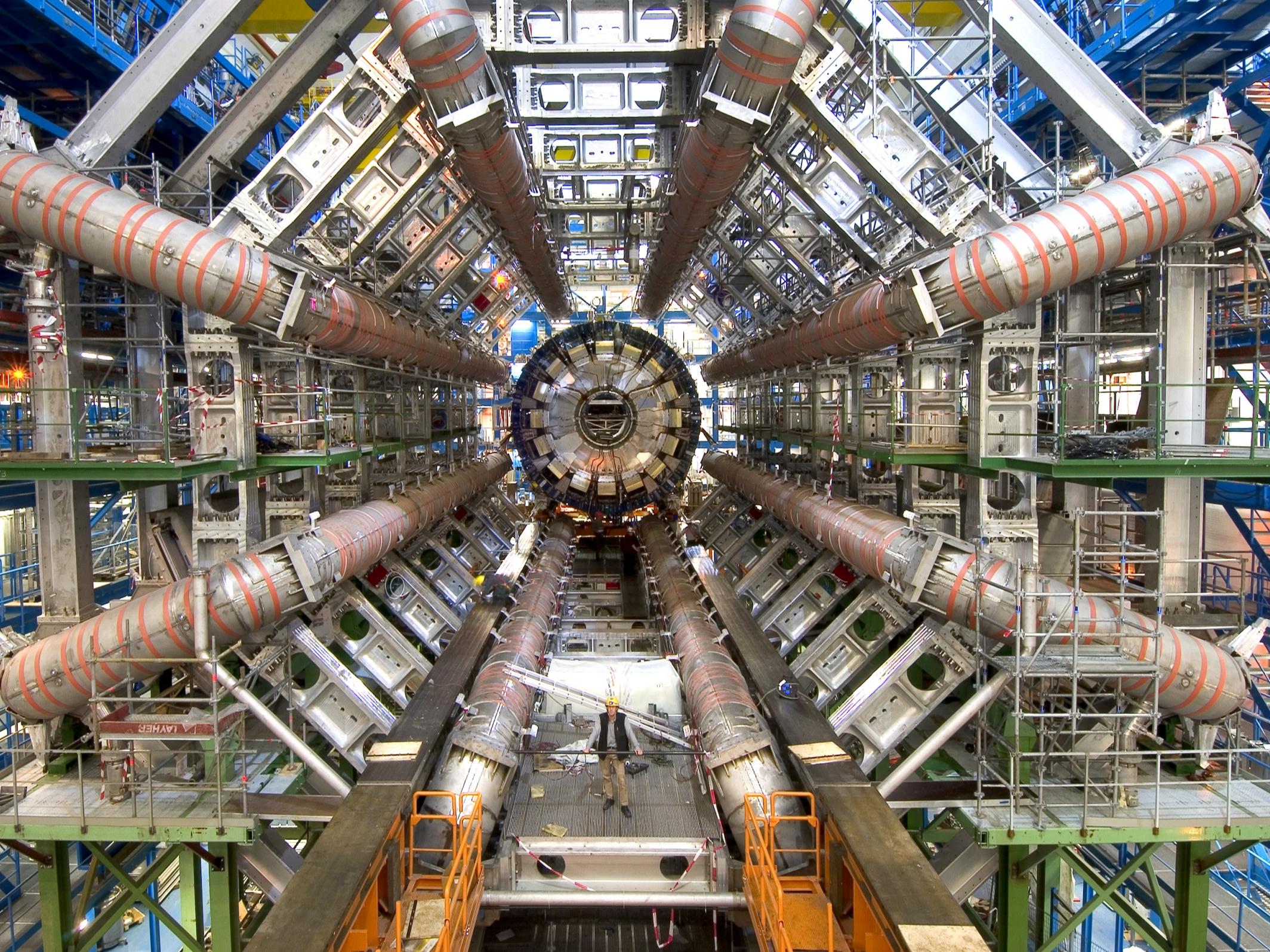
Toroid Magnets

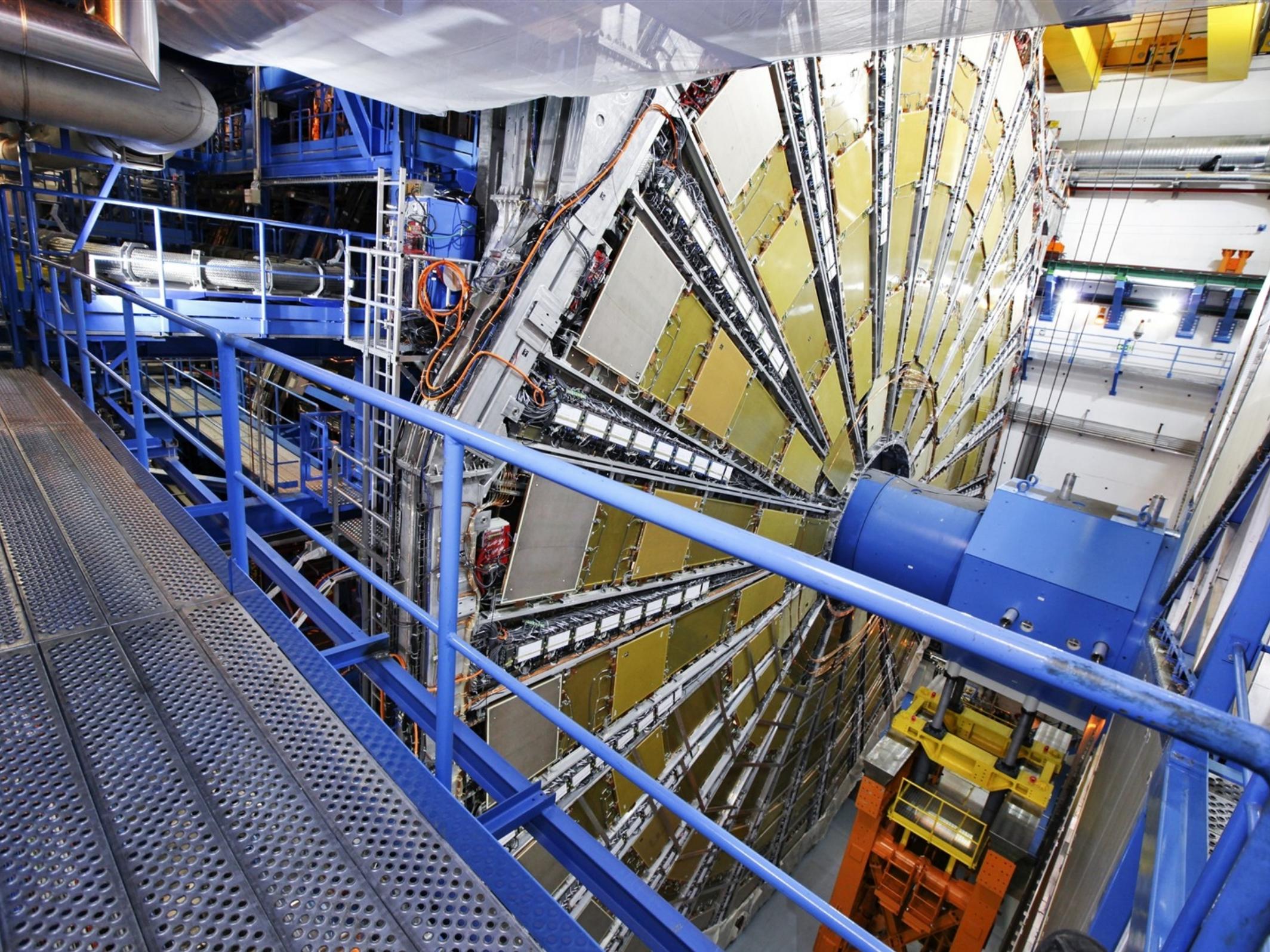
Solenoid Magnet

SCT Tracker

Pixel Detector

TRT Tracker







Muon Spectrometer

Muon

Neutrino

Hadronic Calorimeter

Proton

Neutron

The dashed tracks are invisible to the detector

Electromagnetic Calorimeter

Electron

Photon

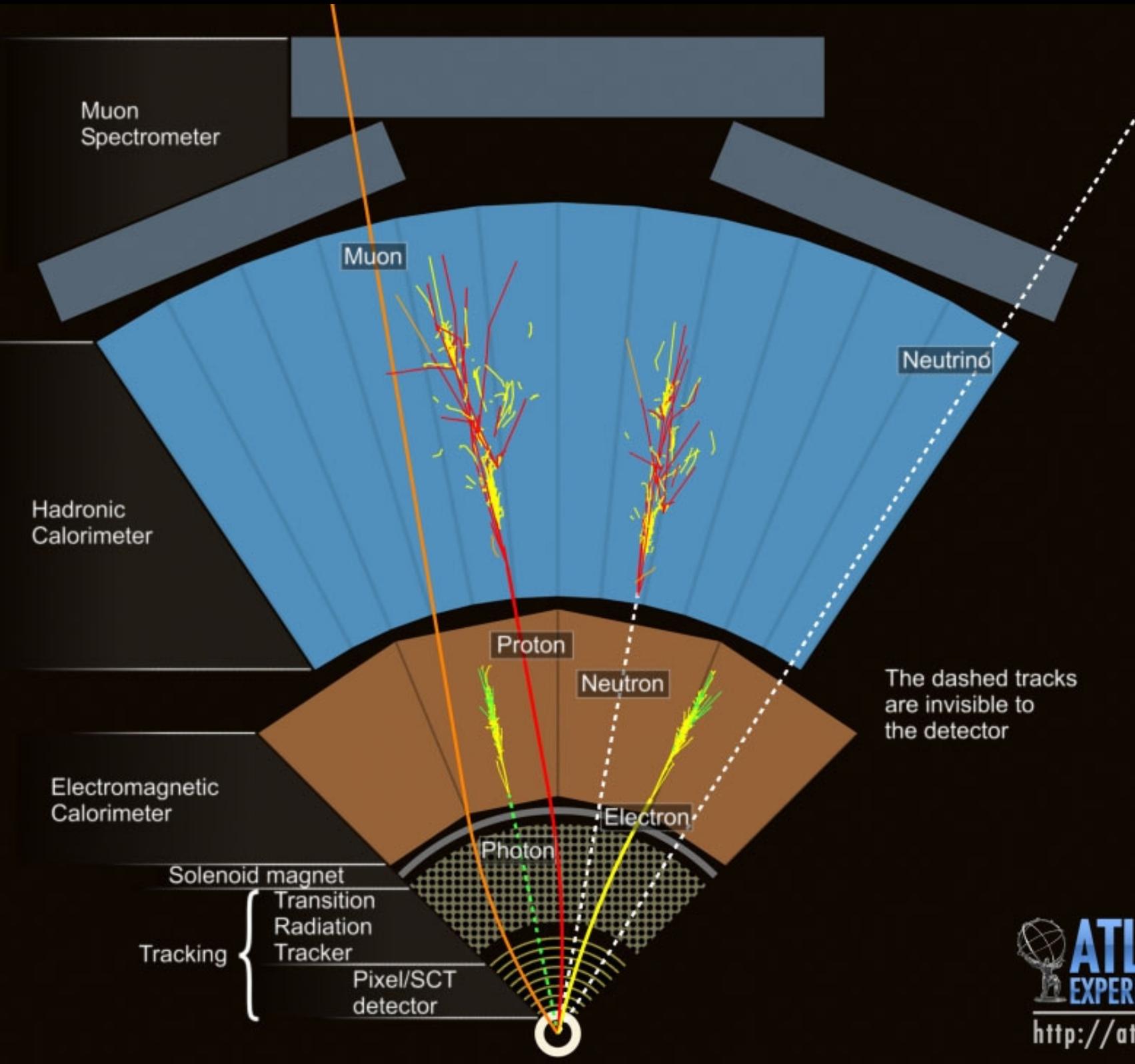
Solenoid magnet

Tracking {
Transition
Radiation
Tracker

Pixel/SCT
detector

 **ATLAS**
EXPERIMENT

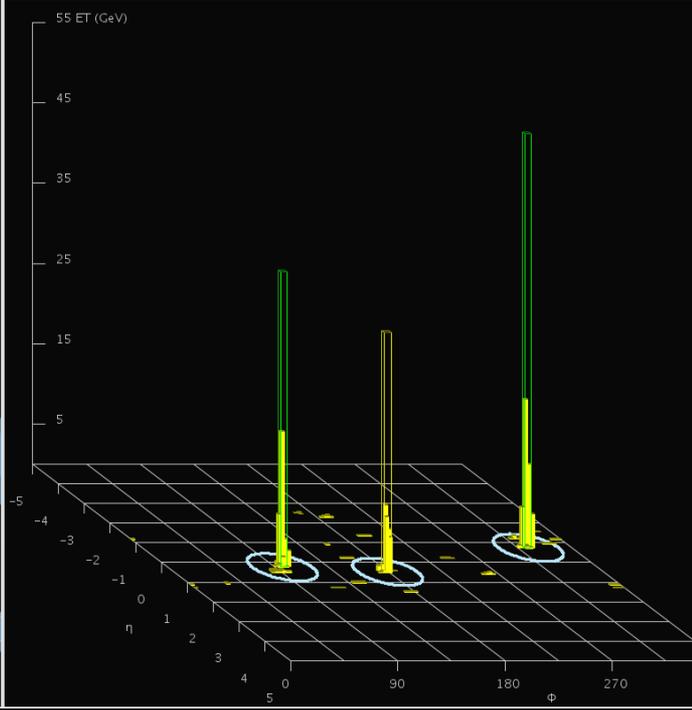
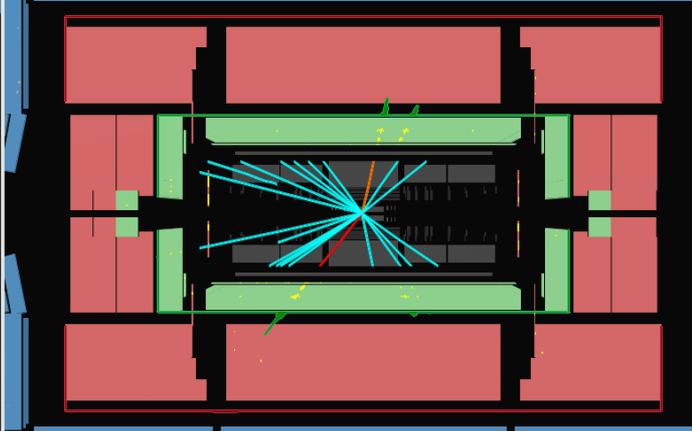
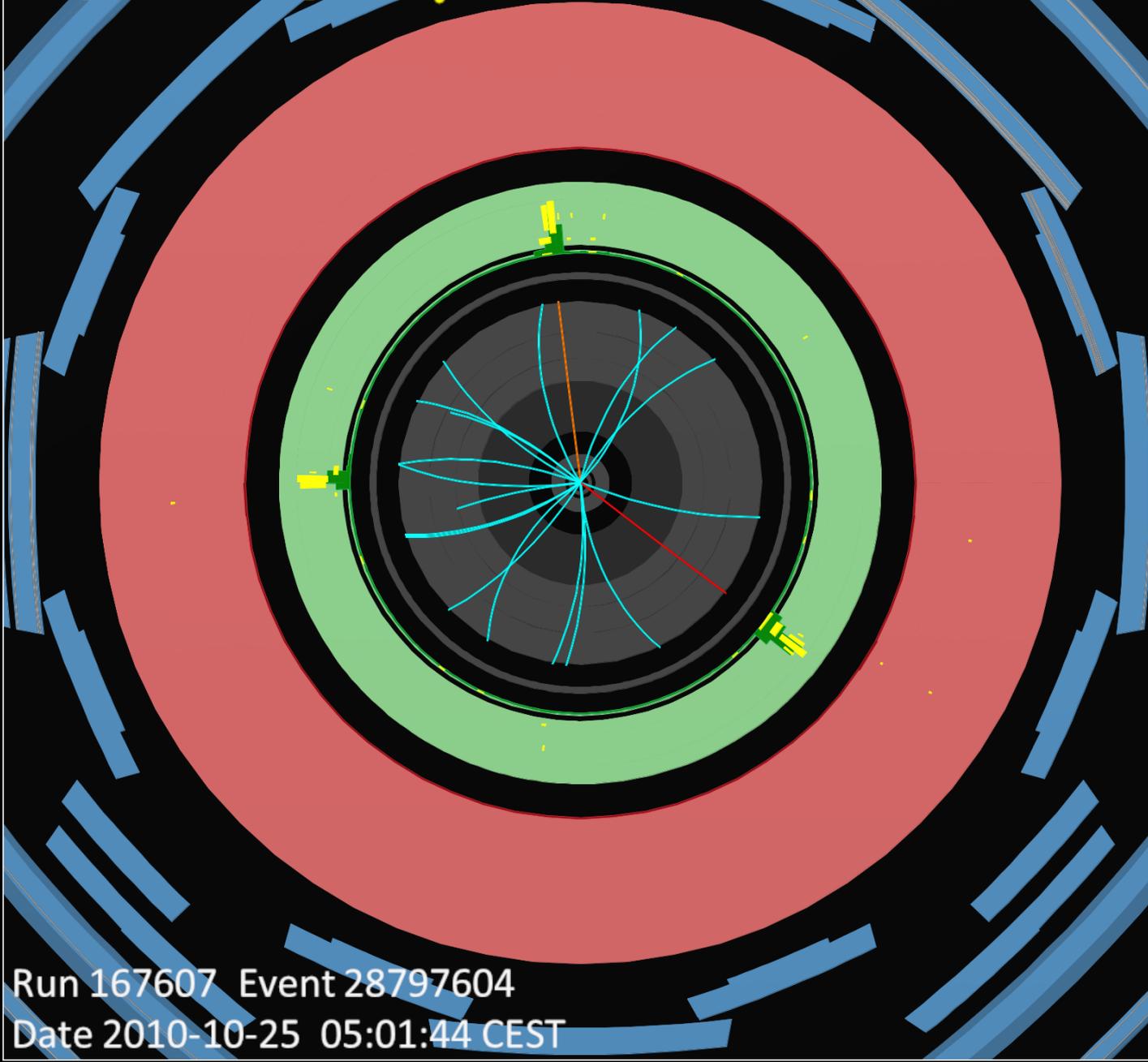
<http://atlas.ch>



