



Physics with Taus at the LHC

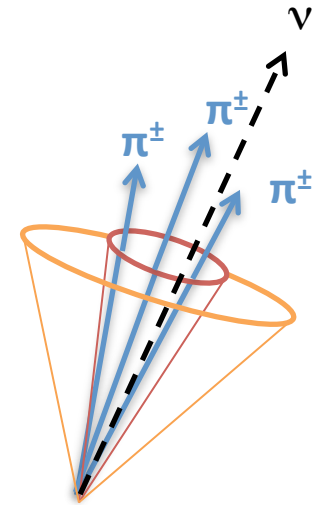
SMU

Monday, February 27, 2012

Sarah Demers
Yale University

Outline

- **WHY taus?**
 - and why not taus . . .
- ATLAS and CMS
 - some of these things are not like the other
- The building blocks
 - Triggering, Reconstructing, Identifying
- Measurements
 - Standard Model, Higgs (SM and MSSM), H^+
- Precision
 - Polarization



What are taus?



T-shirt design by Neil Davies

What are taus?



Taus are leptons



T-shirt design by Neil Davies

What are taus?



Taus are leptons

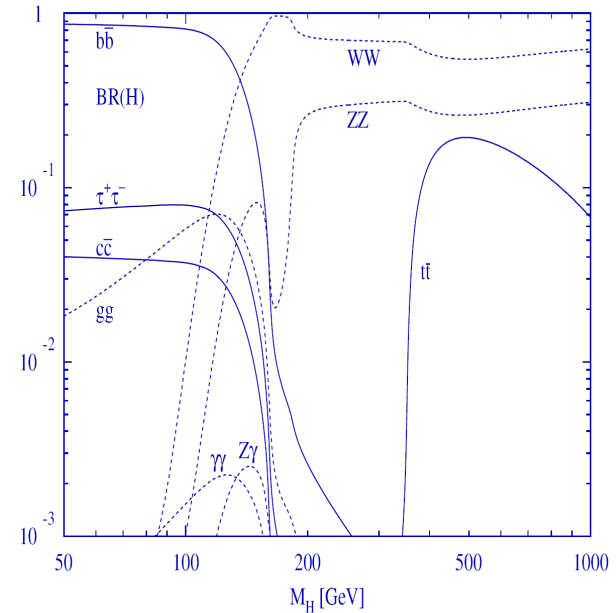
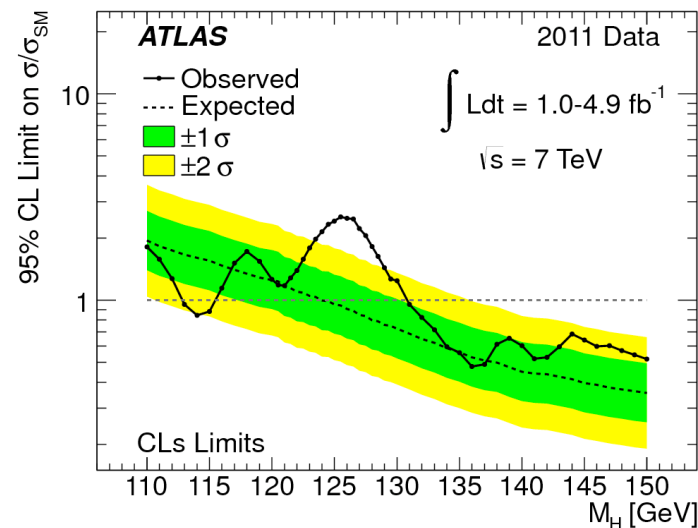
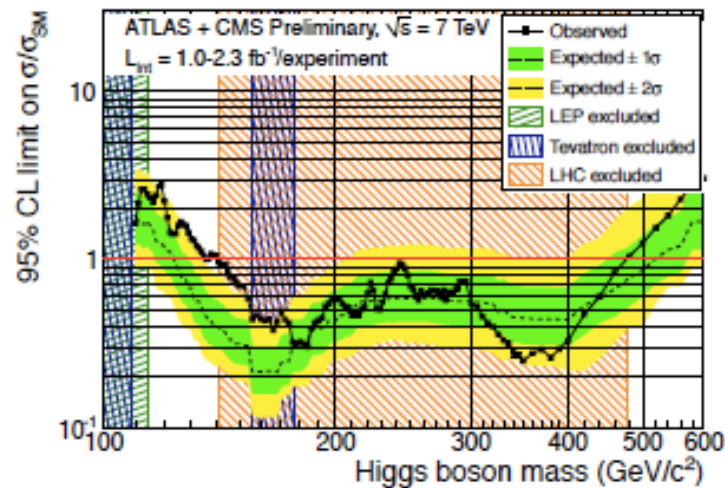
with $|\text{charge}| = 1$

in the 3rd generation



T-shirt design by Neil Davies

Taus as Probes for “New” Physics



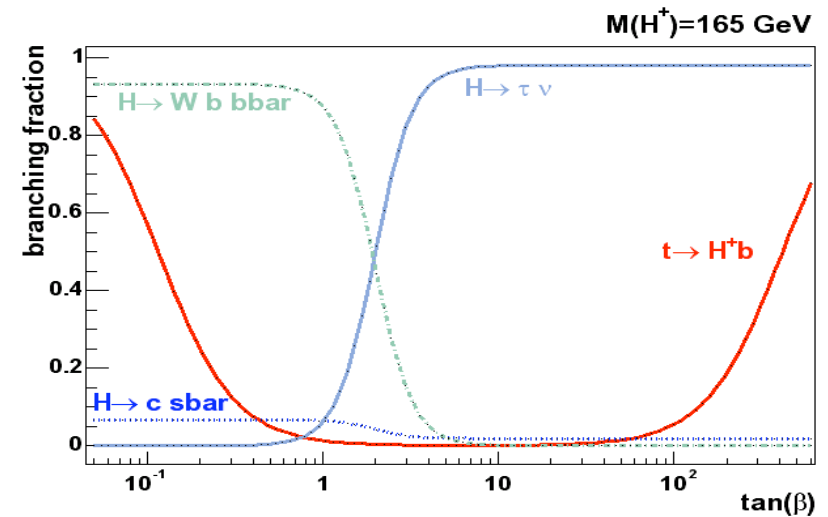
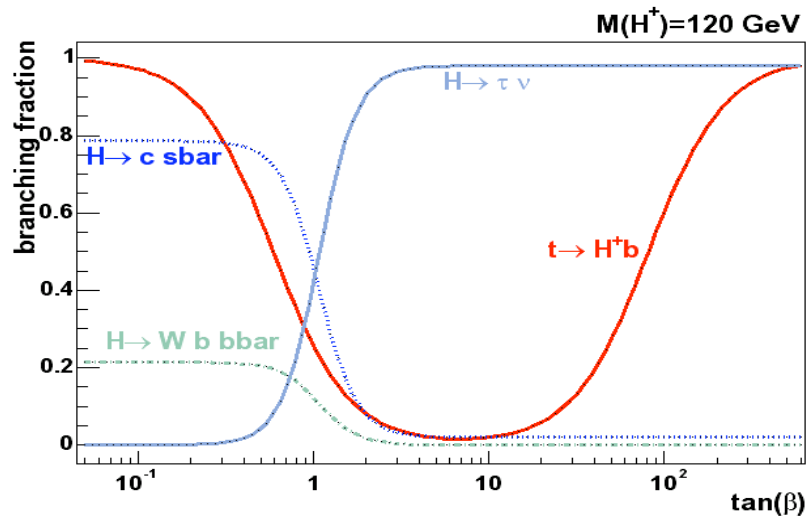
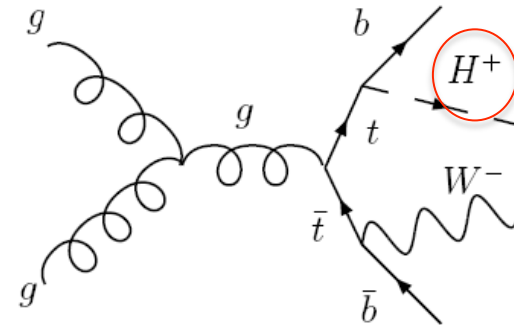
The tau, the heaviest lepton, couples strongly to the Higgs, and is relevant for a 5σ discovery in the important, but challenging, range of 115 – 130 GeV. This analysis could rely (at least partially) on tau triggers.

Channel analyzed by Stephen and Aidan

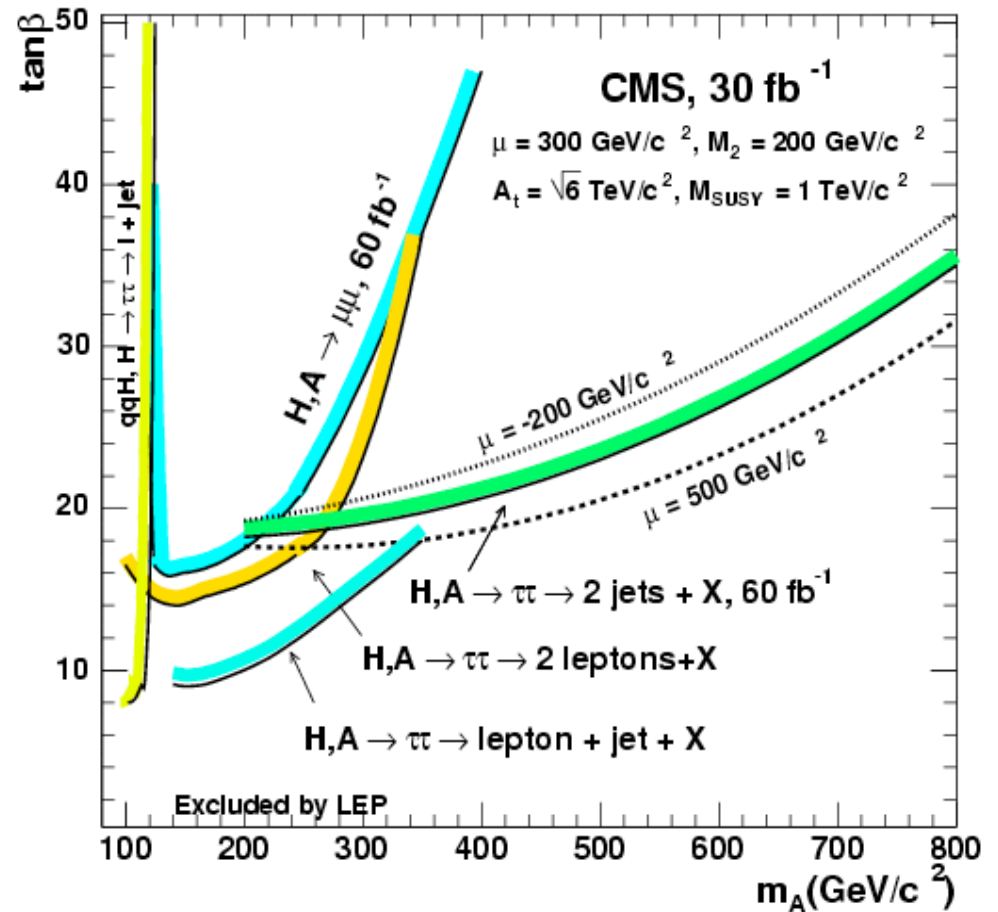
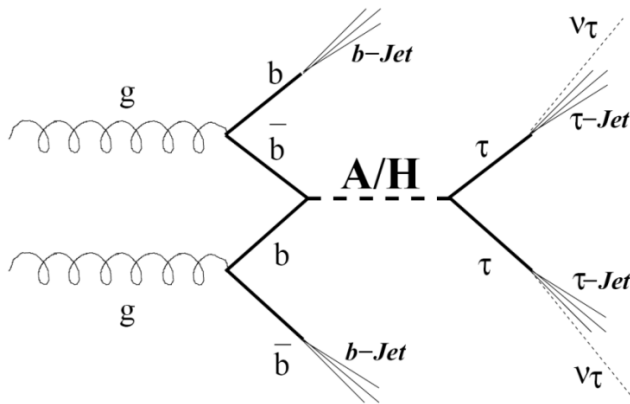


Charged Higgs (SUSY)

*With $\tan\beta$ over ~ 3 , H^+ decays
almost exclusively to taus*

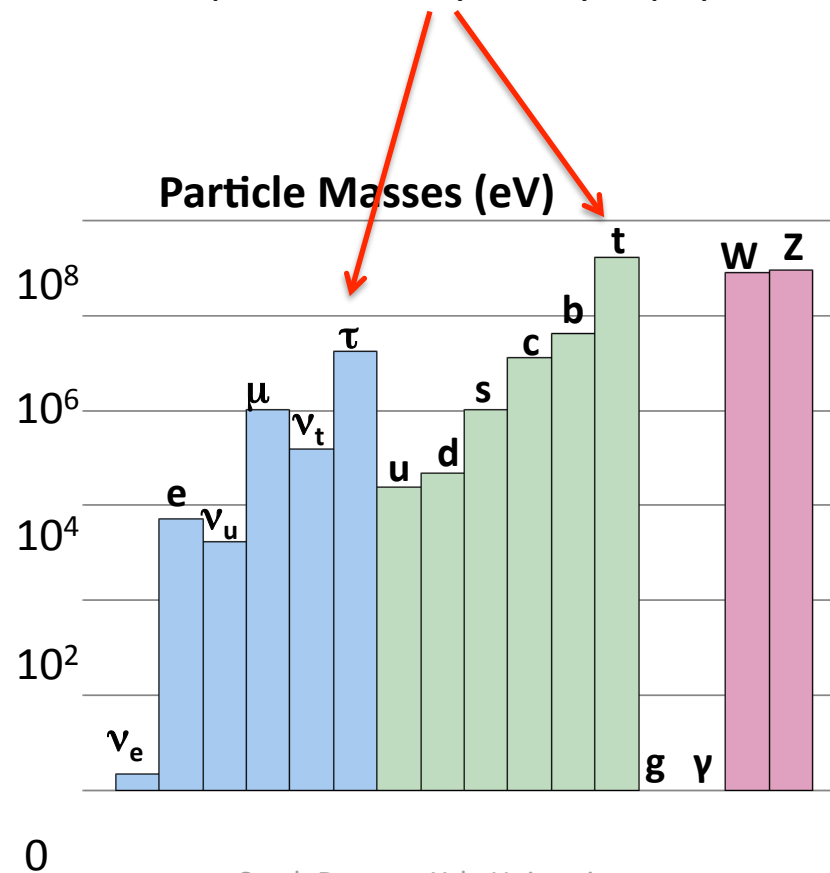


Neutral Higgs (SUSY)



Life without the Higgs?

Many Higgsless theories include new particles with preferential couplings to the third generation, often motivated by trying to explain the very heavy top quark



Understanding New Phenomena



WE WANT Taus!

All decay modes should be explored to understand what we
find at the LHC

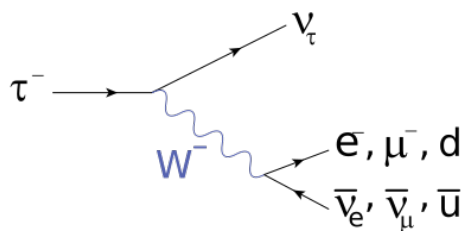
Not just electrons and muons!

Tau decays can carry information about the polarization of the
object that
decays into them

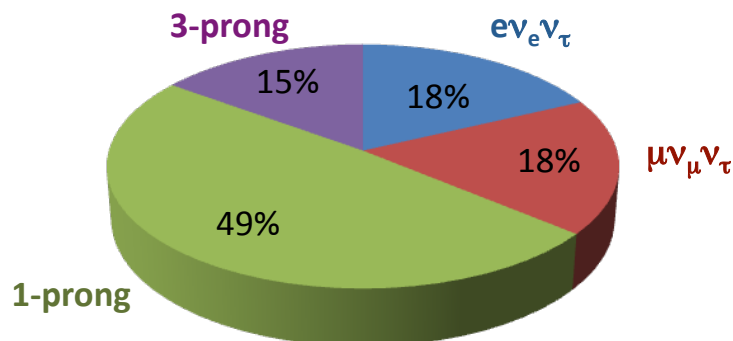
Why NOT taus?

The tau lifetime is 2.9×10^{-13} s.

We detect the decay products of the tau, not the tau itself.



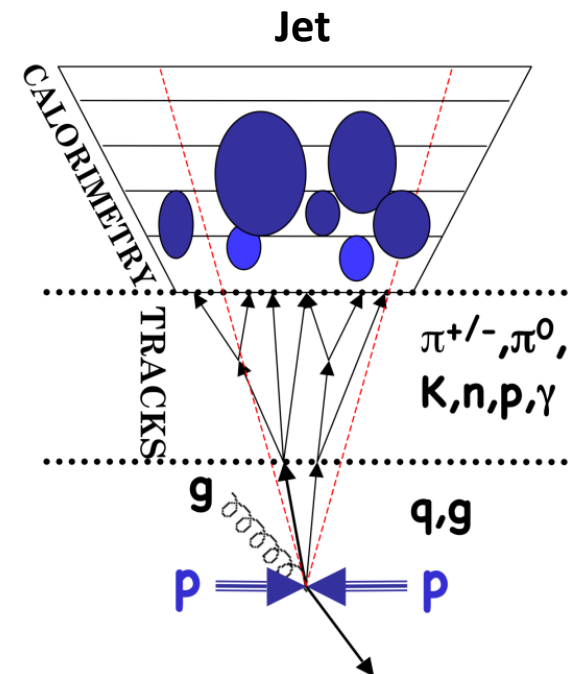
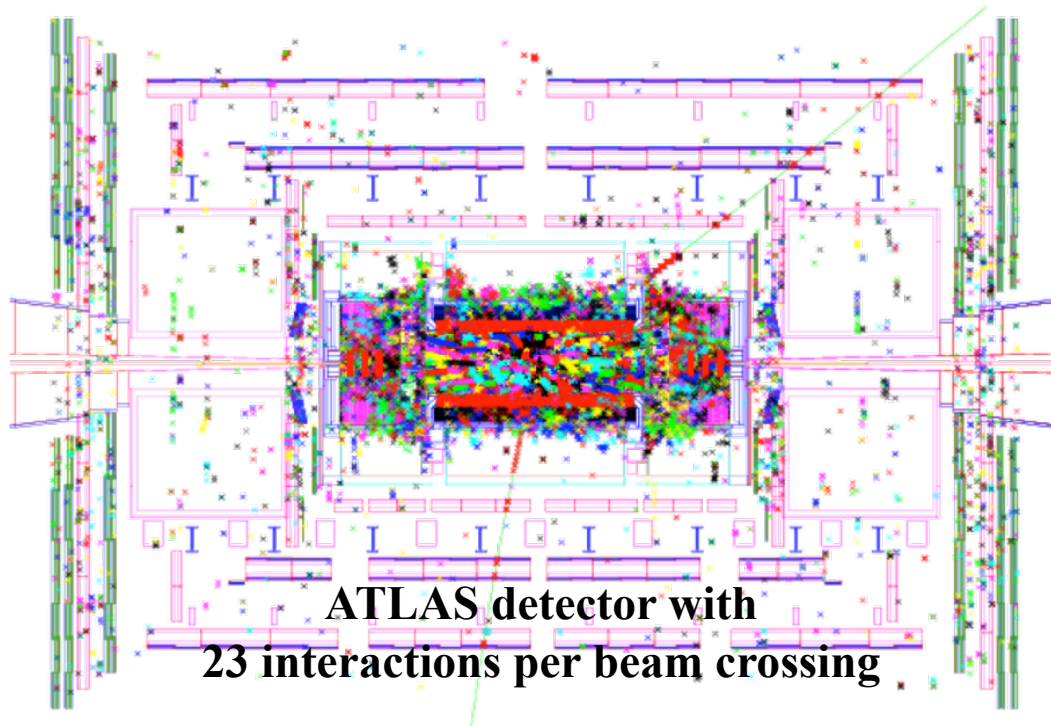
Tau Decay Modes

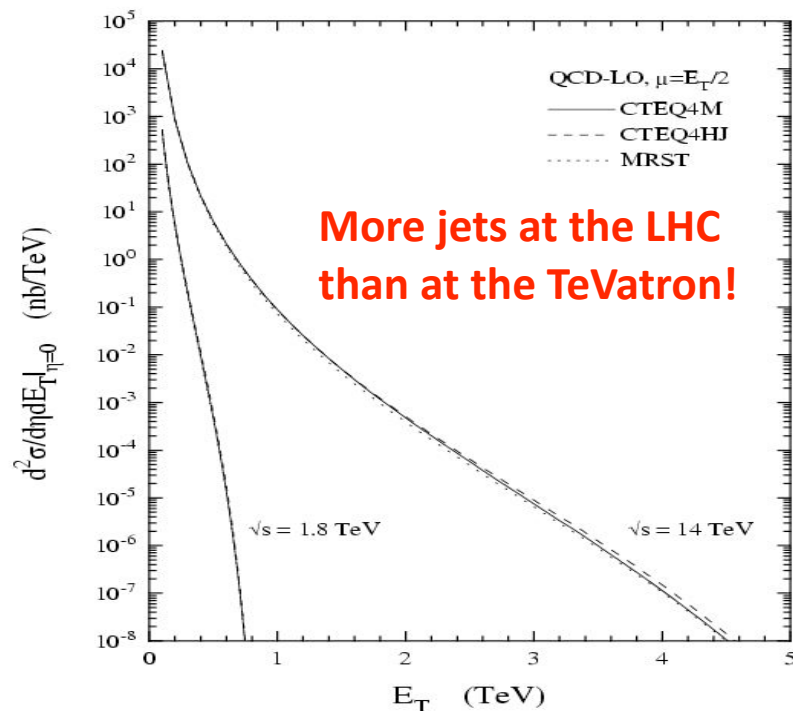


Tau branching ratios from
 10^8 simulated $Z \rightarrow \tau\tau$ events

Decay modes	TAUOLA-CLEO
$\tau \rightarrow e\nu_e\nu_\tau$	17.8 %
$\tau \rightarrow \mu\nu_\mu\nu_\tau$	17.4 %
$\tau \rightarrow h^\pm \text{ neutr. } \nu_\tau$	49.5 %
$\tau \rightarrow \pi^\pm \nu_\tau$	11.1 %
$\tau \rightarrow \pi^0 \pi^\pm \nu_\tau$	25.4 %
$\tau \rightarrow \pi^0 \pi^0 \pi^\pm \nu_\tau$	9.19 %
$\tau \rightarrow \pi^0 \pi^0 \pi^0 \pi^\pm \nu_\tau$	1.08 %
$\tau \rightarrow K^\pm \text{ neutr. } \nu_\tau$	1.56 %
$\tau \rightarrow h^\pm h^\pm h^\pm \text{ neutr. } \nu_\tau$	14.57 %
$\tau \rightarrow \pi^\pm \pi^\pm \pi^\pm \nu_\tau$	8.98 %
$\tau \rightarrow \pi^0 \pi^\pm \pi^\pm \pi^\pm \nu_\tau$	4.30 %
$\tau \rightarrow \pi^0 \pi^0 \pi^\pm \pi^\pm \pi^\pm \nu_\tau$	0.50 %
$\tau \rightarrow \pi^0 \pi^0 \pi^0 \pi^\pm \pi^\pm \pi^\pm \nu_\tau$	0.11 %
$\tau \rightarrow K_S^0 X^\pm \nu_\tau$	0.90 %
$\tau \rightarrow (\pi^0) \pi^\pm \pi^\pm \pi^\pm \pi^\pm \pi^\pm \nu_\tau$	0.10 %
other modes with K	1.30 %
others	0.03 %

Why NOT taus?





Jets, jets and more JETS!

Taus were challenging
at the TeVatron

We find ourselves even
more overwhelmed
by jets at the LHC

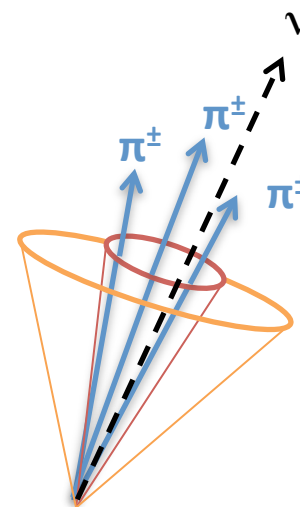
Sample	cross section X branching ratio	#events/8 hours (10^{31})
dijets (p_T 8 – 17 GeV)	1.7×10^{10} pb	5×10^9
dijets (p_T 17 – 35 GeV)	1.4×10^9 pb	4×10^8
dijets (p_T 35 – 70 GeV)	9.3×10^7 pb	3×10^7
$W \rightarrow \tau \nu$, $\tau \rightarrow had$	1.1×10^4 pb	3200
$Z \rightarrow \tau \tau$, $1\tau \rightarrow had$	1.55×10^3 pb	450

Definitions

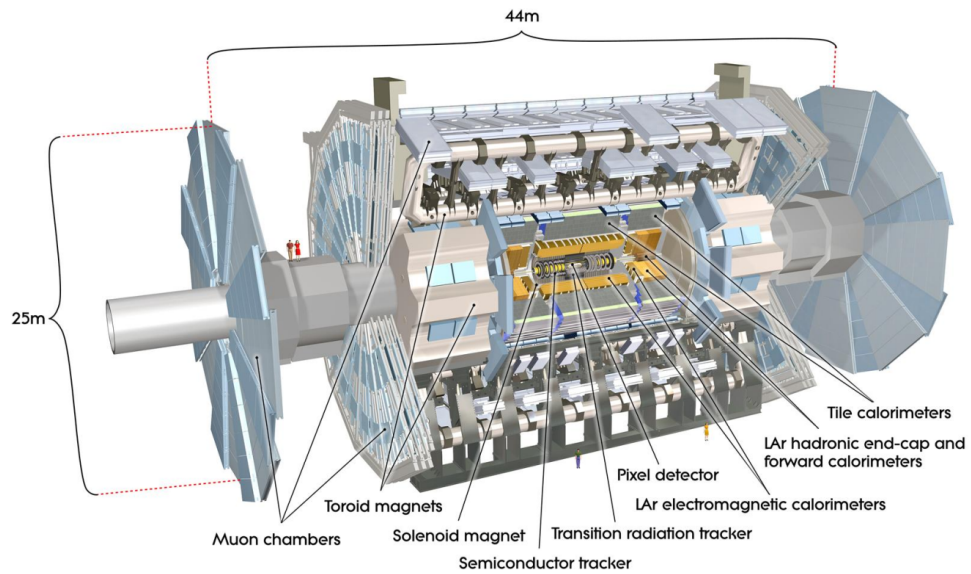
E_T = transverse energy
 p_T = transverse momentum
 $\sqrt{s}$ = center of mass energy

Introduction

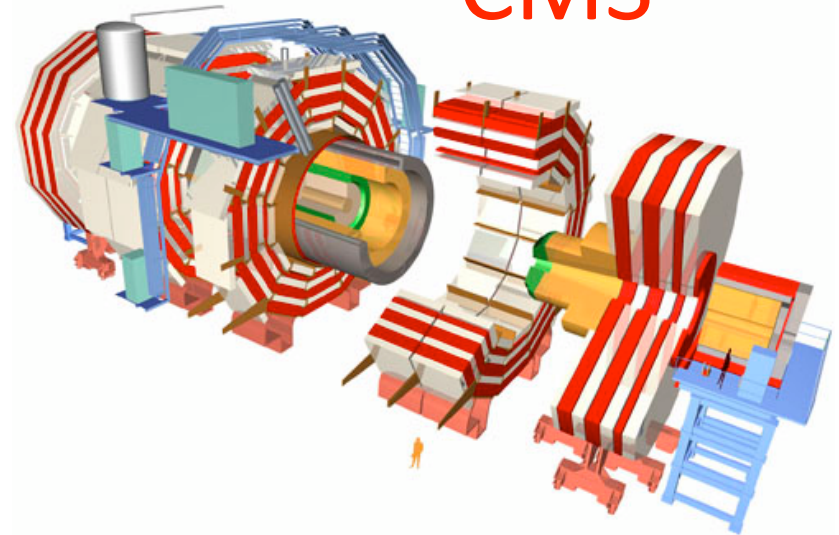
- *WHY* taus?
 - and why not taus . . .
- **ATLAS and CMS**
 - some of these things are not like the other



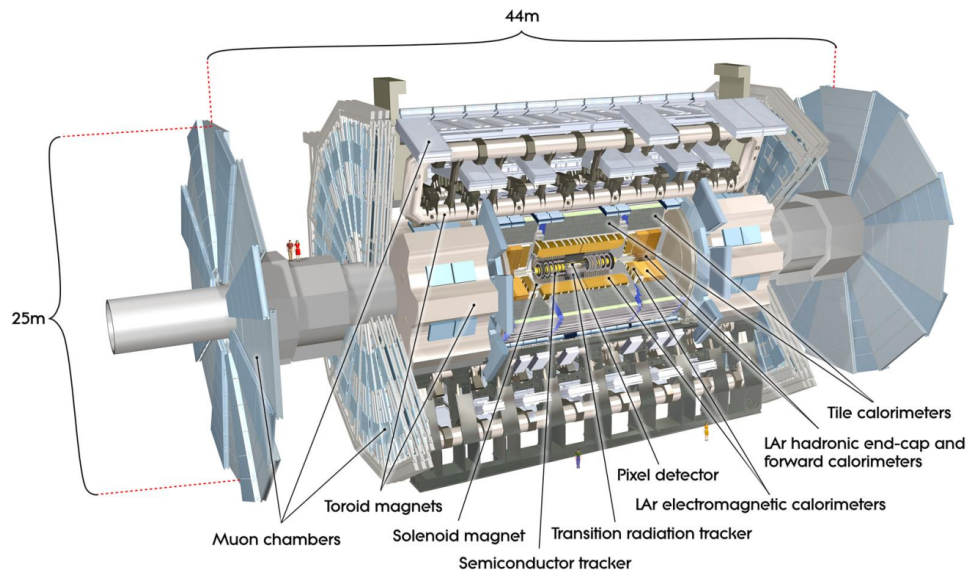
ATLAS



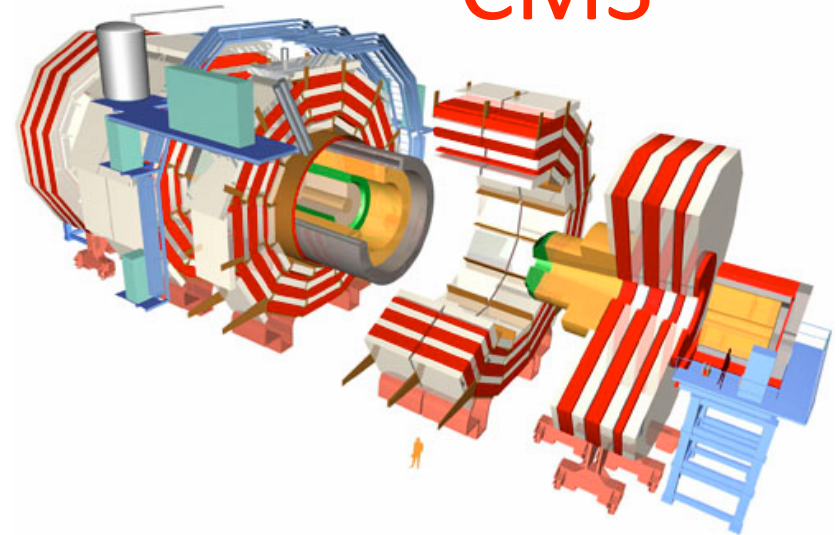
CMS



ATLAS



CMS



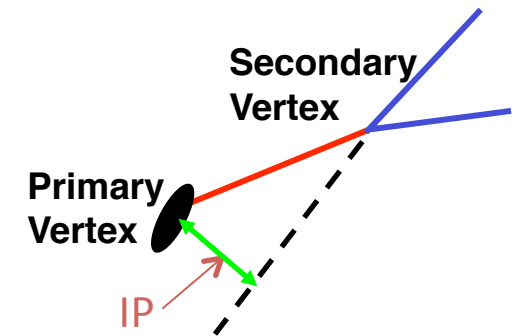
2 T

Magnetic Field

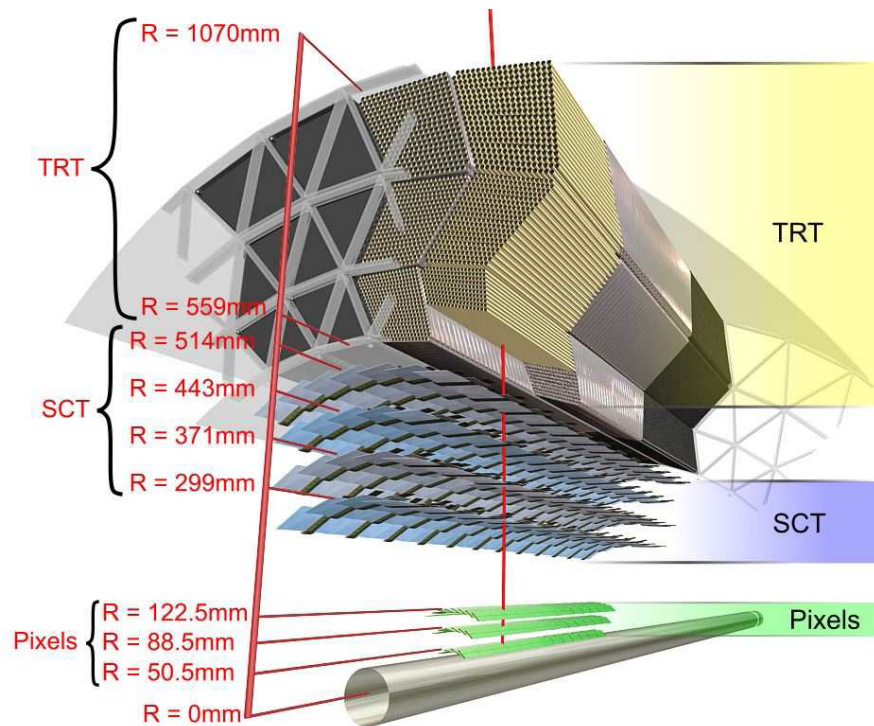
4 T



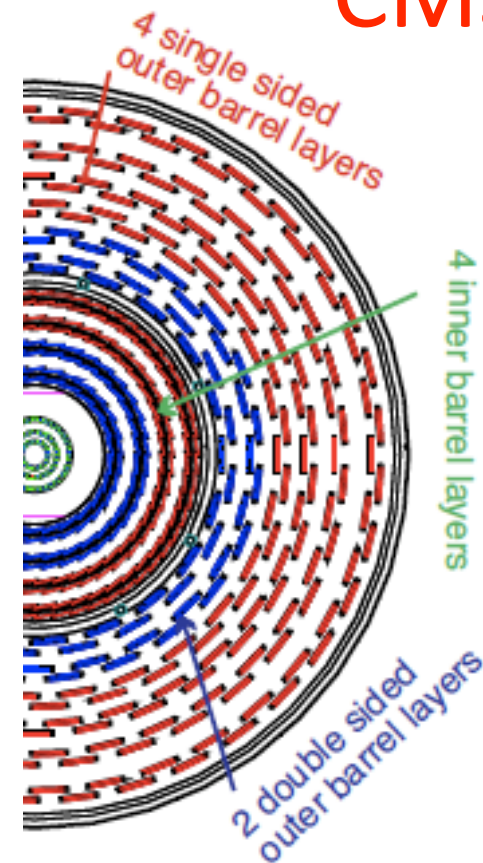
Tracking Chambers



ATLAS

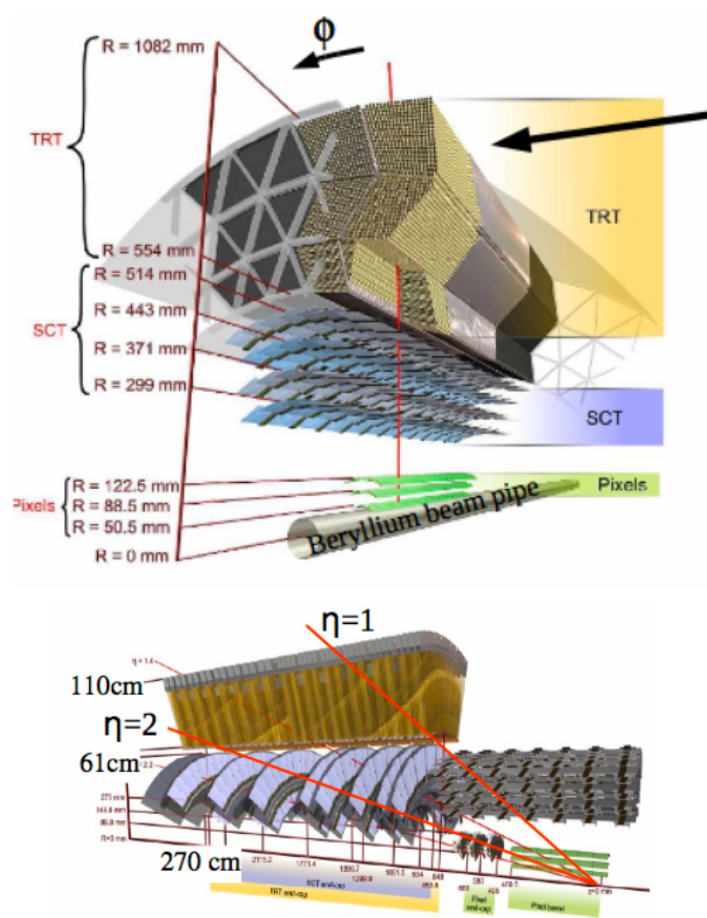


CMS

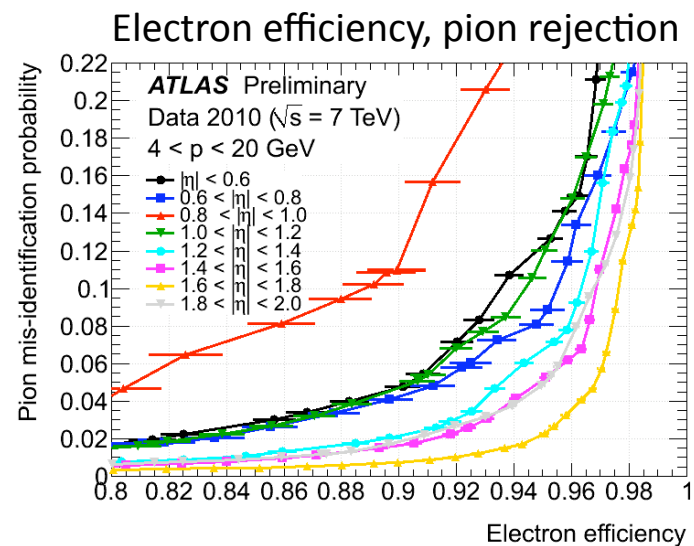
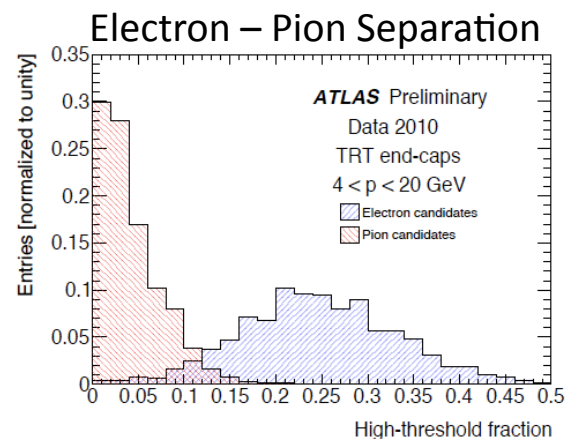