Z($\rightarrow ee$) + γ Candidate



ATLAS EXPERIMENT



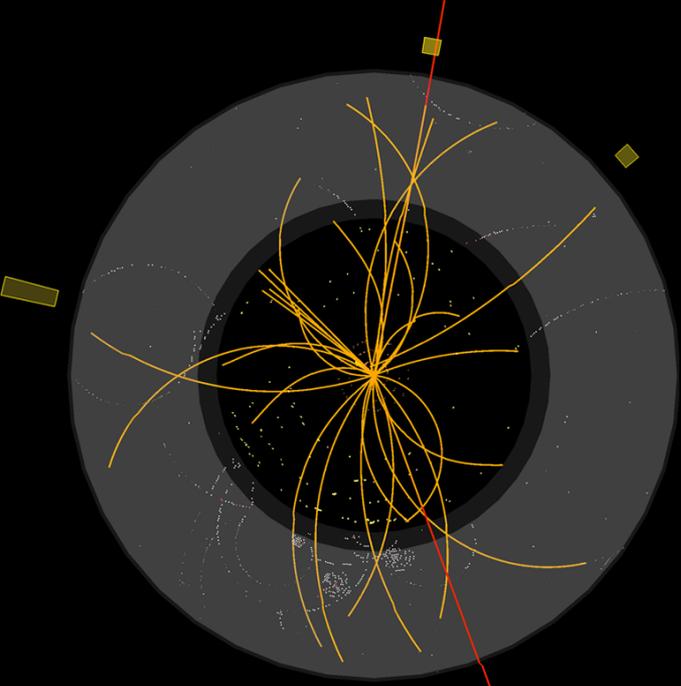
Run 167607 Event 28797604
Date 2010-10-25 05:01:44 CEST

Electron and Photon reconstruction and identification



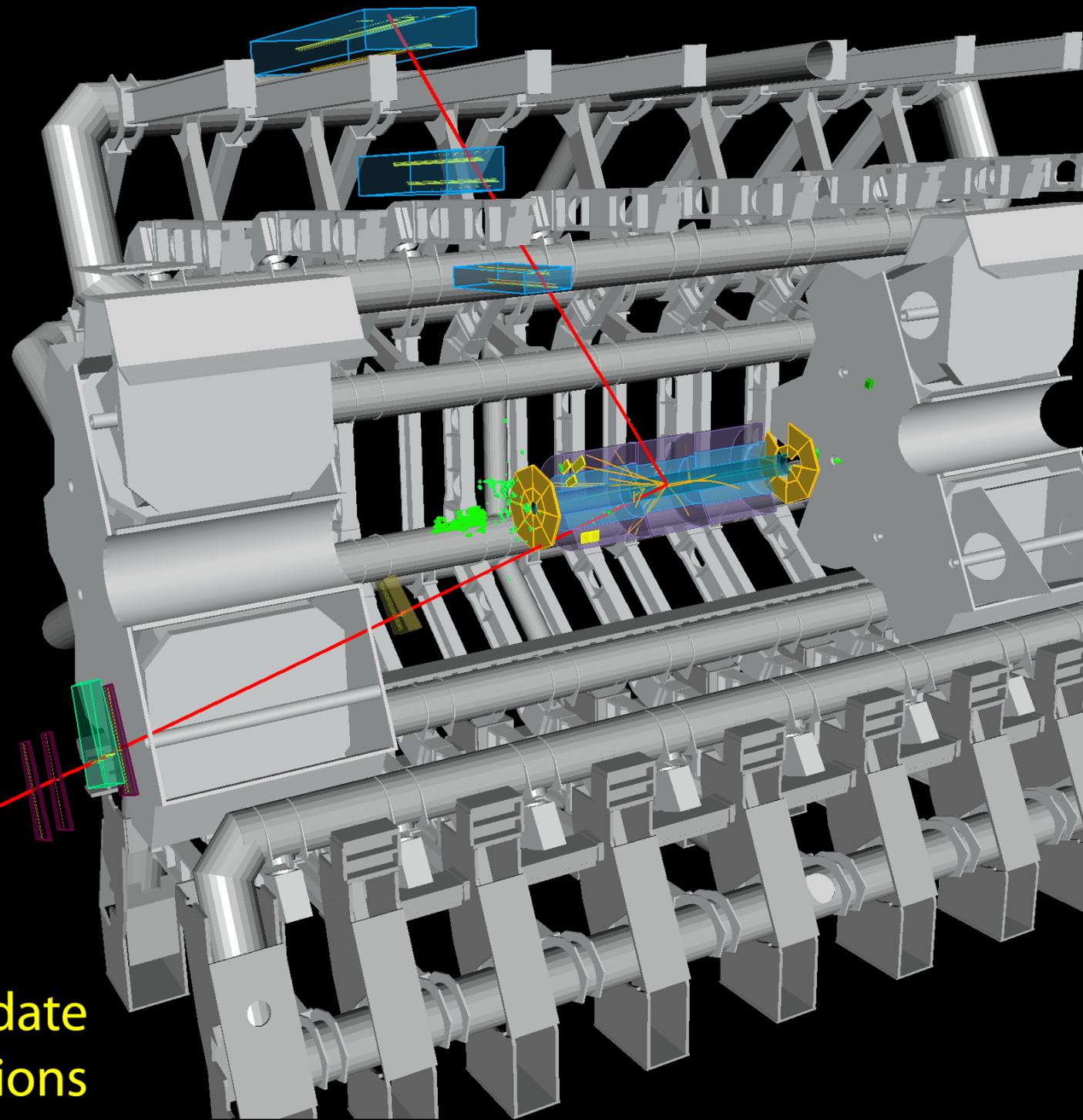
ATLAS EXPERIMENT

Run: 154822, Event: 14321500
Date: 2010-05-10 02:07:22 CEST



$p_T(\mu^-) = 27 \text{ GeV}$ $\eta(\mu^-) = 0.7$
 $p_T(\mu^+) = 45 \text{ GeV}$ $\eta(\mu^+) = 2.2$
 $M_{\mu\mu} = 87 \text{ GeV}$

**Z $\rightarrow\mu\mu$ candidate
in 7 TeV collisions**



Muon reconstruction and identification

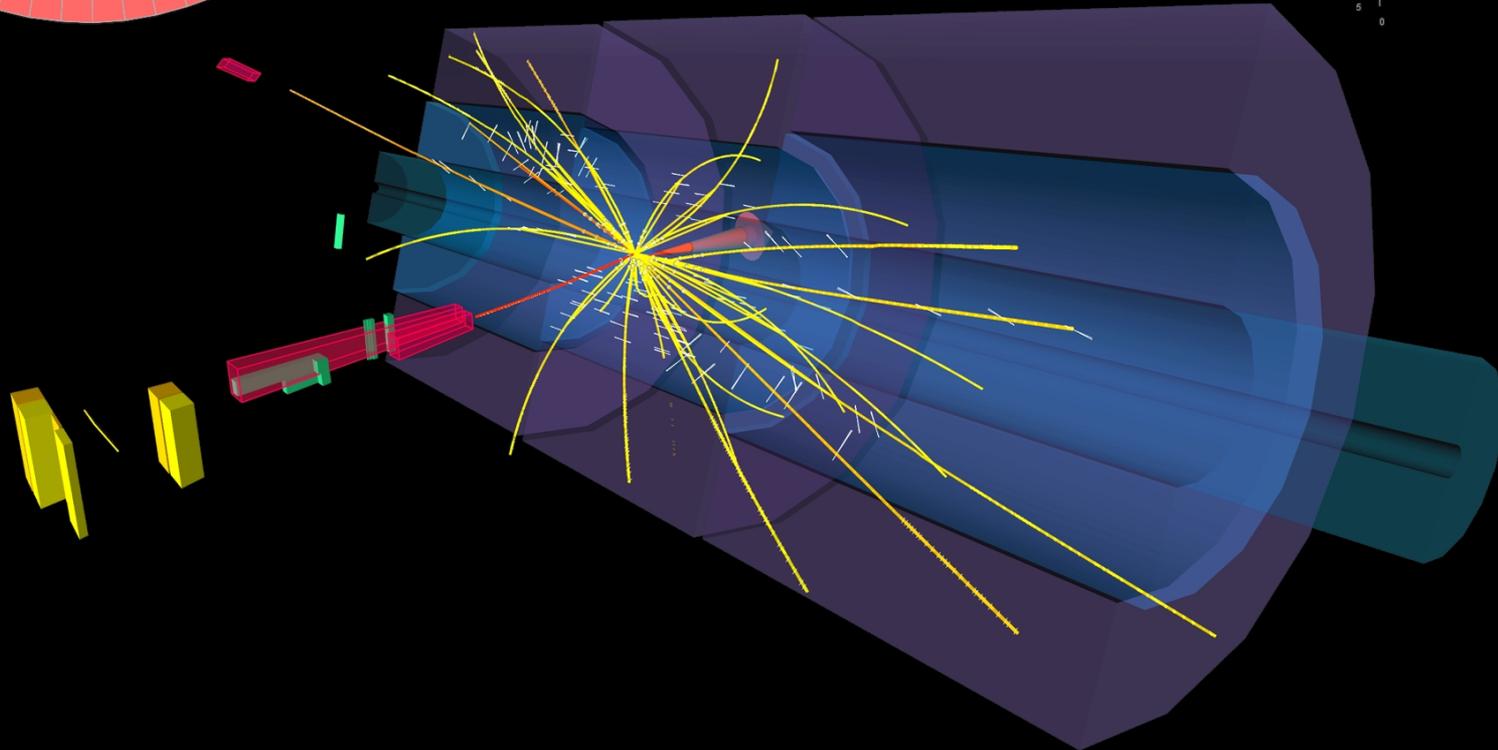
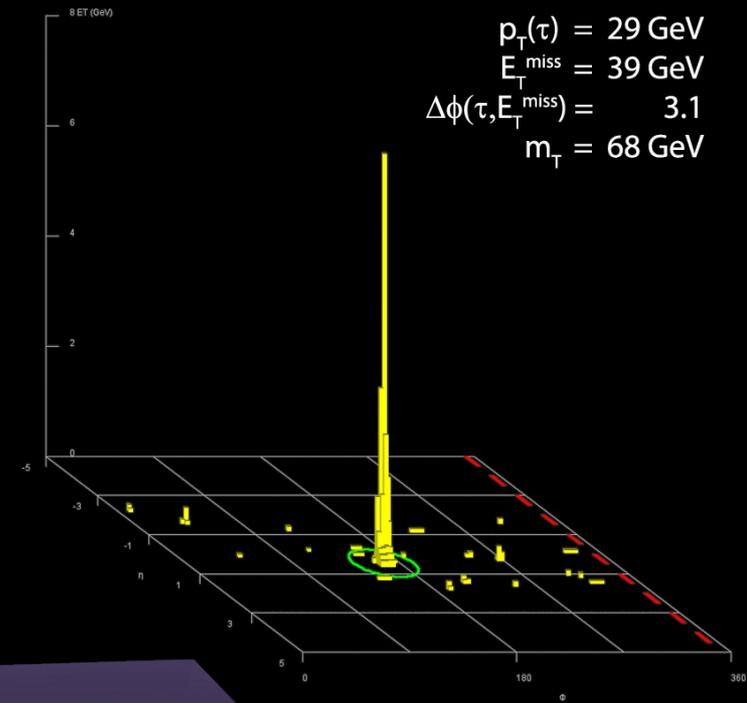
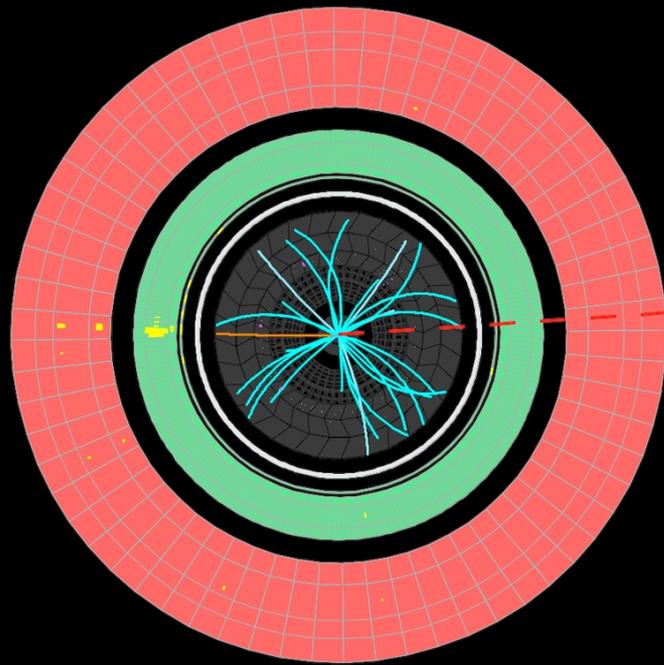


ATLAS EXPERIMENT

Run 155697, Event 6769403

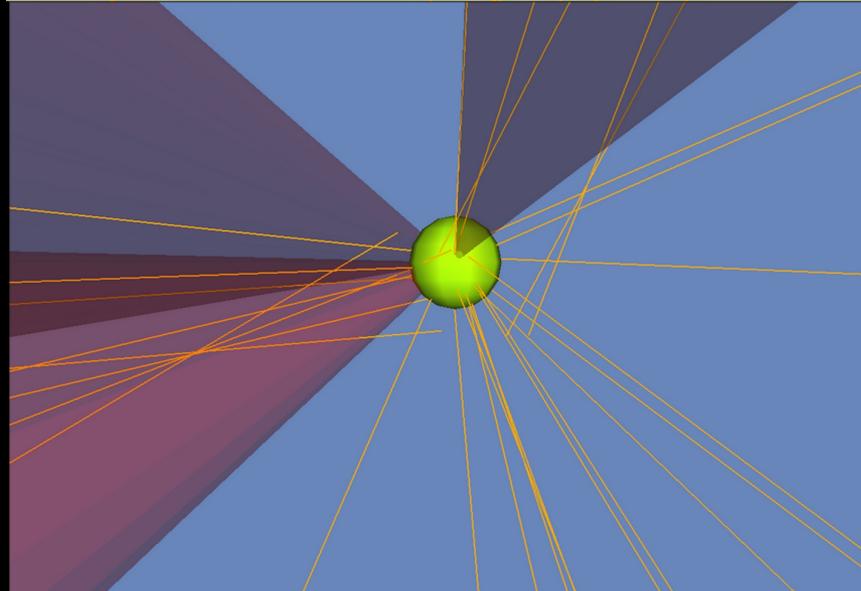
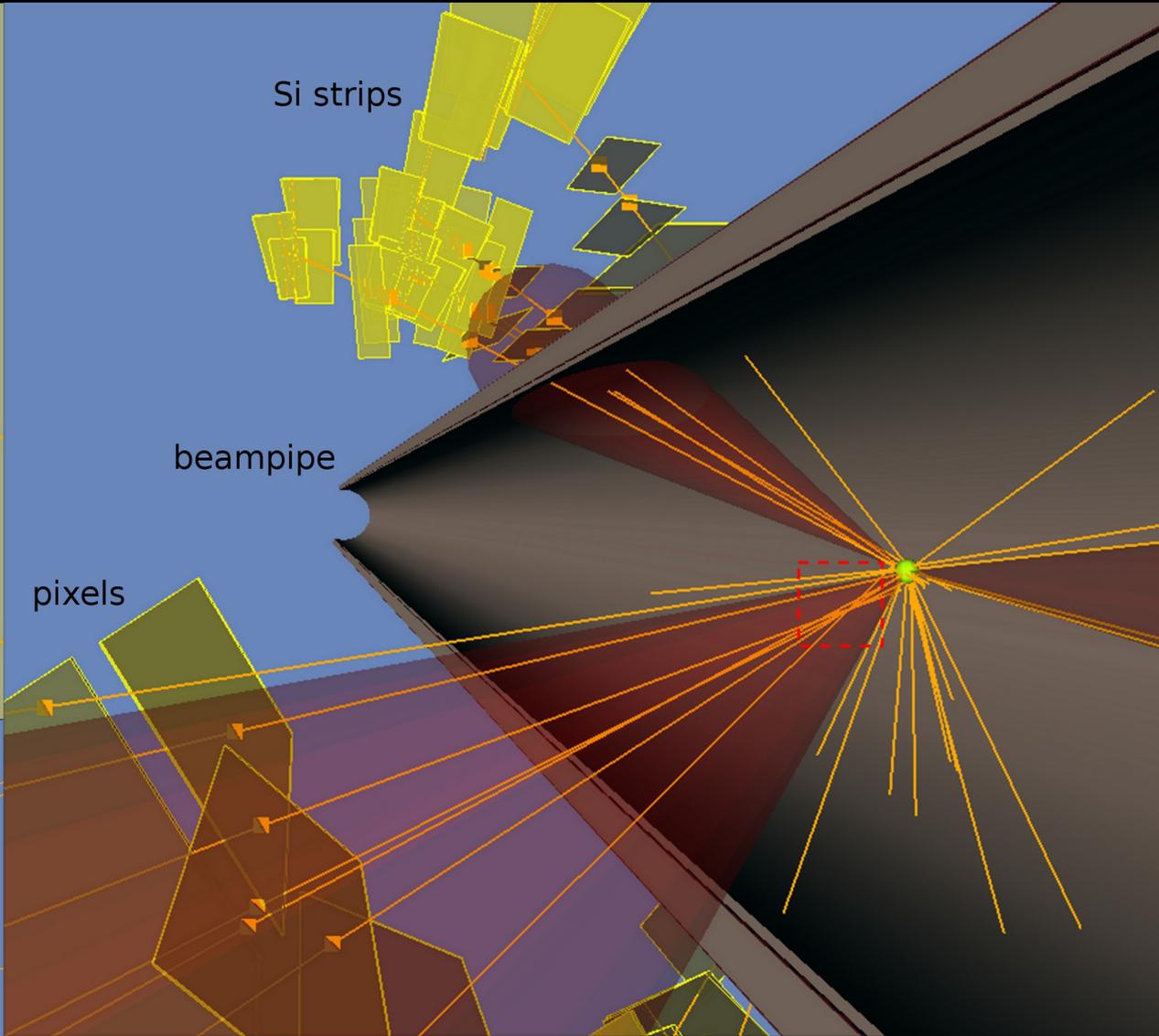
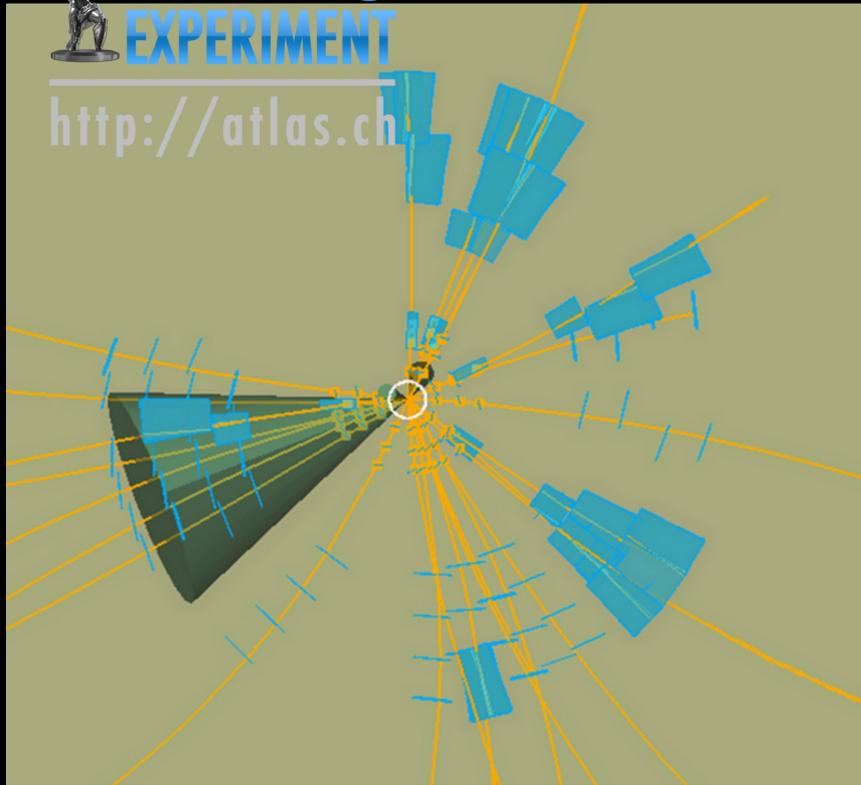
Time 2010-05-24, 17:38 CEST