ATLAS TRT: e/π separation

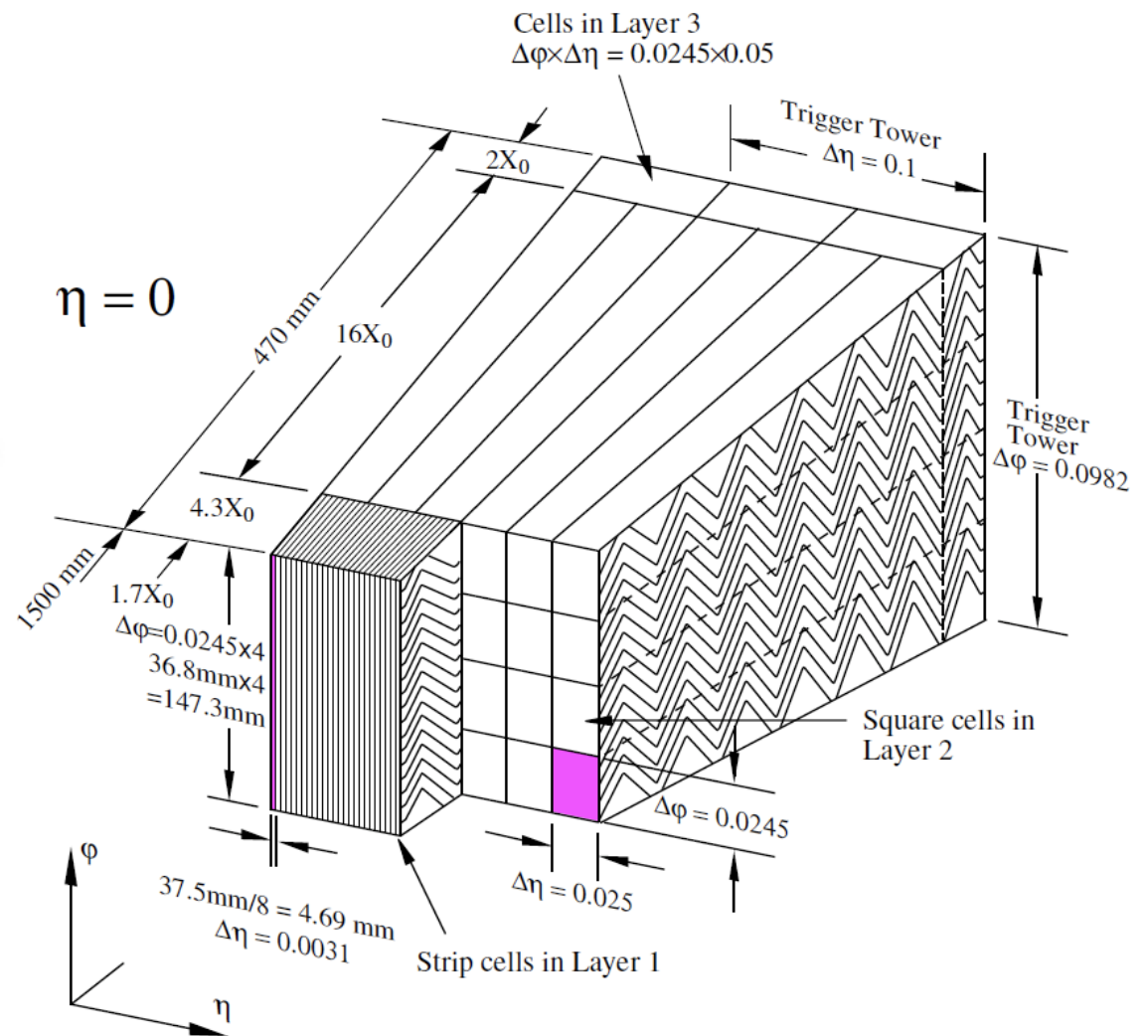
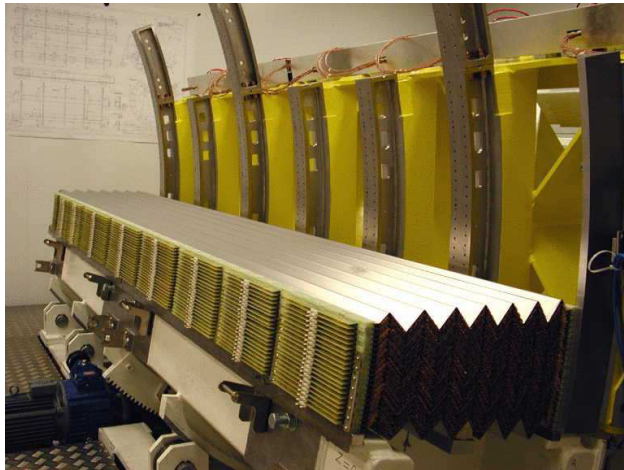
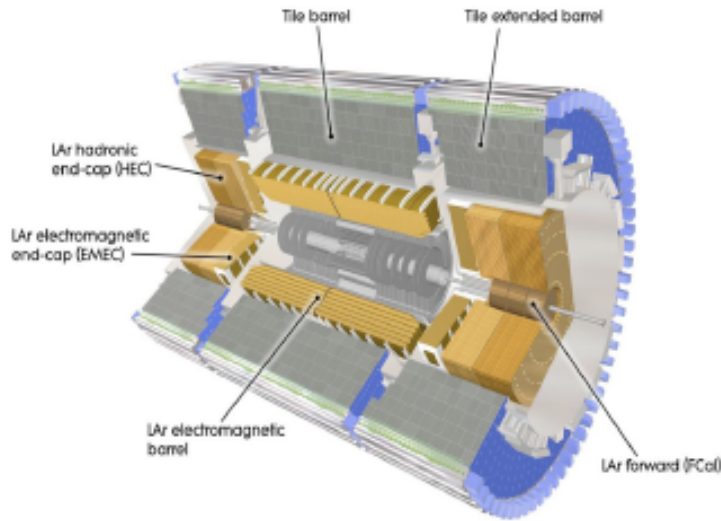
Particle Identification Performance of the ATLAS TRT Tracker: ATLAS-CONF-2011-128



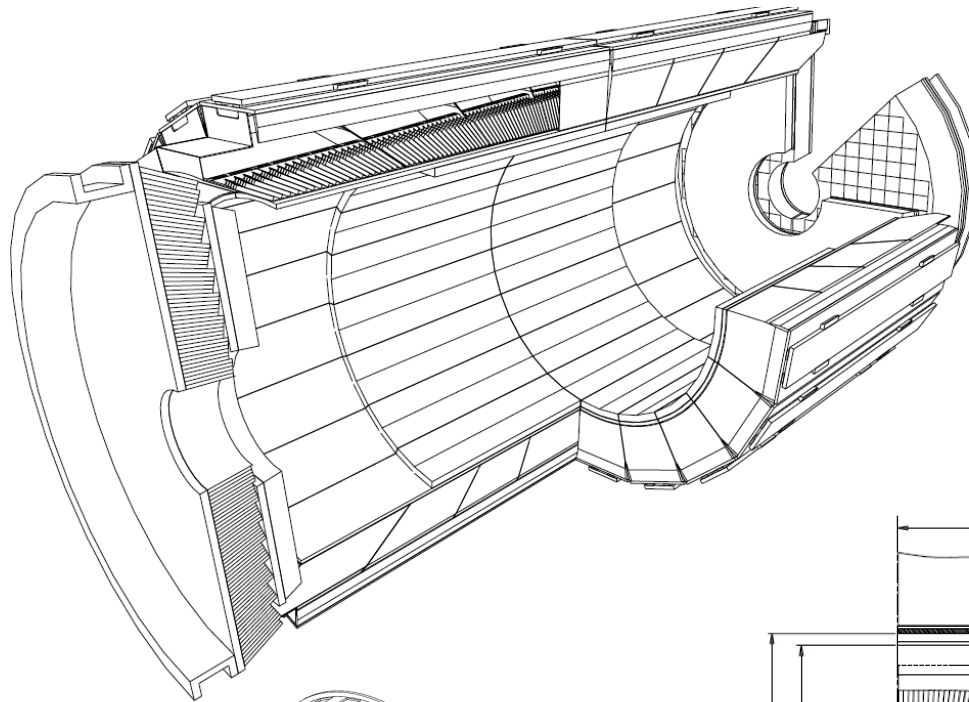
300,000 straw tubes
4 mm diameter
Radiator material
between straws



ATLAS EM Calorimeter

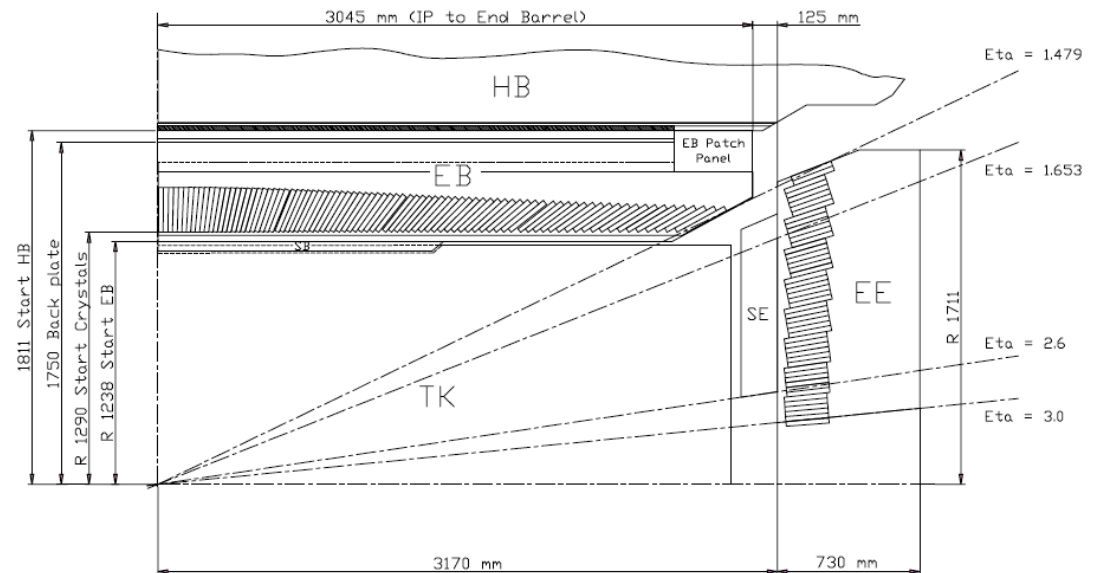
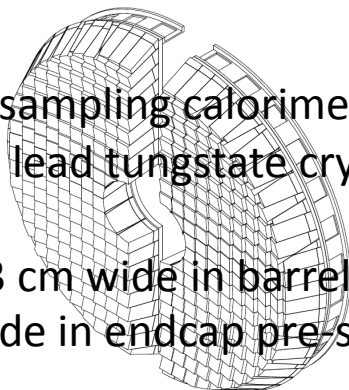


CMS EM Calorimeter



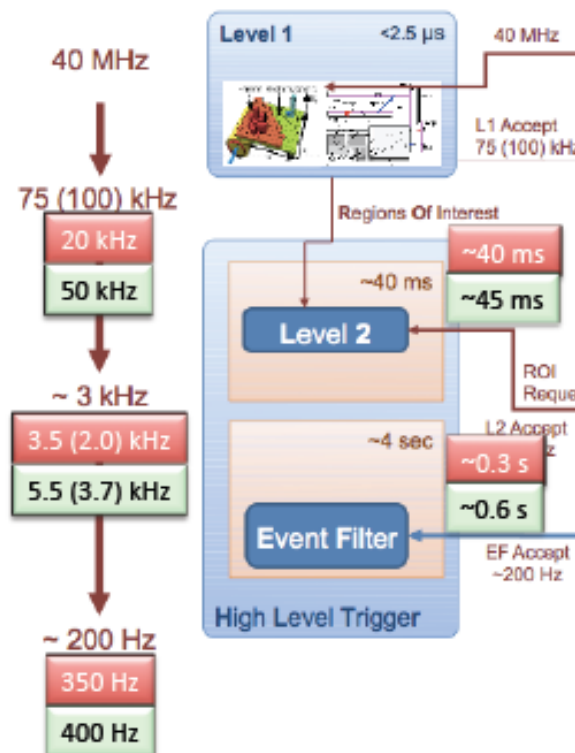
Non-sampling calorimeter
78,000 lead tungstate crystals

~3 cm wide in barrel
2 mm wide in endcap pre-shower

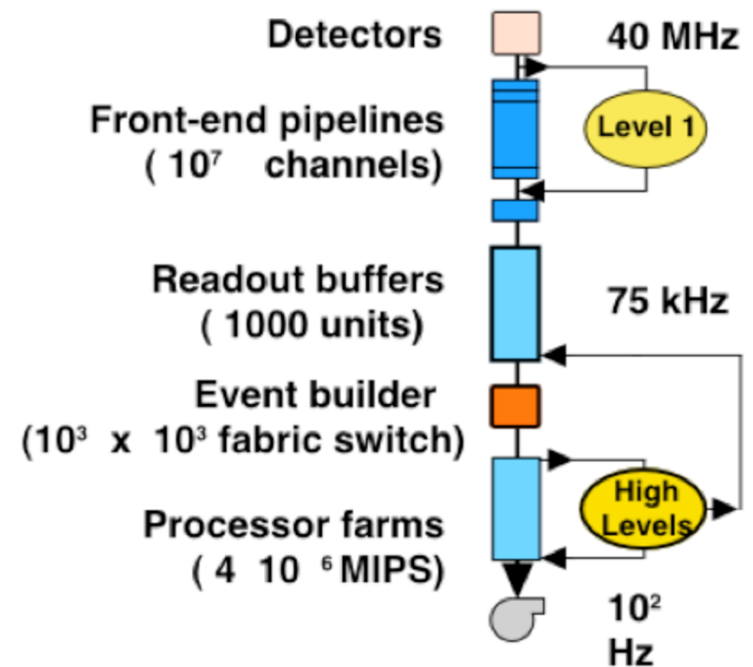


Triggers

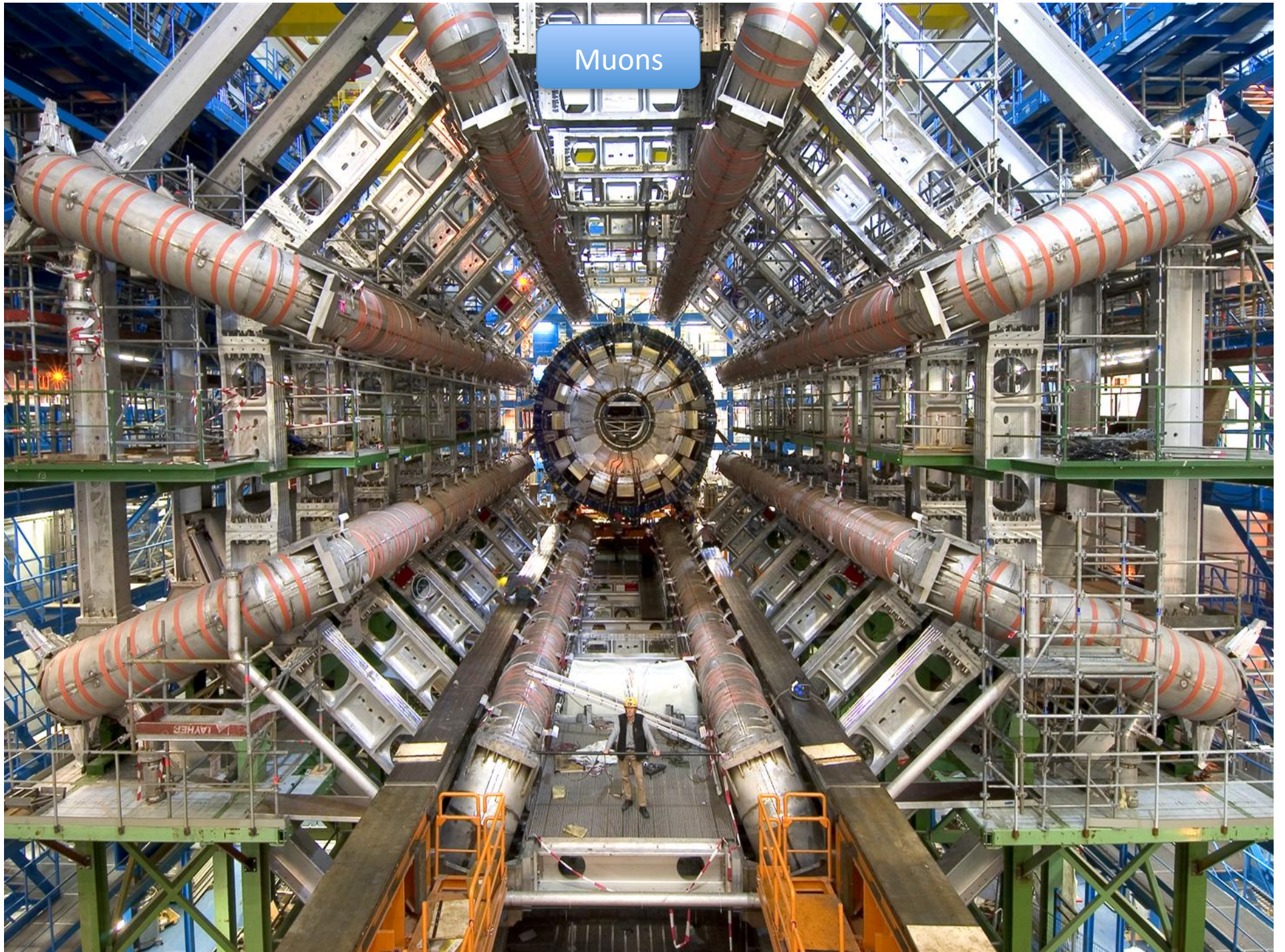
ATLAS: 3 Levels



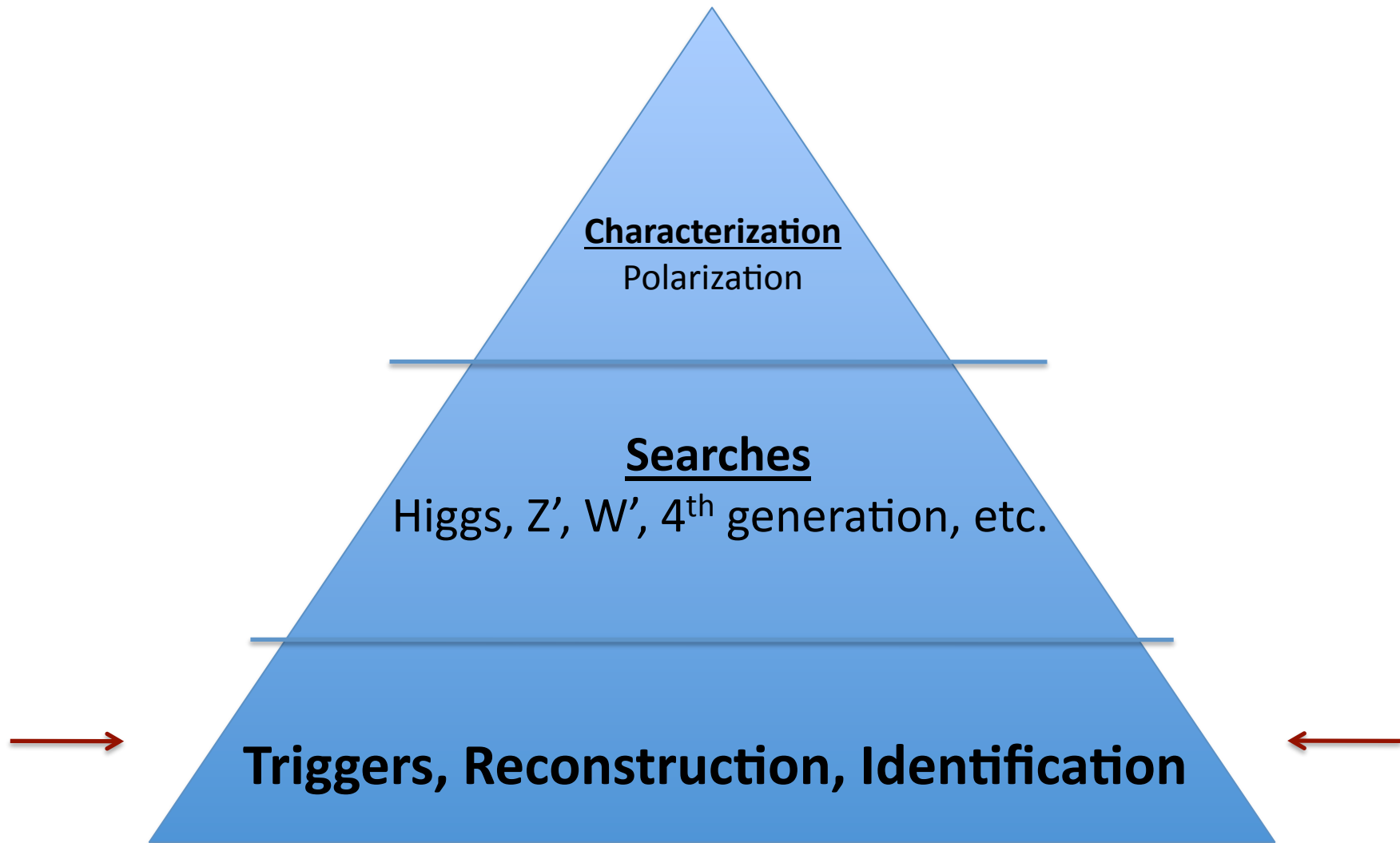
CMS: 2 Levels



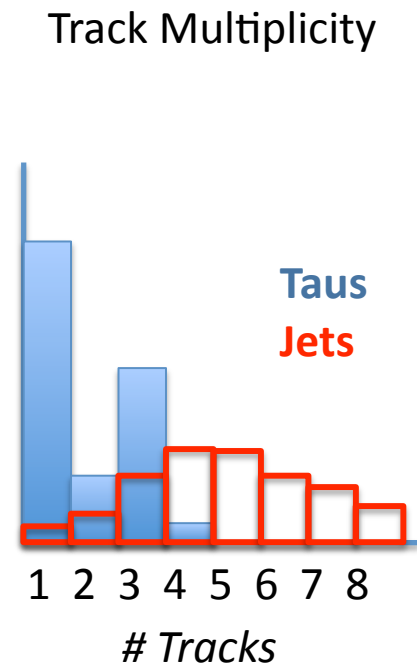
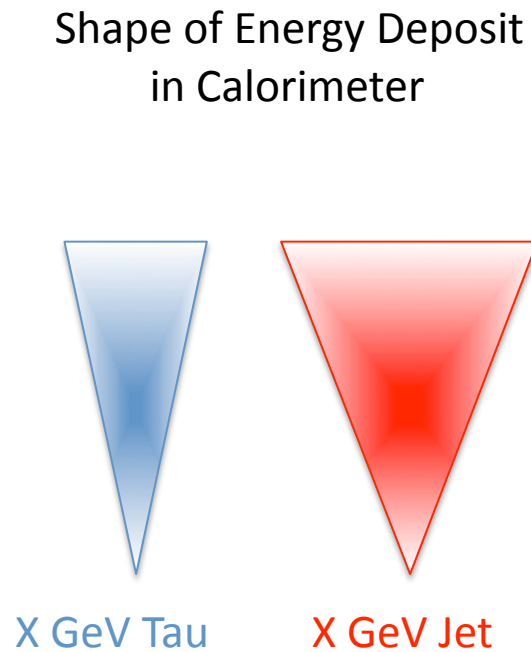
Muons



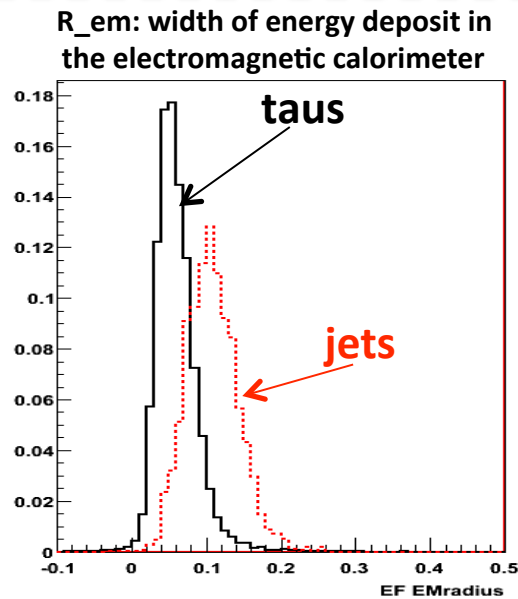
Taus at the LHC



Triggering, Reconstructing, Identifying

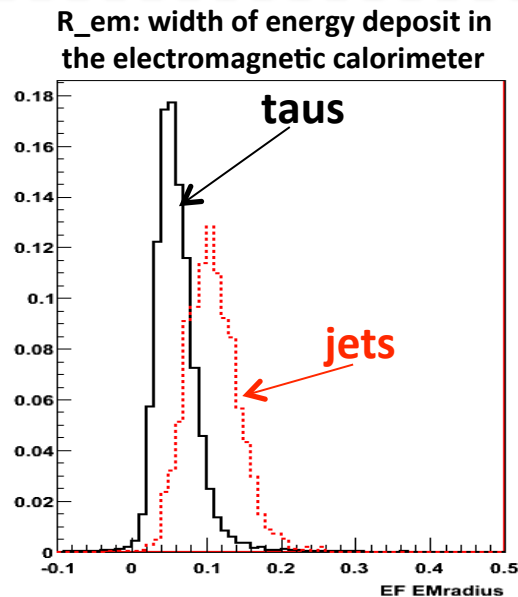


Triggering, Reconstructing, Identifying

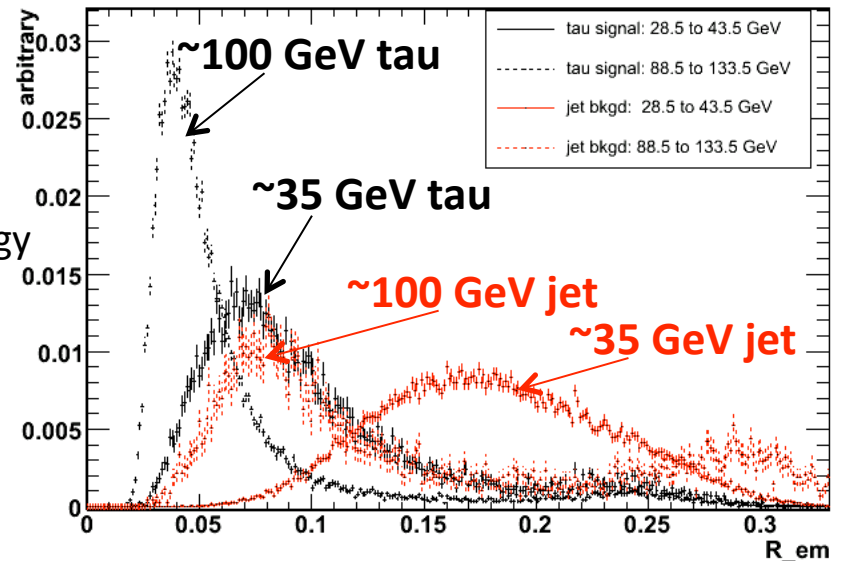


$$R_{em} = \frac{\sum_{i=1}^n E_{Ti} \sqrt{(\eta_i - \eta_{cluster})^2 + (\phi_i - \phi_{cluster})^2}}{\sum_{i=1}^n E_{Ti}}$$

Triggering, Reconstructing, Identifying



Separate in bins of energy



$$R_{em} = \frac{\sum_{i=1}^n E_{Ti} \sqrt{(\eta_i - \eta_{cluster})^2 + (\phi_i - \phi_{cluster})^2}}{\sum_{i=1}^n E_{Ti}}$$

also
isolation in tracking chamber

Tau Triggers

Level 1

Hardware

No Inner tracker information

Rely on calorimeters with low granularity

High Level Trigger

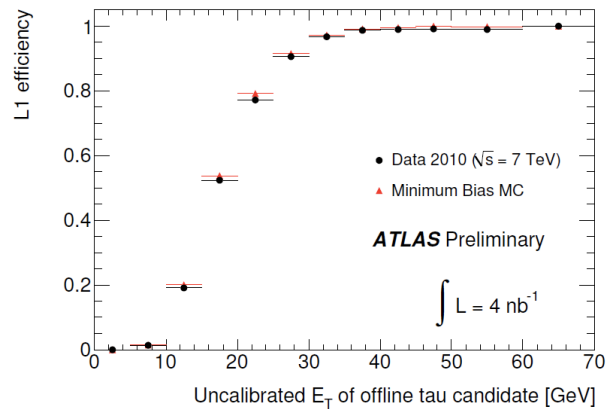
Software

Use “offline” algorithms wherever possible

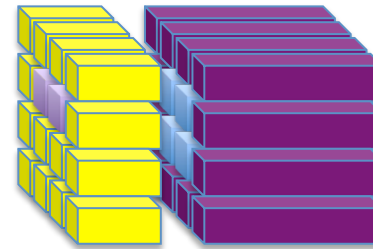


ATLAS Level 1 Tau Trigger

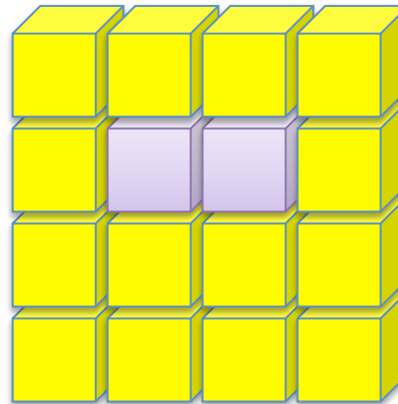
5 GeV L1 Tau Trigger



EM
Trigger Towers

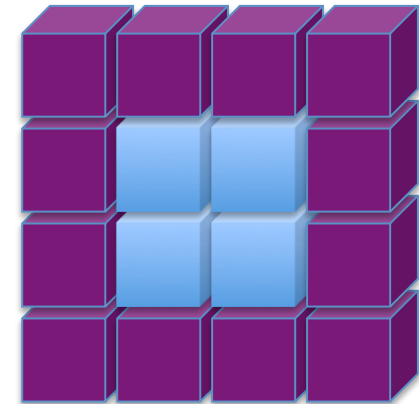


HAD
Trigger Towers



2 Adjacent EM
Trigger Towers

+

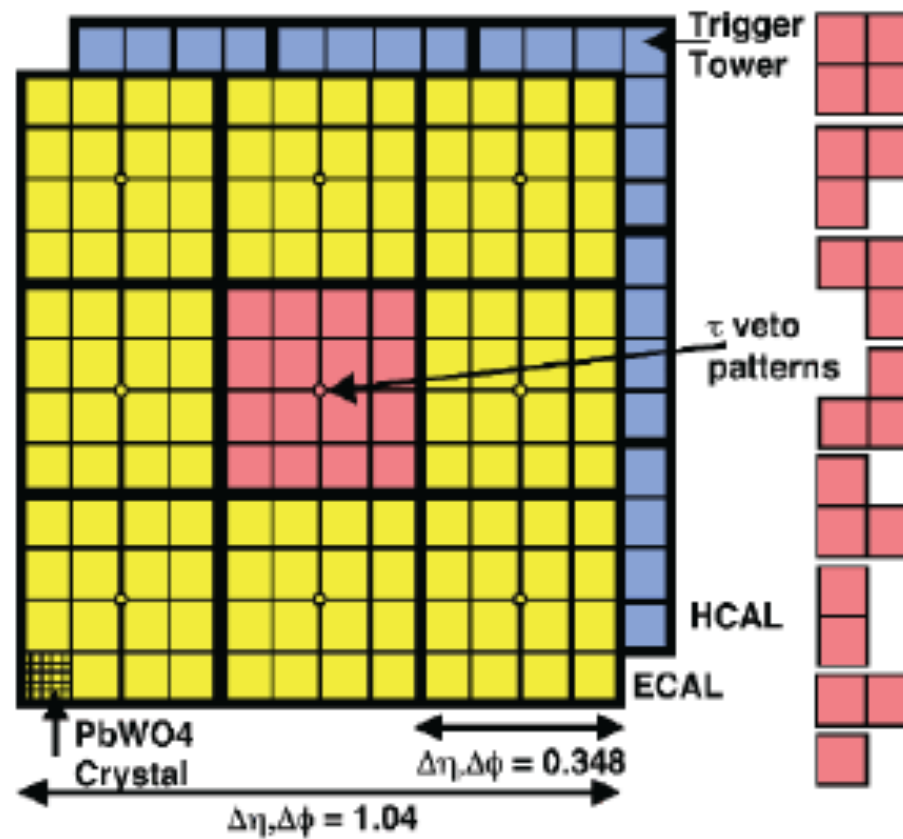


2x2 HAD Towers
behind them

Threshold
Calculation
=

$$\Delta\eta \times \Delta\varphi = 0.1 \times 0.1$$

CMS Level 1 Tau Trigger



Reconstruction and Identification

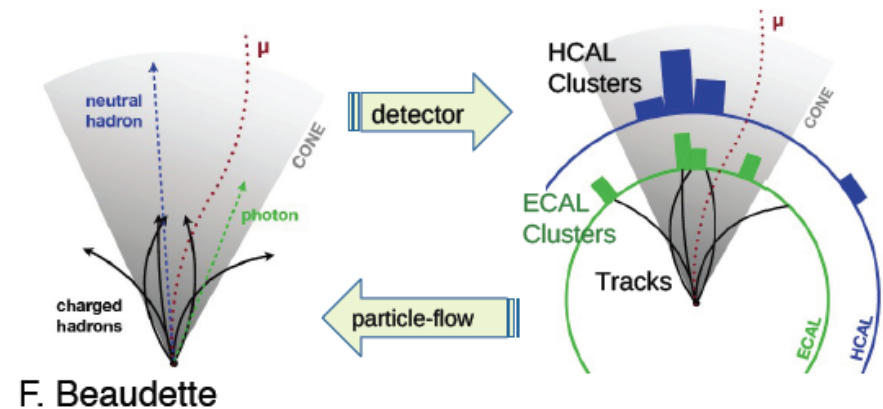
Major Difference: CMS uses particle flow

Reconstruction and Identification

Major Difference: CMS uses particle flow

Particle Flow

Clusters and links signals from sub-detectors
Produces list of particle candidates
To user: looks like Monte Carlo

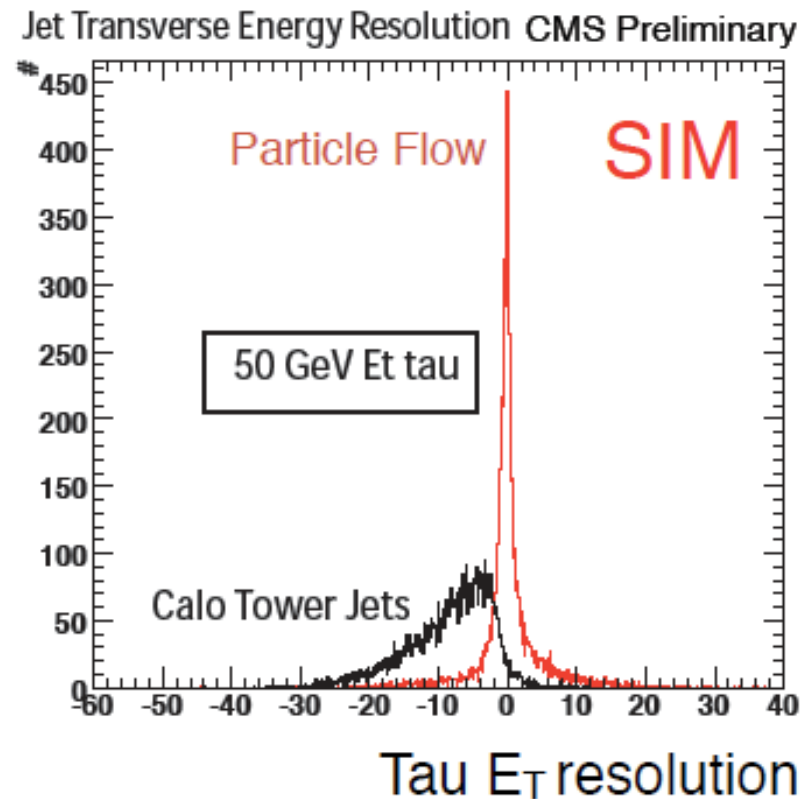


Reconstruction and Identification

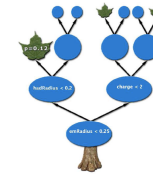
Major Difference: CMS uses particle flow

Particle Flow

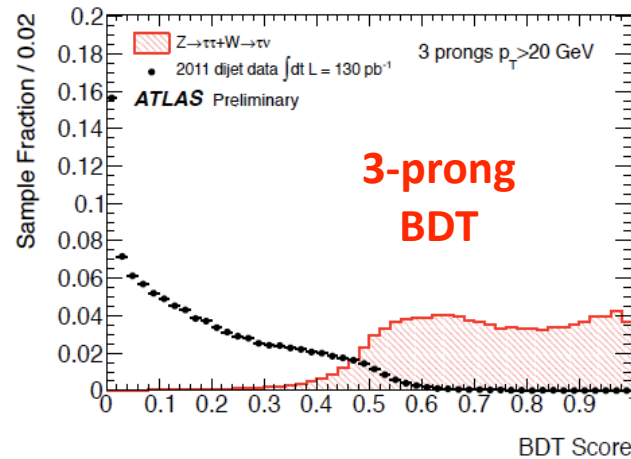
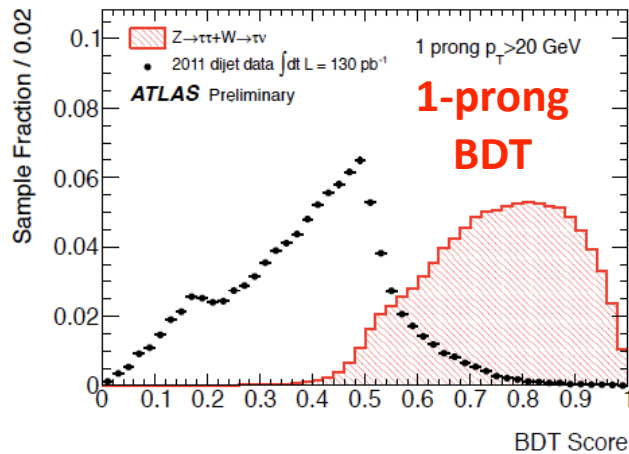
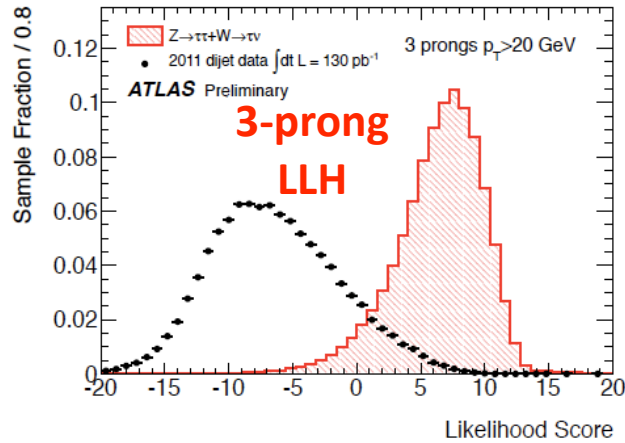
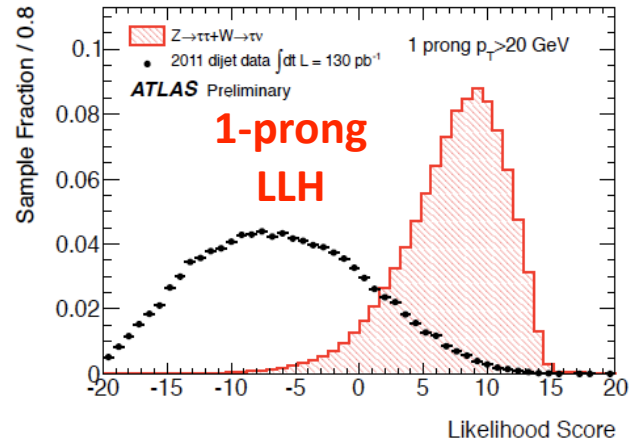
Clusters and links signals from sub-detectors
Produces list of particle candidates
To user: looks like Monte Carlo