$W \rightarrow \tau \nu$ candidate in
7 TeV collisions



Tau lepton reconstruction and identification

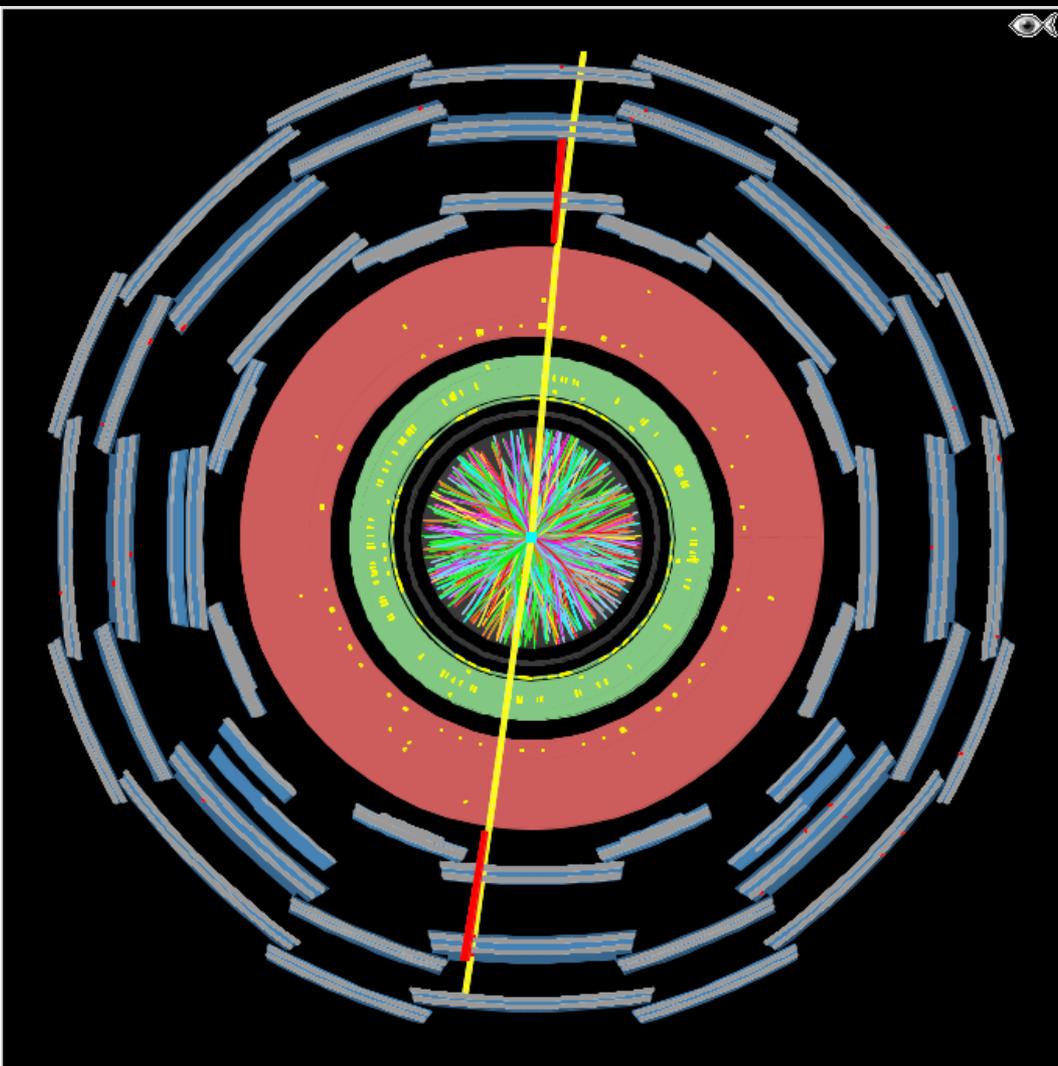
<http://atlas.ch>



jet
 $p_T = 19$ GeV (measured at electromagnetic scale)

4 b-tagging quality tracks in the jet

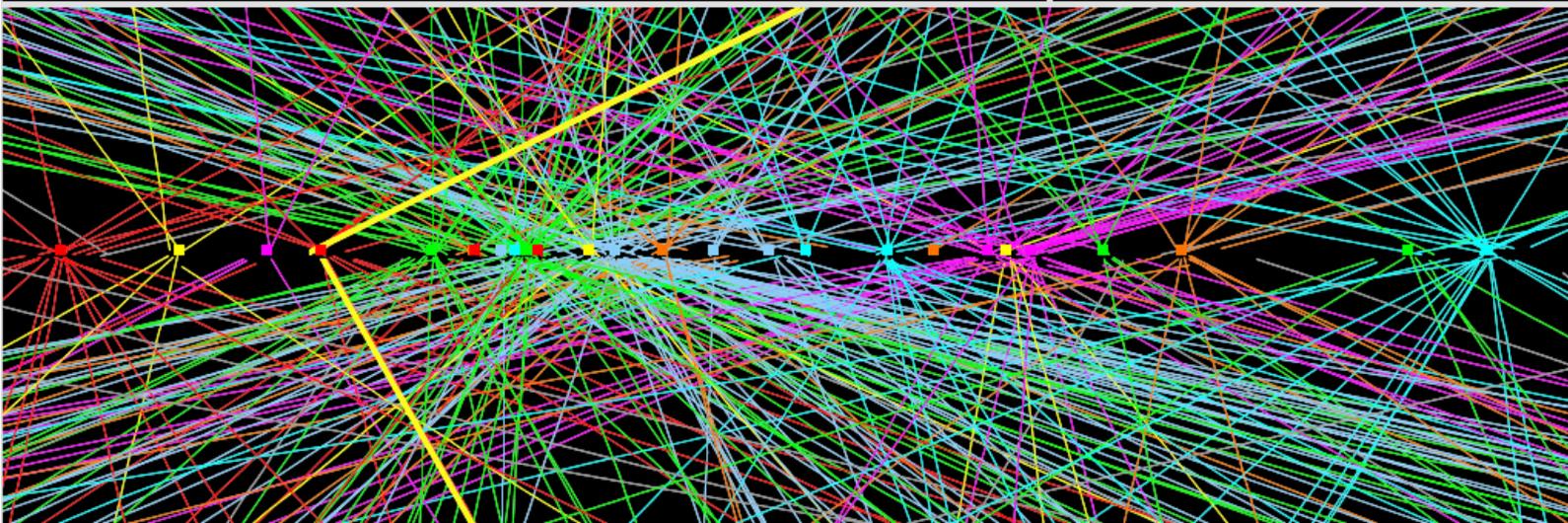
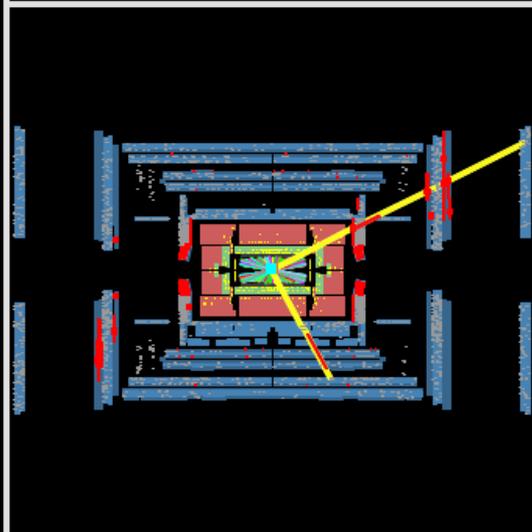
Bottom quark jet (b-jet) reconstruction and identification



ATLAS EXPERIMENT

Run Number: 201289, Event Number: 24151616

Date: 2012-04-15 16:52:58 CEST



Experimental Sidenote:
*How to avoid fooling yourself or
being fooled by nature*

First, you need to establish the question you are trying to answer:

“What is the chance that a fair coin, one which is not preferential to either heads or tails, can yield 14 or more HEADS results in 20 flips?”

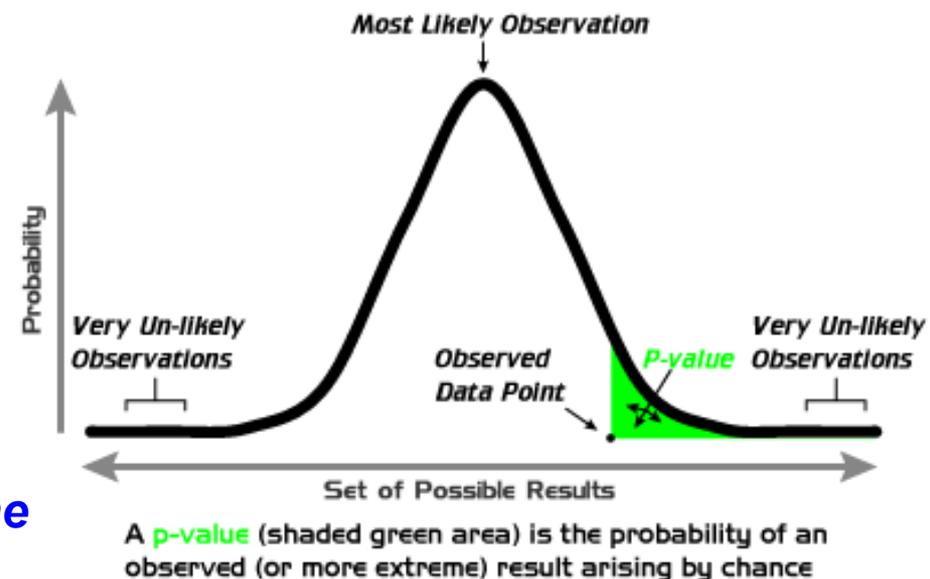
Analogy in the Higgs search: “If I see an excess in data over what is expected from the Standard Model, how probable is it for the Standard Model to have produced that many or more events in data?”

Use binomial probabilities to estimate the chance that a fair coin can yield 14 or more HEADS in 20 coin flips:

$$\begin{aligned} & \text{Prob}(14 \text{ heads}) + \text{Prob}(15 \text{ heads}) + \dots + \text{Prob}(20 \text{ heads}) \\ &= \frac{1}{2^{20}} \left[\binom{20}{14} + \binom{20}{15} + \dots + \binom{20}{20} \right] = \frac{60,460}{1,048,576} \approx 0.058 \end{aligned}$$

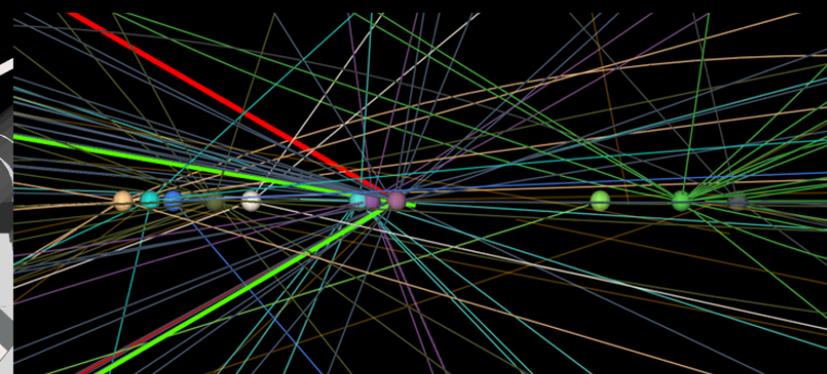
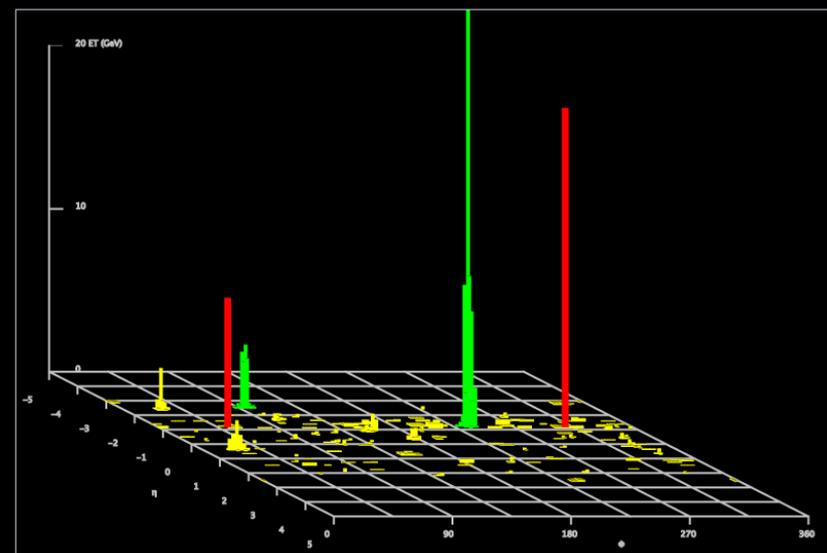
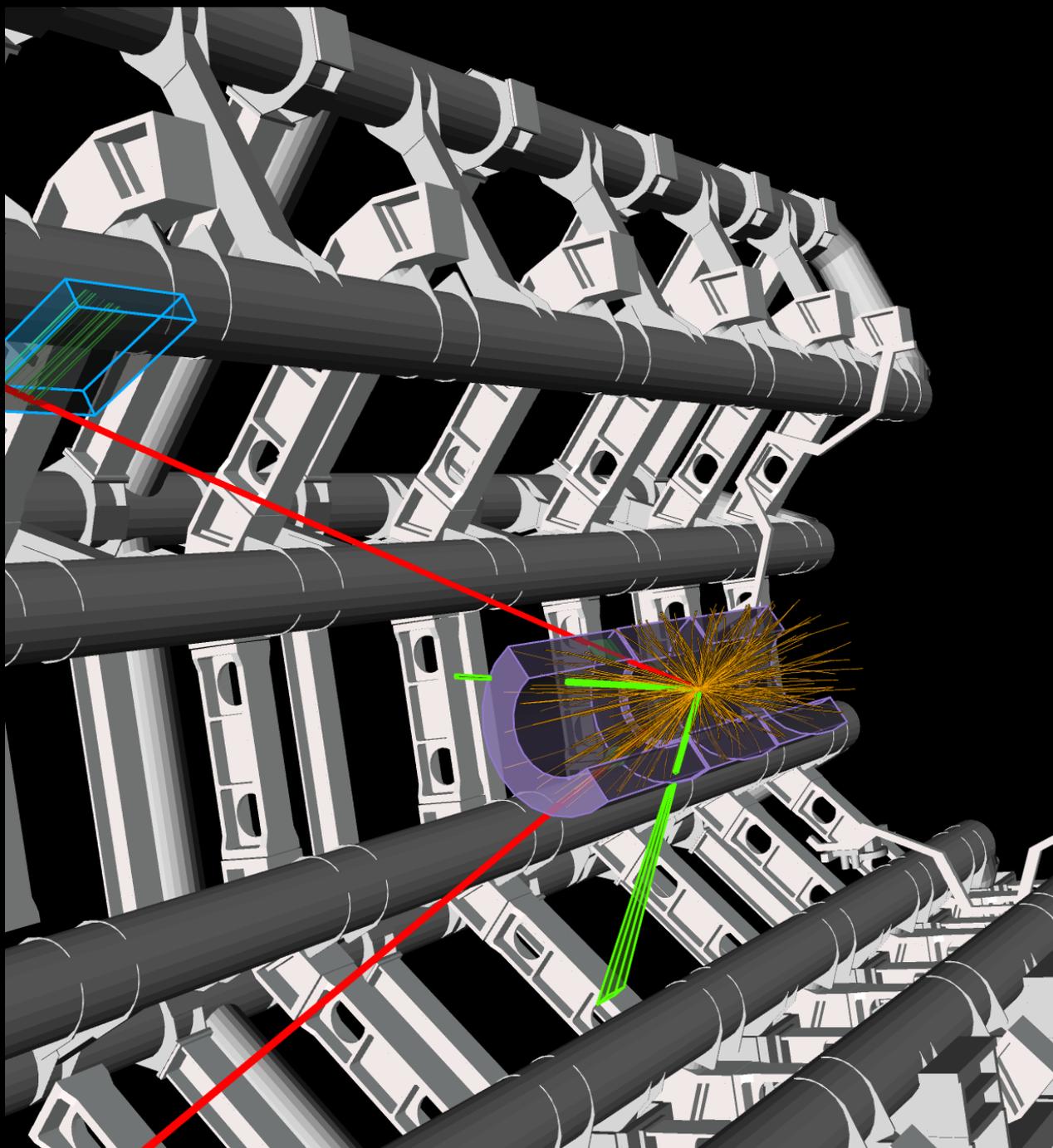
Here, binomial probability theory is the analogy to the Standard Model and how it's used to predict outcomes in an experiment.