ATLAS Tau Identification



Cuts, log-likelihood,
& boosted decision trees



Variable	Eqn.	Jet discriminants					
		Cut		LLH		BDT	
		l	m	l	m	l	m
R_{track}	11	•	•	•	•	•	•
f_{track}	12	•	•			•	•
f_{core}	13			•	•	•	•
$N_{\text{track}}^{\text{iso}}$		•	•	•		•	•
R_{Cal}	14			•		•	•
f_{iso}	15						
$m_{\text{eff. clusters}}$	16					•	•
m_{tracks}	18				•		•
$S_{\text{T}}^{\text{flight}}$	19		•		•		•
$S_{\text{lead track}}$	20					•	•
$f_2^{\text{lead clusters}}$				•			
$f_3^{\text{lead clusters}}$						•	•
ΔR_{max}					•		•
f_{EM}	21						
f_{HT}	22						
$f_{\text{Had}}^{\text{track}}$	23						
$E_{\text{T,max}}^{\text{strip}}$							
$f_{\text{EM}}^{\text{track}}$	24						
R_{Had}	25						
$E_{\text{T,corr}}^{\text{iso}}$	26	•	•				

CMS Tau Identification

Two Identification Methods

Hadrons Plus Strips (HPS)

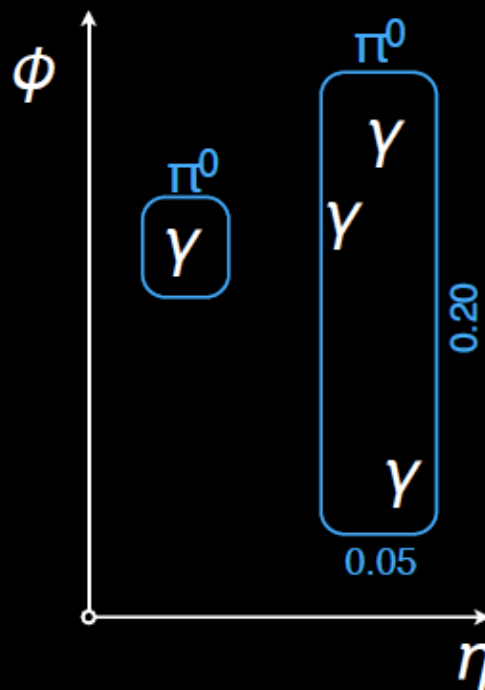
Tau Neural Classifier (TaNC)

Slide from Evan Friis

Hadrons Plus Strips Algorithm

build signal components combinatorially

cluster gammas into π^0
candidates using η - ϕ strips



build all possible taus
that have a 'tau-like' multiplicity
from the seed jet

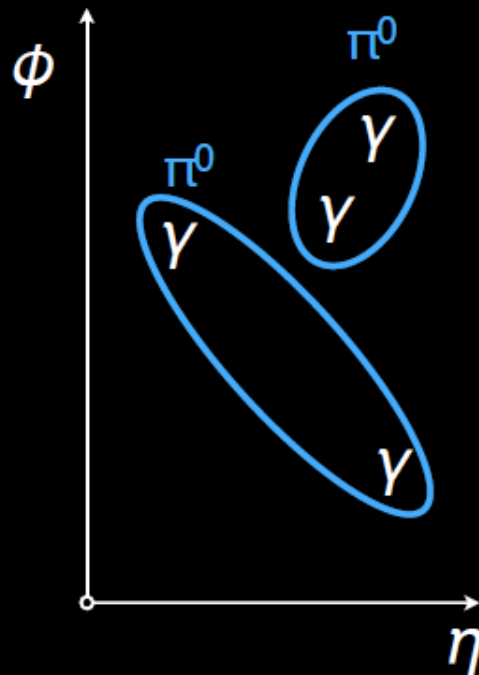
π^+
 $\pi^+ \pi^0$
 $\pi^+ \pi^+ \pi^-$

tau that is 'most isolated'
with compatible m_{vis}
is the final tau candidate
associated to the seed jet

Tau Neural Classifier

a neural network for each decay mode

cluster gammas into π^0
candidates by combinatoric
pairs compatible with m_{π^0}



signal objects are defined
using shrinking cone

depending on decay mode

π^+

$\pi^+ \pi^0$

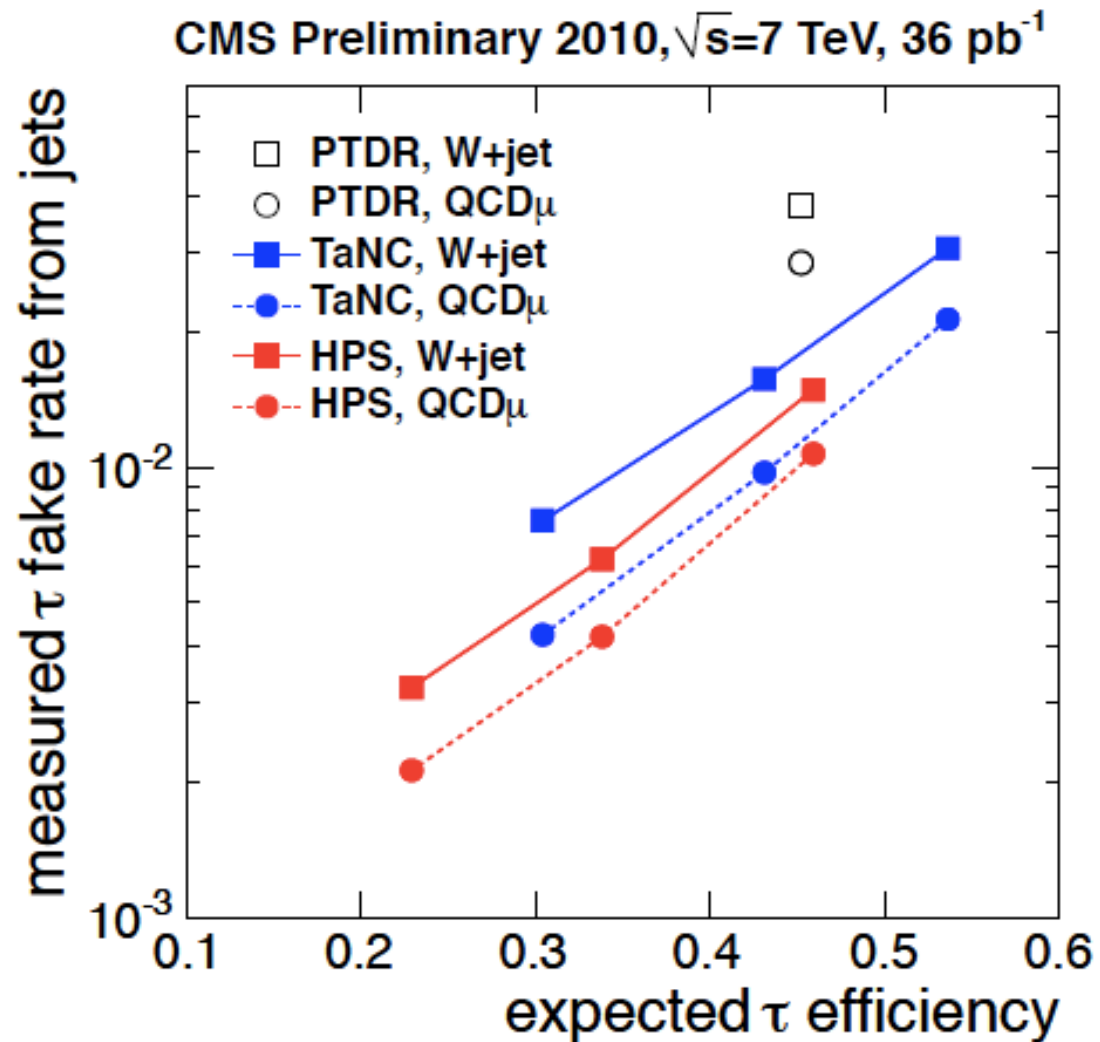
$\pi^+ \pi^0 \pi^0$

$\pi^+ \pi^+ \pi^-$

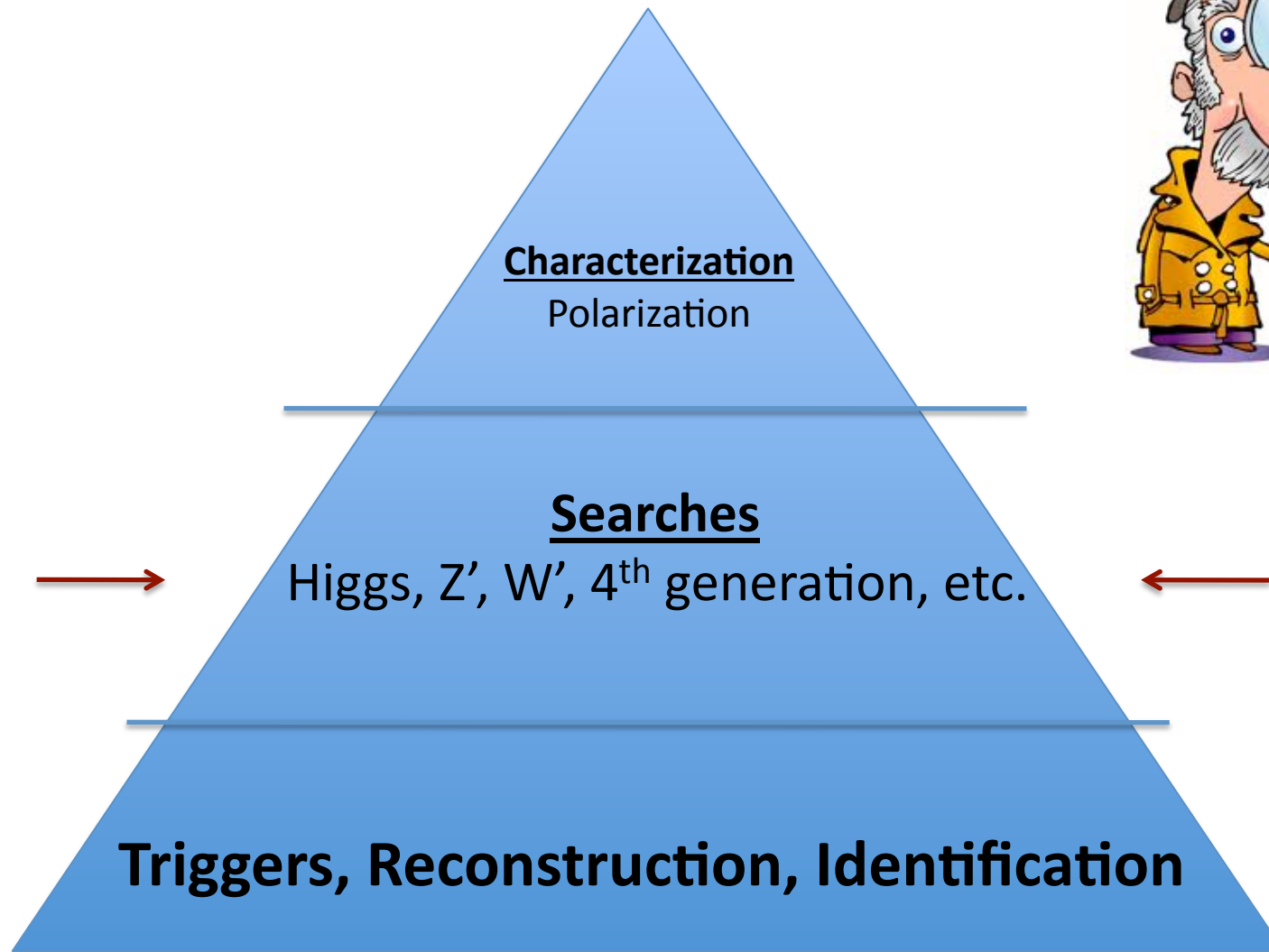
$\pi^+ \pi^+ \pi^- \pi^0$

a different neural network
is applied!

Performance (2010)

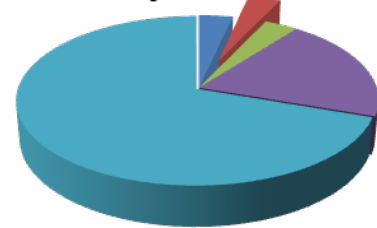


Taus at the LHC



Standard Model Measurements

Z decays



Interest

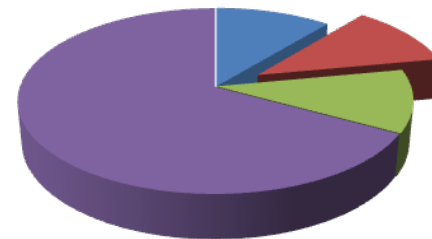
Background in Higgs Search

Challenges

Missing Energy (missing info for reconstructing mass)

Trigger

W decays



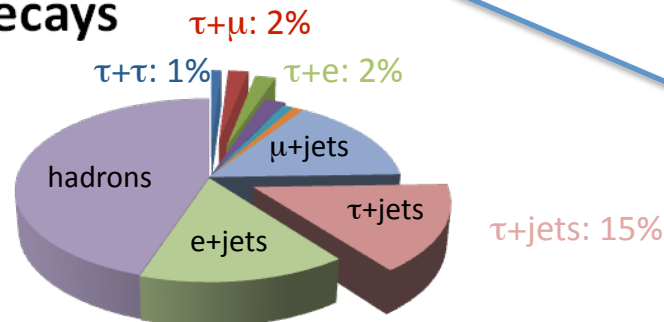
Interest

Early Tau Observation

Challenges

Trigger

$t\bar{t}$ decays



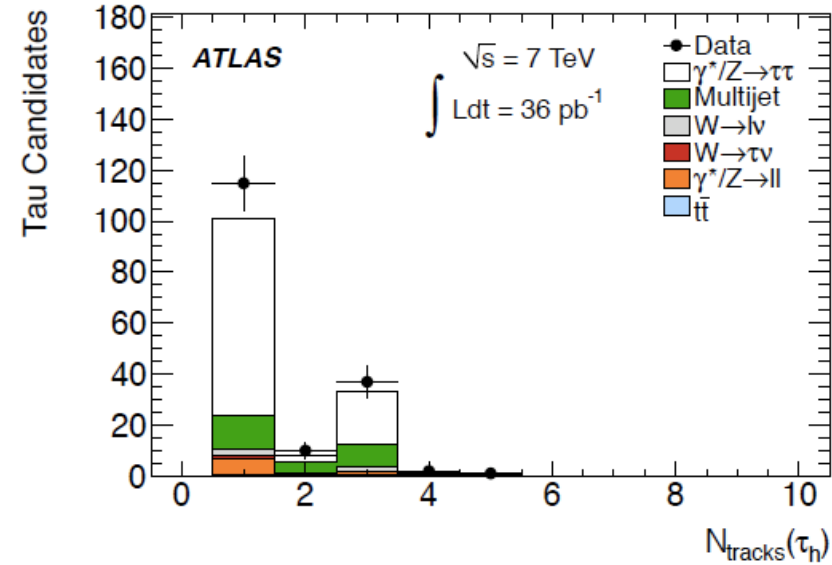
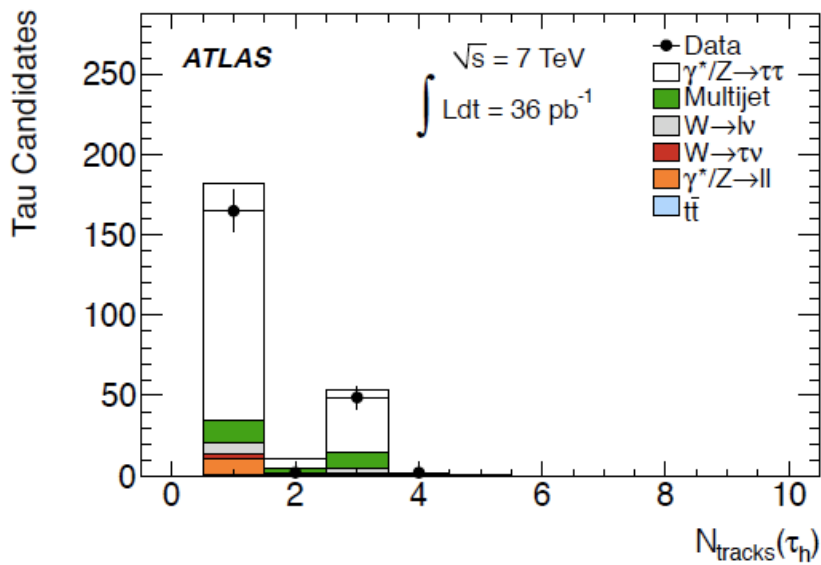
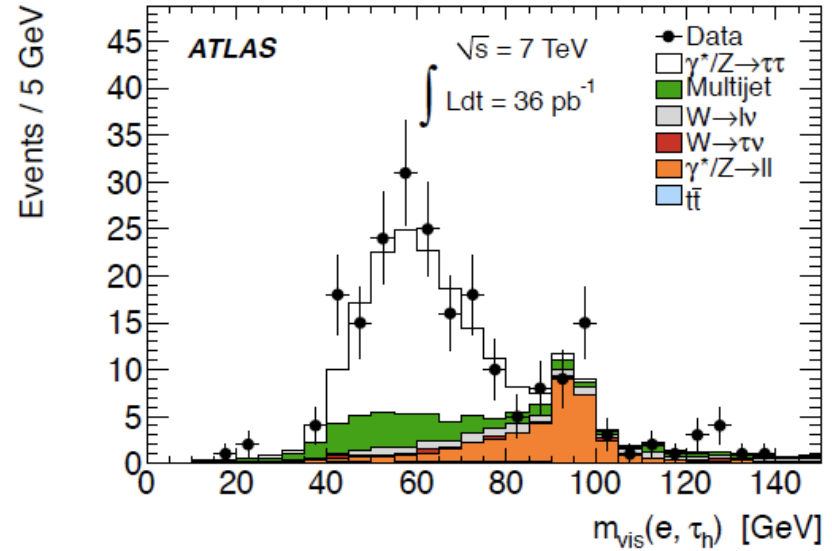
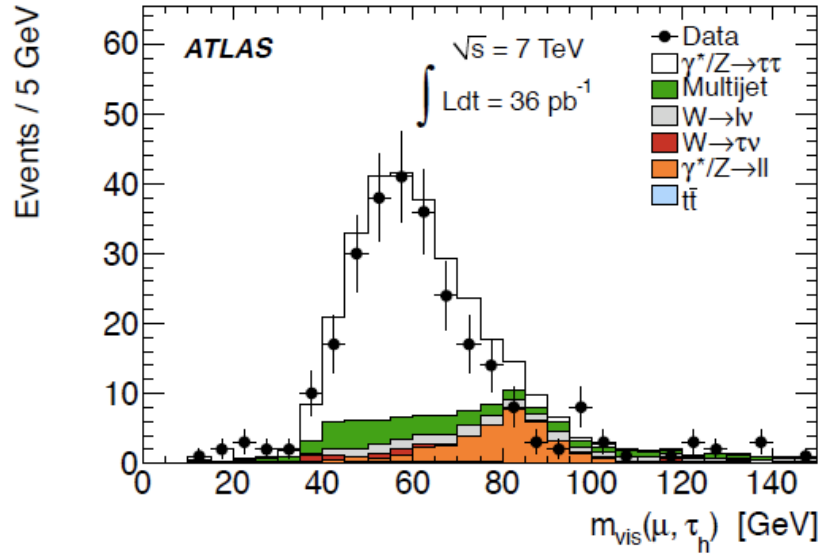
Interest

Charged Higgs Search
Calibration Sample

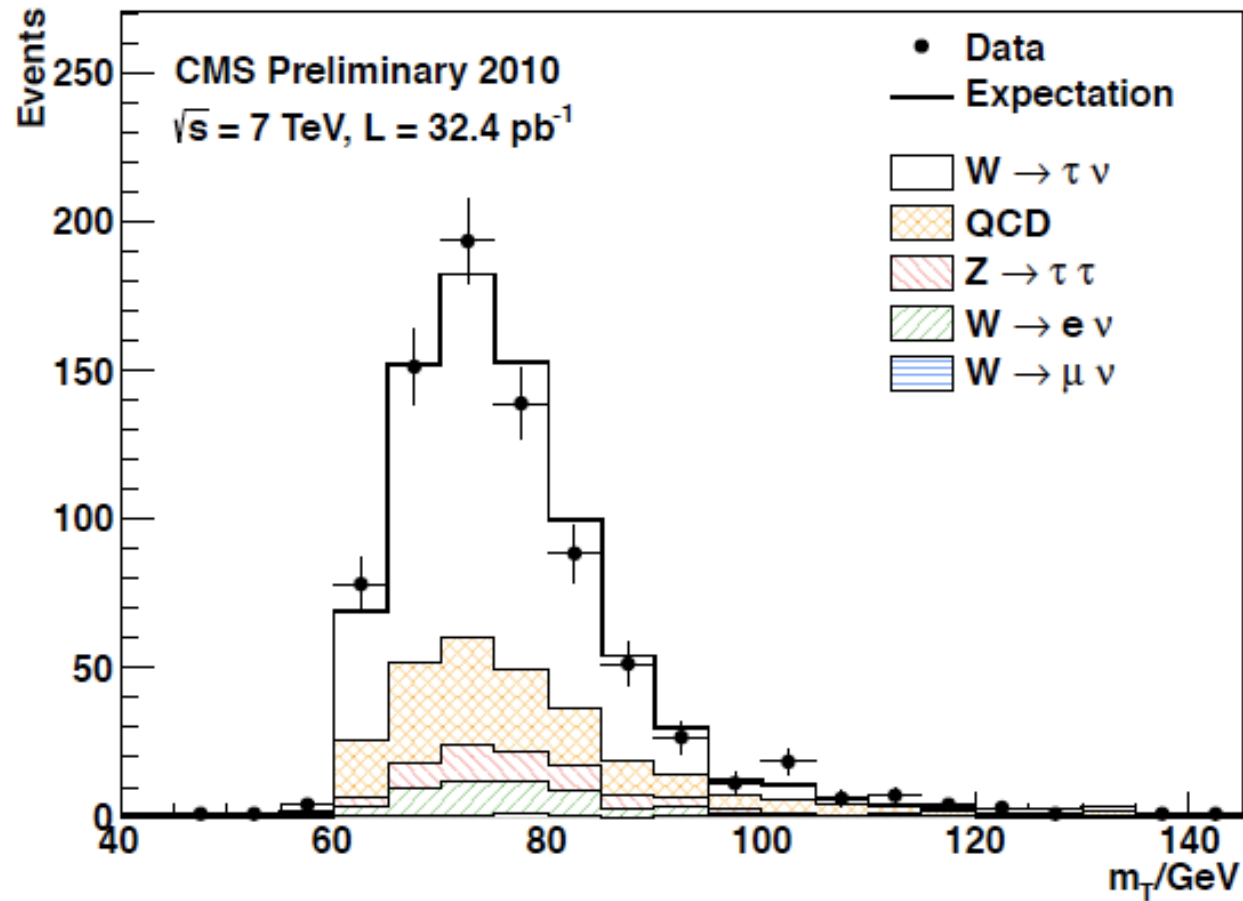
Challenges

Signal/Background

SM Measurements: $Z \rightarrow \tau_l \tau_h$

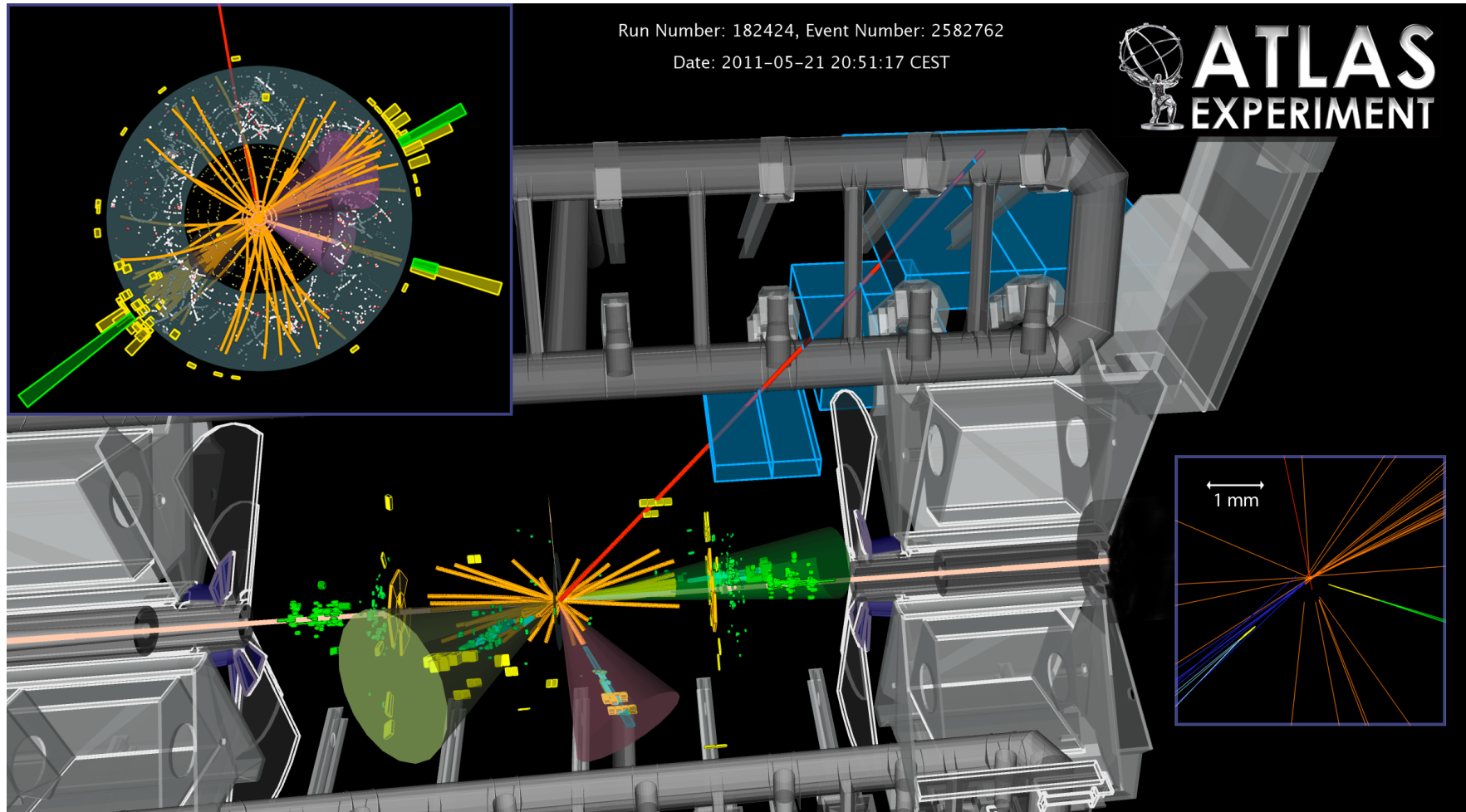


SM Measurements: $W \rightarrow \tau_h \nu$

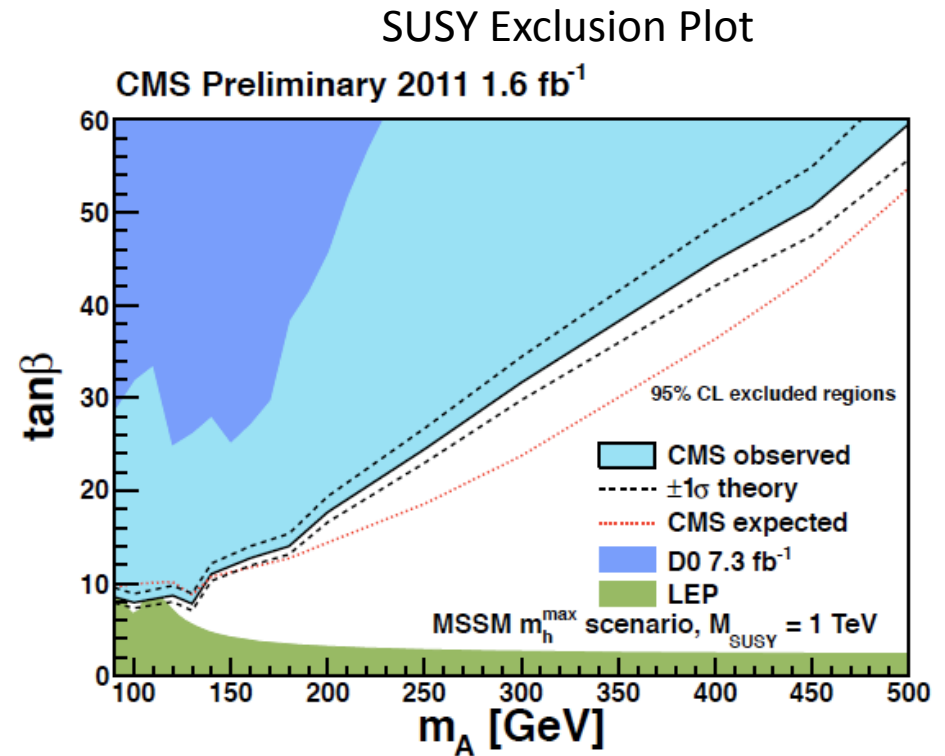
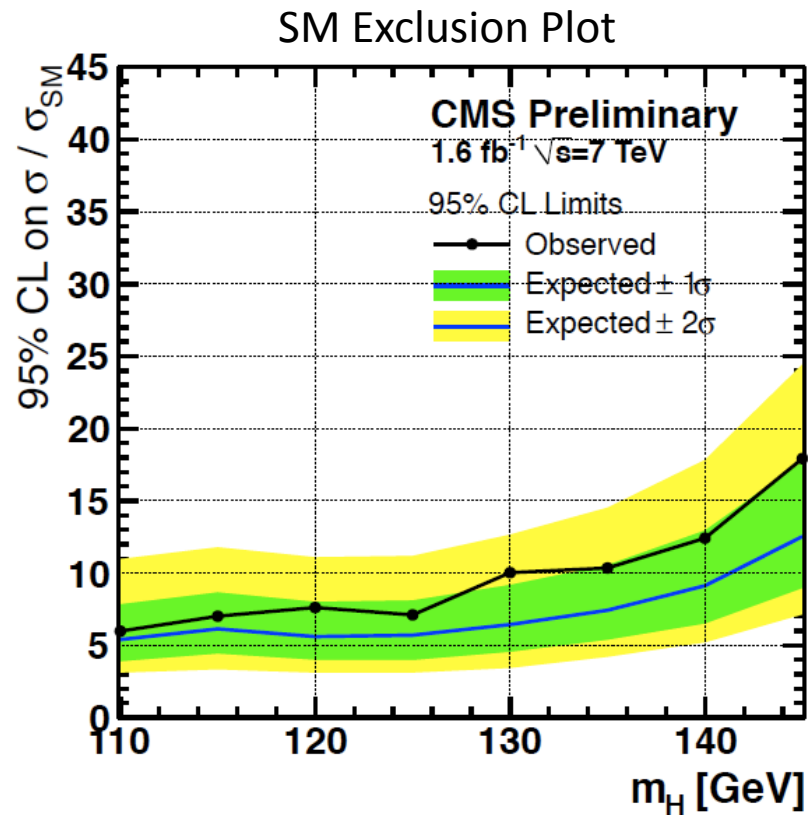


Top events with taus

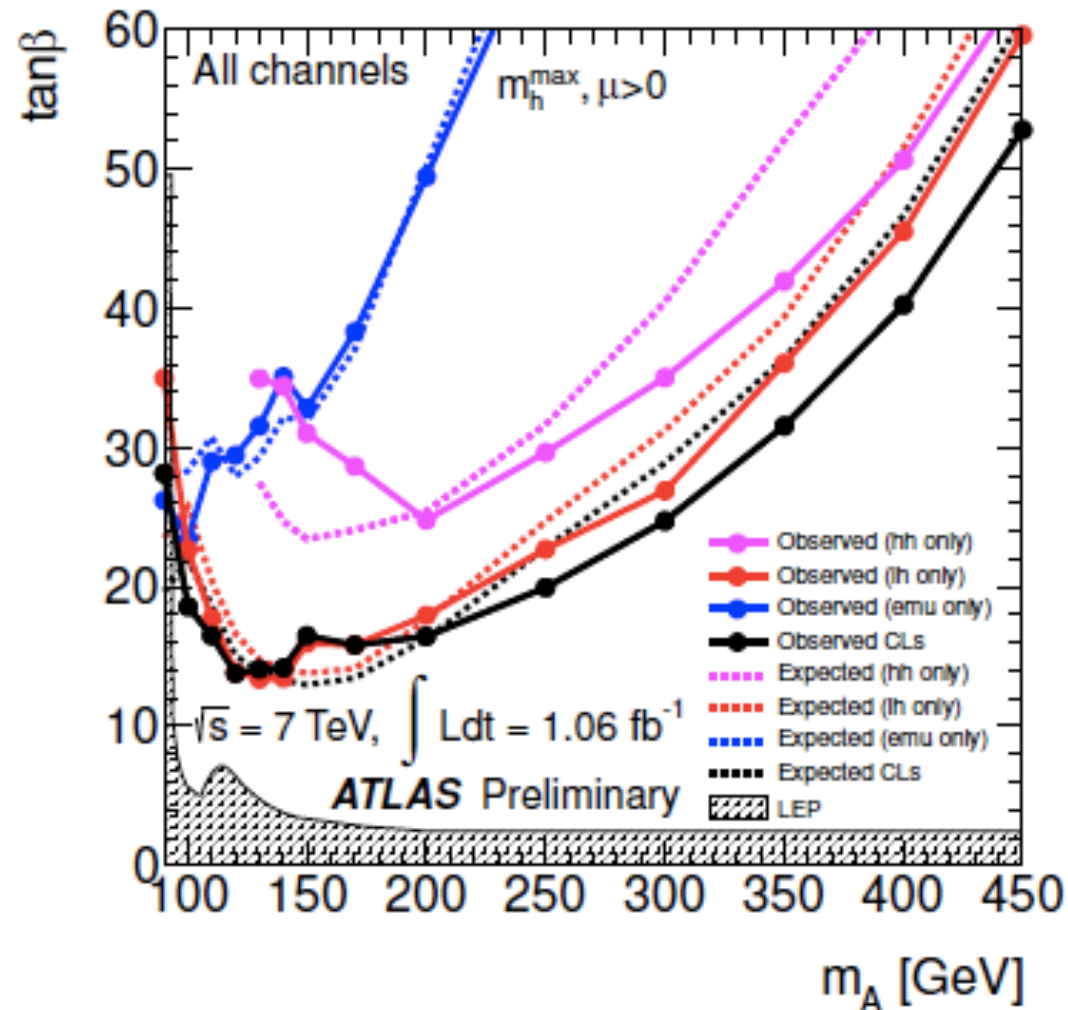
tau-muon dilepton candidate



CMS: Higgs $\rightarrow \tau\tau$



ATLAS MSSM $H \rightarrow \tau\tau$



Search for the Charged Higgs

If

- *the charged Higgs exists and has a mass less than the mass of the top quark*

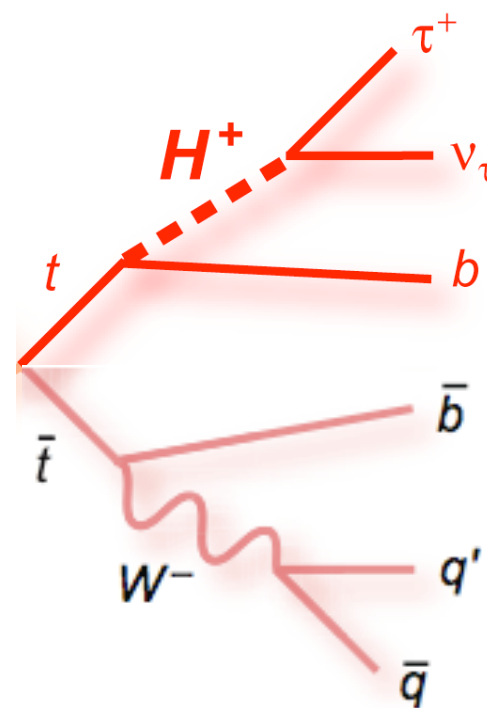
Then

- *the primary production mechanism for the charged Higgs will be via top quarks*

Since

- *the primary source of top quarks is within $t\bar{t}$ production*

Top quark pairs provide a nice playground for searching for and constraining the allowed masses of the charged Higgs



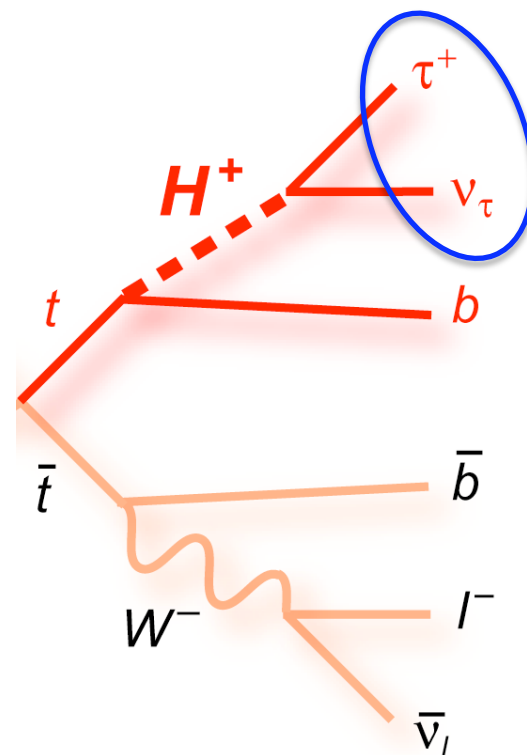
Search for the Charged Higgs

When $\tan\beta$ is greater than ~ 3 , the H^+ decays almost exclusively to $\tau\nu$

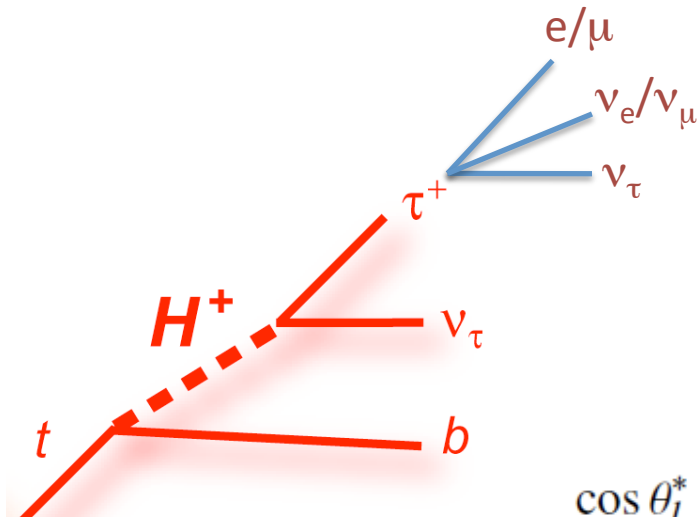
This search was performed with 1.03 fb^{-1} of 2011 ATLAS data in both the **single lepton** and **dilepton** channels