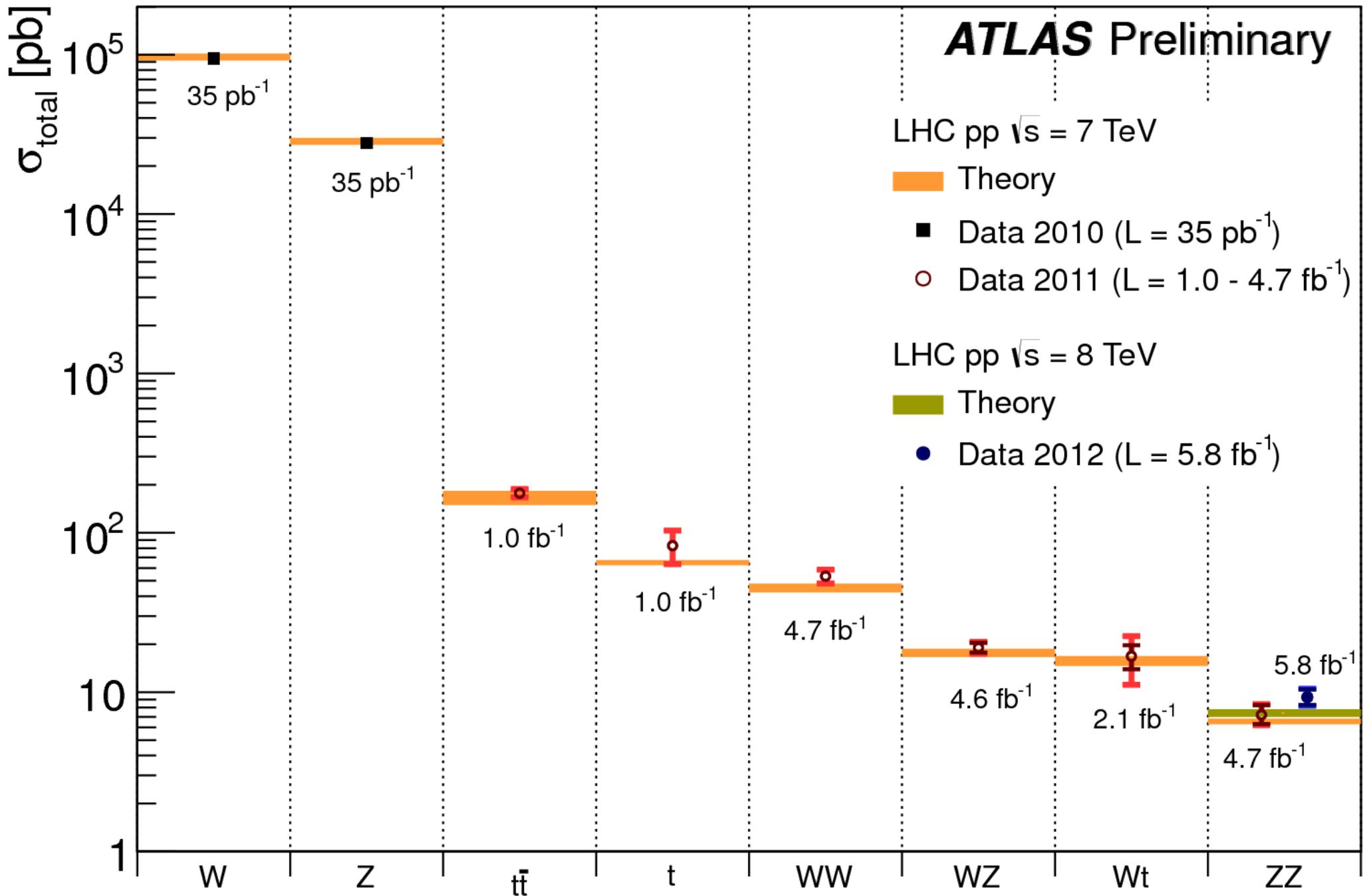
This number, 0.058, is called the “p-value” - the probability that random chance alone, not the presence of an unfair coin, could yield 14 or more HEADS from 20 flips of a coin.



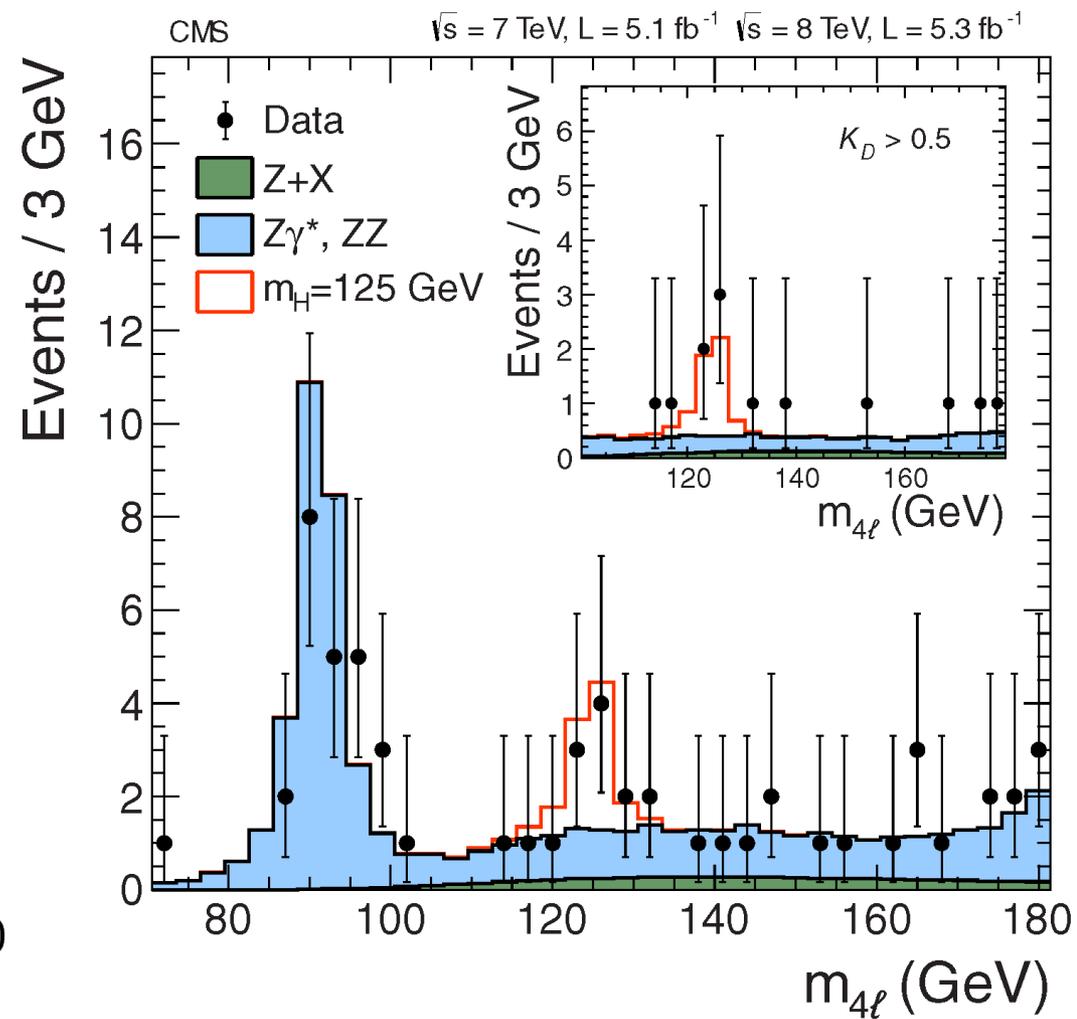
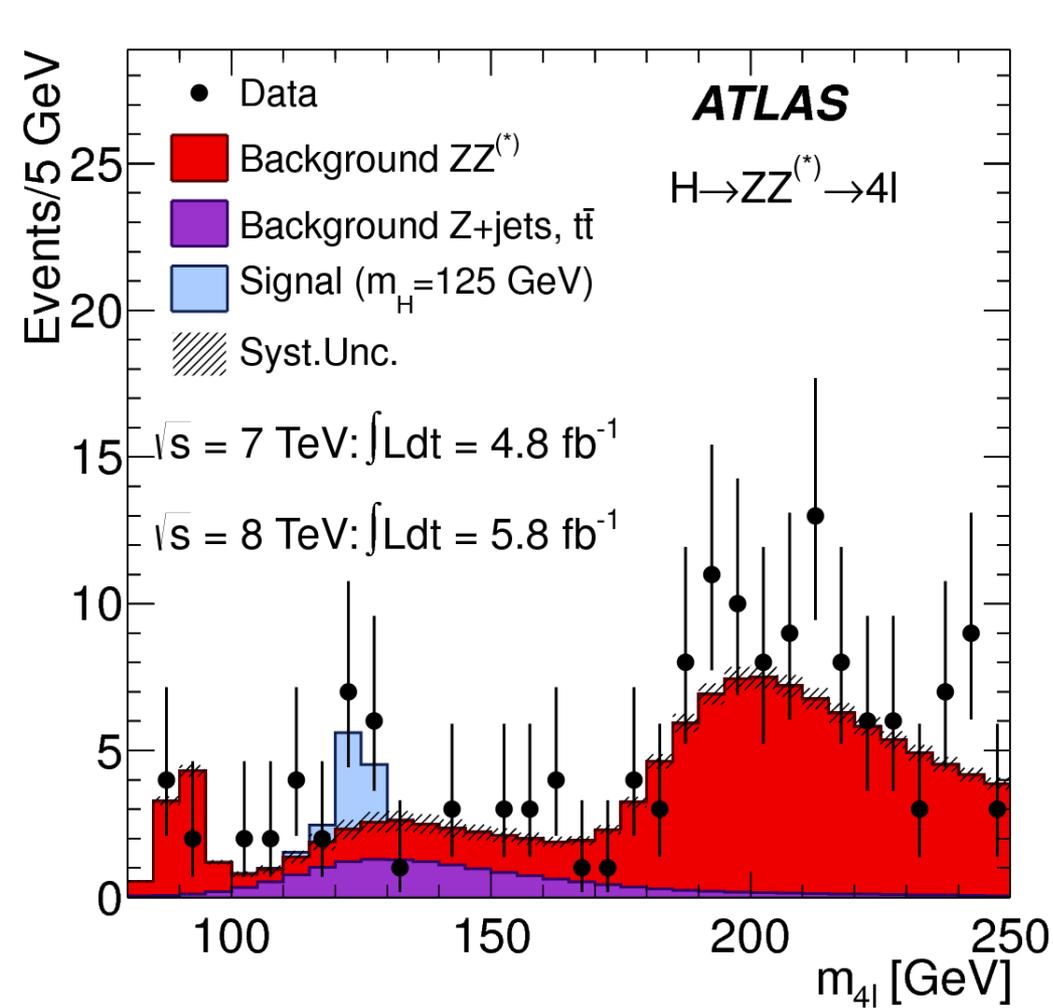
Experimental Results in the Search for a Higgs Boson Candidate

$H \rightarrow ZZ^{(*)} \rightarrow 4l$
The “Golden Mode”