Presence of H^+ ends up giving excess of leptons:
 $B(H^+ \rightarrow \tau\nu \rightarrow l + N\nu) \approx 35\%$ while $B(W \rightarrow l + N\nu) \approx 25\%$

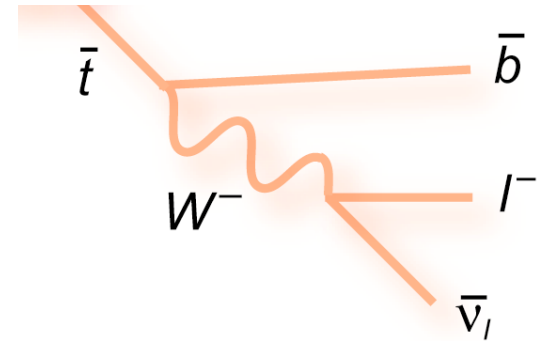
but additional handles are needed:
 $\cos\theta_l^*$ and transverse mass of the H^+



Discriminating variables: $\cos\theta_l^*$

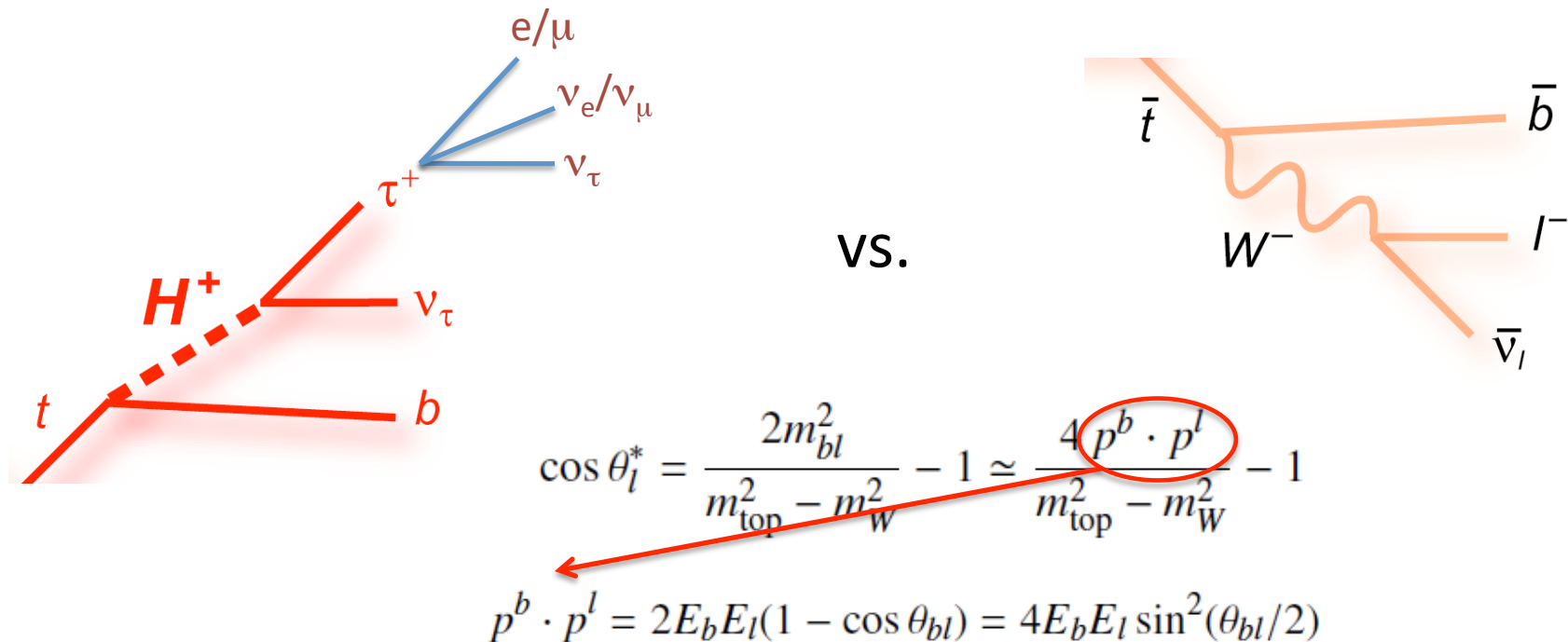


vs.



$$\cos\theta_l^* = \frac{2m_{bl}^2}{m_{\text{top}}^2 - m_W^2} - 1 \simeq \frac{4 p^b \cdot p^l}{m_{\text{top}}^2 - m_W^2} - 1$$

Discriminating variables: $\cos\theta_l^*$

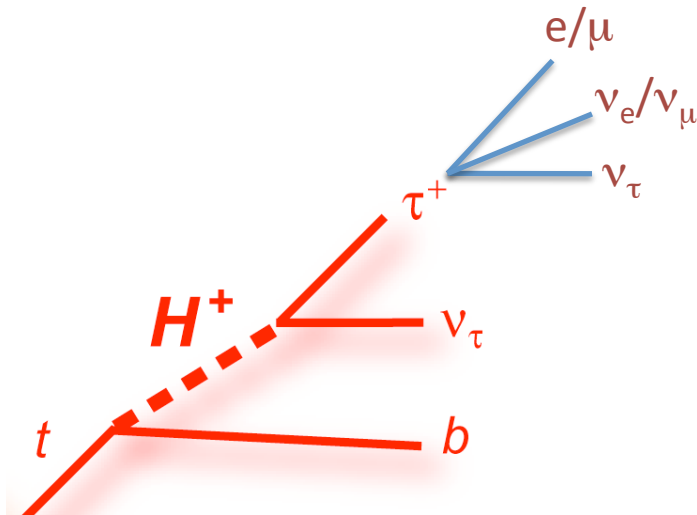


Variable commonly used to measure polarization of W in top quark events

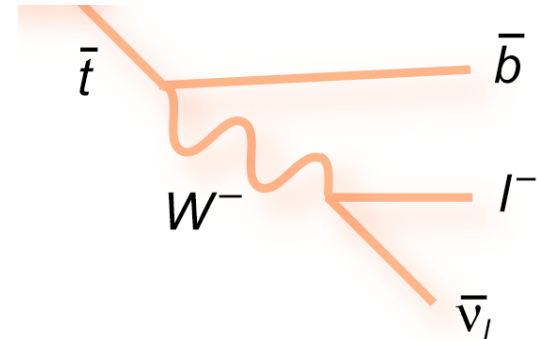
Here, with H^+ heavier than the W, the b-quark tends to have less momentum in H^+ events
also, leptons from the tau decays tend to have less momentum than from Ws

All of this works to push $\cos\theta_l^*$ toward the value of -1

Discriminating variables: transverse mass of Higgs



VS.



Define H^+ transverse mass

For single lepton channel, maximize the lepton + p_T^{miss} :

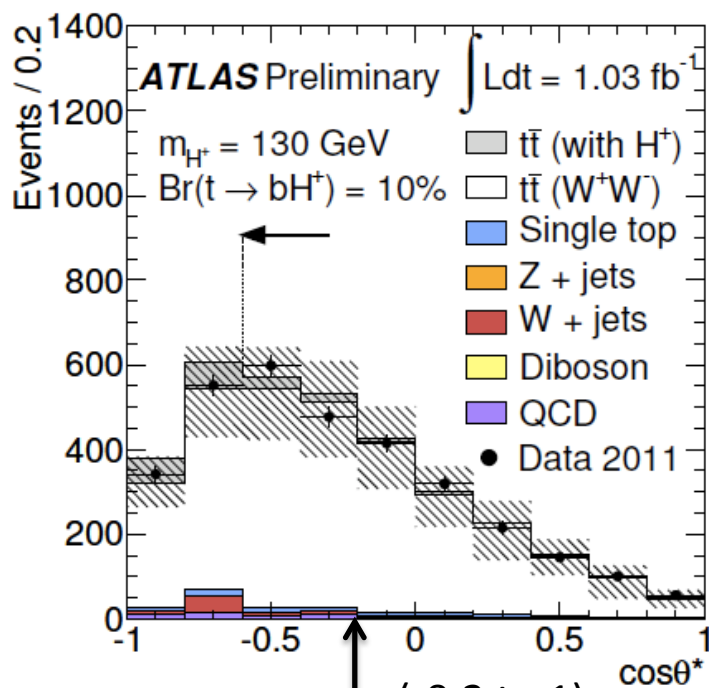
$$(m_T^H)^2 = \left(\sqrt{m_{\text{top}}^2 + (\vec{p}_T^l + \vec{p}_T^b + \vec{p}_T^{\text{miss}})^2} - p_T^b \right)^2 - (\vec{p}_T^l + \vec{p}_T^{\text{miss}})^2$$

$$(\vec{p}^{\text{miss}} + \vec{p}^l + \vec{p}^b)^2 = m_{\text{top}}^2$$

Situation is slightly more complicated in dilepton case

Charged Higgs: Single Lepton Results

Signal region, with additional requirement
that “transverse mass of W” < 60 GeV
to enhance charged boson
decays through $\tau \rightarrow \nu_l \nu_\tau$



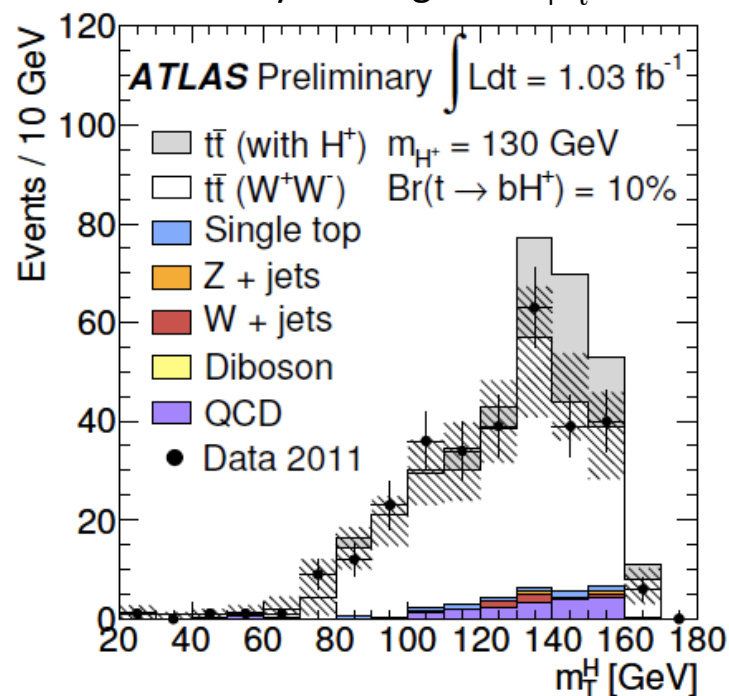
Signal region

(-0.2 to 1)

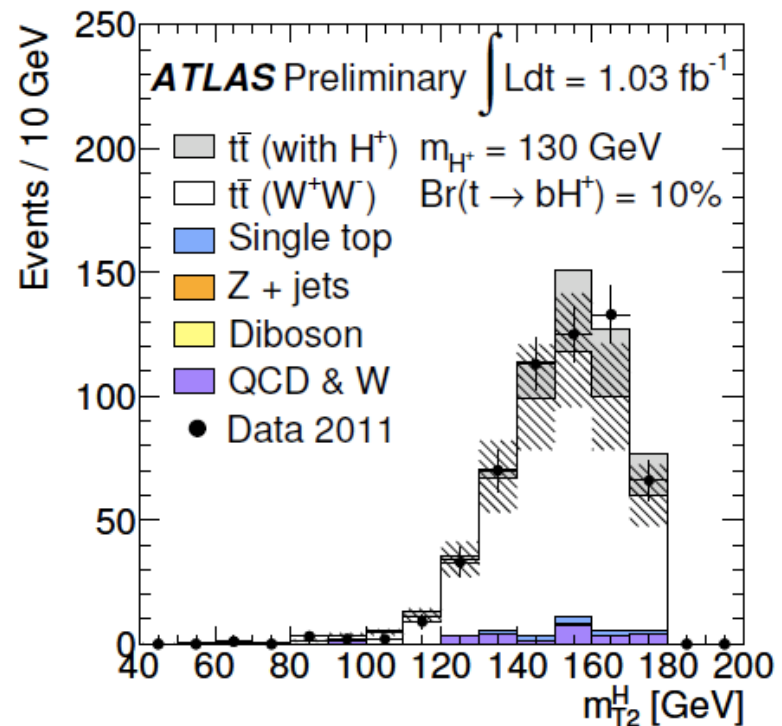
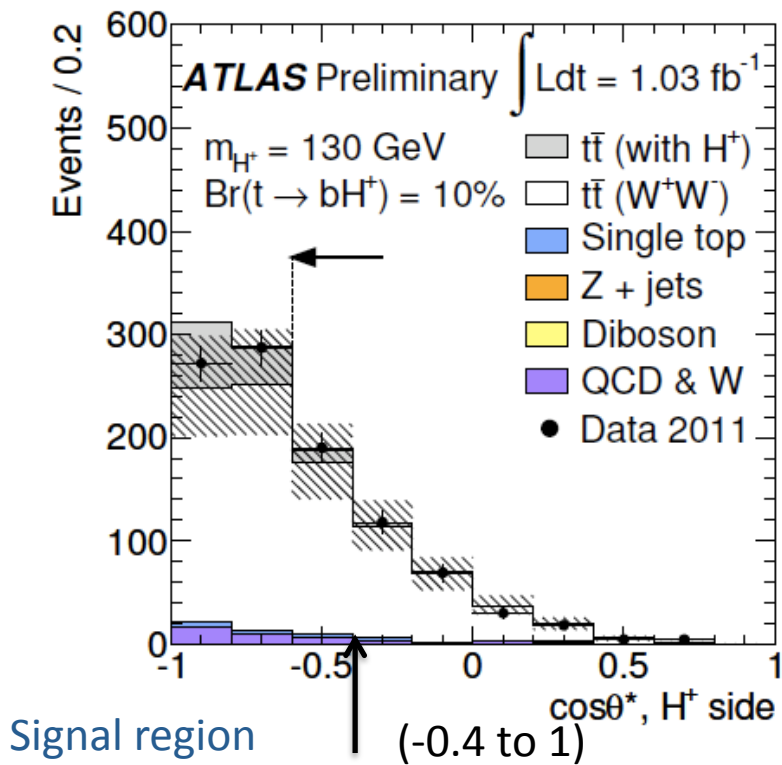
SM enriched region

for fitting to cross section: $t\bar{t} \rightarrow b\bar{b}W^+W^-$

165.1 pb



Charged Higgs: Dilepton Results

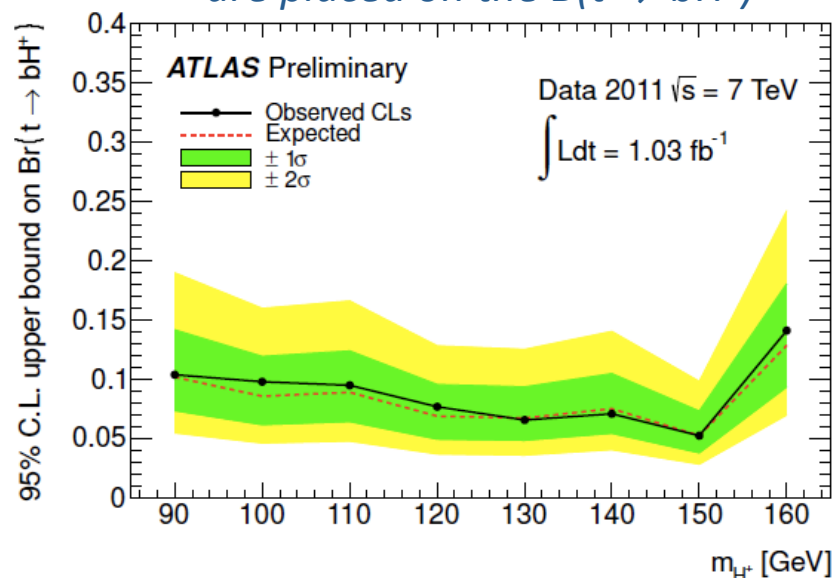


SM enriched region
 for fitting to cross section: $t\bar{t} \rightarrow b\bar{b}W^+W^-$
 150.4 pb

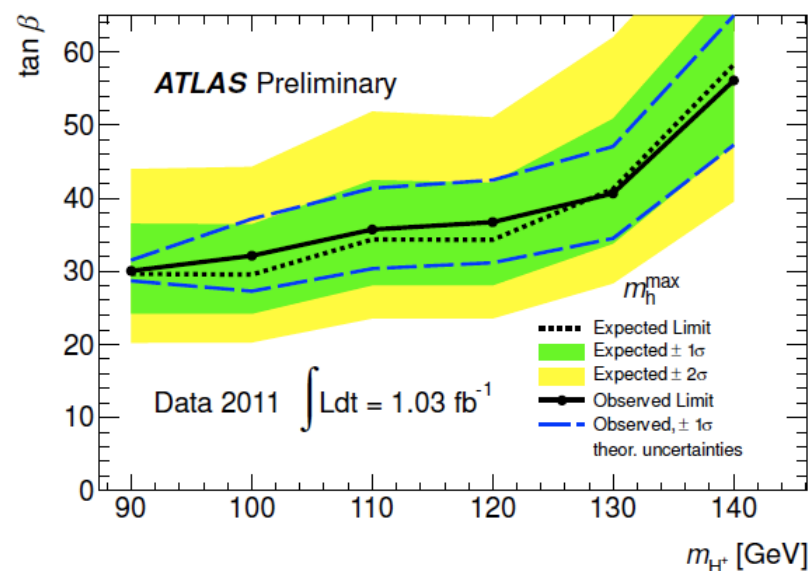
Combined Limits

Assuming $B(H^+ \rightarrow \tau\nu) = 1$,
place upper limits on the branching fraction $B(t \rightarrow bH^+)$

*In the 90 – 160 GeV range,
upper limits between 5.2% and 14.1%
are placed on the $B(t \rightarrow bH^+)$*