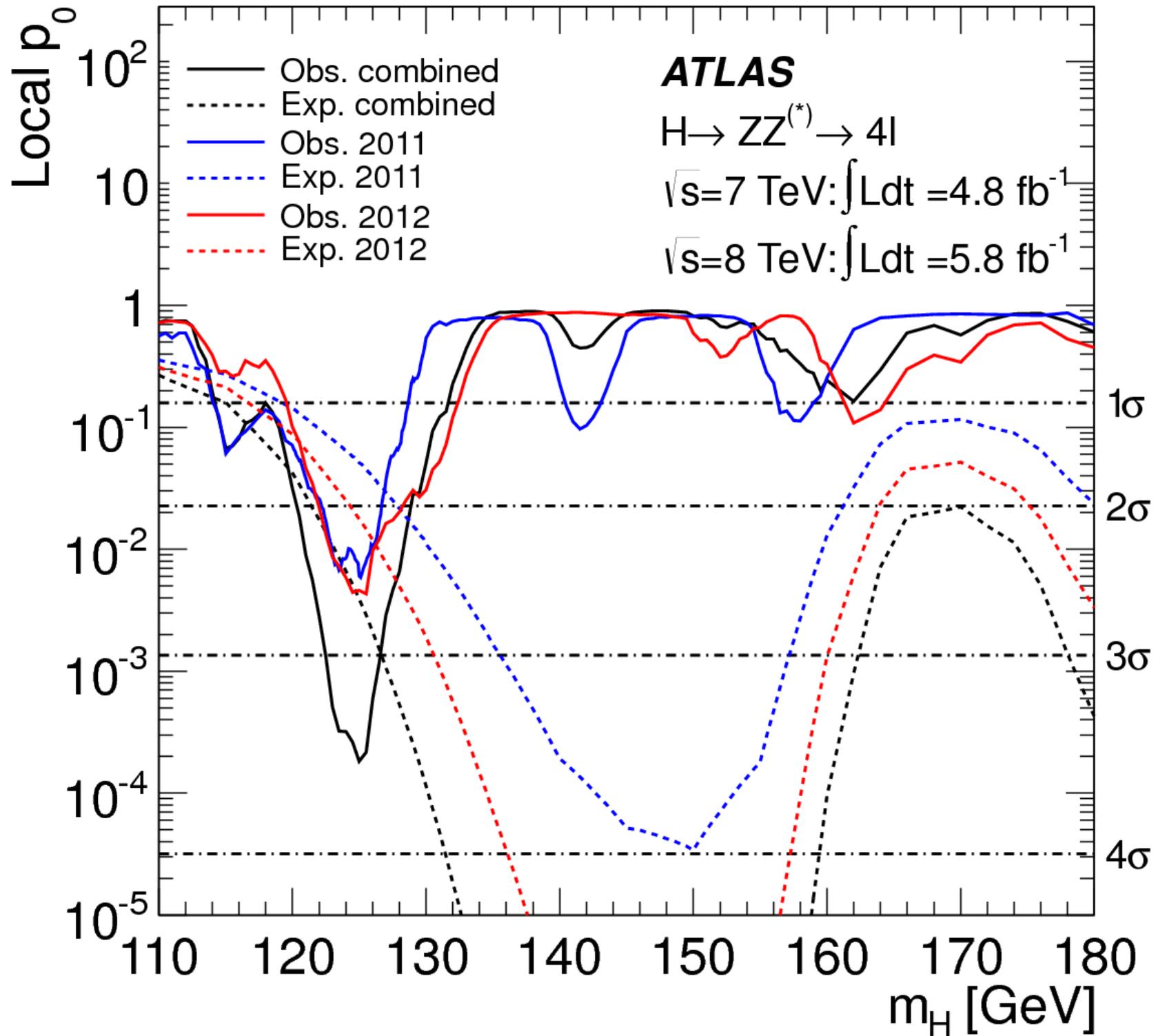


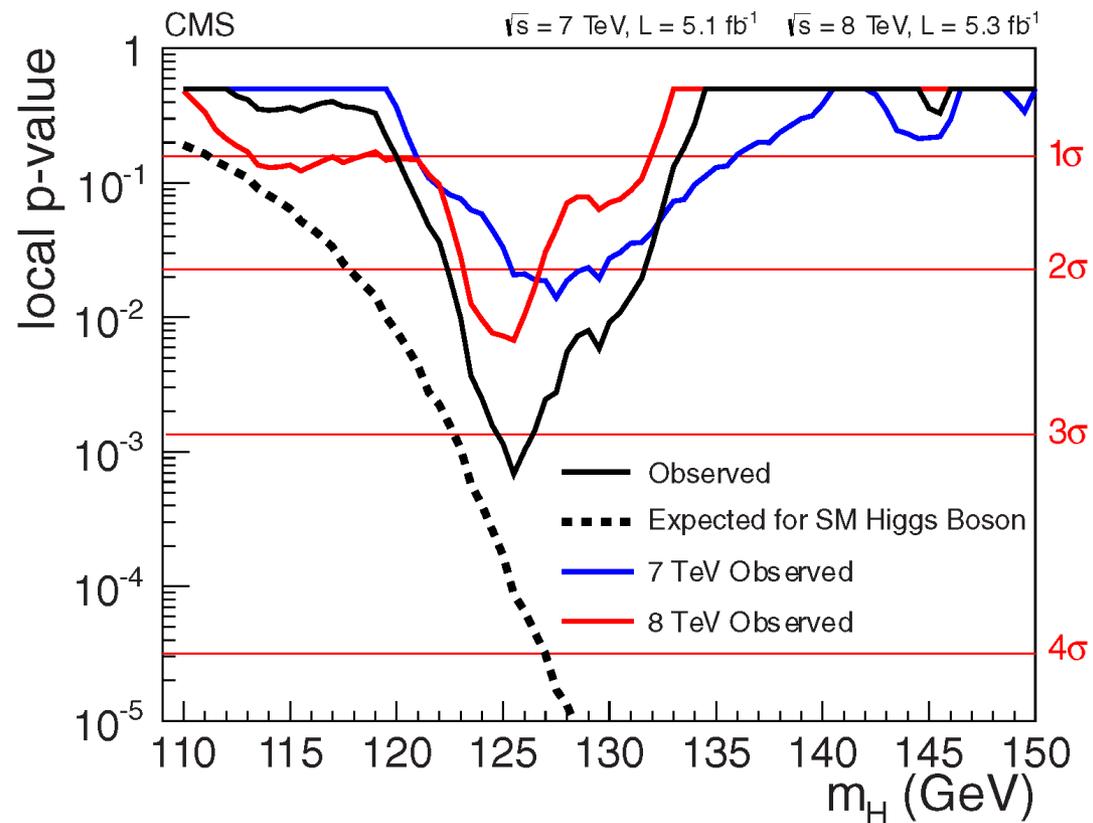
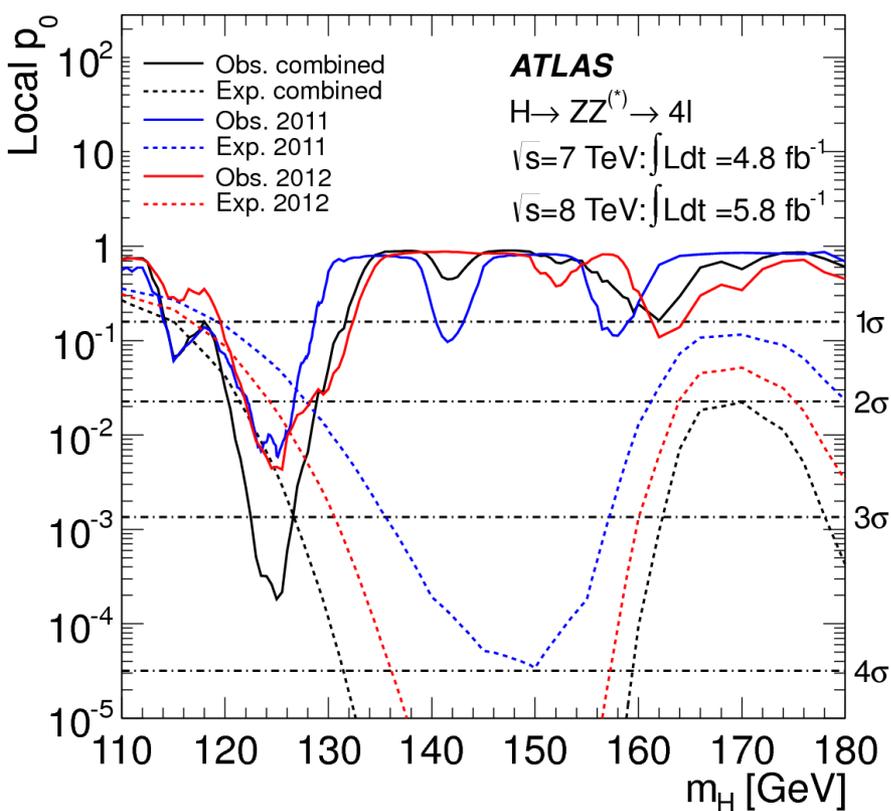


For comparison, the product of Higgs production cross-section and the branching fraction to ZZ(*) and then to four leptons is between 2×10^{-4} pb and 6×10^{-3} pb for Higgs masses between 100-200 GeV/c²



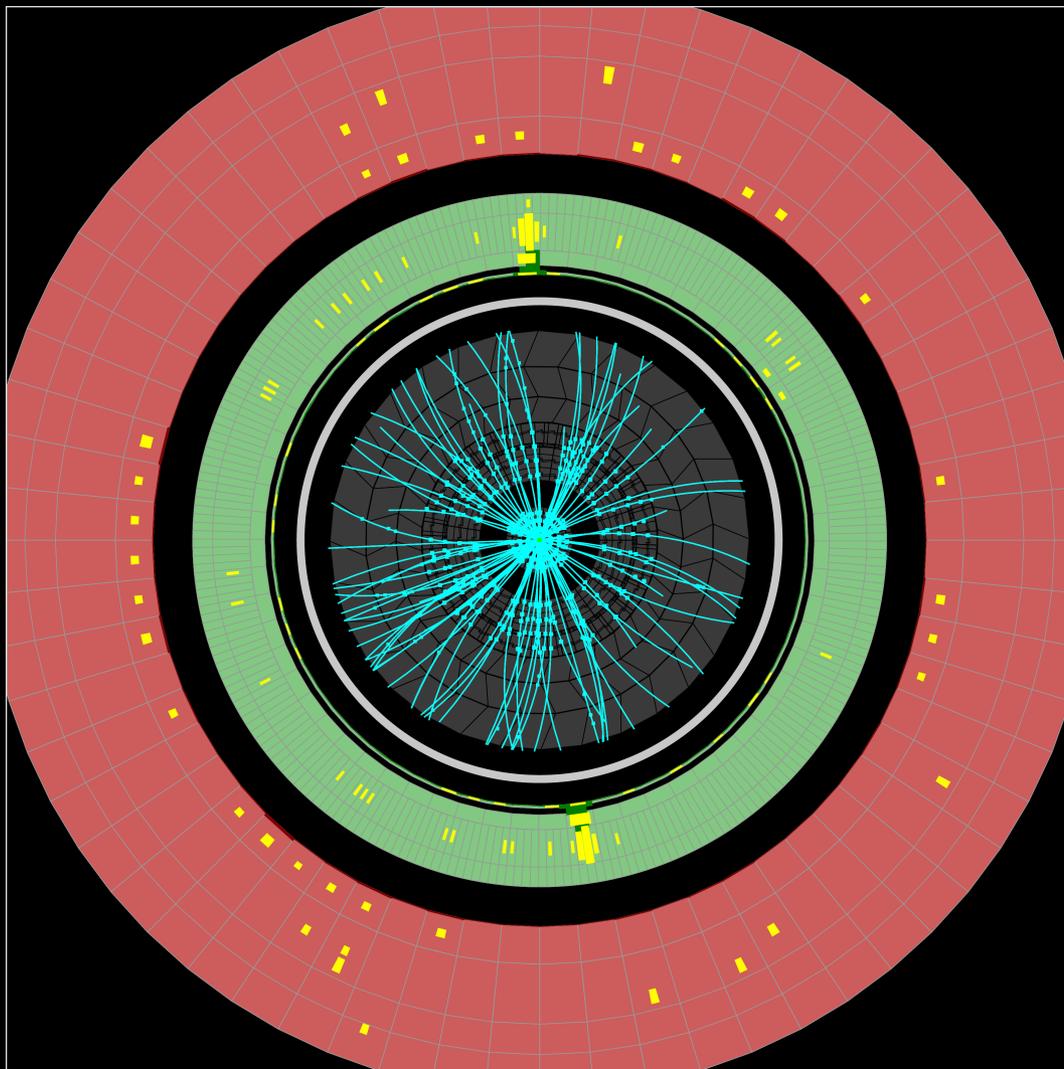
Both experiments independently observe an excess above the expected contributions from non-Higgs “background” events. The regions with the largest excesses are consistent and centered in the region around 125 GeV/c 2 .





From the perspective of a detailed statistical analysis, both experiments independently observe an excess above the expected contributions from non-Higgs “background” events. The regions with the largest excesses are fairly consistent and centered in the region around 124-126 GeV/c^2 .

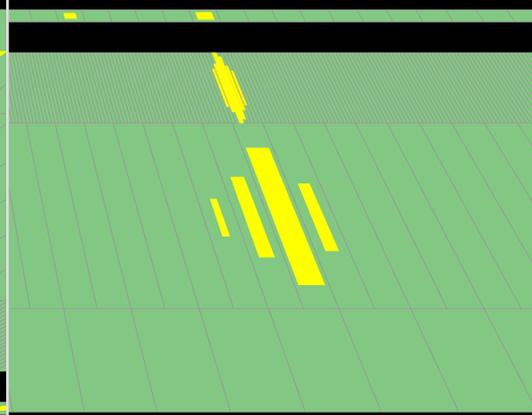
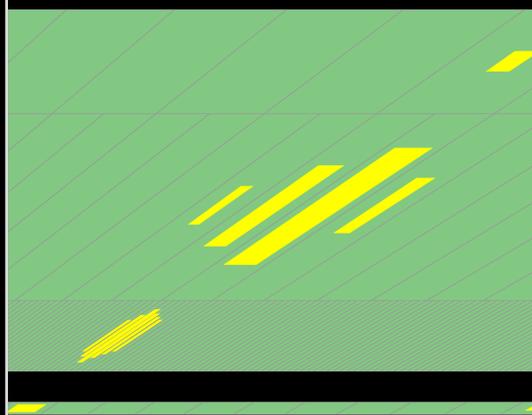
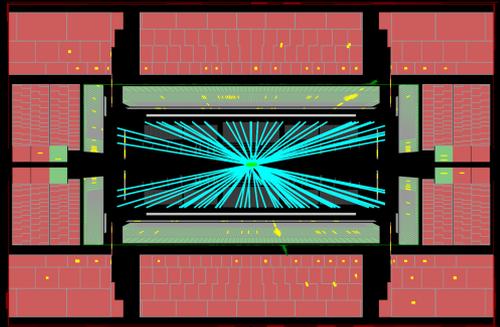
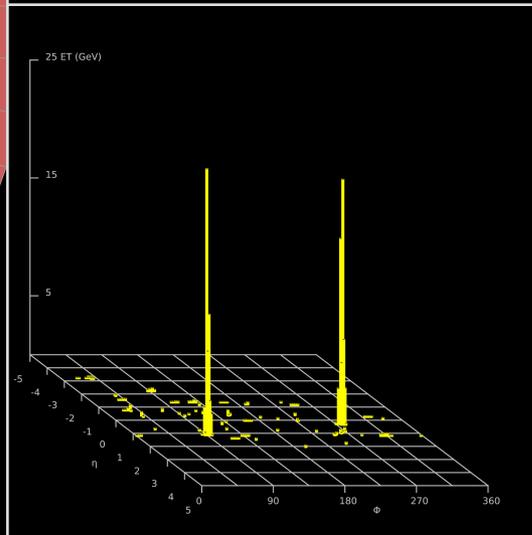
H → YY

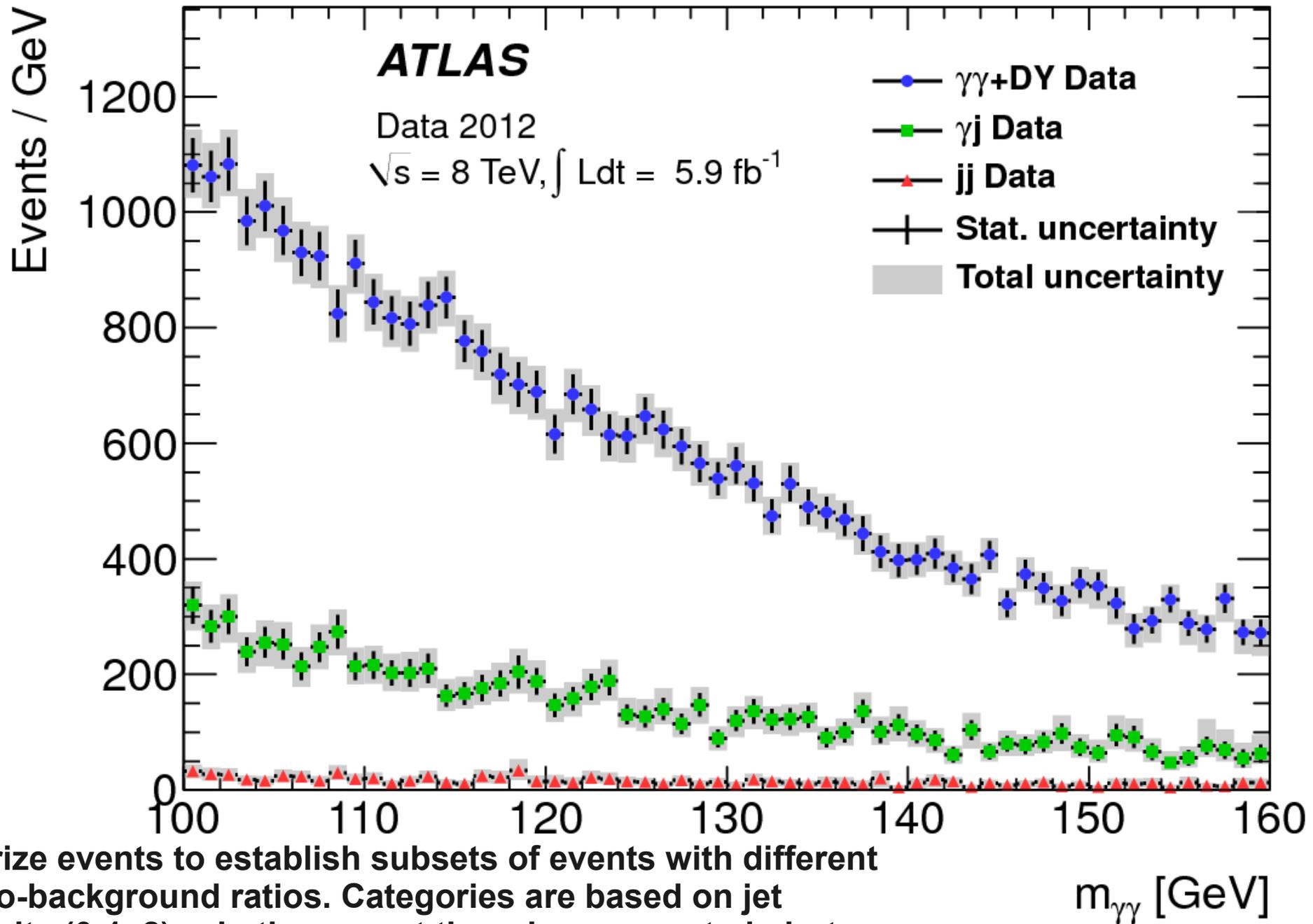


ATLAS EXPERIMENT

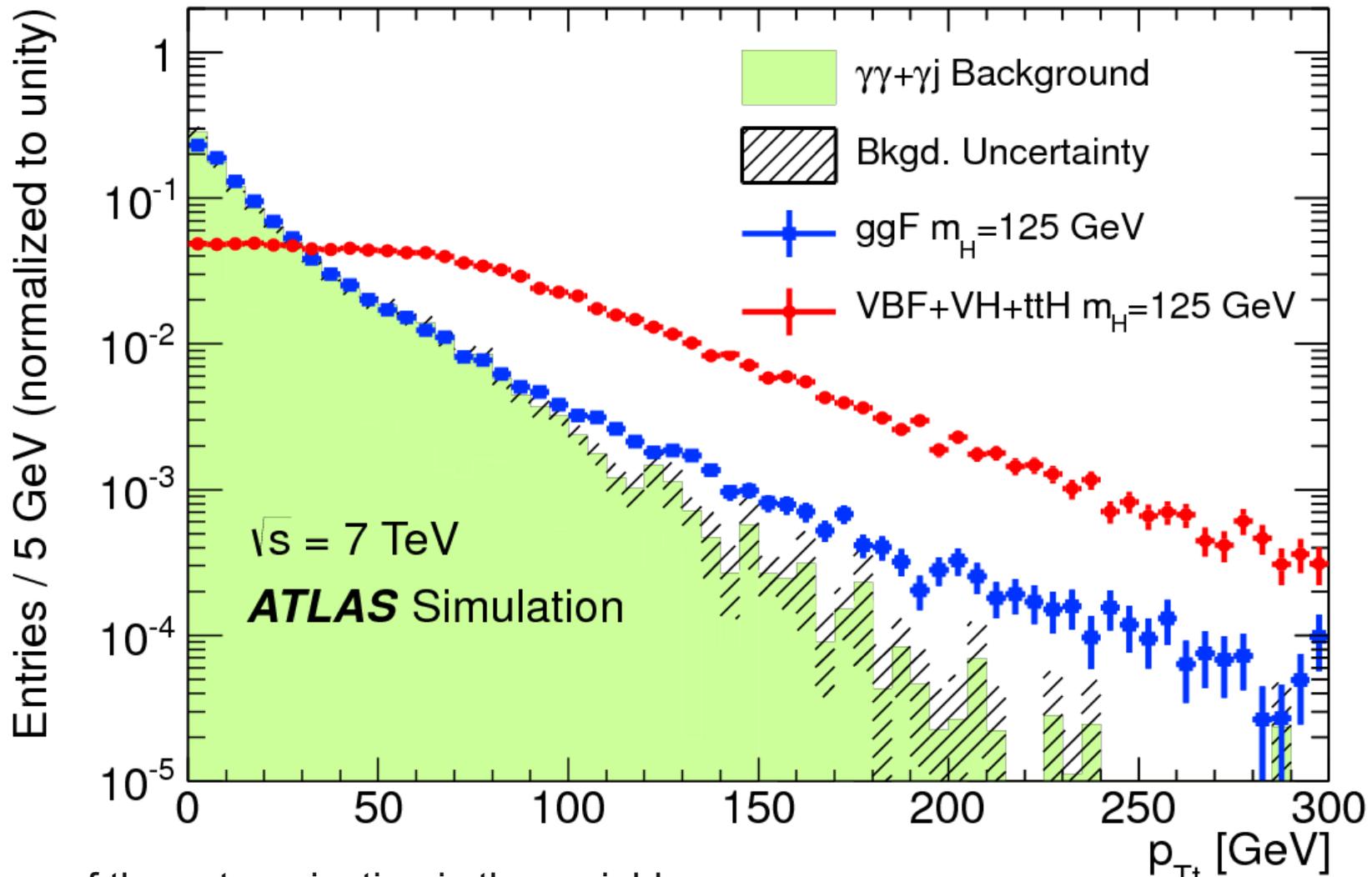
Run Number: 203779, Event Number: 56662314

Date: 2012-05-23 22:19:29 CEST

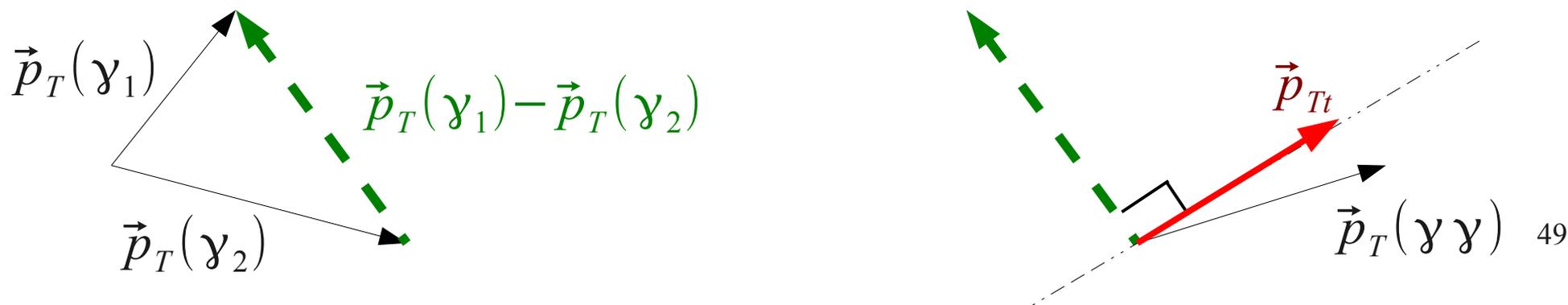


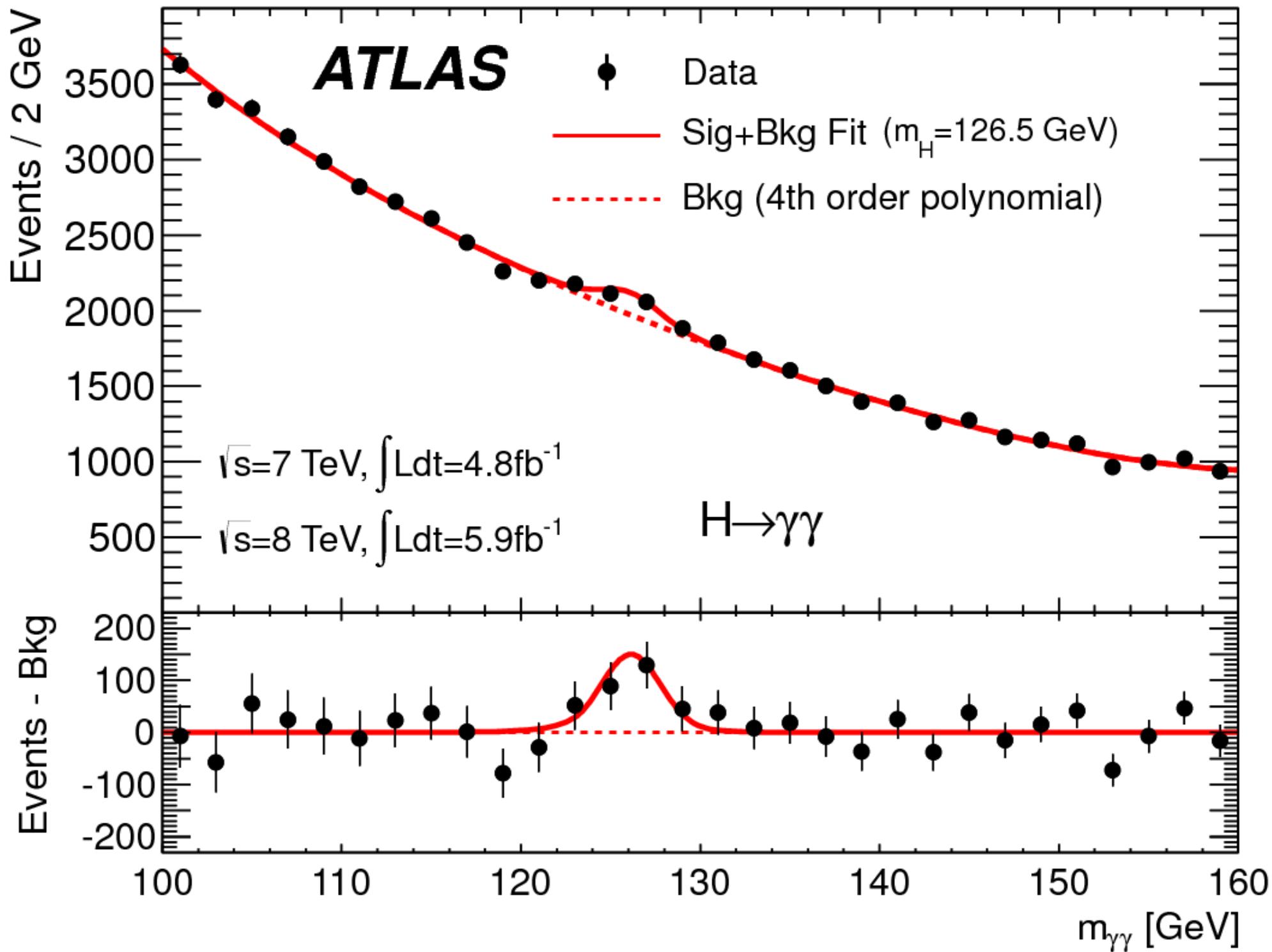


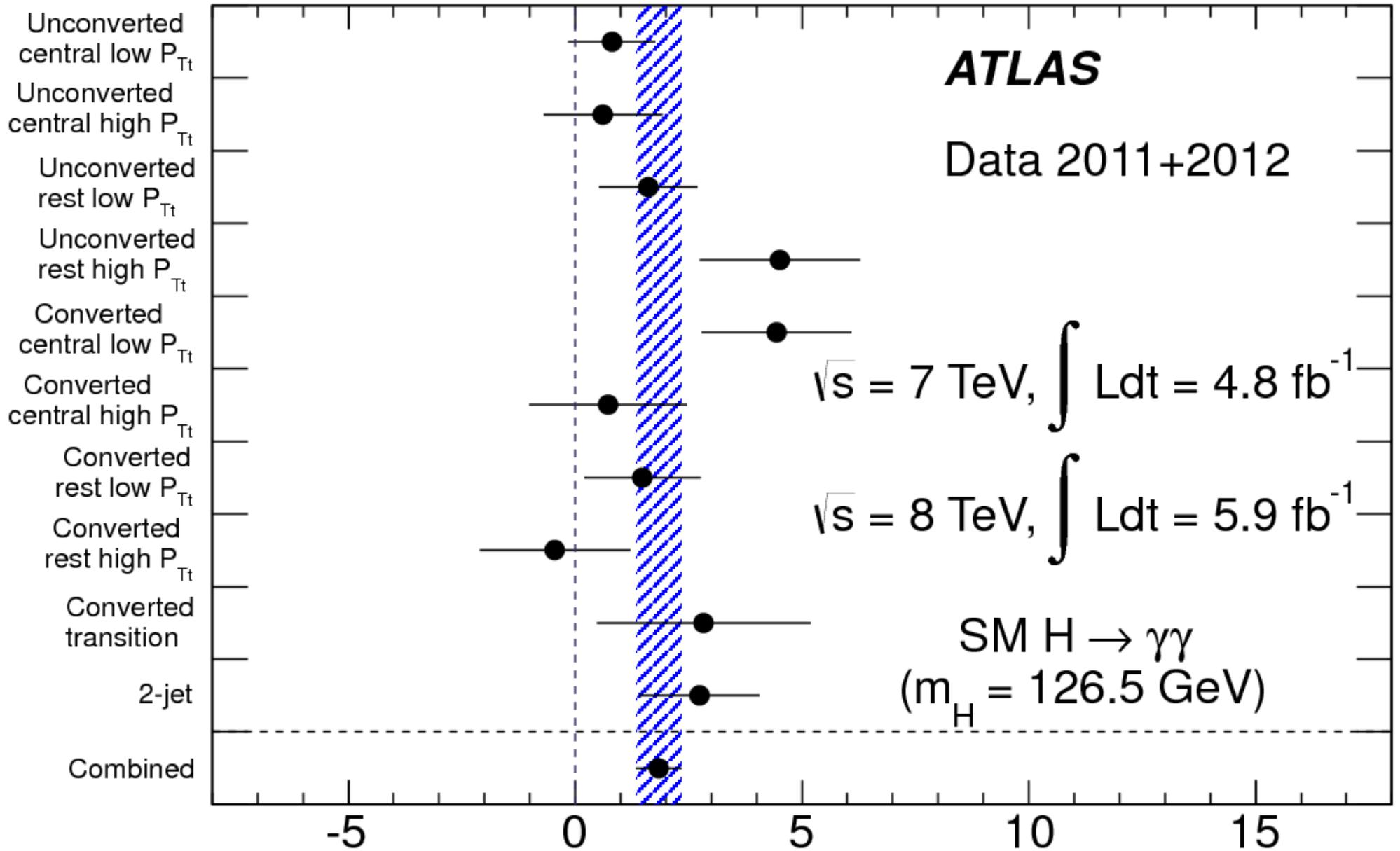
Categorize events to establish subsets of events with different signal-to-background ratios. Categories are based on jet multiplicity (0-1, 2), whether or not there is a converted photon, the polar angle of the photons, and the angular relationships of the photons . . .



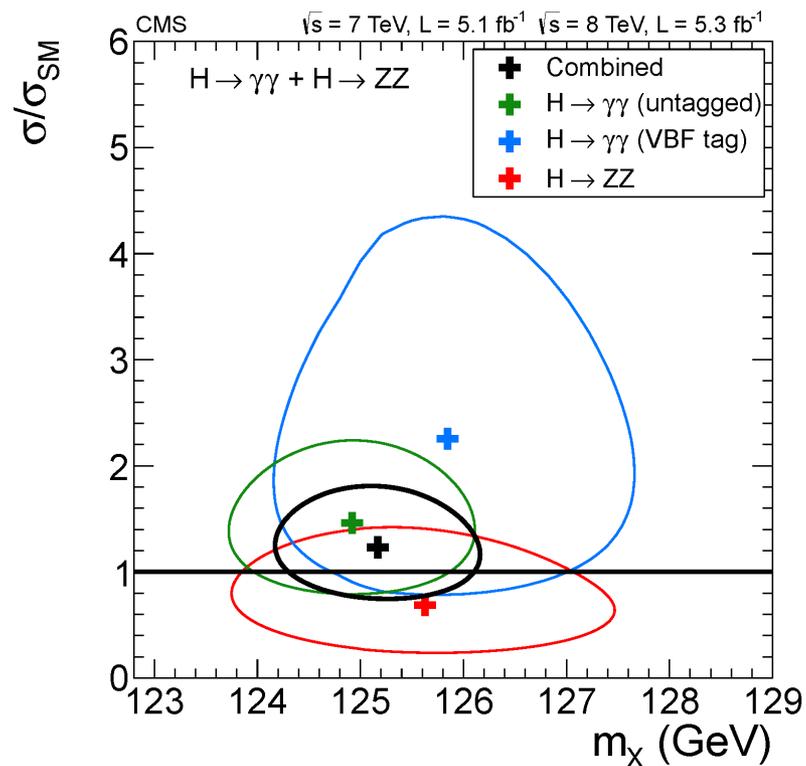
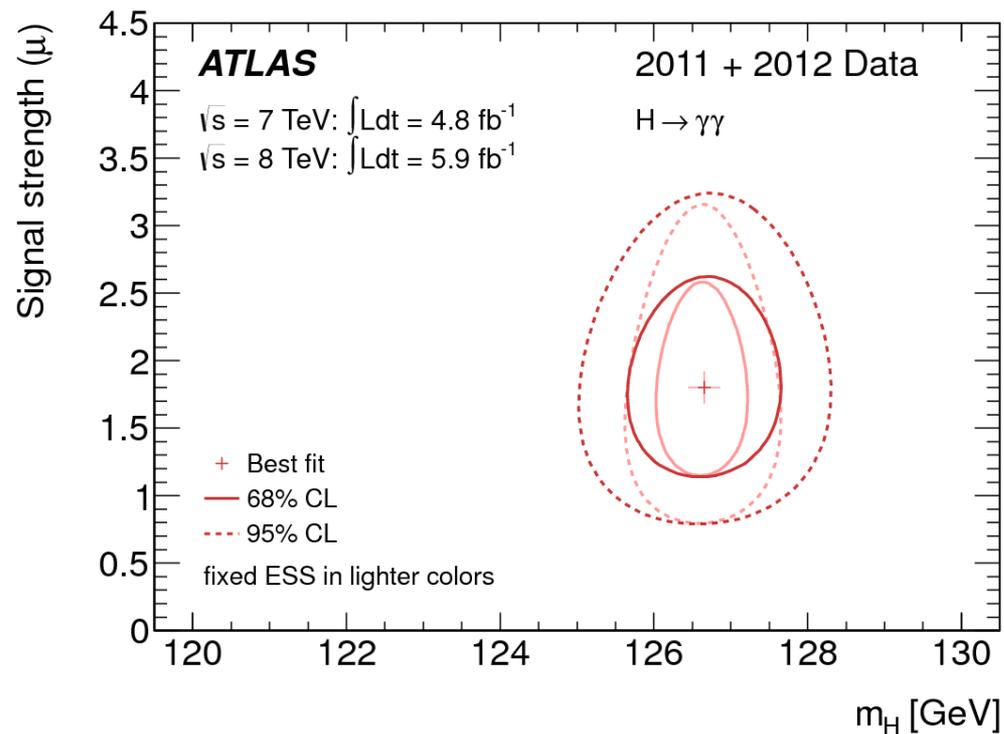
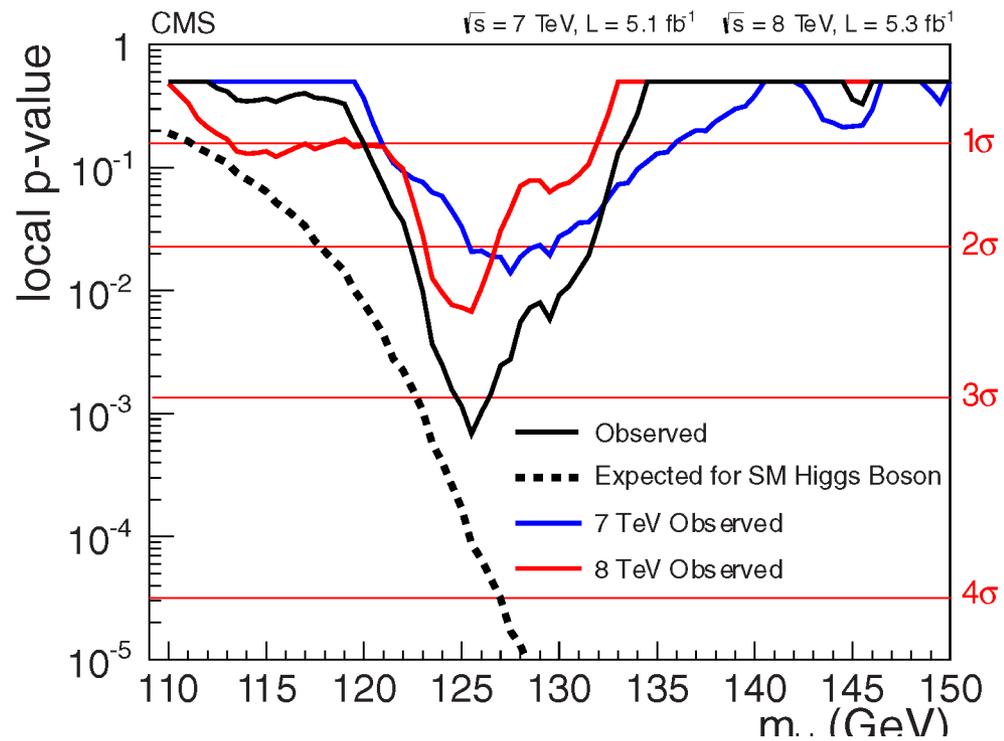
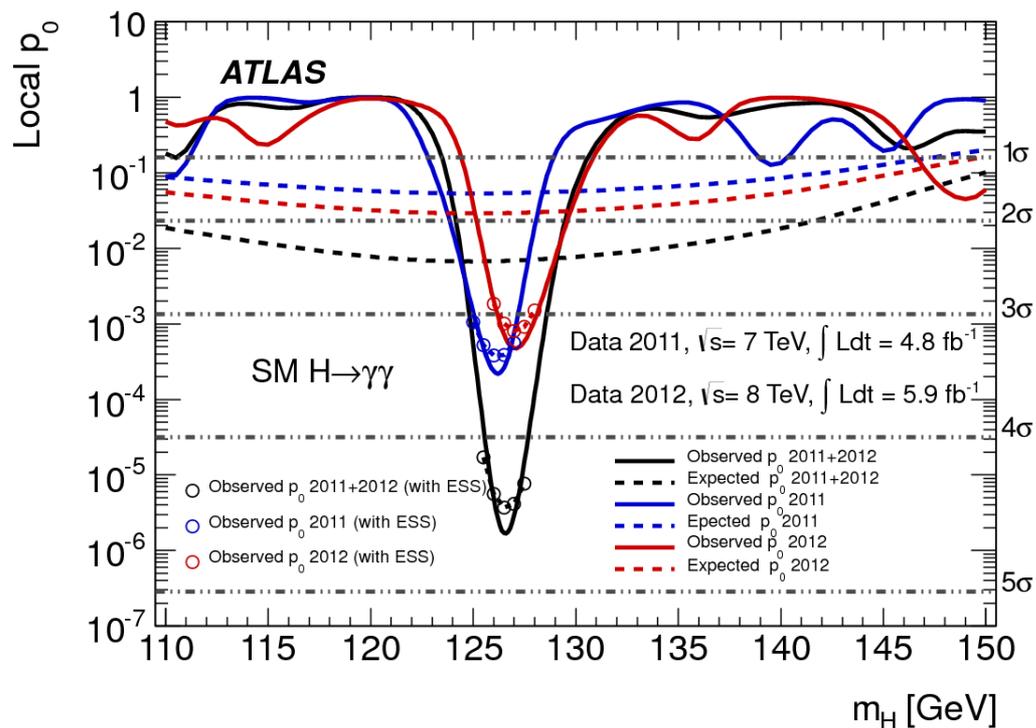
One of the dimensions of the categorization is the variable p_{Tt} :