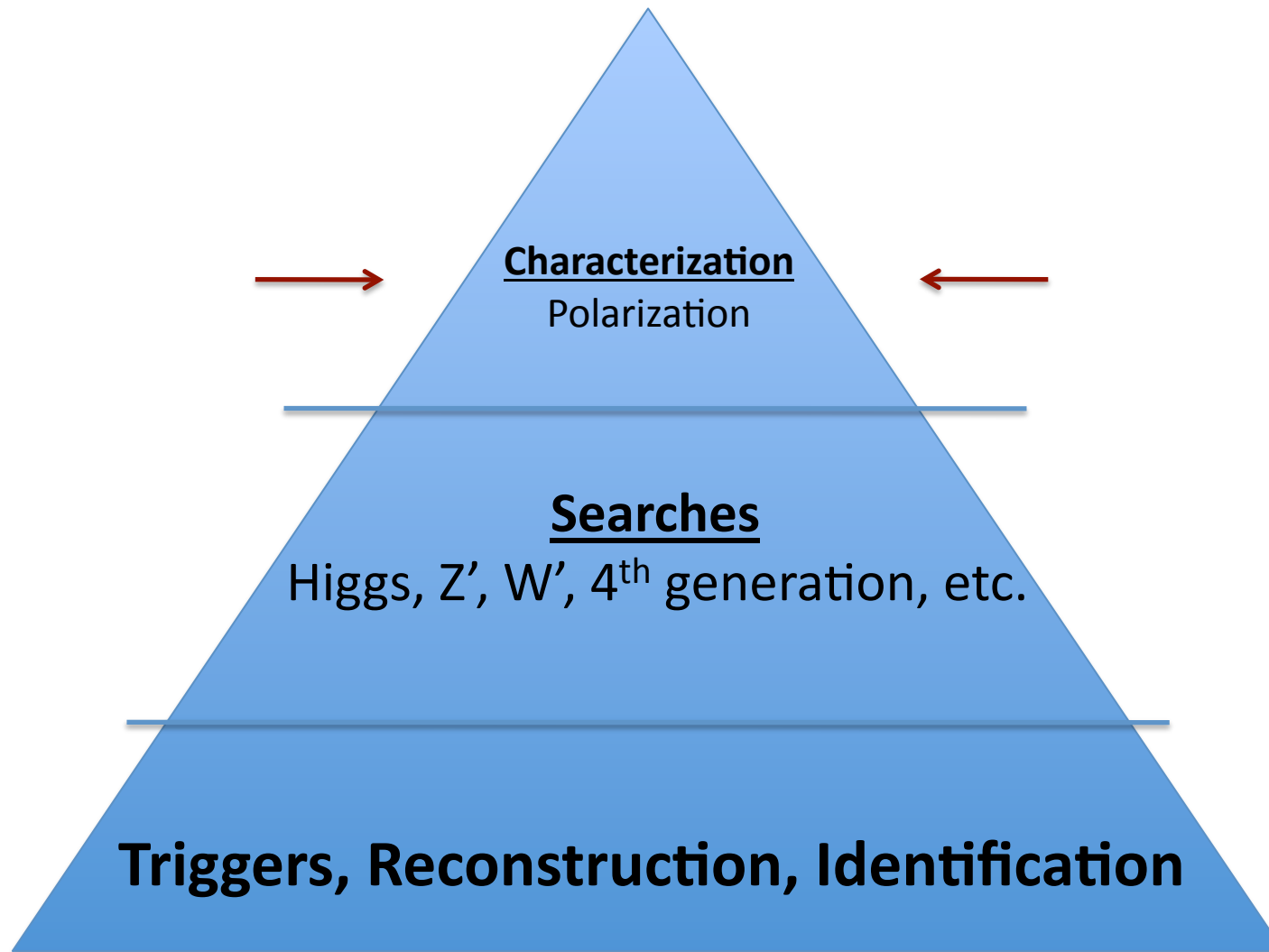
*In $m_h^{\max}$ scenario of MSSM
 $\tan \beta > 30$ –56 is excluded for H^+ masses
 $90 \text{ GeV} < m_{H^+} < 140 \text{ GeV}$*



Theoretical Uncertainty Band

missing one-loop EW corrections
 missing two-loop QCD corrections
 uncertainties related to running b-quark mass

Taus at the LHC



Polarization

$$P_{\tau} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}$$

*Relative cross section of
right- and left-handed taus*

Process	$W \rightarrow \tau_h \nu_{\tau}$	$Z \rightarrow \tau\tau$	$H \rightarrow \tau\tau$	$H^- \rightarrow \tau\nu$
P_{τ}	-1	≈ -0.15	0	+1

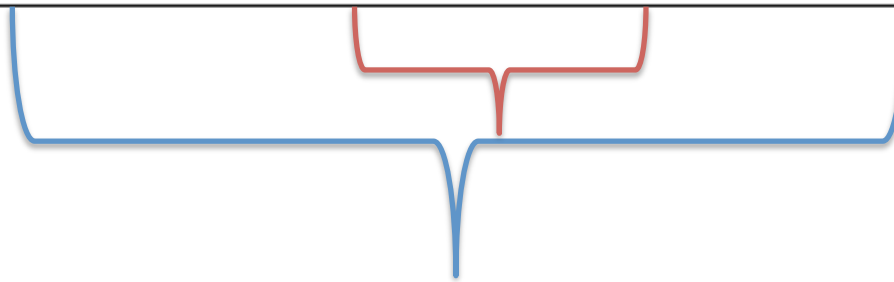
The ability to access tau polarization would give us information regarding the resonance that decays to taus

Polarization

$$P_{\tau} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}$$

*Relative cross section of
right- and left-handed taus*

Process	$W \rightarrow \tau_{\text{h}} \nu_{\tau}$	$Z \rightarrow \tau\tau$	$H \rightarrow \tau\tau$	$H^{-} \rightarrow \tau\nu$
P_{τ}	-1	≈ -0.15	0	+1



The ability to access tau polarization would give us information regarding the resonance that decays to taus

How to access Polarization?

Neutrino has *intrinsic parity*: its angular momentum vector points opposite to its velocity vector, regardless of reference frame.

Neutrinos are left-handed, Anti-neutrinos are right-handed

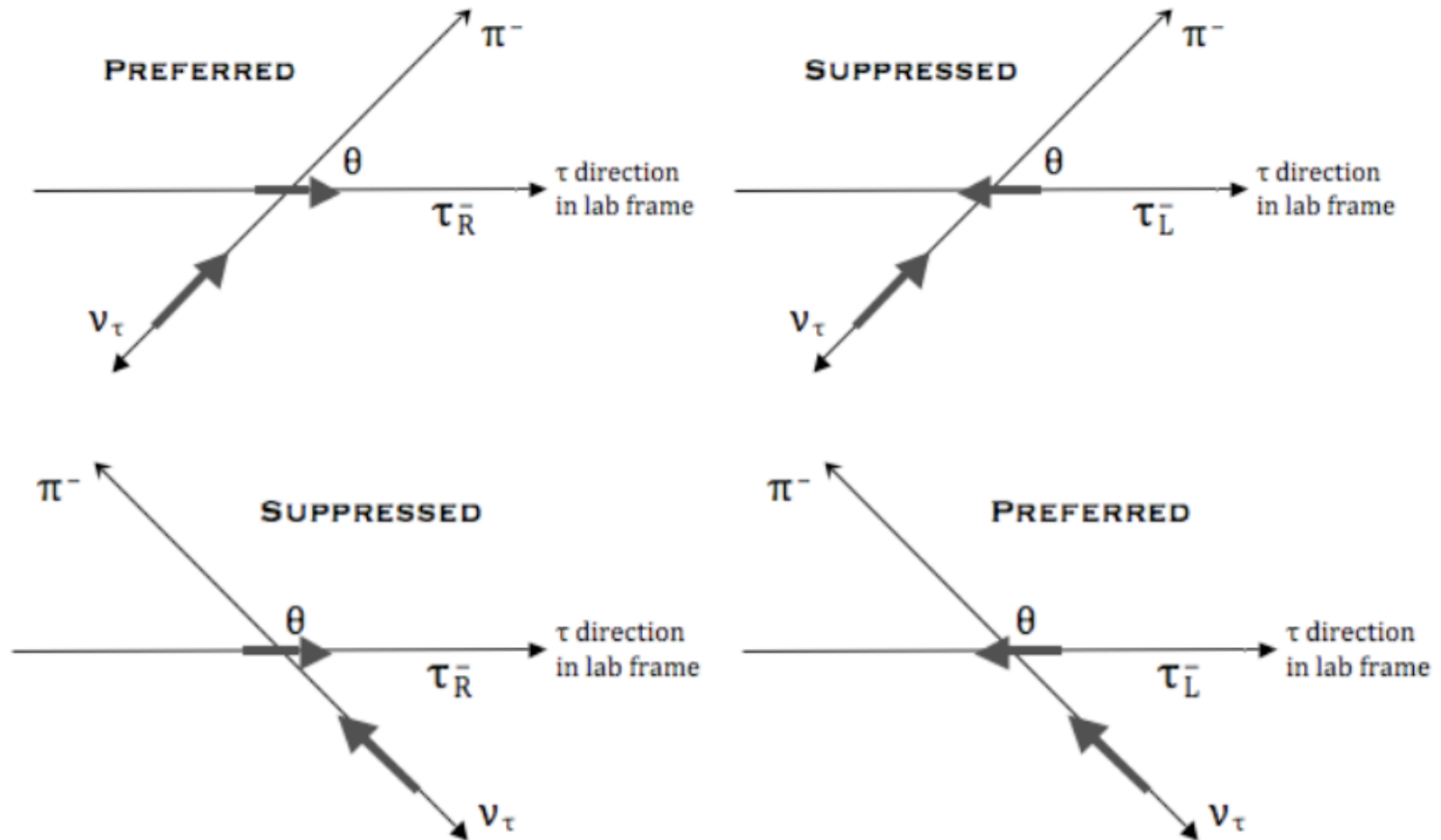
Angular momentum is conserved -> angle between tau decay products depends on the handed-ness of the tau

$$\cos \theta = \frac{2x - 1 - y^2}{1 - y^2}$$

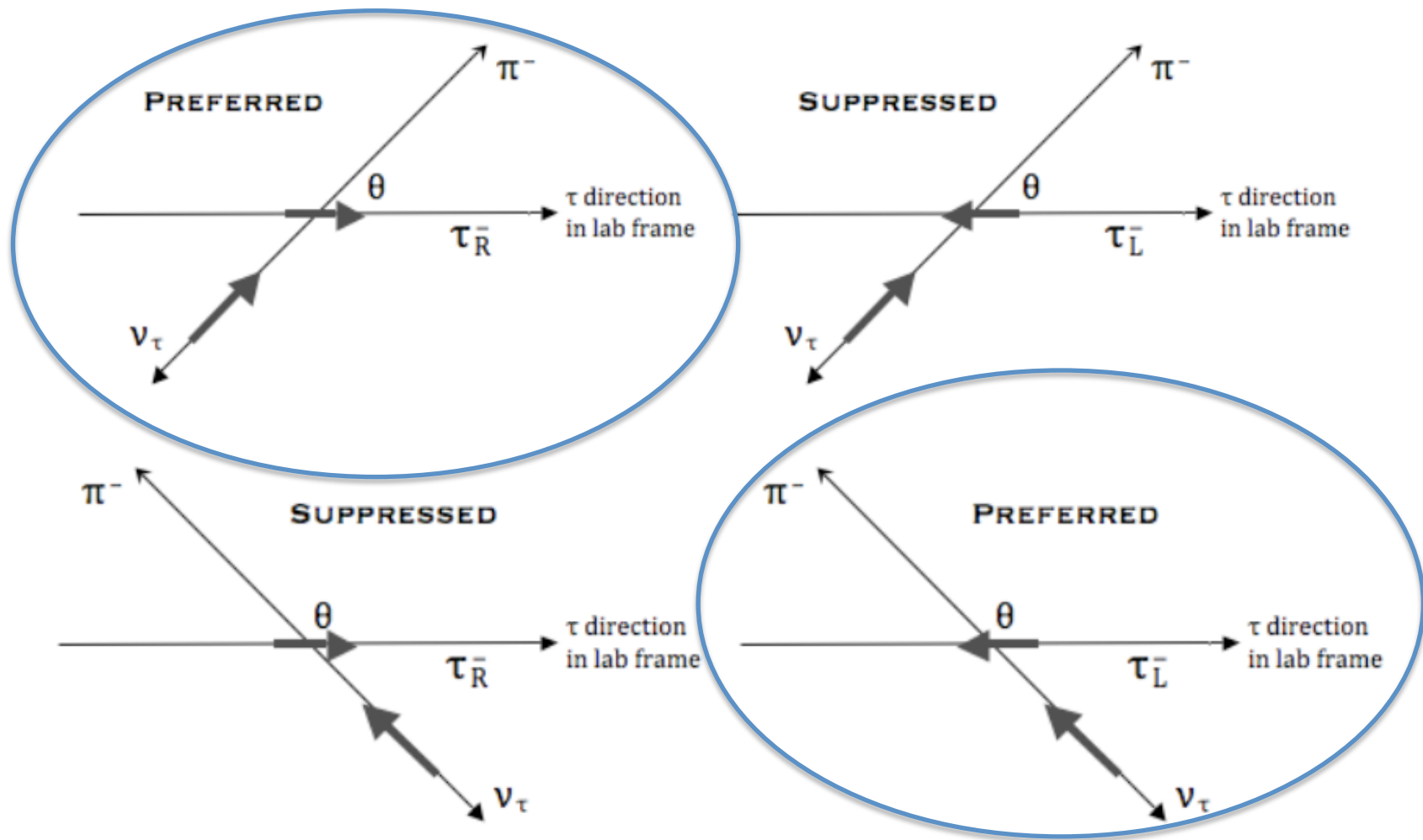
$$x = E_h / E_\tau$$

y refers to the fraction of the tau mass in the hadronic system m_h / m_τ

Polarization: taus decay in detector



Polarization: taus decay in detector



How to access Polarization?

Neutrino has *intrinsic parity*: its angular momentum vector points opposite to its velocity vector, regardless of reference frame.

Neutrinos are left-handed, Anti-neutrinos are right-handed

Angular momentum is conserved -> angle between tau decay products depends on the handed-ness of the tau

$$\cos \theta = \frac{2x - 1 - y^2}{1 - y^2}$$

$$x = E_h / \textcircled{E_\tau}$$

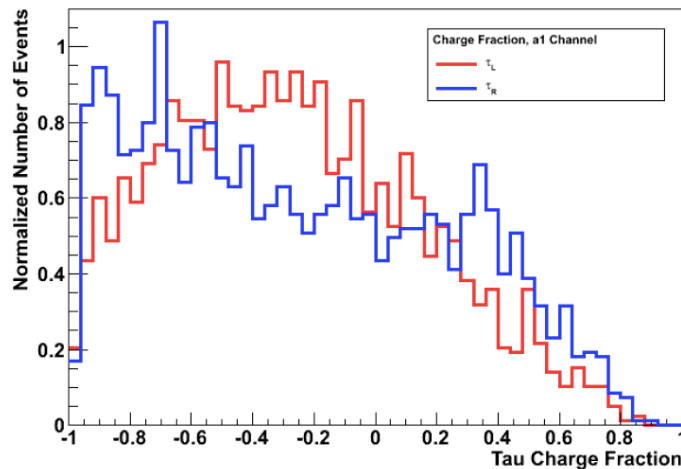
But we don't know the total energy of the tau

y refers to the fraction of the tau mass in the hadronic system m_h/m_τ

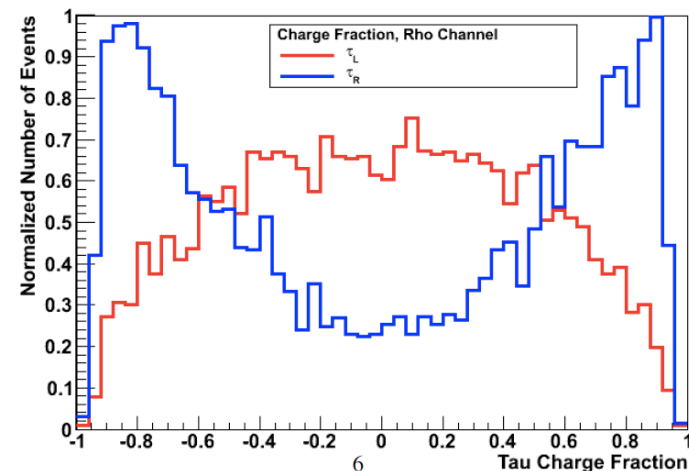
How to access Polarization

$$\frac{E_{\text{T}}^{\pi^-} - E_{\text{T}}^{\pi^0}}{p_{\text{T}}} \approx 2 \frac{p_{\text{T}}^{\text{trk}}}{p_{\text{T}}} - 1 = \gamma.$$

1. $\tau^\pm \rightarrow \pi^\pm \nu$, **the $\pi^\pm$ channel²** (11.6%)
2. $\tau^\pm \rightarrow \rho^\pm \nu \rightarrow \pi^\pm \pi^0 \nu$, **the $\rho^\pm$ channel³** (26.0%)
3. $\tau^\pm \rightarrow a_1^\pm \nu \rightarrow \rho^\pm \pi^0 \nu \rightarrow \pi^\pm 2\pi^0 \nu$, **the $a_1^\pm$ channel** (10.8%)

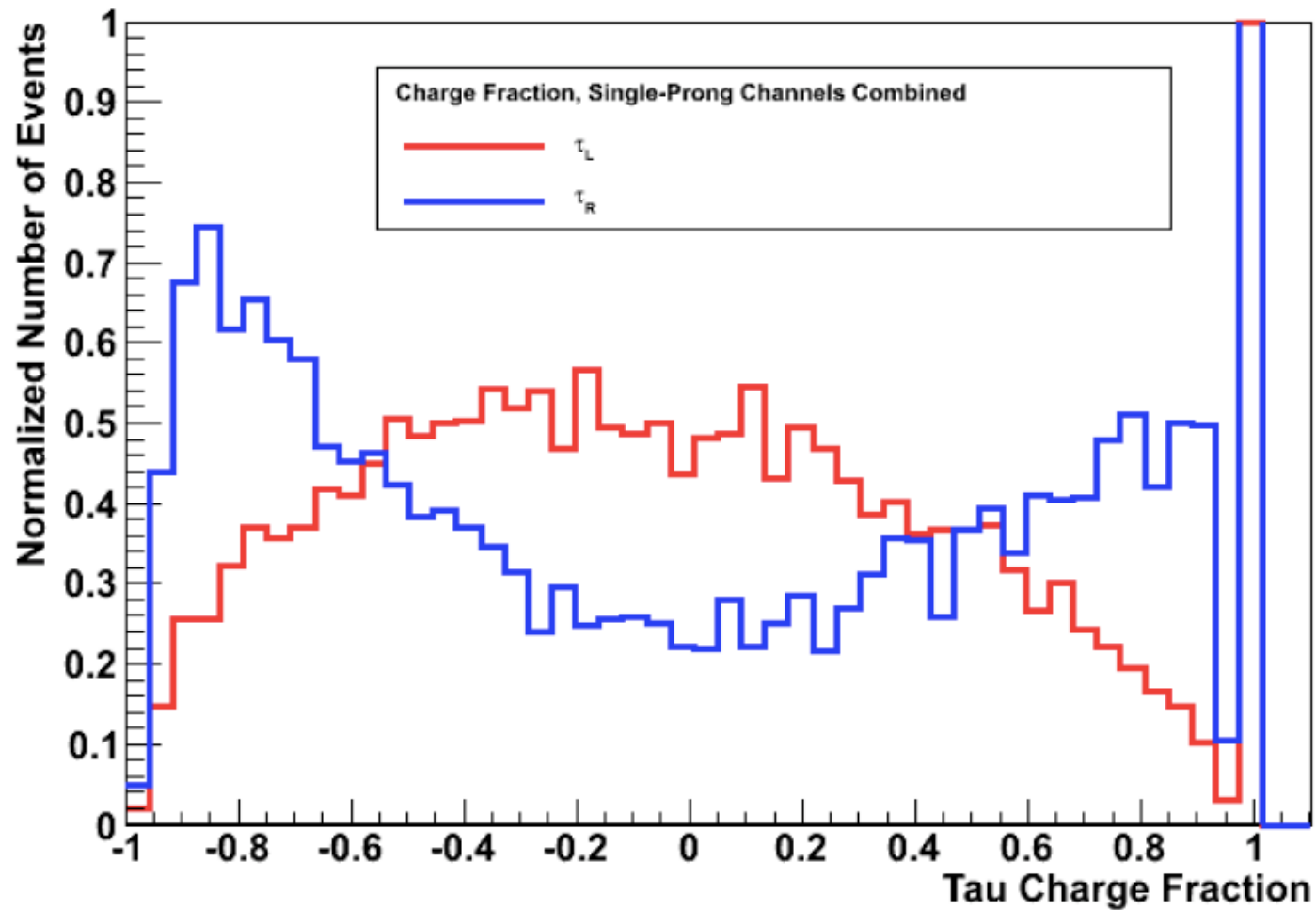


(b) Charge Fraction in the a_1 Channel



(c) Charge Fraction in the Rho Channel

Generator-Level Study



Polarization: Next Steps

This is where I would have shown you results...

Keep your eyes open for Moriond EW results!

Conclusions

Watch for results in
Moriond!



The physics motivation for using taus as probes to discover and/or understand new phenomena at the LHC is compelling.

The tools are in place and rapidly improving to maximize the reach of the tau physics program.

In spite of the promise and the progress, we remain in many ways the few, the proud, and the crazy. We are always looking for company!

The addition of polarization to our arsenal of tools is critical for the future of characterizing new physics at the LHC.

THANKS very much for the opportunity to talk about taus!