$$\sigma_{measured} = \mu \cdot \sigma_{predicted}$$



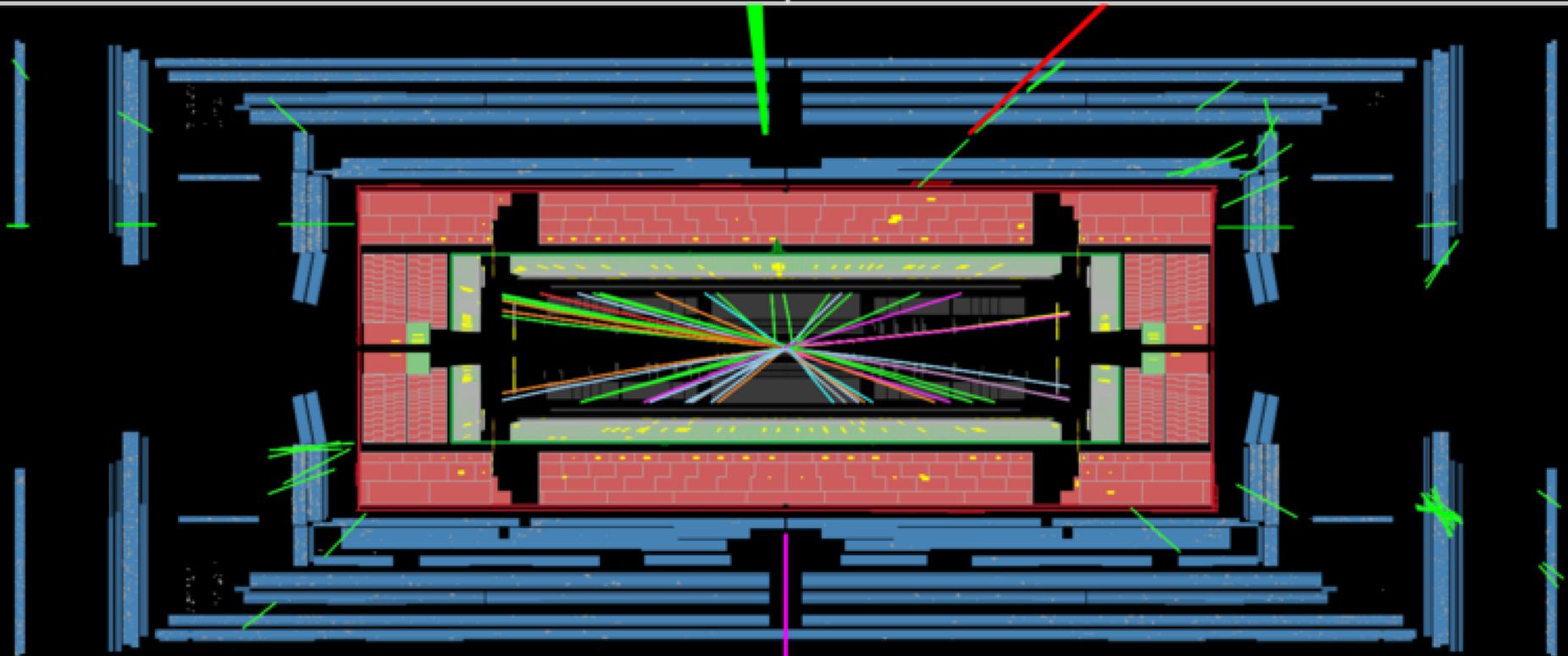
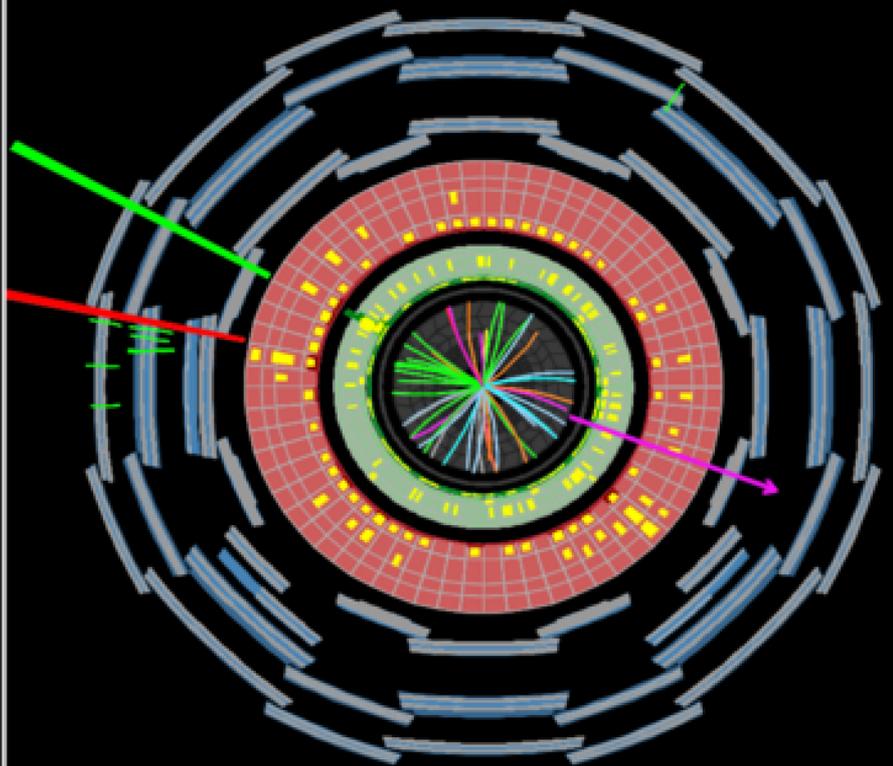
$$H \rightarrow W^+W^-$$



ATLAS EXPERIMENT

Run Number: 204026, Event Number: 33133446

Date: 2012-05-28 07:23:47 CEST



Events / 10 GeV

ATLAS

$\sqrt{s} = 8 \text{ TeV}, \int \text{Ldt} = 5.8 \text{ fb}^{-1}$

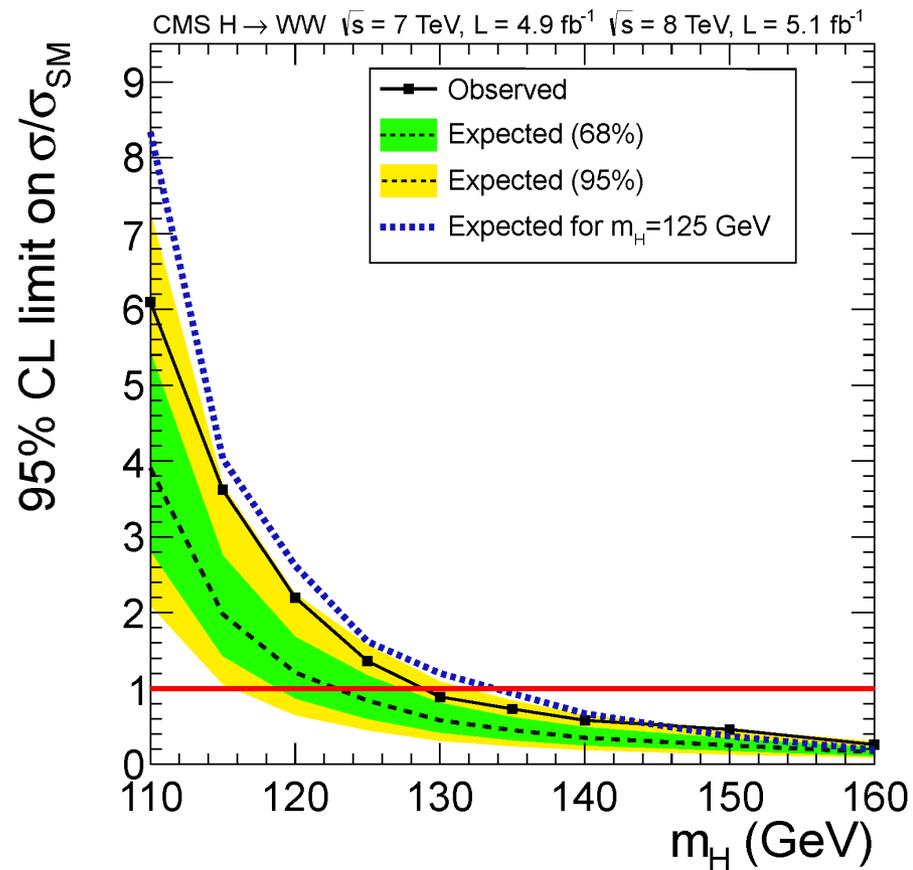
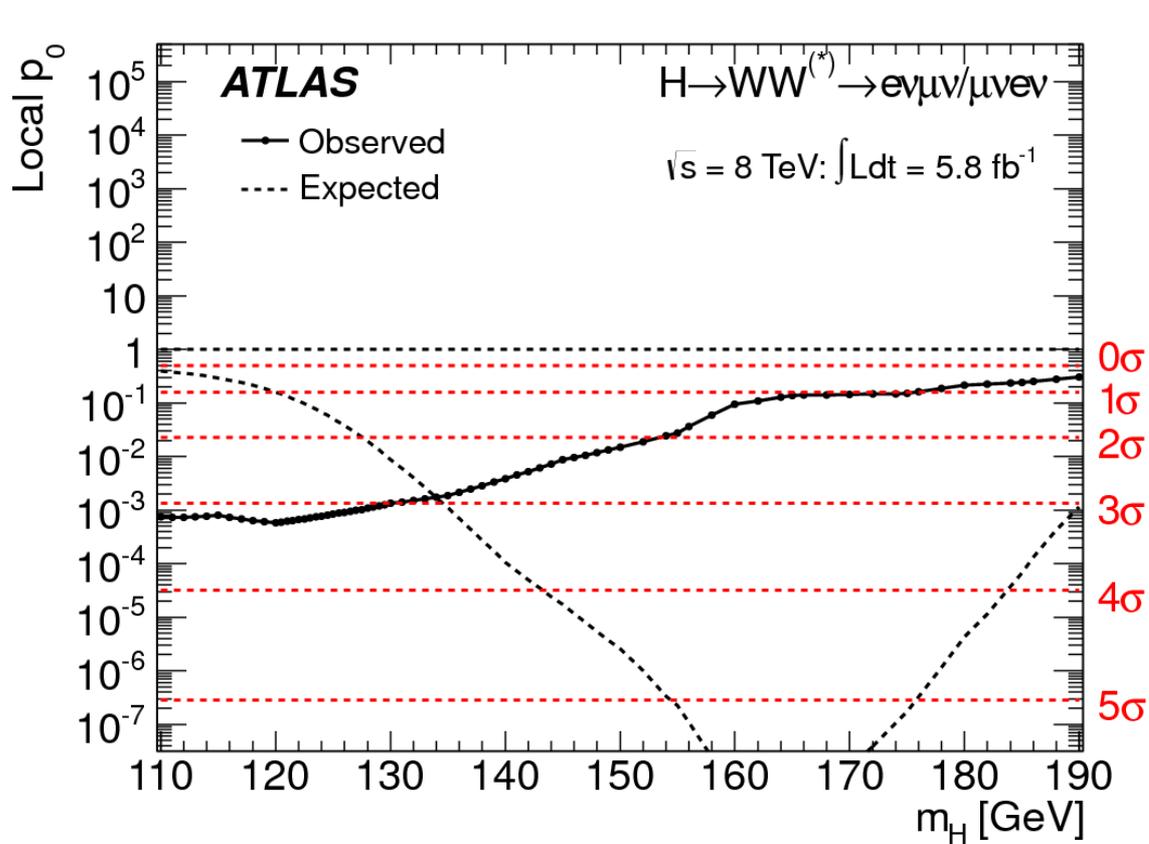
$H \rightarrow WW^{(*)} \rightarrow e\nu\mu\nu/\mu\nu e\nu + 0/1 \text{ jets}$



$$m_T = \sqrt{(E_T^{\text{ll}} + E_T^{\text{miss}})^2 - |\vec{p}_T^{\text{ll}} + \vec{p}_T^{\text{miss}}|^2}$$

This variable has a physical threshold at the mass of the resonance, but is broad below the threshold (unlike a pure invariant mass). Experimentally, it is further broadened/smeared by reconstruction effects.



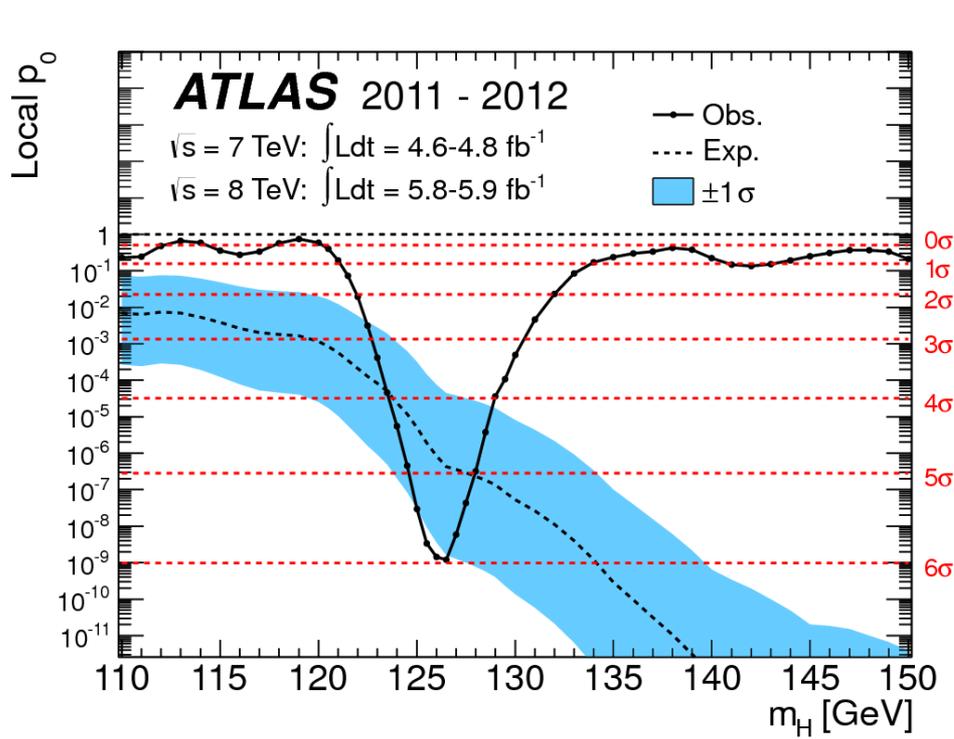


The excesses observed by both ATLAS and CMS are “broad,” as expected from this channel. But both experiments independently observe such an excess in their data.

Experimental Summary

ATLAS Publication:
Phys. Lett. B 716 (2012) 1-29

CMS Publication:
Phys. Lett. B 716 (2012) 30-61



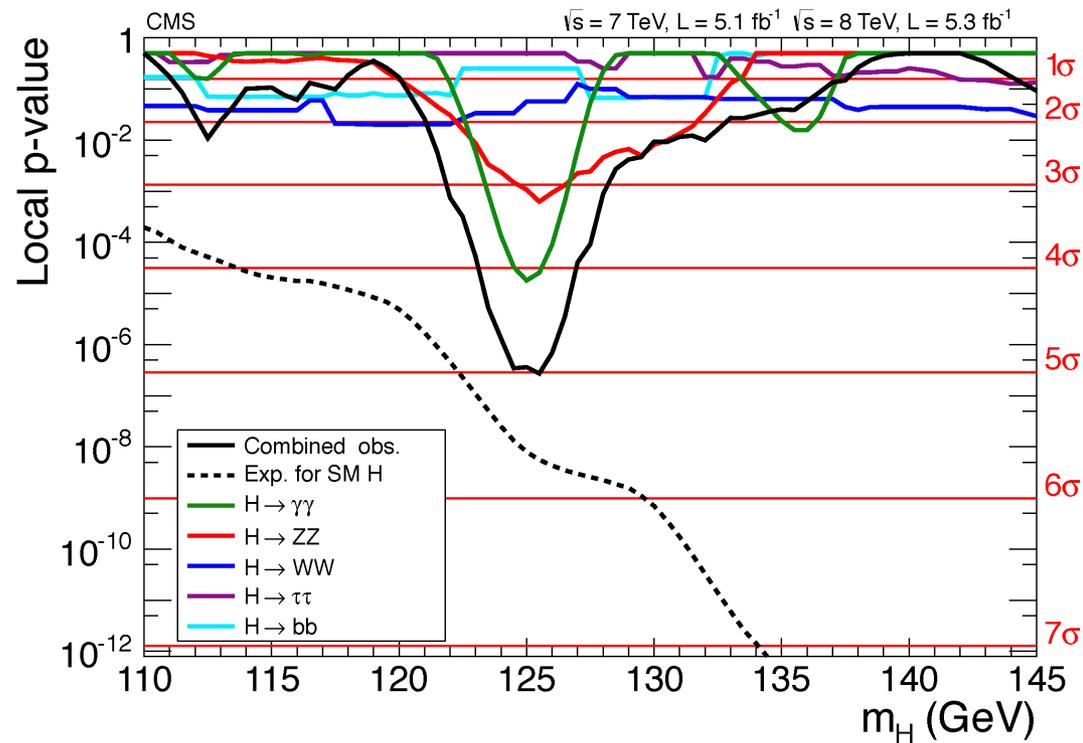
Strength of the ATLAS excess

The maximum value of the local significance is 6.0σ .

Taking into account systematic effects from energy measurement resolution, that maximum reduces to 5.9σ .

Taking into account the fact that the search actually looks over a wide range of masses (and fluctuations can occur when you look in more places), the global significance is:

5.1σ (global significance)

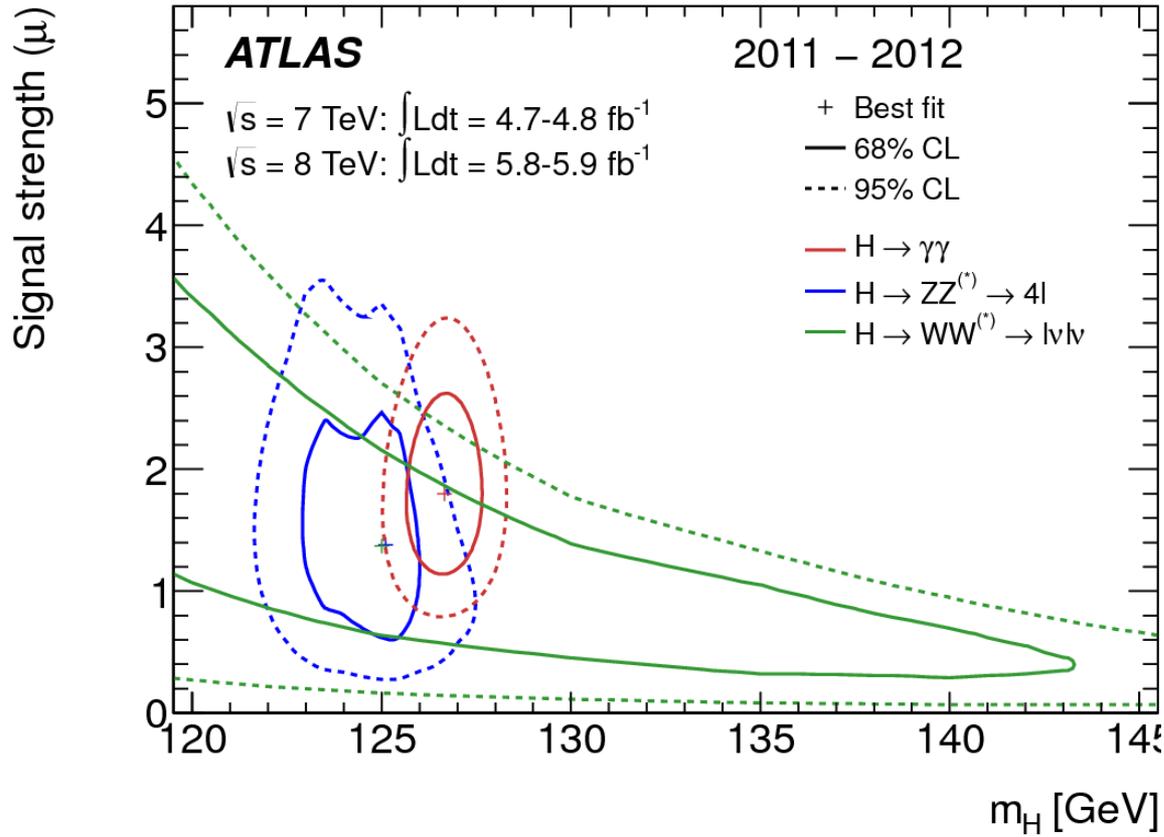


Strength of the CMS excess

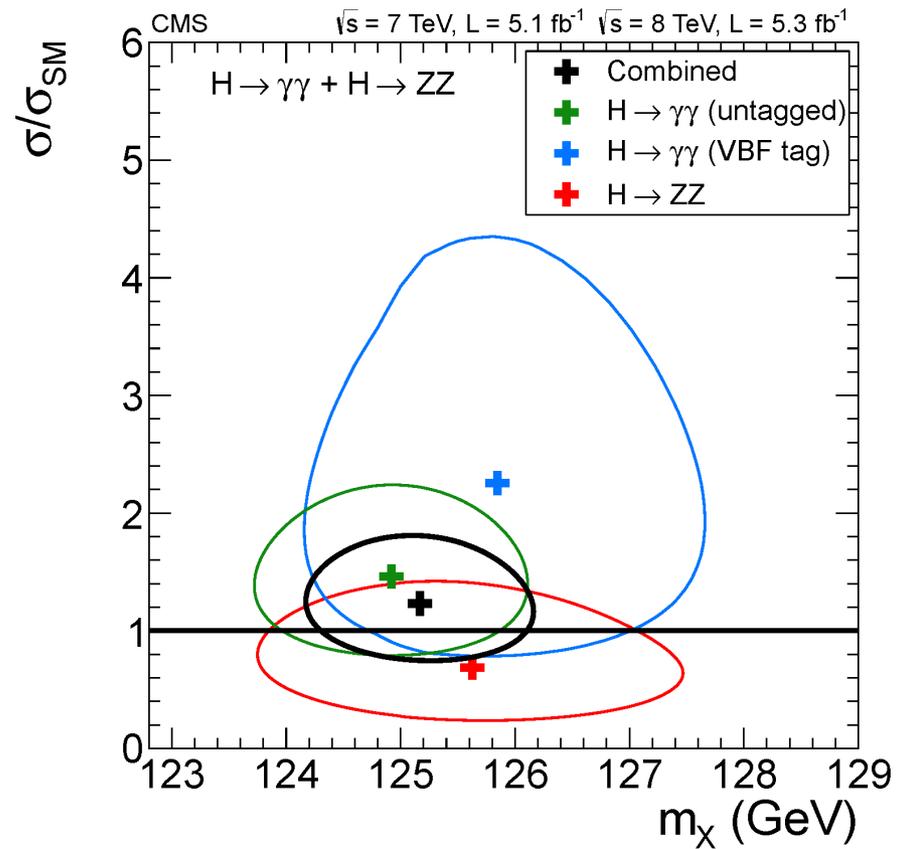
The maximum value of the local significance is 5.0σ .

Taking into account the fact that the search actually looks over a wide range of masses (and fluctuations can occur when you look in more places), the global significance is:

4.5σ (global significance)



ATLAS Mass Measurement

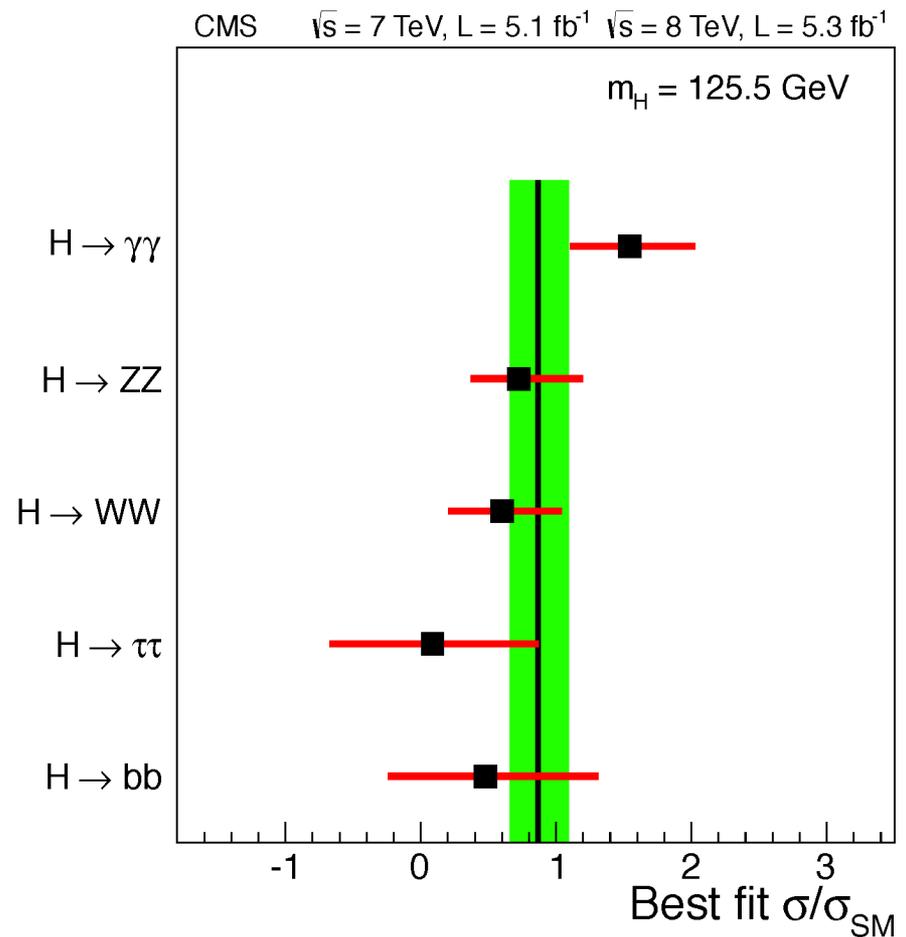
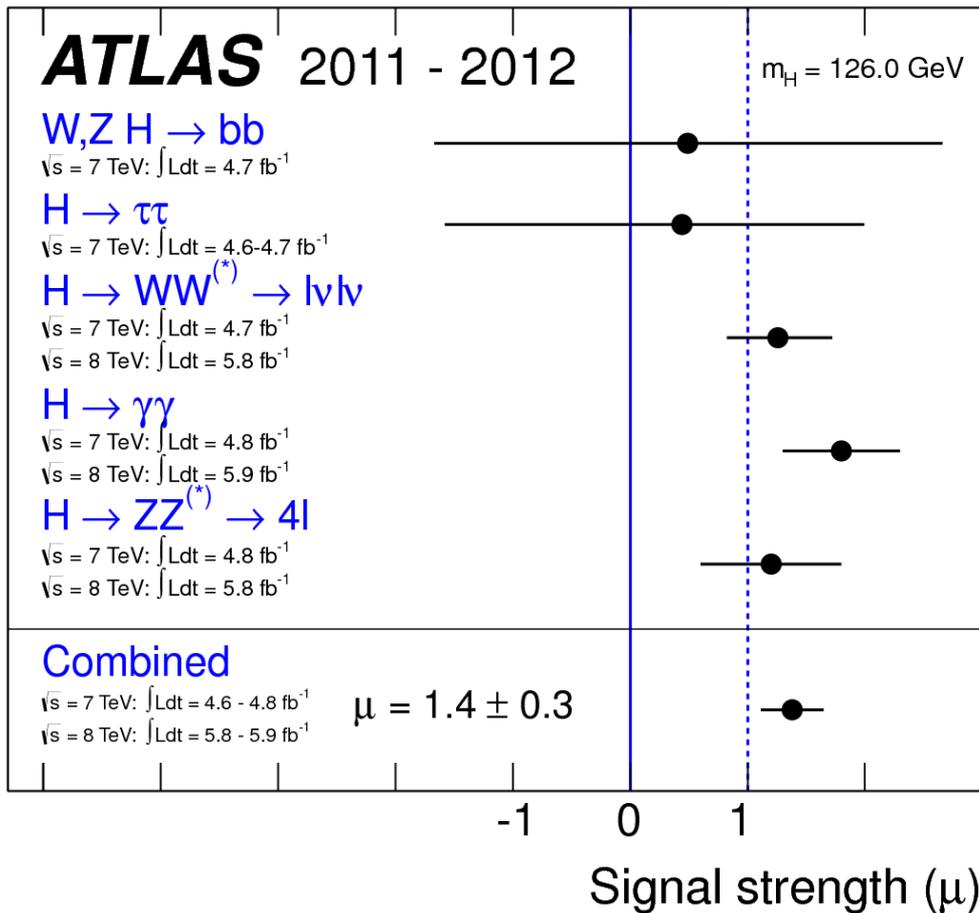


CMS Mass Measurement

$$M_H = 126.0 \pm 0.4 (\text{stat.}) \pm 0.4 (\text{syst.}) \text{ GeV}/c^2$$

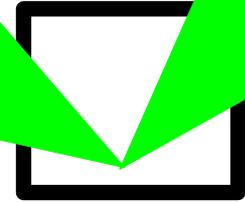
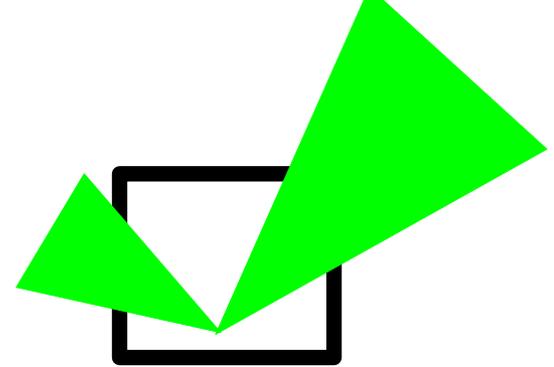
$$M_H = 125.3 \pm 0.4 (\text{stat.}) \pm 0.5 (\text{syst.}) \text{ GeV}/c^2$$

The mass is consistent between the two experiments.

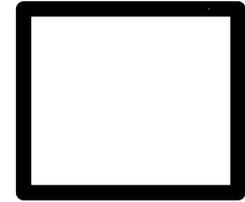


Within currently published statistics, this particle is so-far consistent with the Higgs boson hypothesis. But we want to reduce the uncertainties here (more statistics, better experimental techniques, better theoretical modeling) to see if that $H \rightarrow \gamma\gamma$ “excess” holds up.

PARTICLE MASS



PARTICLE SPIN

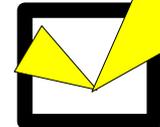
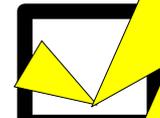
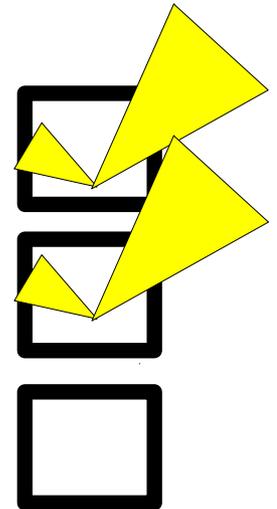


PARTICLE CP

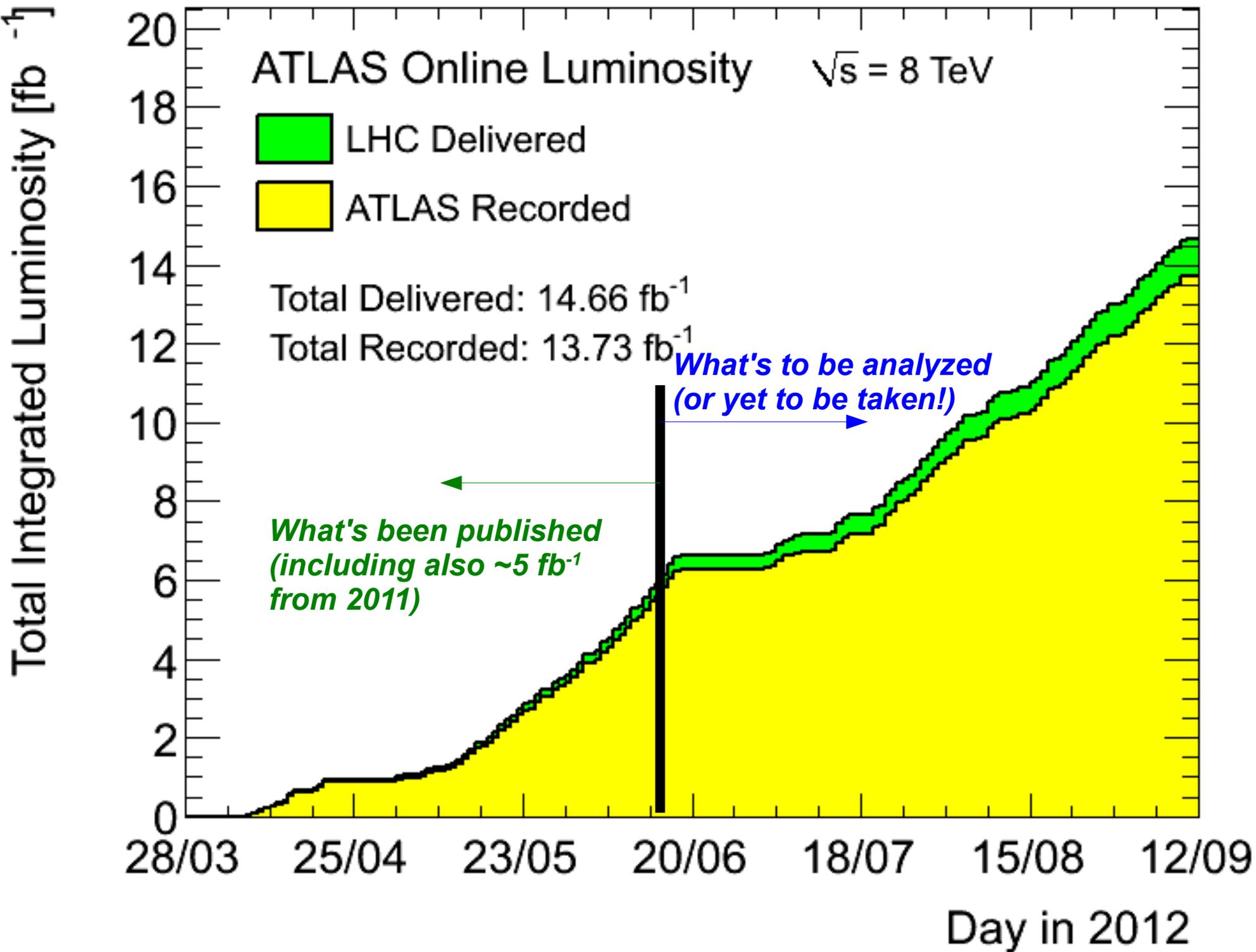


PARTICLE COUPLINGS

*[for instance:
rate of production
rate of decay]*



The Next Steps



Outlook

- Properties of the Higgs
 - Efforts underway to determine spin and CP quantum numbers of this new particle
 - Efforts continue to look at existing and new decay modes of the new particle
- Beyond the Higgs
 - Efforts are expanding to look for cousins of this new particle (e.g. additional Higgs particles that might give mass to things like dark matter or neutrinos - “Godzilla Particles”)
 - Efforts continue to generally look for signs of physics not predicted in the Standard Model (extra dimensions of space, supersymmetric partners of the Standard Model particles, a connection between gravity and the Standard Model, etc.)

All figures and numbers in this talk were obtained from public ATLAS and CMS repositories. For copies, or more plots and numbers than are dreamed of in your philosophy, check out these links:

<https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/HIGG-2012-27/>
<https://twiki.cern.ch/twiki/bin/view/CMSPublic/Hig12028TWiki>

Event Displays, etc.:
<https://twiki.cern.ch/twiki/bin/view/AtlasPublic>